

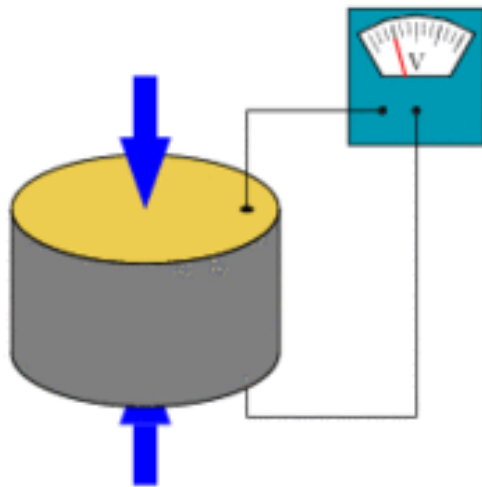
PSI Master School 2017

Introducing photons, neutrons and
muons for materials characterization

Lecture 6: Neutron Cross Sections
with Matter

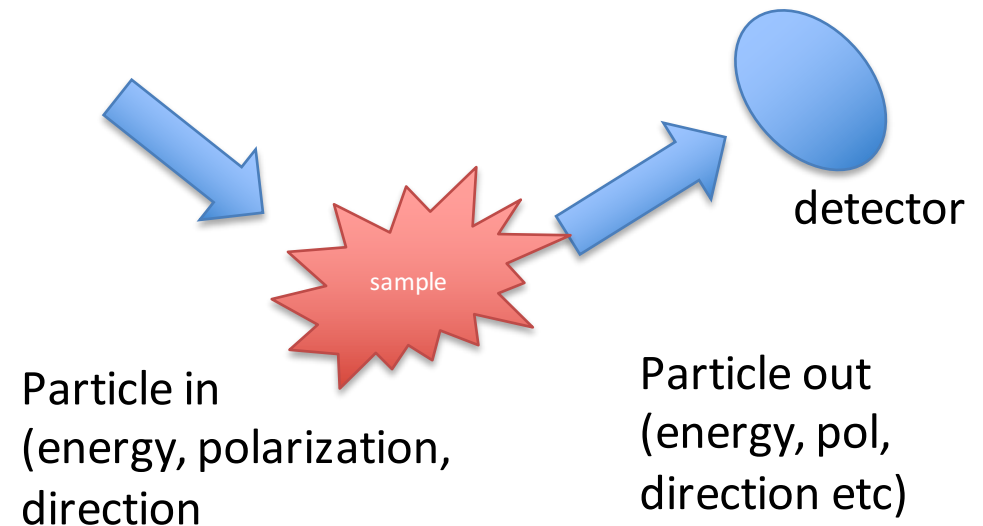
Sensor and particle probe measurements

Sensor measurement



Sensor experiments:
Sensor is sensitive to physical property and converts the signal to something observable

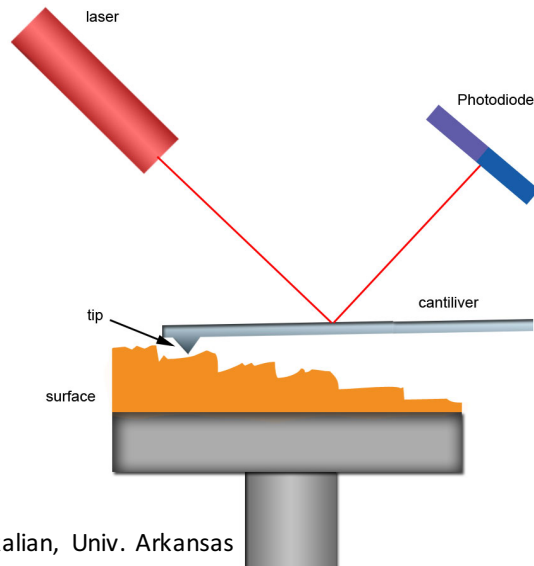
Particle probes for measurements



Particle probe experiments:
Particles injected and ejected in sample probe the materials properties of the material

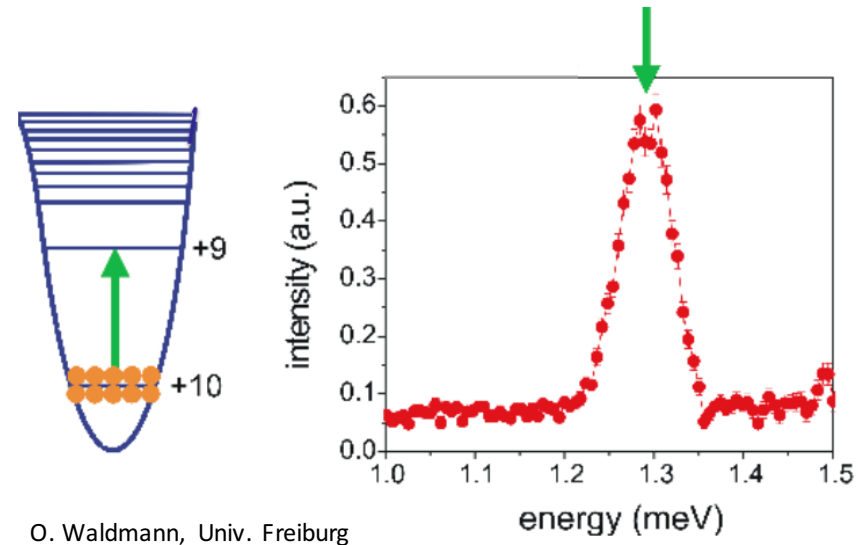
Examples: Sensor vs particle probe measurements

Sensor measurement



Example: Atomic force microscopy
AFM acts as sensor, whose position
can be measured with a laser

Particle probes for measurements



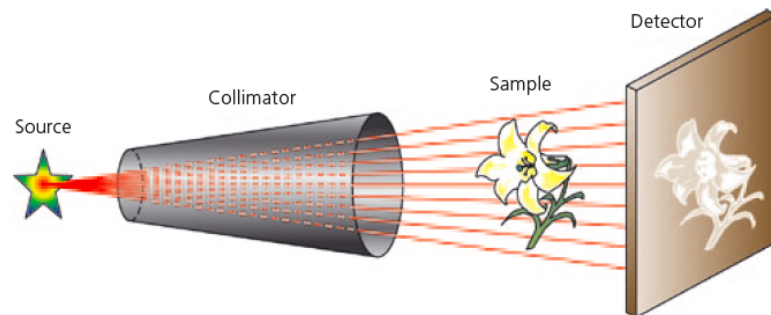
Example: Inelastic neutron scattering
Neutrons transfer momentum and
energy to excitations in the material,
these changes can be detected by
the scattered neutron

Types of Photon/Neutron Experiments

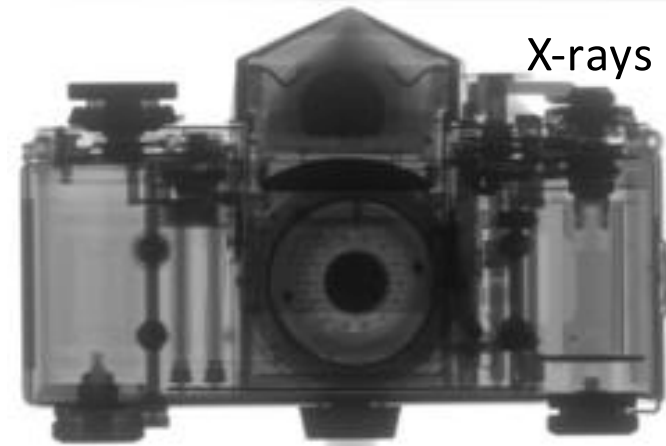
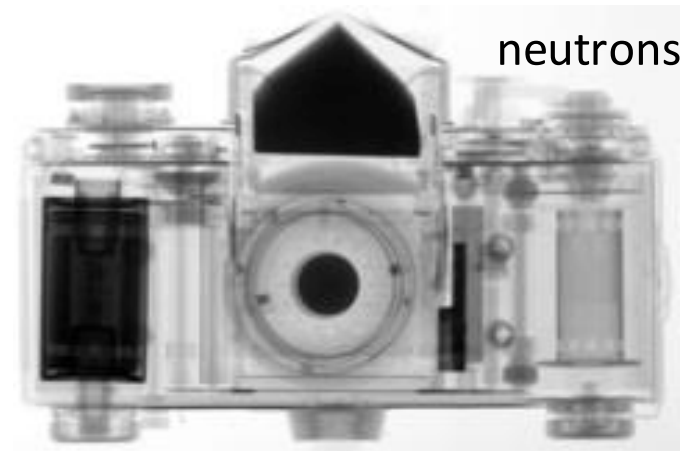
- Absorption: absorption of particles in transmission
 - Neutron and X-ray Imaging
 - X-ray absorption spectroscopy
- Scattering: radiation (light, particles) change direction is used to probe materials
 - Coherent scattering experiment
 - Reflection
- Emission experiments (Photon in-photon out, Neutron in-Photon out; Photon in-Electron out)
 - Neutron prompt gamma experiments
 - ARPES (angular resolved photoemission)

Photons vs neutrons

Example of neutron imaging

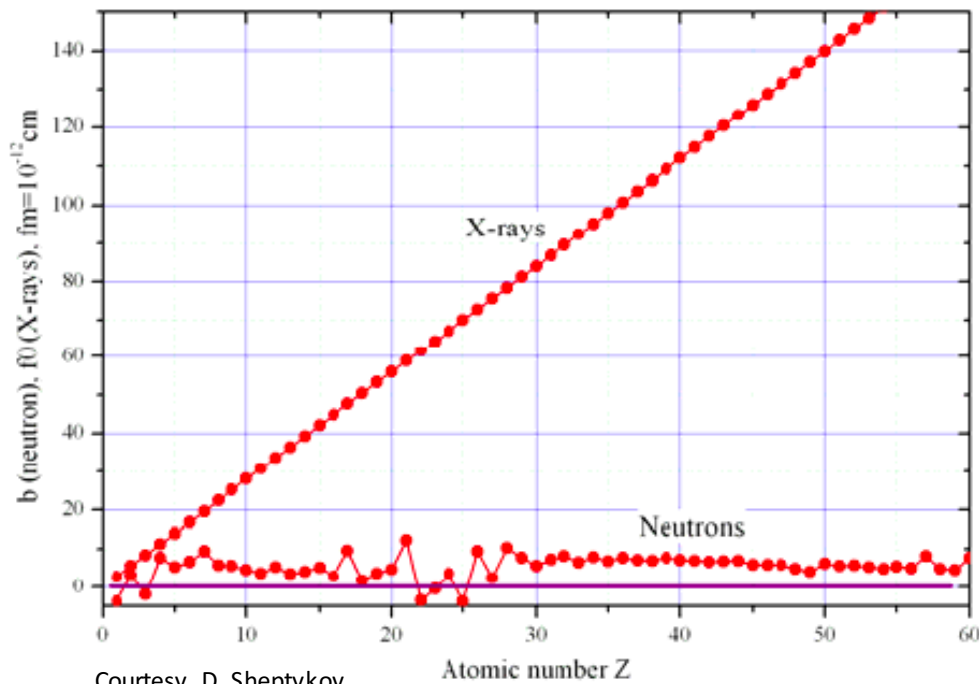


X-rays and neutrons have different cross-sections with matter

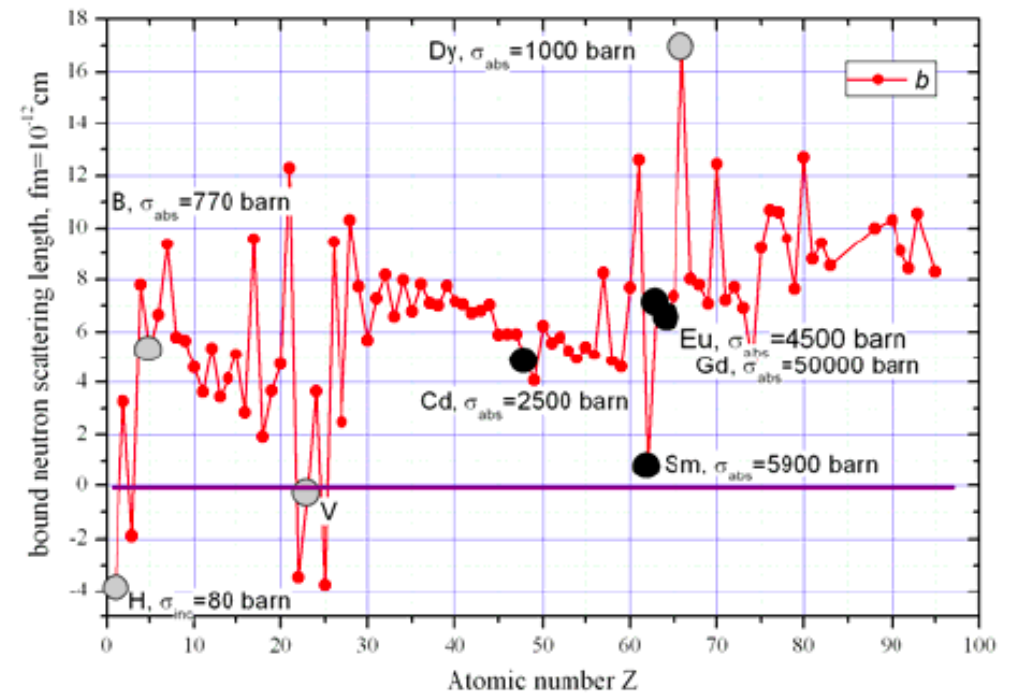


Cross-sections photons and neutrons

*X-ray scattering amplitudes
and neutron scattering lengths:*

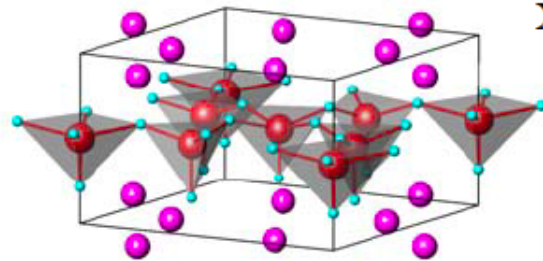


*(Coherent, elastic) bound neutron scattering lengths
for the natural occurrences of the elements:*



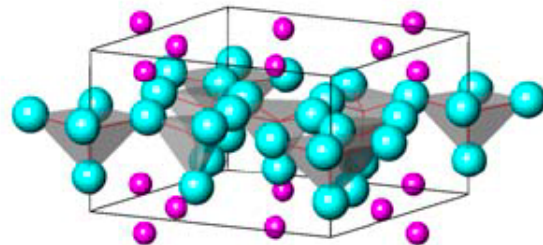
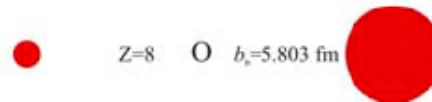
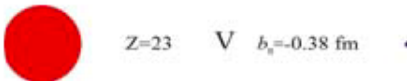
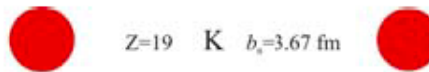
X-rays interact with electron shell, nuclear neutron scattering interacts with the nuclei

Example: X-ray vs. neutron diffraction

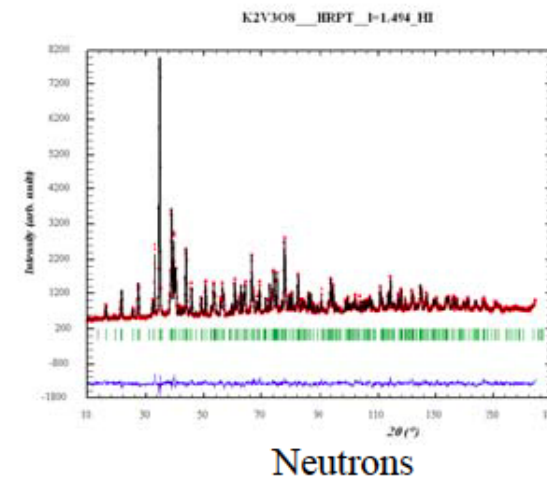
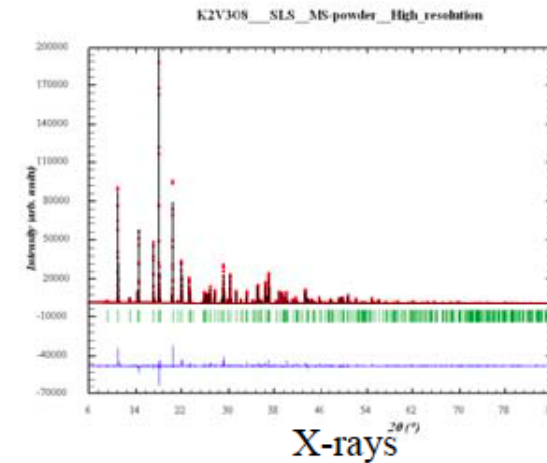
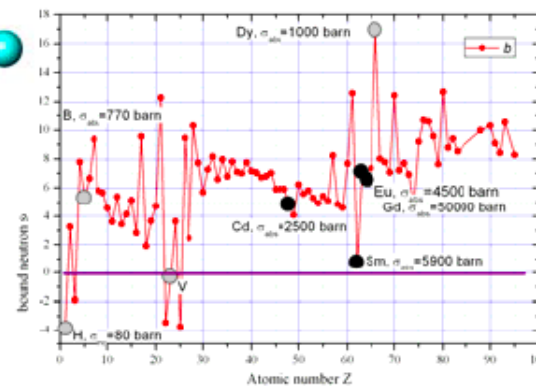


X-rays

X-rays Scattering amplitudes Neutrons



Neutrons



Courtesy D. Sheptykov

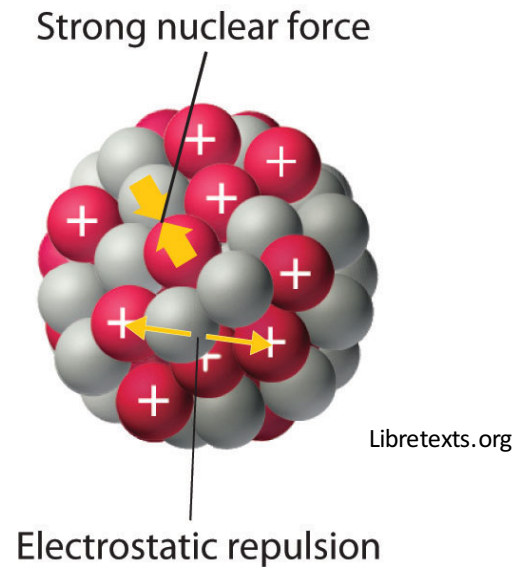
Neutron interactions with matter

- **Strong nuclear interaction**

length scale 10fm

→ scattering

→ absorption



- **Magnetic interaction**

power law decay

Write potentials for two interactions

$$V(\mathbf{R}) = -\boldsymbol{\mu}_n \cdot \mathbf{B}(\mathbf{R})$$

Nuclear scattering length

- Origin of nuclear scattering
 - Strong interaction of neutron with nuclei is complex
 - Wave-length of thermal neutrons is $\sim 0.1\text{nm} \gg 10^{-6}\text{nm}$ of nuclear force
 - Nuclear interactions can be approximated by point like potential

$$V(\mathbf{r}) = \frac{2\pi\hbar^2}{m_N} b\delta(\mathbf{r})$$

- b is nuclear scattering length (units of a length)
- b depends on nuclei and isotope
- Cross section (units of barn or cm^2)

$$\sigma_{\text{tot}} = 4\pi b^2$$

Nuclear spin scattering

- Some nuclei possess nuclear spin
- This lead to additional cross-section, and a change of the scattering length

$$b = b_c + \frac{2b_i}{\sqrt{I(I+1)}} \mathbf{s} \cdot \mathbf{I}$$

- Very anisotropic scattering length
- Nuclear spin is not ordered, so in average scattering is isotropic
- Because scattering length can be different on same atomic position, this can lead to incoherent scattering

Nuclear length from different isotopes

- Atoms are stable with different number of neutrons
→ isotopes
- Interactions between neutron and nuclei depends on number of neutrons
- Neutron scattering length depends on isotope

Example: H

Isotope	conc	Coh b	Inc b
H	---	-3.7390	---
1H	99.985	-3.7406	25.274
2H	0.015	6.671	4.04
3H	(12.32 a)	4.792	-1.04

Example: He

Isotope	conc	Coh b	Inc b
He	---	3.26(3)	---
3He	0.00014	5.74-1.483i	-2.5+2.568i
4He	99.99986	3.26	0

Example: Isotope with different nuclear spin

Case: nucleus with one isotope with nuclear spin I

Two states: $I+1/2, I-1/2$

Examples: hydrogen $I=1/2$: $b^+=10.85$ fm, $b^-=-47.50$ fm
 deuterium $I=1$: $b^+=9.53$ fm, $b^-=0.98$ fm

How to calculate the incoherent scattering?

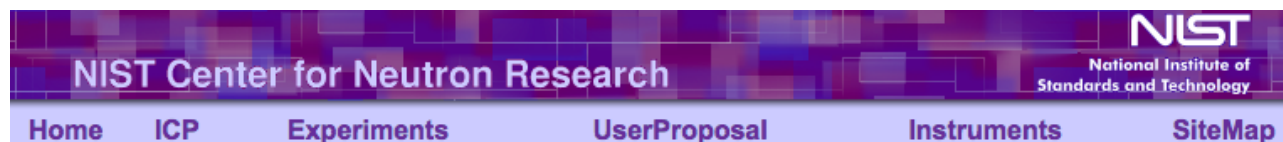
The number of states are:

$$\begin{array}{ll} \text{For } I+1/2 : & 2I+2 \\ \text{For } I-1/2 : & 2I \end{array} \quad \begin{array}{l} f^+ = \frac{2I+2}{4I+2} \\ f^- = \frac{2I}{4I+2} \end{array}$$

Total states: $4I+2$

$$\bar{b} = \frac{1}{2I+1} ((I+1)b^+ + Ib^-)$$

Neutron scattering length and cross sections



Neutron scattering lengths and cross sections

A periodic table of elements where each element's box contains its symbol. The table is color-coded: light blue for H, He, Li, Be, B, C, N, O, F, Ne, Na, Mg, Al, Si, P, S, Cl, Ar, K, Ca, Sc, Ti, V, Cr, Mn, Fe, Co, Ni, Cu, Zn, Ga, Ge, As, Se, Br, Kr, Rb, Sr, Y, Zr, Nb, Mo, Tc, Ru, Rh, Pd, Ag, Cd, In, Sn, Sb, Te, I, Xe, Cs, Ba, La, Hf, Ta, W, Re, Os, Ir, Pt, Au, Hg, Tl, Pb, Bi, Po, At, Rn, Fr, Ra, Ac, and dark blue for the remaining elements (Ce, Pr, Nd, Pm, Sm, Eu, Gd, Tb, Dy, Ho, Er, Tm, Yb, Lu, Th, Pa, U, Np, Pu, Am, Cm, Bk, Cf, Es, Fm, Md, No, Lr).

Example Cu

Neutron scattering lengths and cross sections							
Isotope	conc	Coh b	Inc b	Coh xs	Inc xs	Scatt xs	Abs xs
Cu	---	7.718	---	7.485	0.55	8.03	3.78
63Cu	69.17	6.43	0.22	5.2	0.006	5.2	4.5
65Cu	30.83	10.61	1.79	14.1	0.4	14.5	2.17

<https://www.ncnr.nist.gov/resources/n-lengths/>

Scattering length (fm units)

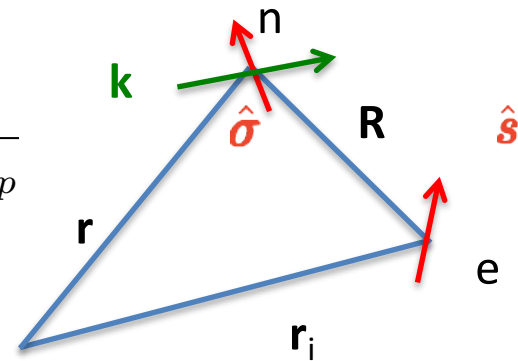
Cross sections (barn= 10^{-24}cm^2 units)

Magnetic neutron scattering

$$\mu_n = 2\gamma\mu_n \frac{\hat{\sigma}}{2}$$

$$\mu_N = \frac{e\hbar}{2m_p}$$

$$\mu_e = -2\mu_B \hat{s}$$



Magnetic field from electron:

$$\mathbf{B}(\mathbf{R}) = \nabla \times \left(\frac{\mu_0}{4\pi} \frac{\boldsymbol{\mu}_e \times \hat{\mathbf{R}}}{R^2} \right)$$

Neutron-electron interaction:

$$V(\mathbf{R}) = -\boldsymbol{\mu}_n \cdot \mathbf{B}(\mathbf{R})$$

Absorption cross-section

- Neutrons can be absorbed by nuclei (neutron capture)
- Very different absorption length than X-rays (for 100keV)

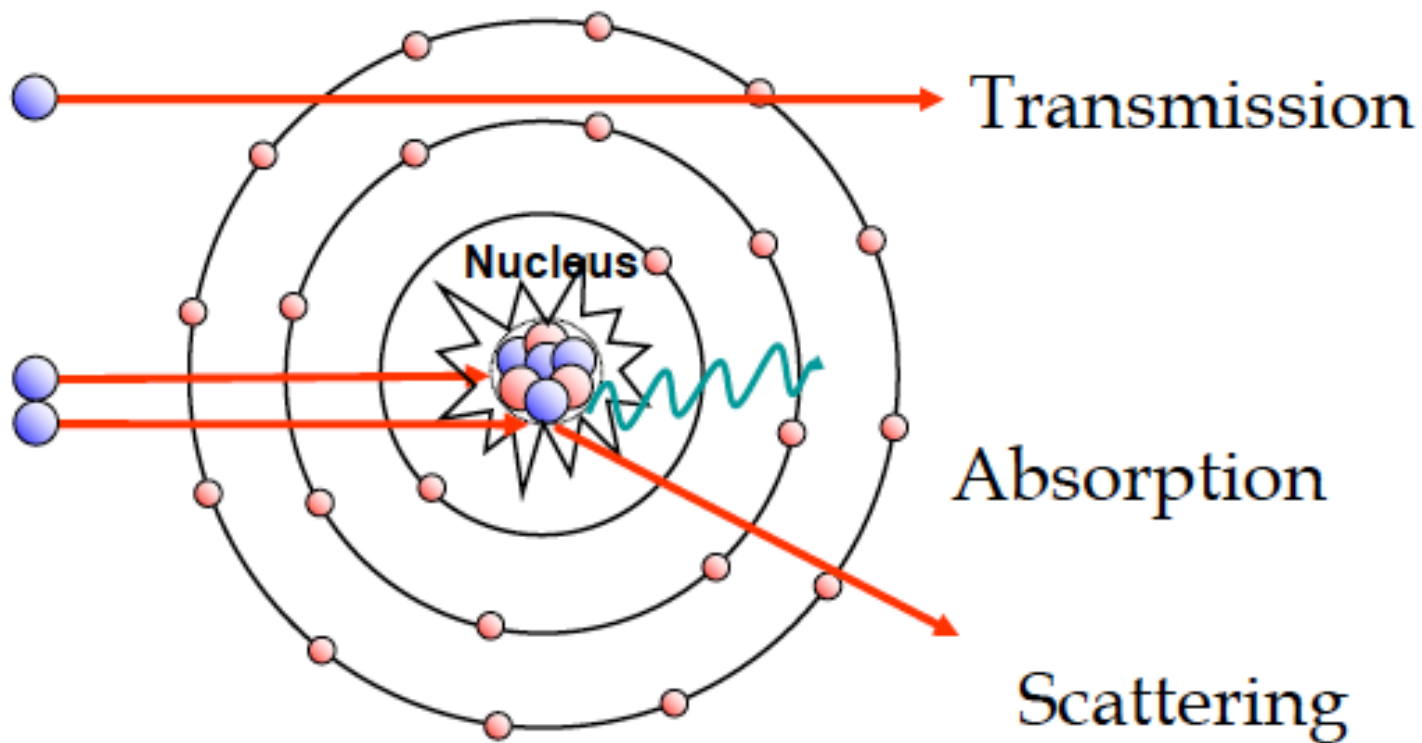
Group →	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18
↓ Period																		
1	H 1.01																	He 4.00
2	Li 6.94	Be 9.01											B 10.81	C 12.01	N 14.01	O 16.00	F 18.99	Ne 20.18
3	Na 22.99	Mg 24.31											Al 26.98	Si 28.09	P 30.97	S 32.07	Cl 35.45	Ar 39.95
4	K 39.10	Ca 40.08	Sc 44.96	Ti 47.88	V 50.94	Cr 52.00	Mn 54.94	Fe 55.85	Co 58.93	Ni 58.69	Cu 63.55	Zn 65.38	Ga 69.72	Ge 72.64	As 74.92	Se 78.96	Br 79.90	Kr 83.80
5	Rb 85.47	Sr 87.62	Y 88.91	Zr 91.22	Nb 92.91	Mo 95.94	Tc 98.91	Ru 101.07	Rh 102.91	Pd 106.42	Ag 107.87	Cd 112.41	In 114.82	Sn 118.71	Sb 121.76	Te 127.60	I 126.91	Xe 131.29
6	Cs 132.91	Ba 137.33		Hf 178.49	Ta 180.95	W 183.84	Re 186.21	Os 190.23	Ir 192.22	Pt 195.08	Au 196.97	Hg 200.59	Tl 204.38	Pb 207.2	Bi 208.98	Po -	At -	Rn -
7	Fr -	Ra 226		Rf -	Db -	Sg -	Bh -	Hs -	Mt -	Ds -	Rg -	Uub -	Uut -	Uuq -	Uup -	Uuh -	Uus -	Uuo -

X-rays 100keV)

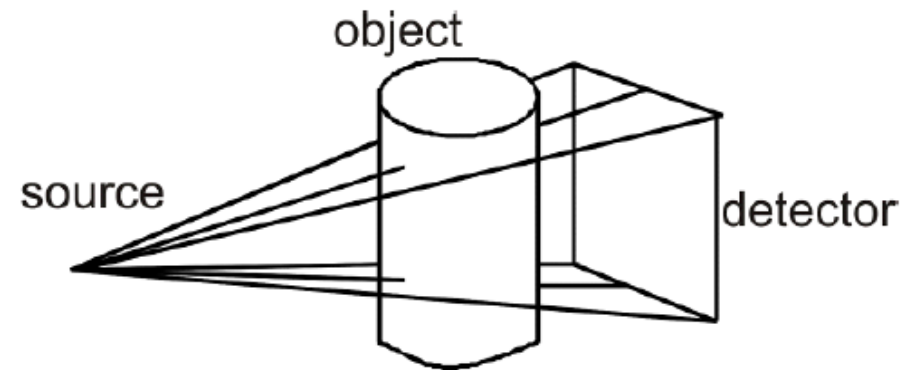
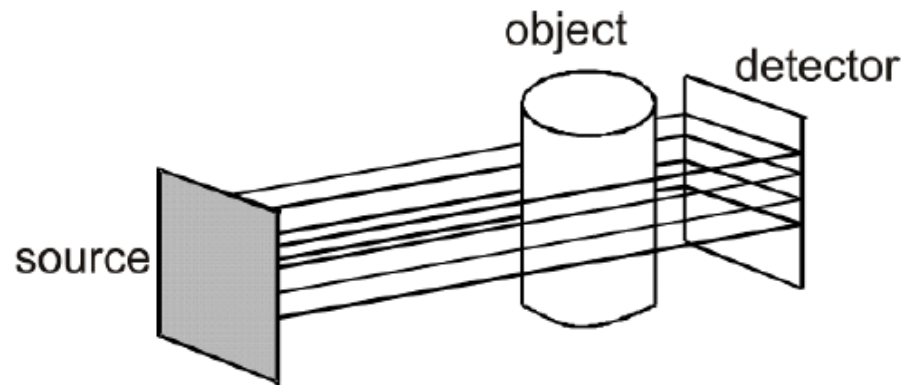
La 0.52	Ce 0.14	Pr 0.41	Nd 1.97	Pm 5.72	Sm 171.47	Eu 94.58	Gd 1479.0	Tb 0.92	Dy 32.42	Ho 2.25	Er 5.49	Tm 3.59	Yb 1.40	Lu 2.75
Ac -	Th 0.59	Pa 8.49	U 0.82	Np 9.80	Pu 60.29	Am 2.89	Cm -	Bk -	Cf -	Es -	Fm -	Md -	No -	Lr -

Neutrons and matter

- Very weak interaction
- Many neutrons are transmitted



Neutron Imaging (neutron radiography)

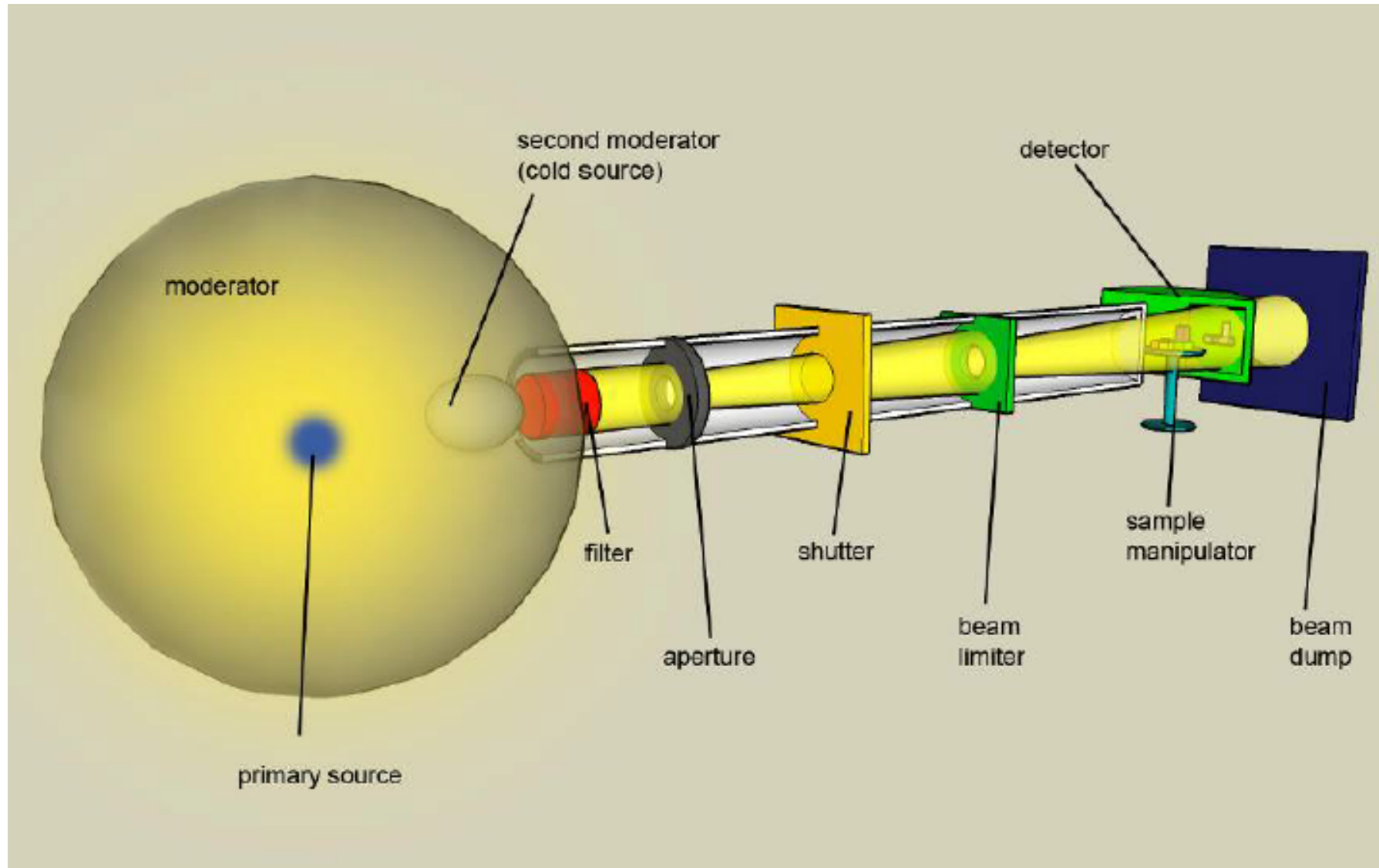


- Objects or materials can be imaged quantitatively by using neutron transmission pictures
- Intensity decreases exponentially with thickness d of the sample

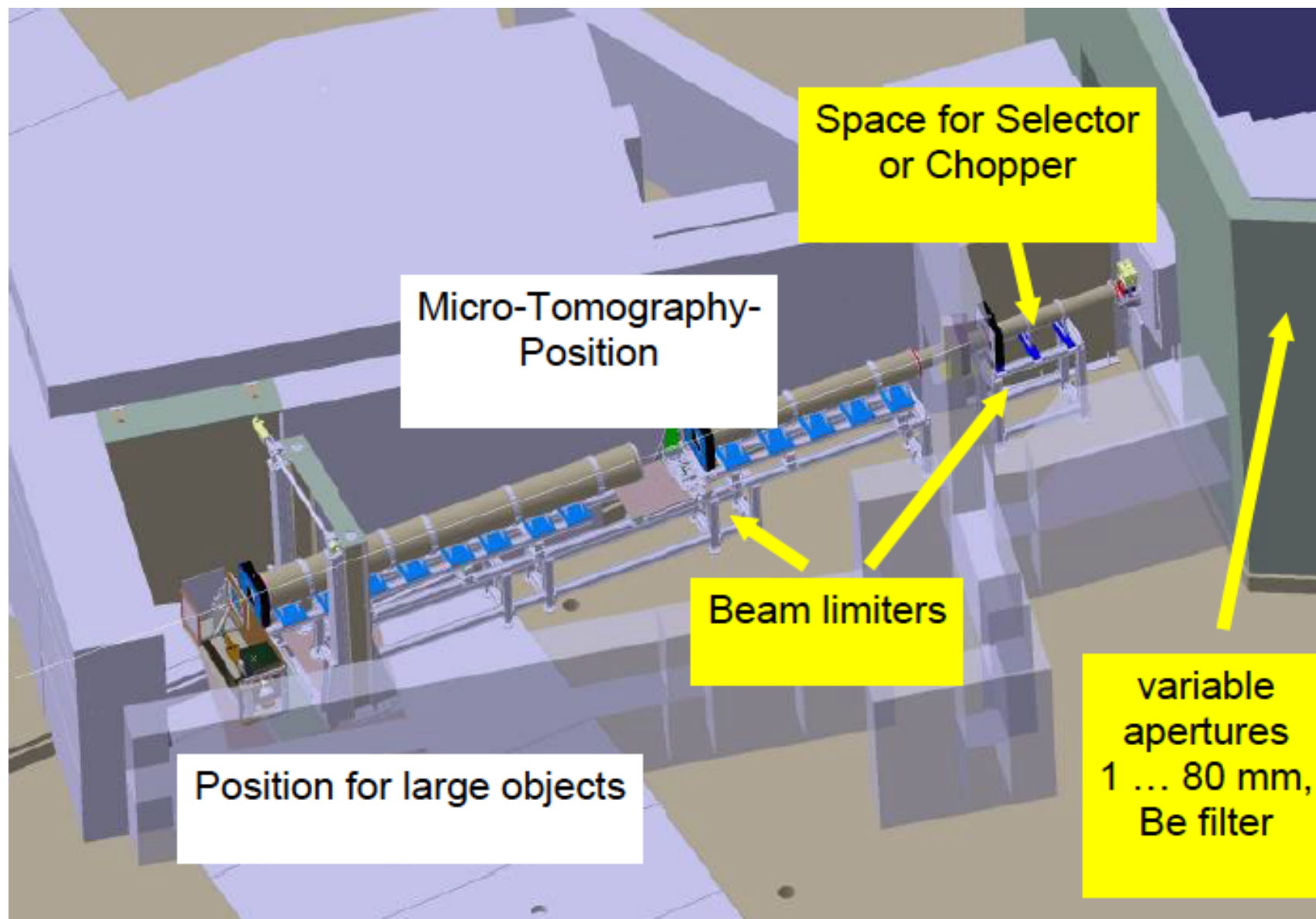
$$I(x,y,E) = I_0(x,y,E) \cdot e^{-\Sigma(x,y,E) \cdot d(x,y)}$$

$\Sigma(x,y,E)$ is the effective attenuation coefficient at x,y and for E

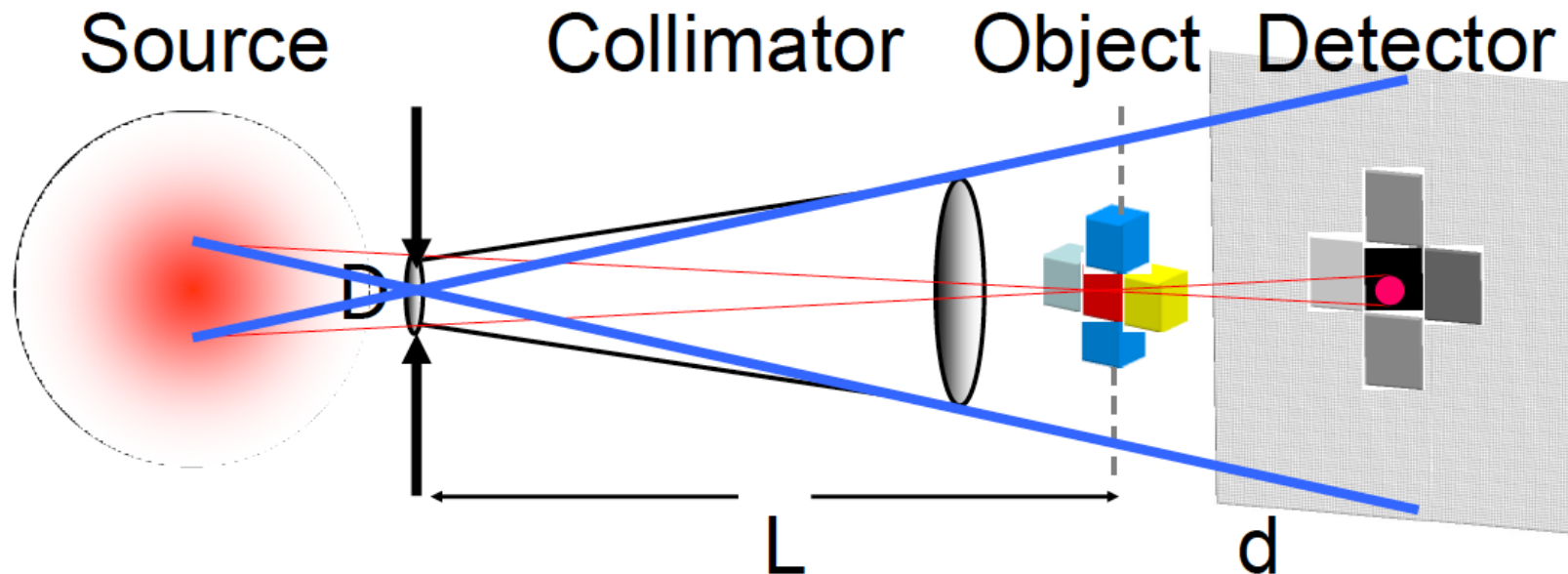
Neutron imaging instrument



ICON neutron imaging beamline



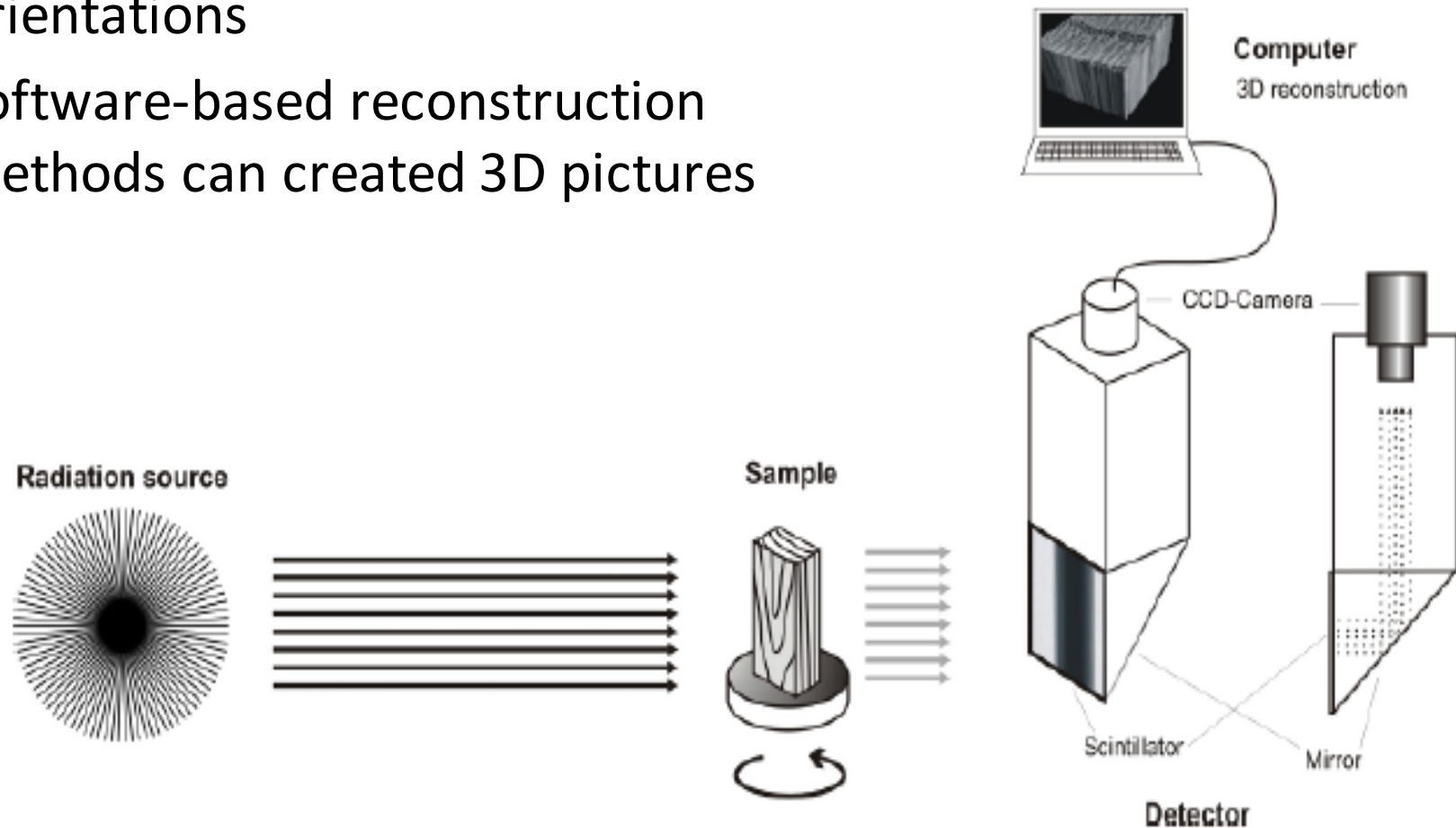
Imaging: spatial resolution



- Collimation ratio $C = L/D$
L: collimator length
D: primary aperture opening
- Geometric blurring $u = d/C$

Tomography Imaging

- neutron radiography in different orientations
- Software-based reconstruction methods can create 3D pictures



Example: neutron imaging

Object



Neutron
radiography



Neutron
tomography



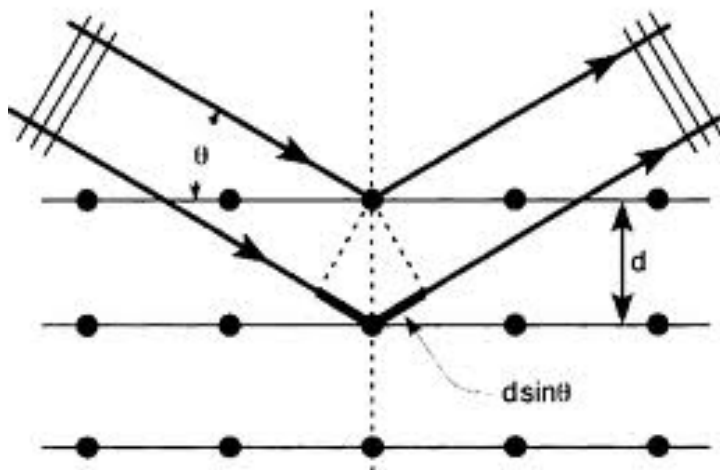
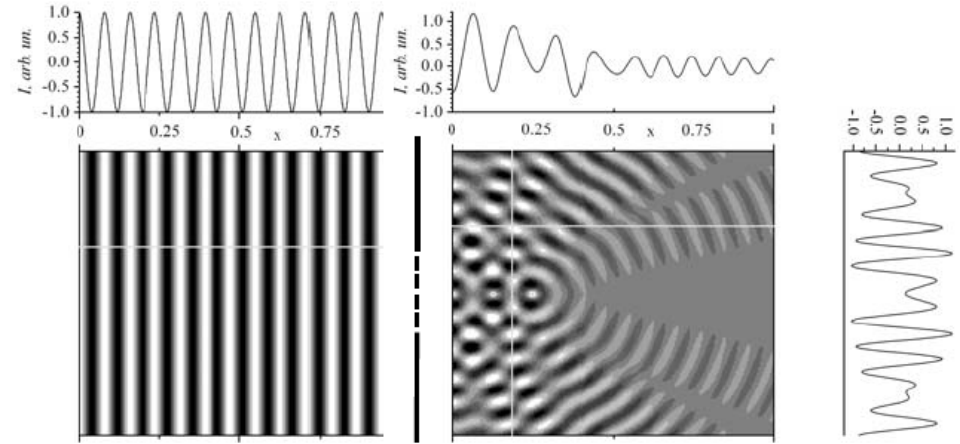
Principle of scattering

Interference of waves at objects

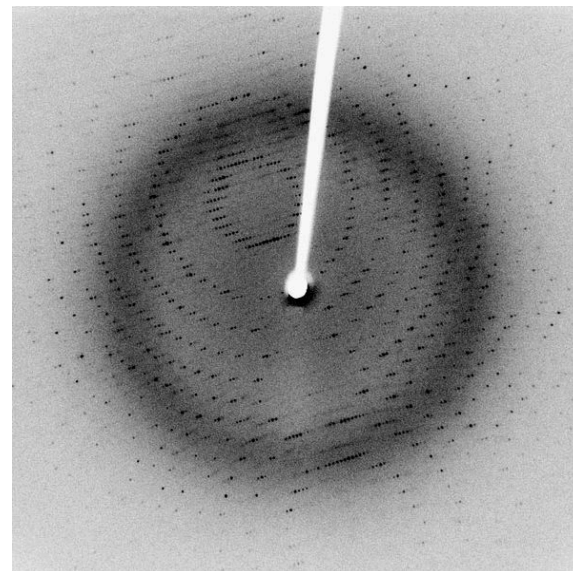
Constructive and destructive interference at two slit object

Bragg's law in crystals

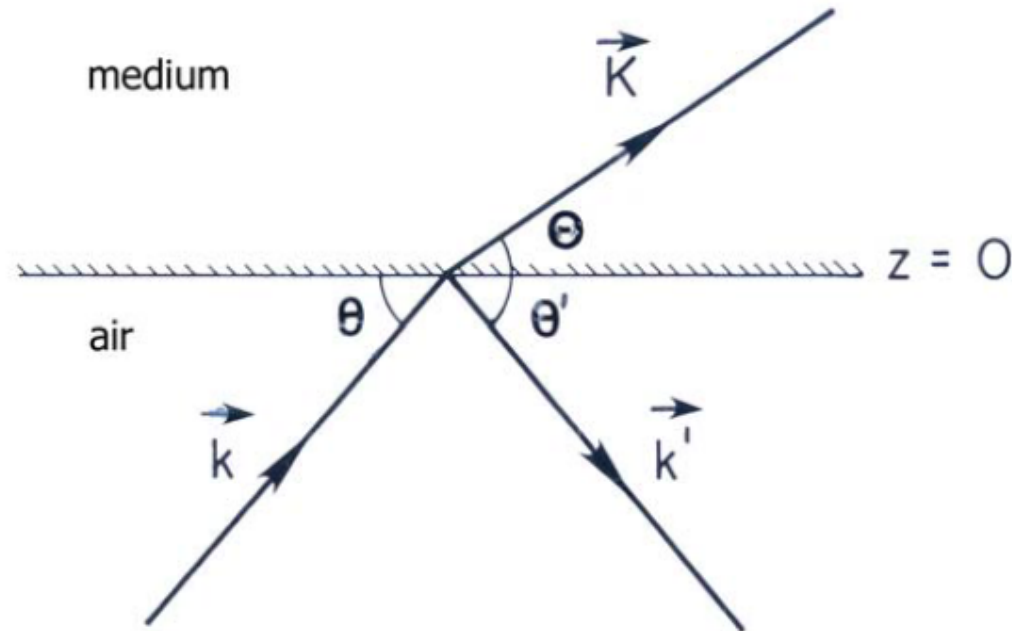
→ Constructive interference at reciprocal wave-vectors



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Wave-function and matter



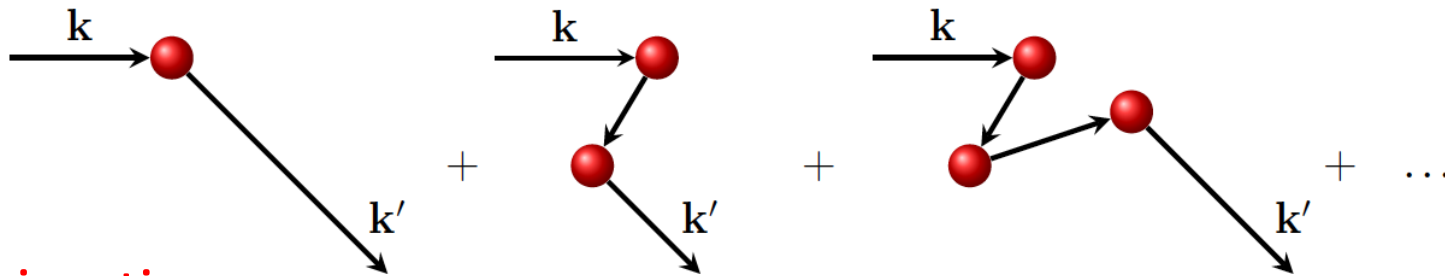
- Wave-function describing neutron particles can change when entering a material
- This can affect the overall interaction of neutron with materials

→ Refraction

→ Total reflection

Kinematic or Born approximation

- Interaction between neutron and matter is very weak
- Often only one scattering even in the material
- the incident wave is in many cases only very weakly affected by the material



Approximation:

- Only one scattering event
- Wave-function in vacuum and in material are the same

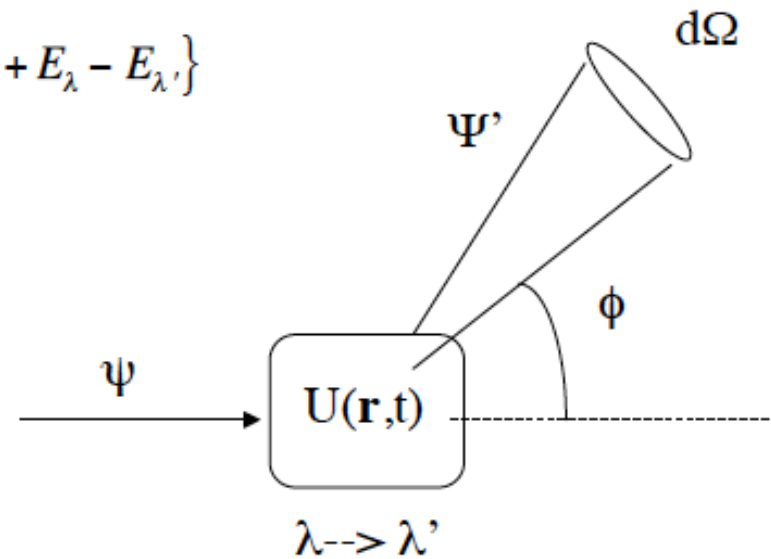
Fermi Golden Rule

- Single scattering event
- Incident wave k , interaction potential U , scattered wave k'

$$\frac{d^2\sigma}{d\Omega d\omega} = \frac{k'}{k} \left(\frac{m}{2\pi\hbar^2} \right)^2 \sum_{\lambda} p_{\lambda} \sum_{\lambda'} \left| \langle \mathbf{k}'\lambda' | \hat{U} | \mathbf{k}\lambda \rangle \right|^2 \delta\{\hbar\omega + E_{\lambda} - E_{\lambda'}\}$$

$$\hat{U}(\mathbf{r}) = \frac{2\pi\hbar^2}{m} \sum_j b_j \delta\{\mathbf{r} - \mathbf{R}_j\}$$

$$b_i \approx 10^{-12} \text{ cm} \rightarrow \frac{d^2\sigma}{d\Omega d\omega} \approx 10^{-24} \text{ cm}^2 = 1 \text{ barn}$$



$$\mathbf{Q} = \mathbf{k}_f - \mathbf{k}_i$$

Calculation of correlation function

$$\frac{d^2\sigma}{d\Omega d\omega} = \frac{k'}{k} \frac{1}{2\pi\hbar} \sum_{\lambda} p_{\lambda} \sum_{j,j'} b_j b_{j'} \int_{-\infty}^{\infty} \langle \lambda | e^{-i\hat{\mathbf{Q}}\hat{\mathbf{R}}_{j'}(0)} e^{i\hat{\mathbf{Q}}\hat{\mathbf{R}}_j(t)} | \lambda \rangle e^{-i\omega t} dt$$

$$\langle \hat{A} \rangle = \sum_{\lambda} p_{\lambda} \langle \lambda | \hat{A} | \lambda \rangle$$

$$\mathbf{Q} = \mathbf{k}_f - \mathbf{k}_i$$

$$\frac{d^2\sigma}{d\Omega d\omega} = \frac{k'}{k} \frac{1}{2\pi\hbar} \sum_{j,j'} b_j b_{j'} \int_{-\infty}^{\infty} \left\langle e^{-i\hat{\mathbf{Q}}\hat{\mathbf{R}}_{j'}(0)} e^{i\hat{\mathbf{Q}}\hat{\mathbf{R}}_j(t)} \right\rangle e^{-i\omega t} dt$$

Correlation functions

Intermediate correlation function:

$$I(\mathbf{Q}, t) = \frac{1}{N} \sum_{j,j'} \left\langle e^{-i\mathbf{Q}\mathbf{R}_{j'}(0)} e^{i\mathbf{Q}\mathbf{R}_j(t)} \right\rangle$$

Pair correlation function:

$$G(\mathbf{R}, t) = \frac{1}{(2\pi)^3} \int_{-\infty}^{\infty} I(\mathbf{Q}, t) e^{-i\mathbf{Q}\mathbf{R}} d\mathbf{Q}$$

Dynamical structure factor:

$$S(\mathbf{Q}, \omega) = \frac{1}{(2\pi\hbar)} \int_{-\infty}^{\infty} I(\mathbf{Q}, t) e^{-i\omega t} dt$$

$$\left(\frac{d^2\sigma}{d\Omega dE'} \right)_{\text{coh}} = \frac{\sigma_{\text{coh}}}{4\pi} \frac{k'}{k} NS(\mathbf{q}, \omega) \quad \sigma_{\text{coh}} = 4\pi b^2$$

Elastic Scattering

$$\frac{d^2\sigma}{d\Omega d\omega} = \frac{k'}{k} \frac{1}{2\pi\hbar} \sum_{j,j'} b_j b_{j'} \int_{-\infty}^{\infty} \langle e^{-i\mathbf{Q}\cdot\hat{\mathbf{R}}_{j'}(0)} e^{i\mathbf{Q}\cdot\hat{\mathbf{R}}_j(t)} \rangle e^{-i\omega t} dt$$

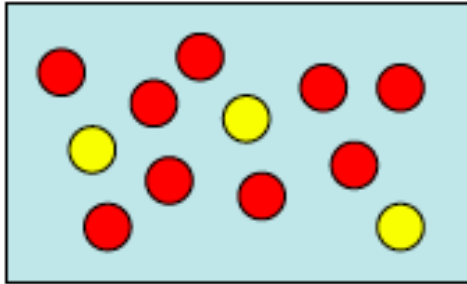
assume there is no time operator time dependence, and integrate over time

$$\frac{d\sigma}{d\Omega} = \int_{-\infty}^{\infty} \left(\frac{d^2\sigma}{d\Omega d\omega} \right) d(\hbar\omega) = \sum_{j,j'} b_j b_{j'} \langle e^{-i\mathbf{Q}\cdot\hat{\mathbf{R}}_{j'}} e^{i\mathbf{Q}\cdot\hat{\mathbf{R}}_j} \rangle$$

Assume fixed atomic positions (replace operators with atomic positions)

$$\begin{aligned} \left(\frac{d\sigma}{d\Omega} \right)_{\text{coh}} &= \langle b \rangle^2 \sum_{j,j'} e^{-i\mathbf{Q}\cdot(\mathbf{R}_{j'} - \mathbf{R}_j)}, \\ \left(\frac{d\sigma}{d\Omega} \right)_{\text{inc}} &= (\langle b^2 \rangle - \langle b \rangle^2) \sum_{j=j'} e^{-i\mathbf{Q}\cdot(\mathbf{R}_{j'} - \mathbf{R}_j)} = N (\langle b^2 \rangle - \langle b \rangle^2) \end{aligned}$$

Incoherent Scattering



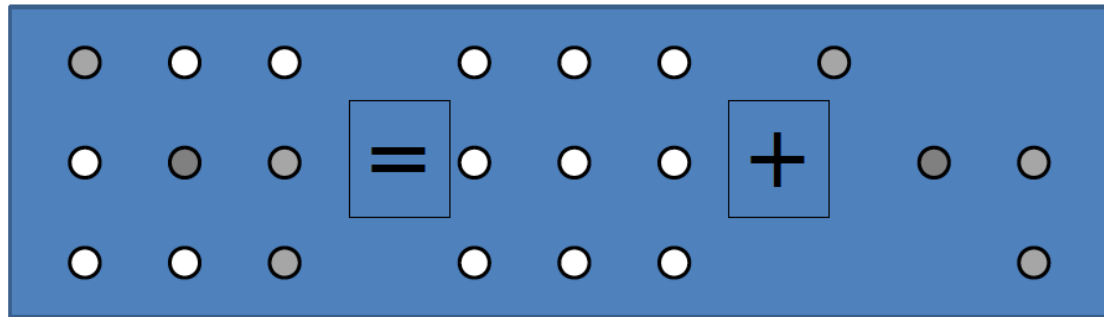
P_α : relative probability for isotope

$$\sum_{\alpha} P_{\alpha} = 1$$

$$\langle b \rangle = \sum_{\alpha} P_{\alpha} b_{\alpha} \quad \langle b^2 \rangle = \sum_{\alpha} P_{\alpha} b_{\alpha}^2$$

$$\sigma_{coh} = 4\pi \langle b \rangle^2$$

$$\sigma_{inc} = 4\pi \left[\langle b^2 \rangle - \langle b \rangle^2 \right]$$



Dynamical scattering theory

- Multiple scattering was previously ignored because these effects are often weak
- Below an critical angle of incidence neutrons are completely reflected, this scattering is not weak

$$\left[-\frac{\hbar^2}{2m} \nabla^2 + V(\mathbf{r}) \right] \psi_{\mathbf{k}}(\mathbf{r}) = E \psi_{\mathbf{k}}(\mathbf{r}) \quad V(\mathbf{r}) = 0 \quad \text{except } \mathbf{r} \in \text{target region } \mathbb{T}$$

$$\psi'_{\mathbf{k}}(\mathbf{r}) = \psi_{\mathbf{k}}(\mathbf{r}) + \int_{\mathbb{T}} d^3 r' G_{\circ}(\mathbf{r}, \mathbf{r}' | E) V(\mathbf{r}') \psi'_{\mathbf{k}}(\mathbf{r}')$$

$$G_{\circ}(\mathbf{r}, \mathbf{r}' | E) = -\frac{2m}{\hbar^2} \frac{1}{4\pi} \frac{e^{ik|\mathbf{r}-\mathbf{r}'|}}{|\mathbf{r}-\mathbf{r}'|} \quad \text{with} \quad k = \sqrt{\frac{2m}{\hbar^2} E}$$

Dynamical scattering theory

- The wave-vector can change when a wave enters a medium

$$\left[\nabla^2 + 2m(E - \bar{V}) / \hbar^2 \right] \psi(r) = 0$$

incident wave: $\psi(r) = e^{i\vec{k}_o \cdot \vec{r}}$ $\bar{V} = \frac{2\pi\hbar^2}{m} \rho$

- Vacuum : $k_0^2 = 2mE / \hbar^2$
- Medium: $k^2 = 2m(E - \bar{V}) / \hbar^2 = k_0^2 - 4\pi\rho$
- Refractive index $n=k/k_0$: $n = 1 - \lambda^2 \rho / 2\pi$

Snell's and Fresnel's Law

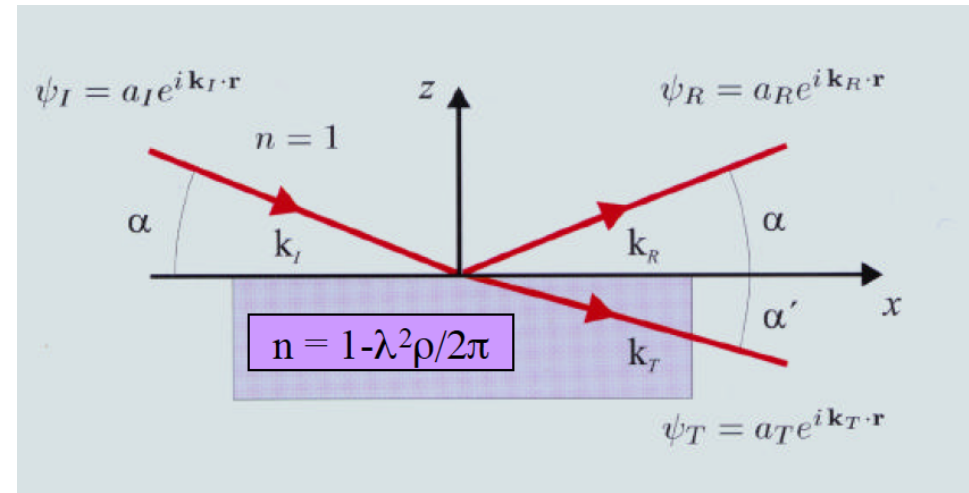
- Continuity of Ψ and Ψ' at surface

$$a_I + a_R = a_T \quad (1)$$

$$a_I \vec{k}_I + a_R \vec{k}_R = a_T \vec{k}_T$$

$$a_I k \cos \alpha + a_R k \cos \alpha = a_T n k \cos \alpha'$$

$$-(a_I - a_R) k \sin \alpha = -a_T n k \sin \alpha'$$



- Snell's law:

$$\cos \alpha = n \cos \alpha'$$

- Fresnel's law:

$$r = a_R / a_I = (k_{Iz} - k_{Tz}) / (k_{Iz} + k_{Tz})$$

Energy/time scales and length scales

