



### Wir schaffen Wissen – heute für morgen

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SUMMARY OF THE RECENT POSITRON ANNIHILATION EXPERIMENTS ON STIP SAMPLES



- WHY WE USE POSITRONS TO PROBE SPALLATION SAMPLES
- HOW DO WE DO THAT
- WHAT INFORMATION (ON HELIUM) CAN WE PROVIDE – OVERVIEW OF EXPERIMENTAL DATA
- SUMMARY AND CONCLUSIONS

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- 1) requirement of reliable accelerator-driven systems (ADS) for an effective transmutation of long-lived radioisotopes in the nuclear fuel;
- 2) requirement of affordable neutron sources for the neutron scattering and imaging techniques;
- 3) requirement of irradiation facilities that enable greatly accelerated irradiation studies with fusion-relevant transmutation rates.



To understand **HELIUM embrittlement** means to understand the helium-vacancy interaction at sub-nm scale level

Problem! Spatial resolution limits (and/or sampling of small area) and low sensitivity to helium at most experimental methods.



(100) direction (au ) 15

## WHY POSITRON ANNIHILATION?

### Self-Seeking

(positron diffuse typically ~100nm in metals and seek for sites with higher positron affinity than bulk i.e. it is attracted by certain type of defects!)

### Non destructive

#### Sensitive

Defect type	Sensitivity range (detection limit vs. saturated trapping)		
neutral vacancies	5×10 <sup>21</sup> 10 <sup>25</sup> m <sup>-3</sup>		
dislocations	10 <sup>12</sup> 5×10 <sup>15</sup> m <sup>-2</sup>		

### Macroscopic samples

Information on sub-nm scale features from a large volume (few mm<sup>3</sup>)

### He - sensitive

Helium presence in defects affect positron lifetime and changes the electron momentum distribution

10102 direction

(100) direction (au)

10107



## HOW DO WE DO THAT

### Positron Annihilation Lifetime

Spectroscopy

Two techniques of positron annihilation spectroscopy, based on different physical principles, have been widely established in the material research. Coincidence Doppler Broadening Spectroscopy



After thermalization (~ 3ps), positron diffuse through the lattice until trapping / annihilation. Diffusion time and trapping rate are a function of the microstructure and they can be measured.

Positrons annihilate mainly with the electrons of the outermost shell due to the repulsion of the nucleus. Such annihilation results in  $E\gamma \cong 511$ keV. But the annihilation occurs also with core electrons (electrons with higher momentum). Several factors determine the increase of the positron–electron annihilation probability at certain momenta.





Positron lifetime increases due to trapping at sites with reduced electron density

E0 = 511keV



### Materials: (EM10, CLAM, Eurofer 97, F82H, Optifer, MA956, ODS Eurofer, T91 ...)

Irradiation doses: (5 – 21dpa)

- Irradiation temperatures: (100 600° C)
- Annealing temperatures: (200 800° C)





results were observed in all samples

V. Krsjak et al, Helium behaviour in irradiated in spallation target, J. Nucl. Mater, DOI: 10.1016/j.jnucmat.2014.10.014















Since we have observed positron lifetimes above 400ps (large vacancy clusters) in all materials, there must be **only small amount of helium** in these defects. As we have not seen by TEM vacancy clusters, they must not be much larger than ~ **30vacancies** 





#### Small defects (lifetime ~ 200ps)



Isolated He atom has an extremely low migration energy;

He mobility drastically reduced in the vicinity of defects



Most of He is expected to be accommodated by defects (at 6dpa ~ 450appm He)

**Figure 4** Correlation between positron lifetime and the number of helium atoms in nano-void (B) 1V+nHe, (D) 2V+nHe, (F) 6V+nHe, (H) 12V+nHe.

- Small defects (lifetime ~ 200ps):

Large defects  $\tau_2 \sim 450$  ps Small defects  $\tau_1 \sim 200$  ps

Small defects  $N_1 \sim 12$  vacancies (high He)

DFT simulations + experiments

Large defects  $N_2 \sim 30$  vacancies (low He)

Trapping model

### Three-state trapping model

$$-\frac{dn(t)}{dt} = \sum_{i=1}^{3} \frac{I_i}{\tau_i} \exp\left(-\frac{t}{\tau_1}\right) \qquad I_1 = 1 - I_2 - I_3$$

$$\tau_1 = \frac{1}{\lambda_1} = \frac{1}{\lambda_b + \kappa_1 + \kappa_2} \qquad \qquad \tau_2 = \frac{1}{\lambda_{D1}} \qquad \tau_3 = \frac{1}{\lambda_{D2}}$$

 $\tau_{1,2,3}$  – lifetime components  $I_{1,2,3}$  – components intensities  $\tau_{\rm b}$  = 1/  $\lambda_{\rm b}$  – lifetime in defect-free bulk  $D_1$ ,  $D_2$  (indexes) – defect 1, defect 2  $\kappa_{D1,D2}$  – positron trapping rate at defect (proportional to defect concentration N and positron trapping coefficient of the given defect  $\mu$ ).

$$\kappa_{1} = \frac{I_{2}I_{3}(\lambda_{D1} - \lambda_{D2}) + I_{2}(\lambda_{b} - \lambda_{D1})}{I_{1}} = \mu_{D1}N_{D1}$$

$$\kappa_2 = \frac{I_2 I_3 (\lambda_{D2} - \lambda_{D1}) + I_3 (\lambda_b - \lambda_{D2})}{I_1} = \mu_{D2} N_{D2}$$

Large defects number density  $\sim 10^{22}$ 

Small defects number density  $\sim 10^{24}$ 





Small defect – 12vacancy cluster with 12 helium atoms (He/V = 1) ~  $10^{24}$  m<sup>-3</sup> Large defect – 30 vacancy cluster with 3 helium atoms (He/V = 0.1) ~  $10^{22}$  m<sup>-3</sup>

Material	dpa / t [°C]	He appm	N <sub>1</sub> [m <sup>-3</sup> ]	N <sub>2</sub> [m <sup>-3</sup> ]	He in defect 1 [m <sup>-3</sup> ]	He in defect 2 [m <sup>-3</sup> ]	He total PALS [m <sup>-3</sup> ]	He total Theory [m <sup>-3</sup> ]
Optifer (STIP 2)	6.1 99	450	2.7×10 <sup>24</sup>	2.3×10 <sup>22</sup>	3.3×10 <sup>25</sup>	5.3×10 <sup>22</sup>	3.3E+25	3.8E+25
Optifer (internal source)	6.1 99	450	3.1×10 <sup>24</sup>	1.5×10 <sup>22</sup>	3.8×10 <sup>25</sup>	4.6×10 <sup>22</sup>	3.8×10 <sup>25</sup>	3.8×10 <sup>25</sup>
EM10 (STIP 2)	6.1 93	450	3.4×10 <sup>24</sup>	5.1×10 <sup>21</sup>	4.0×10 <sup>25</sup>	1.2×10 <sup>22</sup>	4.0×10 <sup>25</sup>	3.8×10 <sup>25</sup>
F82H (STIP 2)	6.2 115	465	3.1×10 <sup>24</sup>	1.5×10 <sup>22</sup>	3.7×10 <sup>25</sup>	3.3×10 <sup>22</sup>	3.7×10 <sup>25</sup>	4.0 ×10 <sup>25</sup>
Eurofer (STIP V)	5.8 116	~ 450	2.5×10 <sup>24</sup>	1.8×10 <sup>22</sup>	3.0×10 <sup>25</sup>	5.5×10 <sup>22</sup>	3.0×10 <sup>25</sup>	3.8×10 <sup>25</sup>

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Lets take low dpa sample and lets assume that all helium is in small vacancy clusters:

- We know the helium concentration (400 500 appm for low-dose STIP sample)
- We know the vacancy clusters number density



What happens to the actual He/V ratio at different temperatures?















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- 10 ferritic/martensitic steels and more than 100 samples have been investigated by means of positron annihilation spectroscopy techniques in the last 2 years. Effect of dpa, irradiation and annealing temperature was studied
- These experiments offer qualitative and quantitative information on sub-nm scale helium-vacancy clusters, which cannot be obtained by any other experimental technique
- Helium peak was clearly identified in the CDBS spectra of annealed spallation samples, which demonstrates its presence in major positron trapping sites (vacancy clusters)
- Small vacancy clusters with the mean size of 12 vacancies contains the majority of helium. He/V ratio was estimated ~1 in most as-irradiated samples and increase with annealing temperature to ~1.5.
- Large vacancy clusters (> 0.8nm), identified in all samples, contains only small amount of helium. Their size increase with irradiation/annealing via absorption of vacancies – increase of positron lifetime and CDBS S-parameter. We assume that these defects evolve eventually to helium bubbles, which can be seen by TEM at elevated temperatures/doses.



# THANK YOU FOR YOUR ATTENTION