

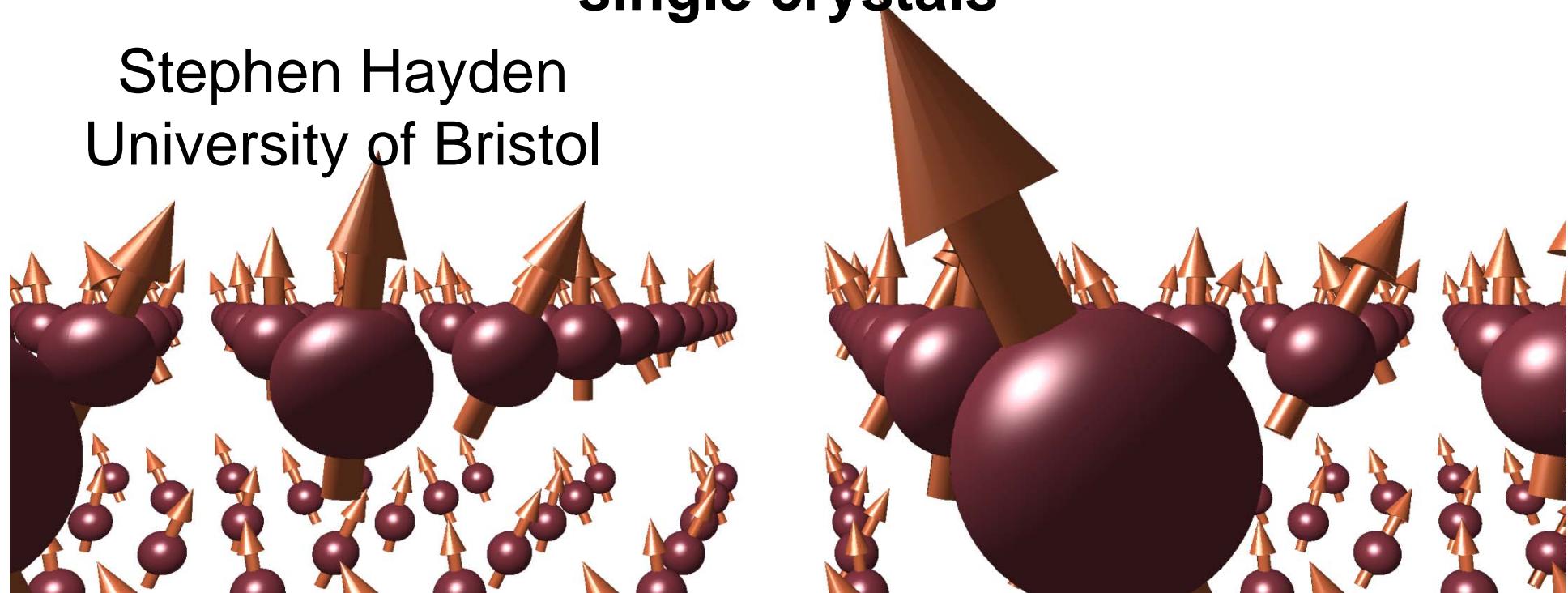


University of  
BRISTOL



# Neutron spectroscopy of collective excitations in single crystals

Stephen Hayden  
University of Bristol



Summer school on condensed matter research, Zuoz, August 2015.

- I. Thinking about resolution (Three axis)
- II. Direct geometry time-of-time Spectroscopy
- III. Magnetic excitations in cuprates
- IV. Magnetic excitations in  $\text{Sr}_3\text{Ru}_2\text{O}_7$

- I. Thinking about resolution (Three axis)
- II. Direct geometry time-of-time Spectroscopy
- III. Magnetic excitations in cuprates
- IV. Magnetic excitations in  $\text{Sr}_3\text{Ru}_2\text{O}_7$

# The story so far.....

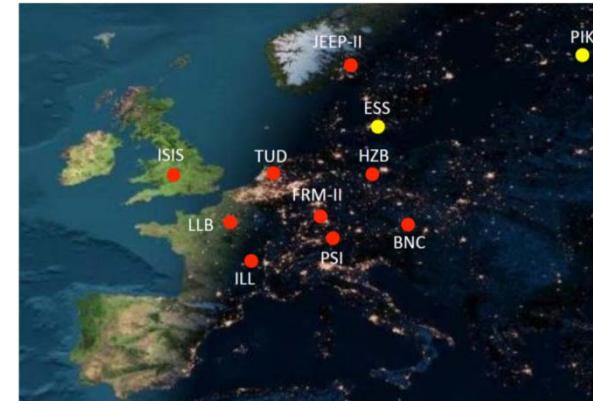
**Summerschool on Condensed Matter Research, Zuoz, 15. August 2015 – 21. August 2015**  
**Spectroscopy with Neutrons, Muons and Photons**

	Sunday 16 Aug 2015	Monday 17 Aug 2015	Tuesday 18 Aug 2015	Wednesday 19 Aug 2015	Thursday 20 Aug 2015	Friday 21 Aug 2015
09:00 – 10:15	Neutron and Muon Sources - (HIPA, target E, SINQ and ESS)  Mike Seidel	Introduction to MuSr practical - scientific problem to be analysed  Rustum Khasanov	Using X-rays – general introduction and science examples  Phil Willmott	Excursion	Practicals (in parallel) 1) Multiplet analysis (F. de Groot) 2) EXAFS analysis (C. Borca)	Neutron spectroscopy of collective excitations in single crystals  Stephen Hayden
10:15 – 10:45	COFFEE	COFFEE	COFFEE		COFFEE	COFFEE
10:45 – 12:00	MuSr - techniques and instrumentation  Hubertus Luetkens	Introduction to Neutron practical - scientific problem to be analysed  Sandor Toth	Generating X-rays – the machines  Riccardo Bartolini		Practicals (in parallel) 1) Multiplet analysis (F. de Groot) 2) EXAFS analysis (C. Borca)	New Developments (neutron optics mm)  Peter Böni
12:15 – 15:30	Lunch and Free Afternoon	Lunch and Free Afternoon	Lunch and Free Afternoon		Lunch and Free Afternoon	Closing remarks Lunch Departure
15:30 – 16:00	COFFEE	COFFEE	COFFEE	COFFEE	COFFEE	
16:00 – 17:15	Neutrons – techniques and instrumentation  Ken Anderson	Practicals (in parallel) 1) Muon Analysis (Z. Guguchia, P. Biswas) 2) Spin wave analysis (C. Rüegg, S. Toth)	Using X-rays – Instrumentation for spectroscopy  Jan Dreiser	Introduction to Exafs practical – scientific problem to be analysed  Camelia Borca	New developments photons: High brightness sources  Christoph Quitmann	
17:15 – 17:30	SHORT BREAK	SHORT BREAK	SHORT BREAK	SHORT BREAK	SHORT BREAK	SHORT BREAK
17:30 – 18:15	CAMEA  Henrik Rønnow	Practicals (in parallel) 1) Muon Analysis (Z. Guguchia, P. Biswas) 2) Spin wave analysis (C. Rüegg, S. Toth)	XFEL – techniques and instrumentation (time-resolved spectroscopy)  Chris Milne	Introduction to Multiplet practical – scientific problem to be analysed  Frank de Groot	Complementary methods, Optical spectroscopy  Leonardo Degiorgi	
18:15 – 18:45	Short poster presentation		Short poster presentation			
19:00 – 20:00	DINNER	DINNER	DINNER	DINNER	Apéro	
20:00 – 21:00	FREE EVENING	Supercomputing  Thomas Schulthess	POSTER SESSION	Monitoring landcover / landuse with satellites and airborne systems  Charlotte Steinmeier	BANQUET	

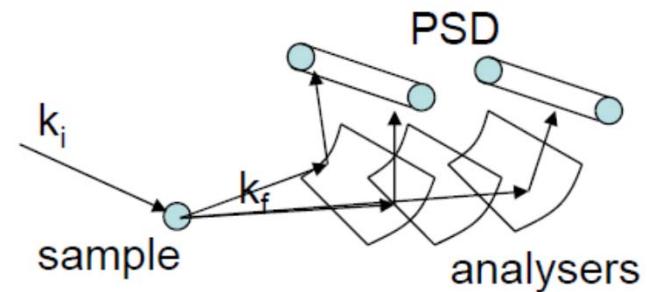
# The story so far.....

16:00 – 17:15	Neutrons – techniques and instrumentation  <b>Ken Anderson</b>
17:15 – 17:30	SHORT BREAK
17:30 – 18:15	CAMEA  <b>Henrik Rønnow</b>

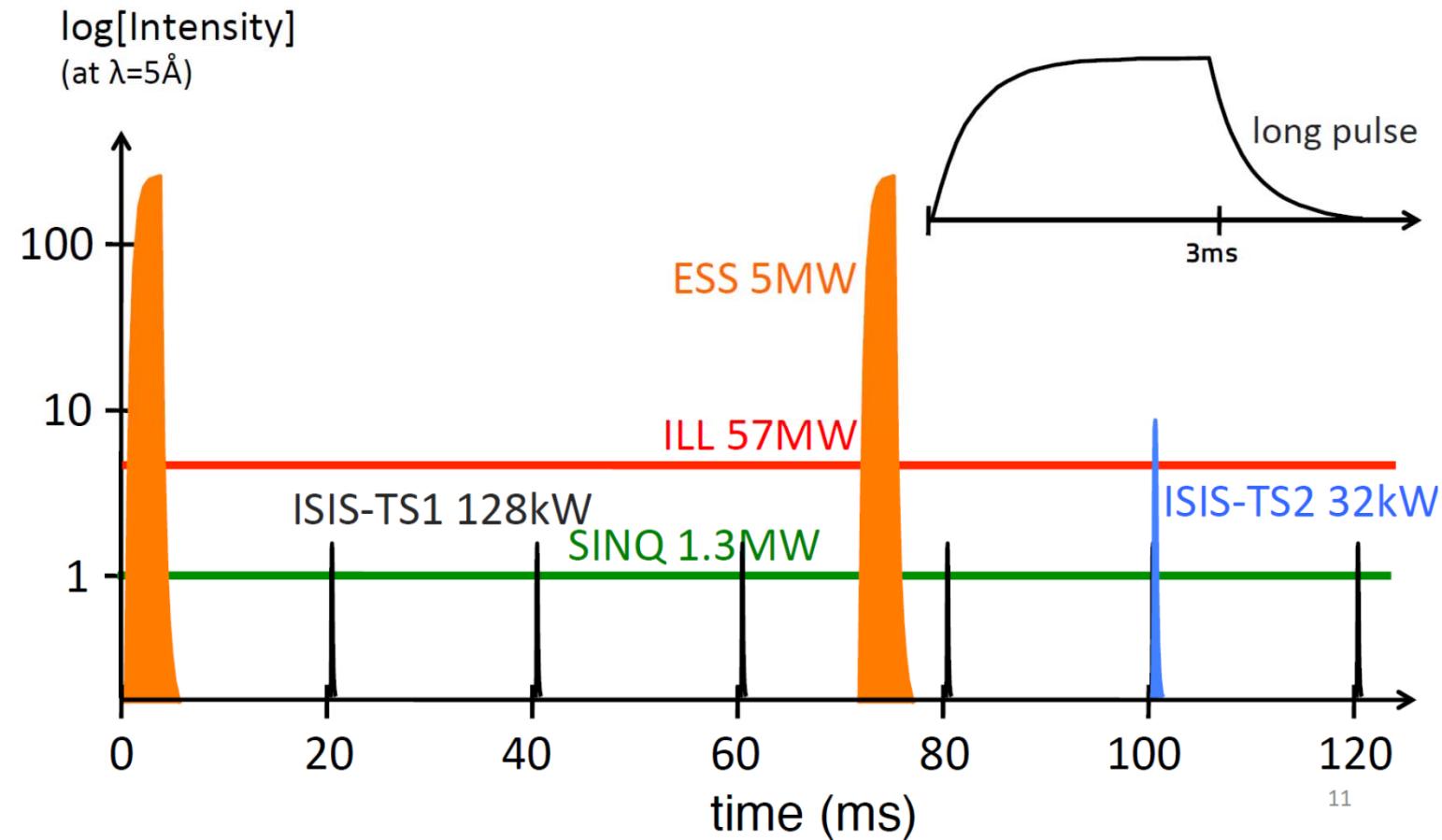
## Neutron Techniques and Instrumentation



# CAMEA



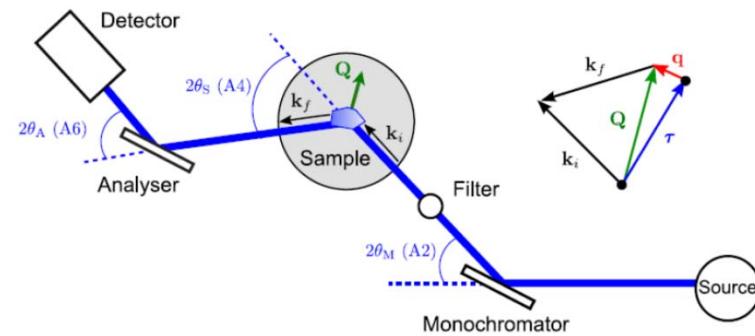
# Time structure of neutron sources



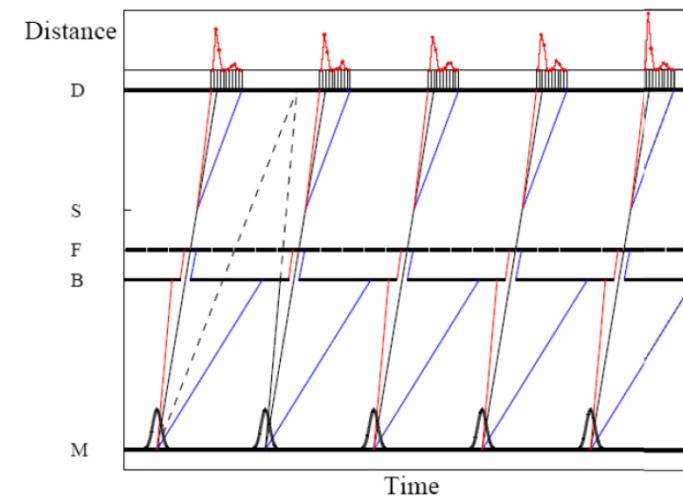
from K Anderson

# TAS and ToF

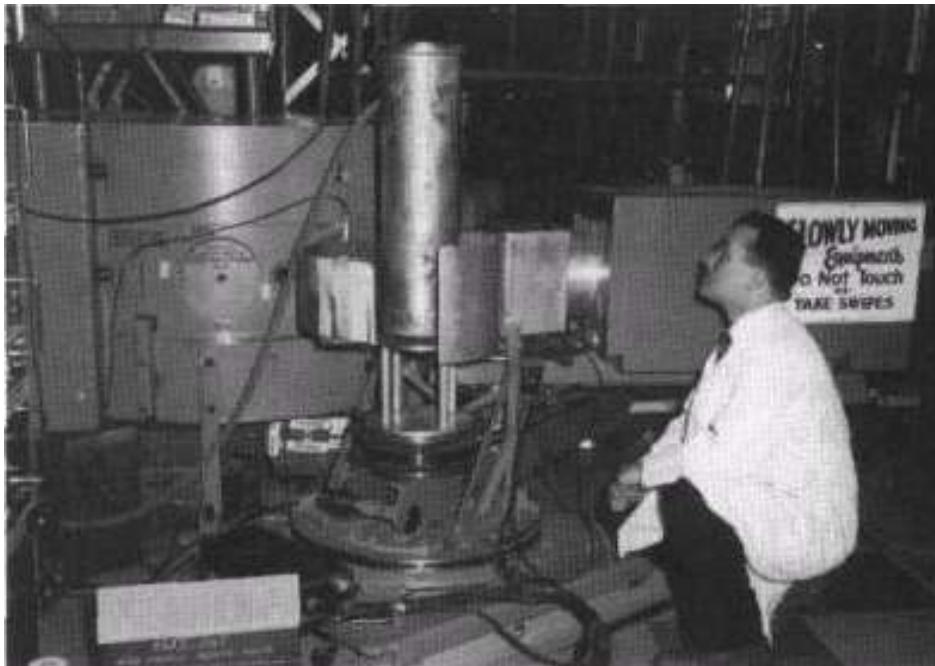
Triple axis spectrometer



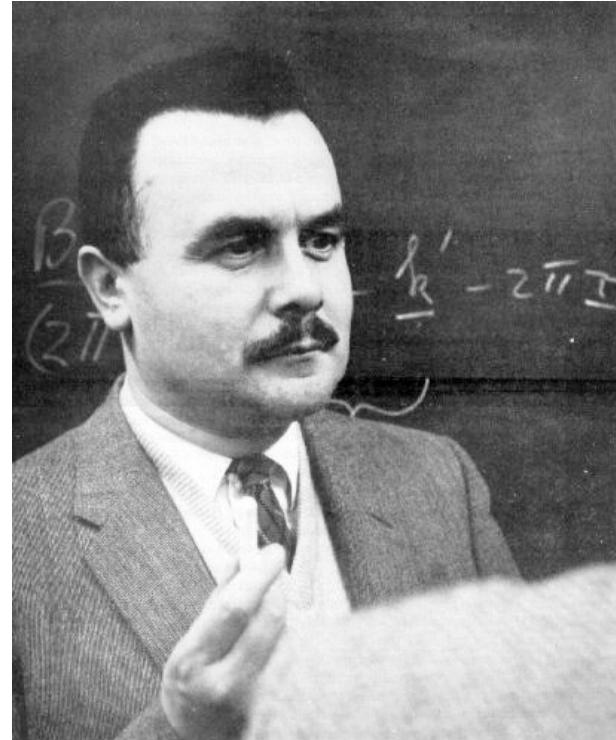
Time of flight spectrometer



# Triple Axis Spectrometer

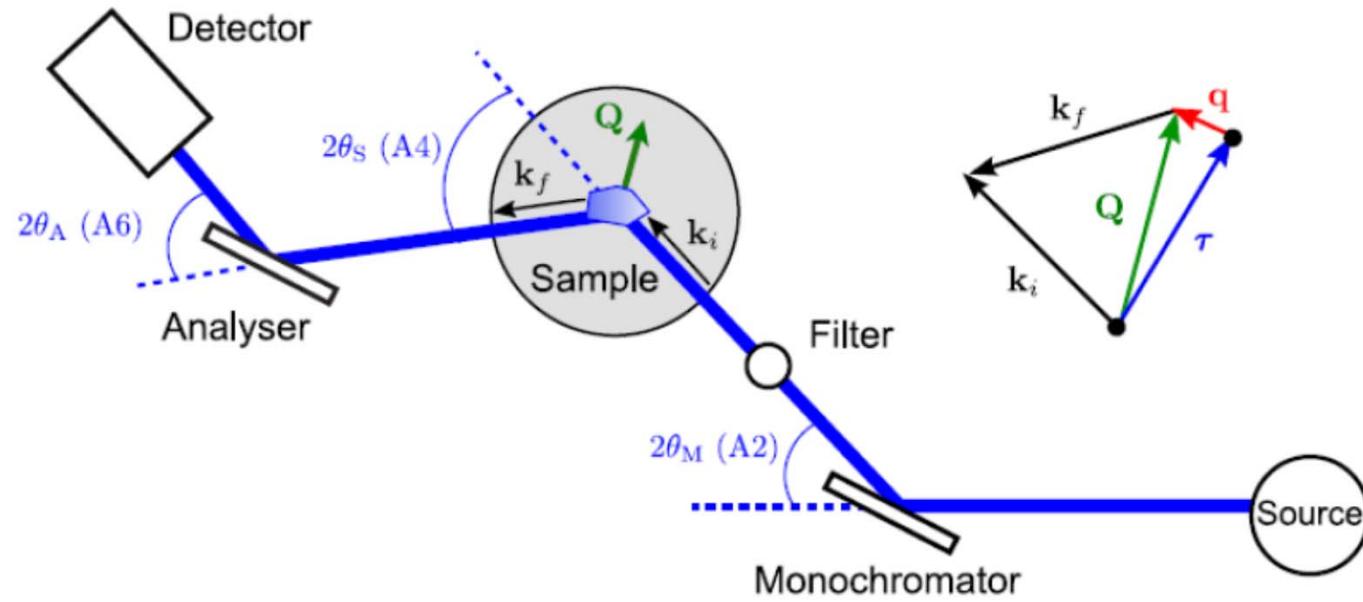


Triple-axis spectrometer installed at NRX in  
Chalk River, 1958



Bert Brockhouse (1918 –2003),  
Noble prize, 1994

# Triple Axis Spectrometer



$$\mathbf{Q} = \mathbf{k}_i - \mathbf{k}_f \quad (\mathbf{Q}, \omega) \text{ refer to the } \textit{excitation}.$$

$$\hbar\omega = E_i - E_f$$

# Typical cross sections

## Phonons vs. magnetic excitations

### Phonons

$$\left( \frac{\partial^2 \sigma}{\partial \Omega \partial E'} \right)_{phonon} \propto \frac{\sigma_{coh}}{4\pi} \frac{k_f}{k_i} \frac{(2\pi)^3}{v_0} \frac{1}{2M} \exp(-2W)x$$

$$\sum_s \sum_{\tau} \frac{(\vec{Q} \cdot \vec{e}_s)^2}{\omega_s} \left[ (n(\omega_s) + 1) \delta(\omega - \omega_s) \delta(\vec{Q} - \vec{q} - \vec{\tau}) + n(\omega_s) \delta(\omega + \omega_s) \delta(\vec{Q} + \vec{q} - \vec{\tau}) \right]$$

Increases with **Q**

Increases with **T**

$$\left( \frac{\partial^2 \sigma}{\partial \Omega \partial E'} \right)_{mag} \propto \frac{k_f}{k_i} \left\{ \frac{1}{2} gF(\vec{Q}) \right\}^2 \sum_{\alpha\beta} (\delta_{\alpha\beta} - \hat{Q}_\alpha \hat{Q}_\beta) \sum_{ll'} \exp(i\vec{Q} \cdot (\vec{l} - \vec{l}')) \langle S_l^\alpha(0) S_{l'}^\beta(t) \rangle \exp(-i\omega t) dt$$

Decreases with **Q**

Bose statistics for simple ordered systems (FM, AFM, etc.).  
Spreads in **Q** and  $\hbar\omega$  at high T  
Limited by sum-rules

+ Neutron polarisation analysis

from H Ronnow

# Scattering Law

We can write the neutron scattering cross section using the general form:

$$\frac{d^2\sigma}{d\Omega dE_f} = \frac{k_f}{k_i} S(\mathbf{q}, \hbar\omega)$$

or

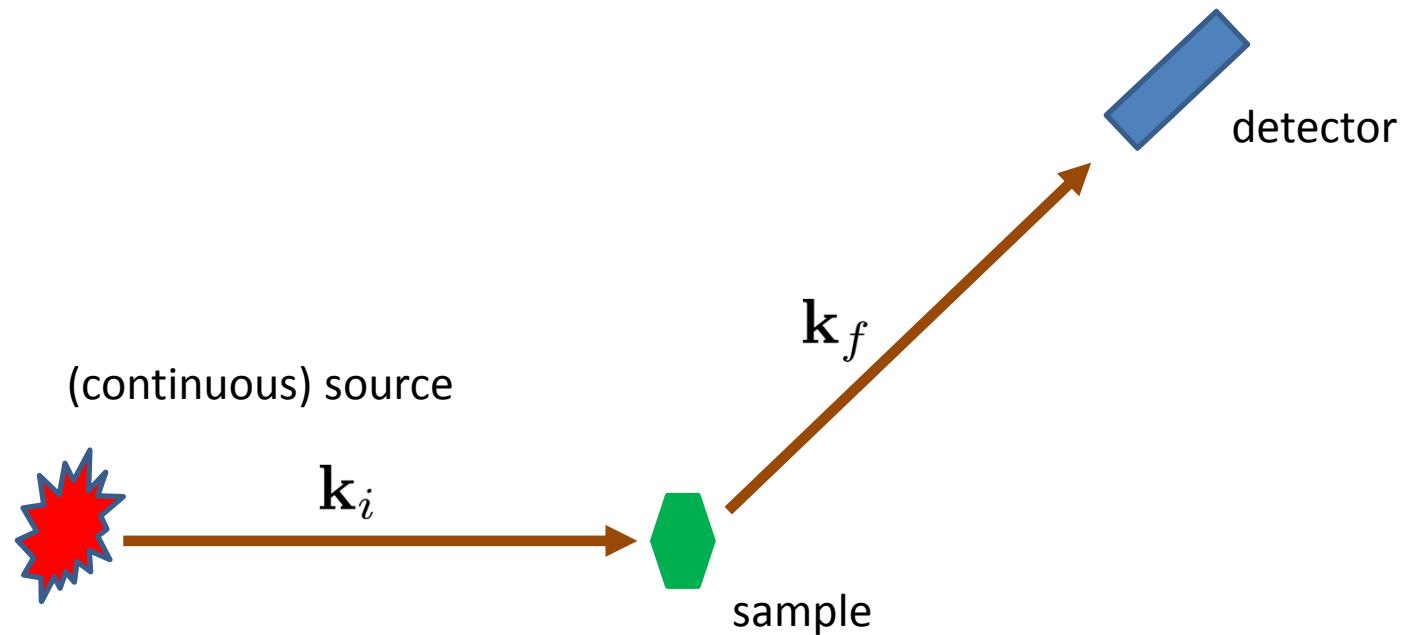
$$\frac{k_i}{k_f} \frac{d^2\sigma}{d\Omega dE_f} = S(\mathbf{q}, \hbar\omega)$$

six coordinates  $(\mathbf{k}_i, \mathbf{k}_f)$

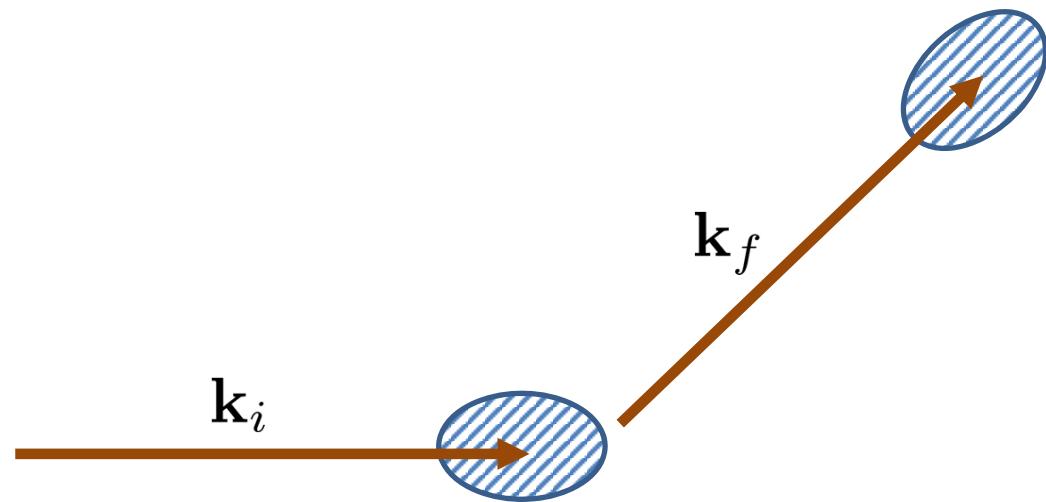
four coordinates  $(\mathbf{Q}, \hbar\omega)$

$$\left. \begin{array}{l} \mathbf{Q} = \mathbf{k}_i - \mathbf{k}_f \\ \hbar\omega = E_i - E_f \end{array} \right\} \text{Transform with Jacobian}$$

# Schematic of inelastic scattering experiment



# Resolutions and Phase Space



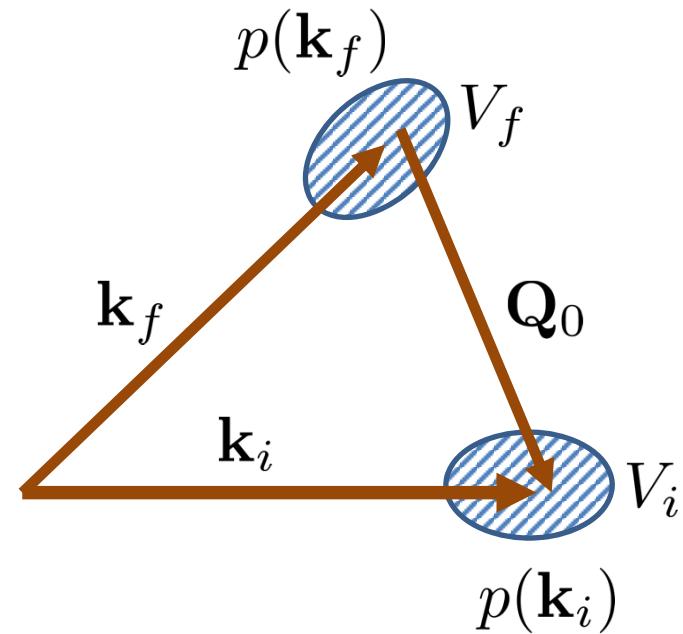
# Definitions

$$d\Omega_f = \frac{dk_{f,x} dk_{f,y}}{k_f^2}; \quad dE_f = \frac{\hbar^2 k_f}{m} dk_{f,z}$$

- $p(\mathbf{k}_i)$  is the probability that a neutron leaves the source and hits the sample with wavevector  $\mathbf{k}_i$
- $p(\mathbf{k}_f)$  is the probability that a neutron with wavevector  $\mathbf{k}_f$  passes the analyer system an is detected.
- $S(\mathbf{Q}, \omega)$  is proportional to the probablity that a neutron is scattered from  $\mathbf{k}_i$  to  $\mathbf{k}_f$

# Resolutions and Phase Space

$$\begin{aligned} I &\propto S(\mathbf{Q}, \omega) p_i(\mathbf{k}_i) p_f(\mathbf{k}_f) dk_{ix} dk_{iy} dk_{iz} dk_{fx} dk_{fy} dk_{ fz} \\ &\propto S(\mathbf{Q}, \omega) p_i(\mathbf{k}_i) p_f(\mathbf{k}_f) dV_i dV_f \end{aligned} \quad \text{small deviations}$$



$$I(\mathbf{Q}_0, \omega_0) = A \int S(\mathbf{Q}, \omega) R(\mathbf{Q} - \mathbf{Q}_0, \omega - \omega_0) d^3 Q d\omega$$

See B. Dorner Acta Cryst. A28, 319 (1972).

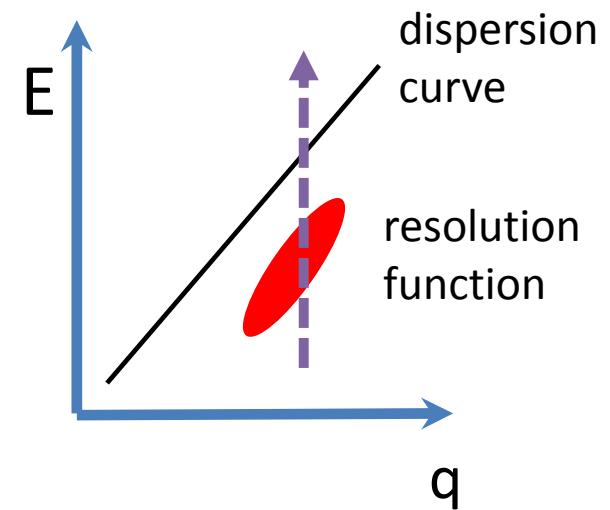
# Qualitative properties of the resolution function

- resolution volumes increase with  $k_i$  and  $k_f$ .  $V_i \propto k_i^3 \cot \theta_M$
- convolution of volumes can lead to focussing
- vertical focussing leads to poor resolution along  $z$ .

Programs to calculate resolution function:

**RESTRAX**      <http://neutron.ujf.cas.cz/restrax/>

**McStas**      <http://www.mcstas.org/>



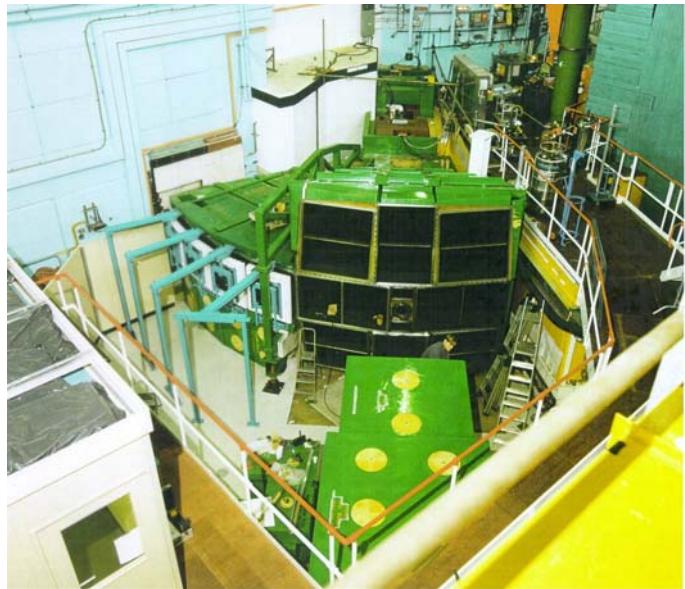
Key references:

MAIER-LEIBNITZ, H. (1966). Nukleonik, 8, 61.

COOPER, M. J. and NATHANS, R. (1967). Acta Cryst. 23, 357.

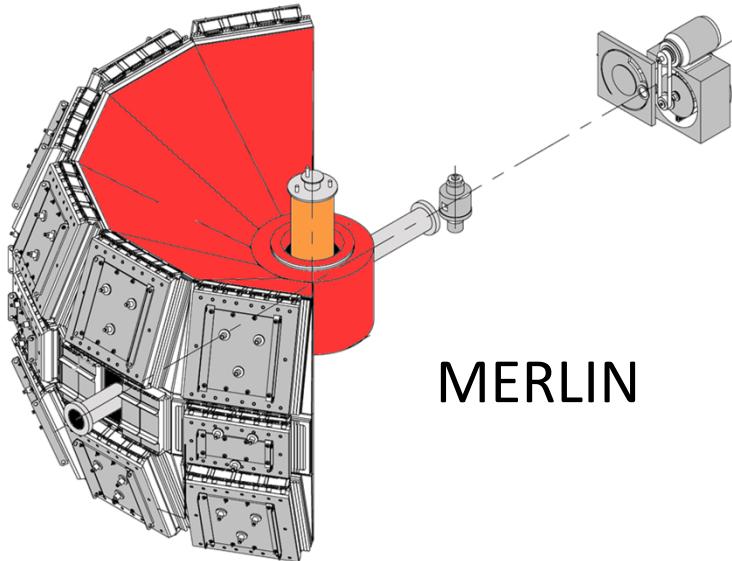
- I. Thinking about resolution (Three axis)
- II. Direct geometry time-of-time Spectroscopy
- III. Magnetic excitations in cuprates
- IV. Magnetic excitations in  $\text{Sr}_3\text{Ru}_2\text{O}_7$

# Direct geometry Time-of-time Spectroscopy



Spectrometers at ISIS

MAPS

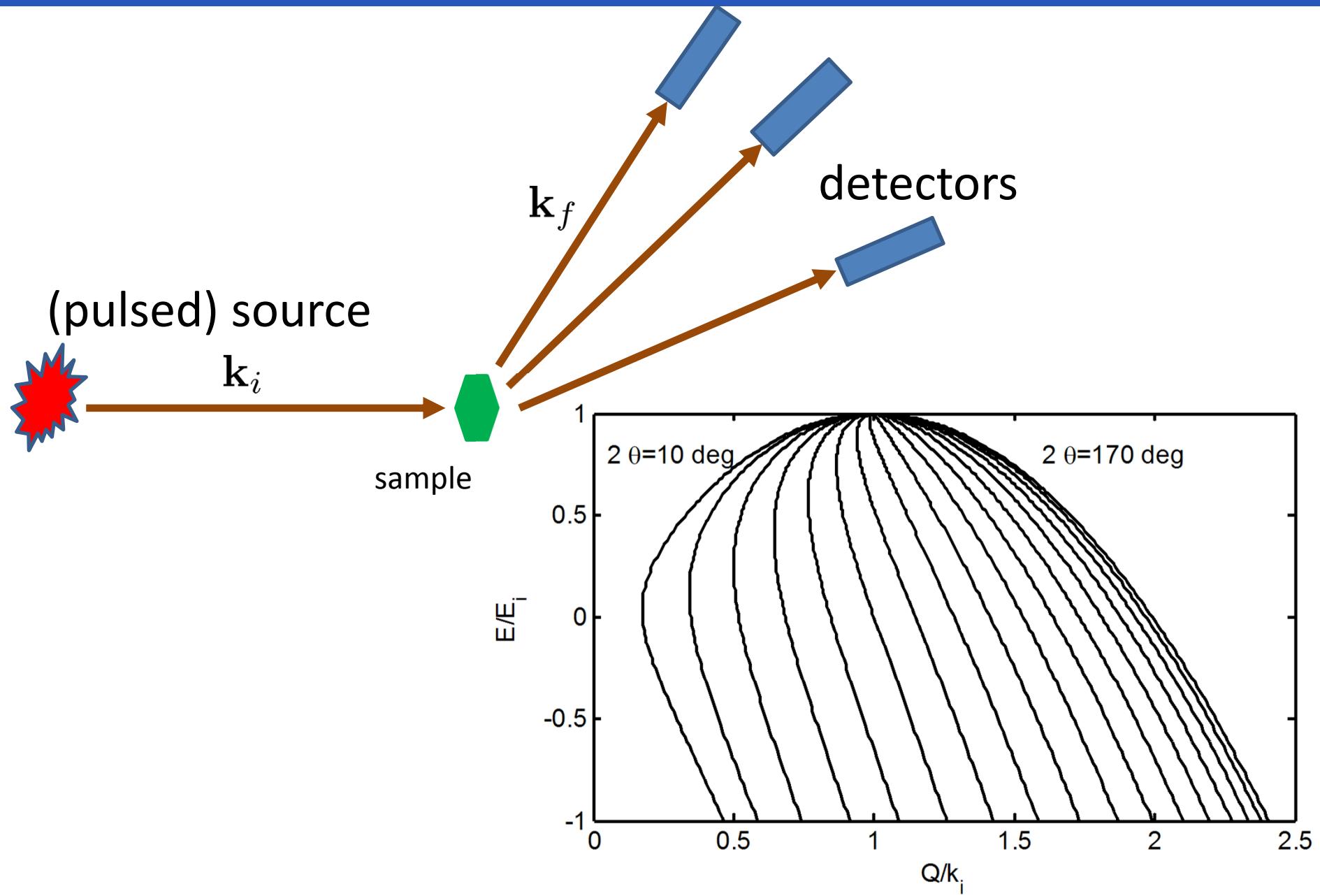


MERLIN



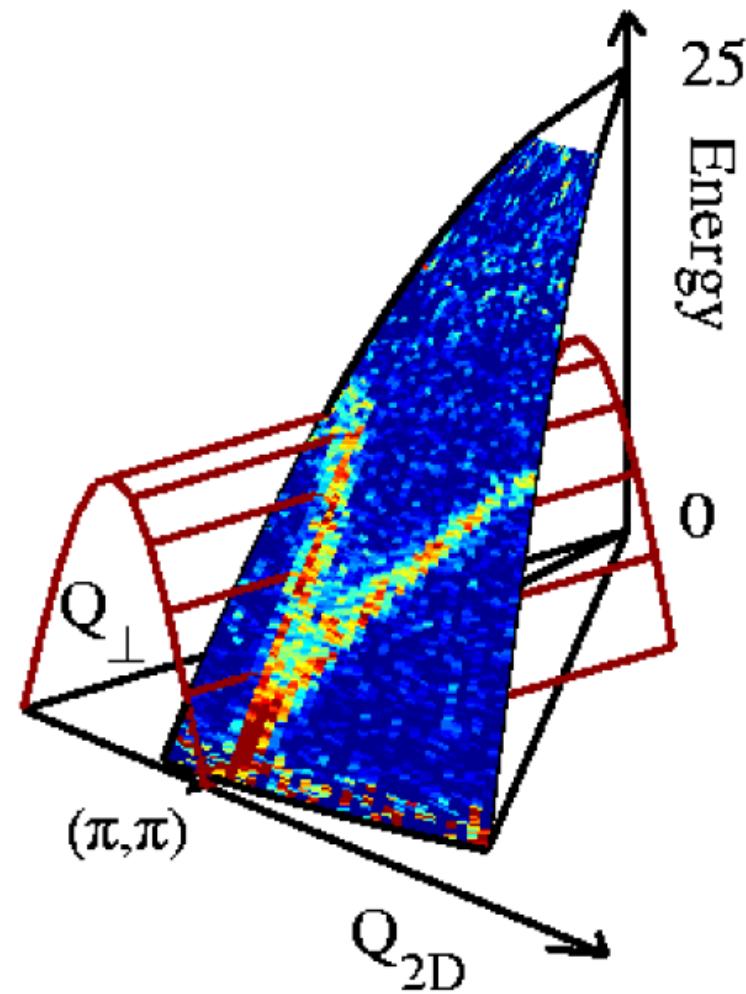
LET

## Direct geometry ToF (one detector or a line)

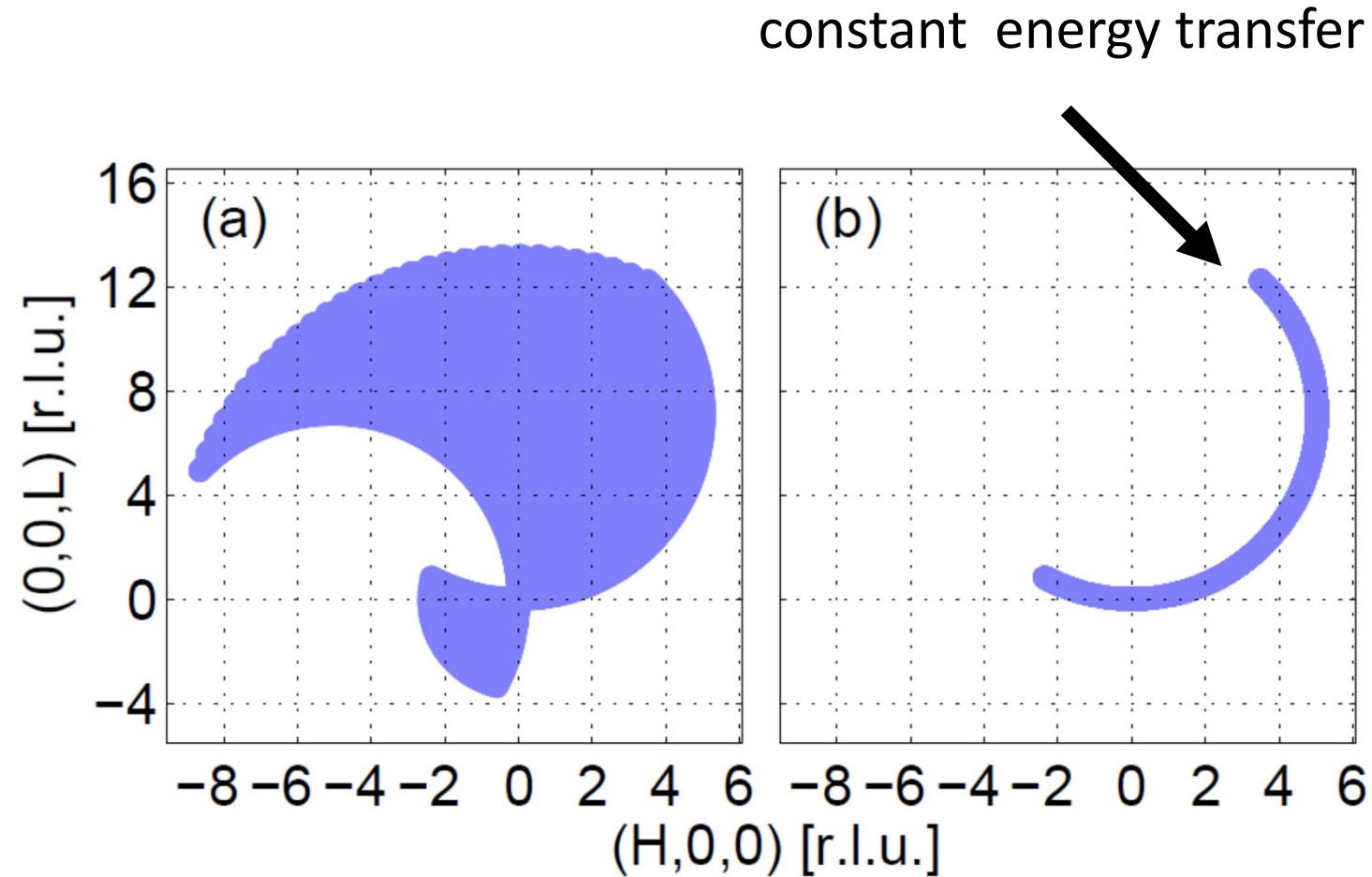


## Direct geometry ToF (line of detectors)

and 1D system...

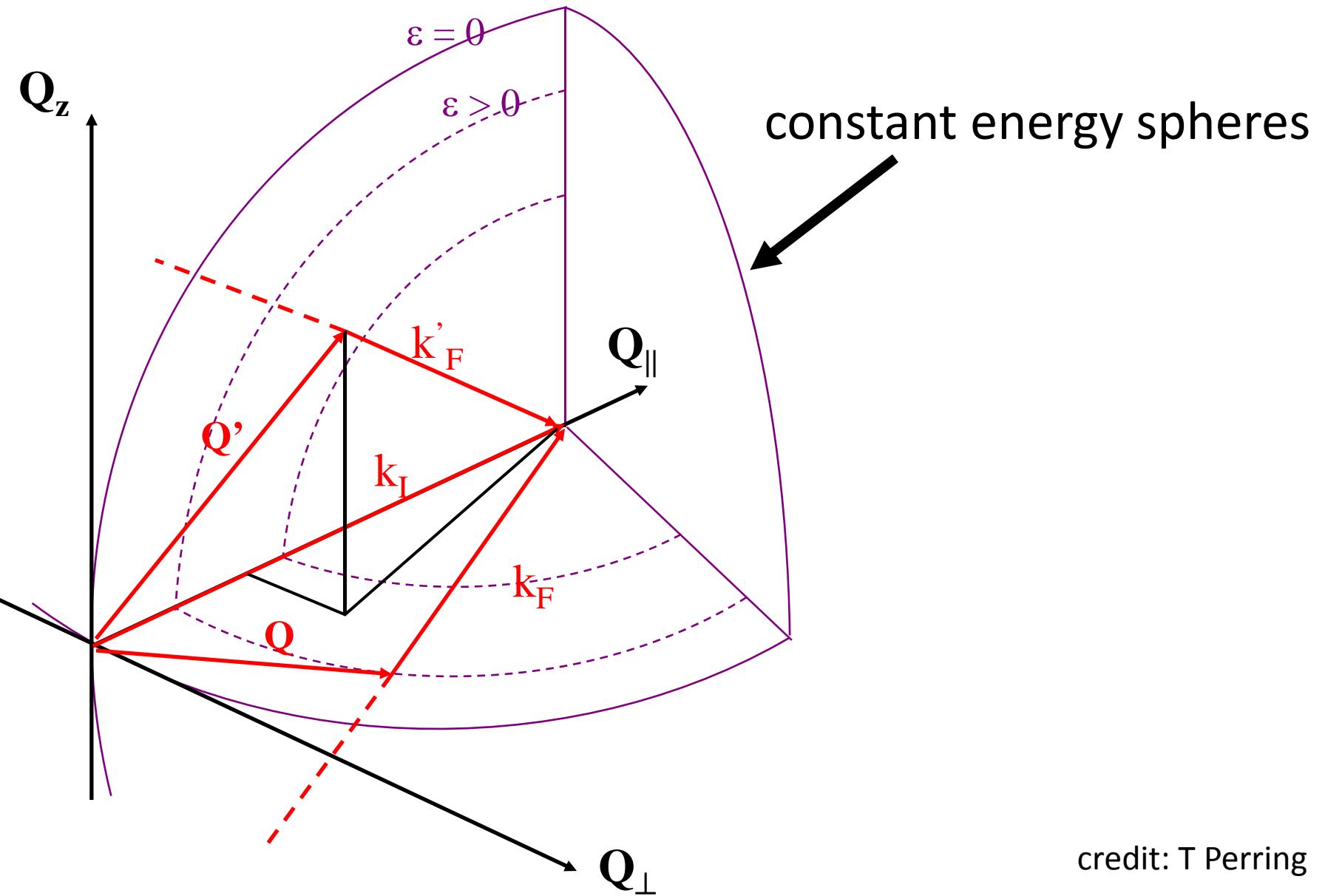


# Rotating sample



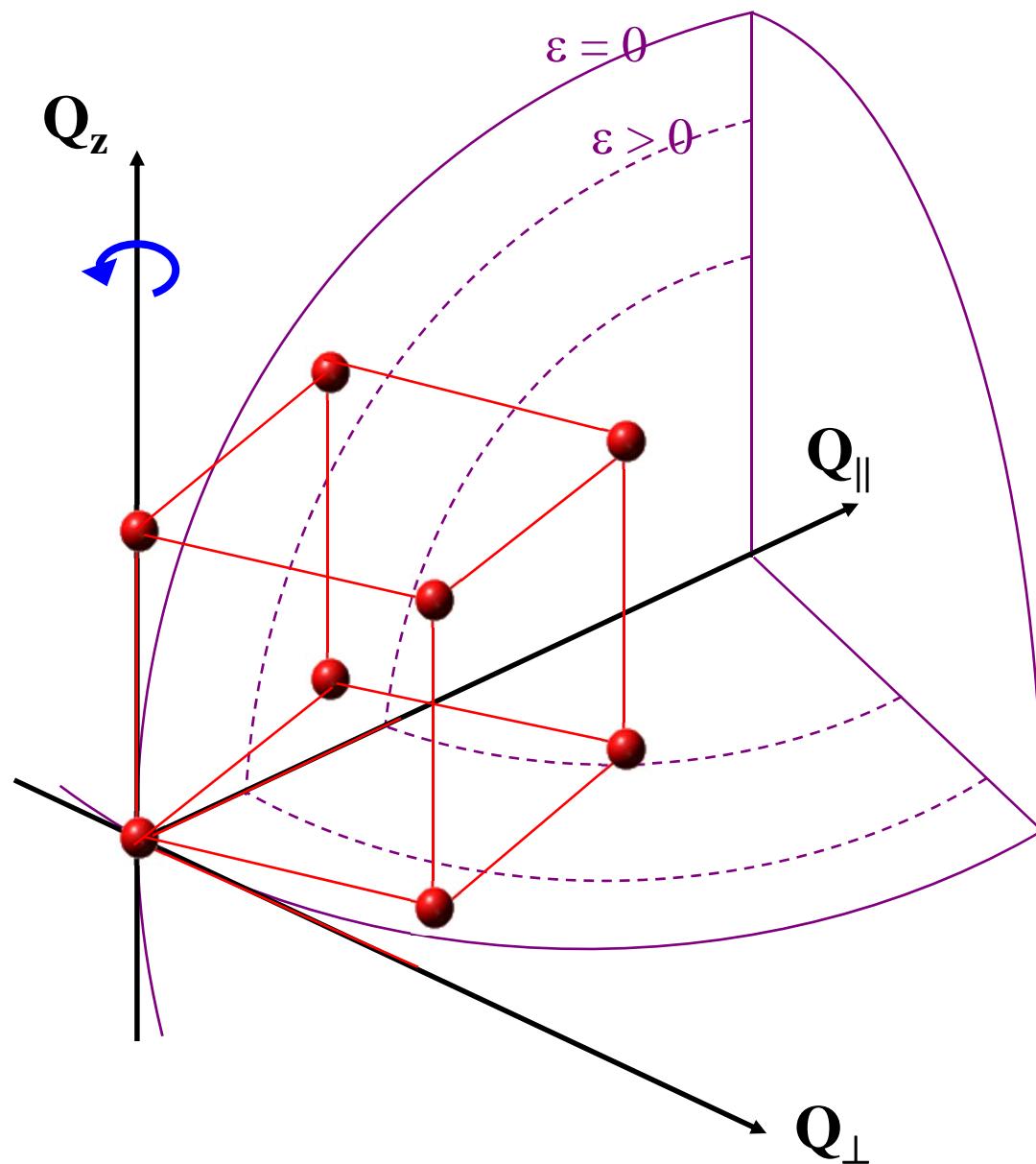
credit: R Ewings

## 2-D (area) of detectors



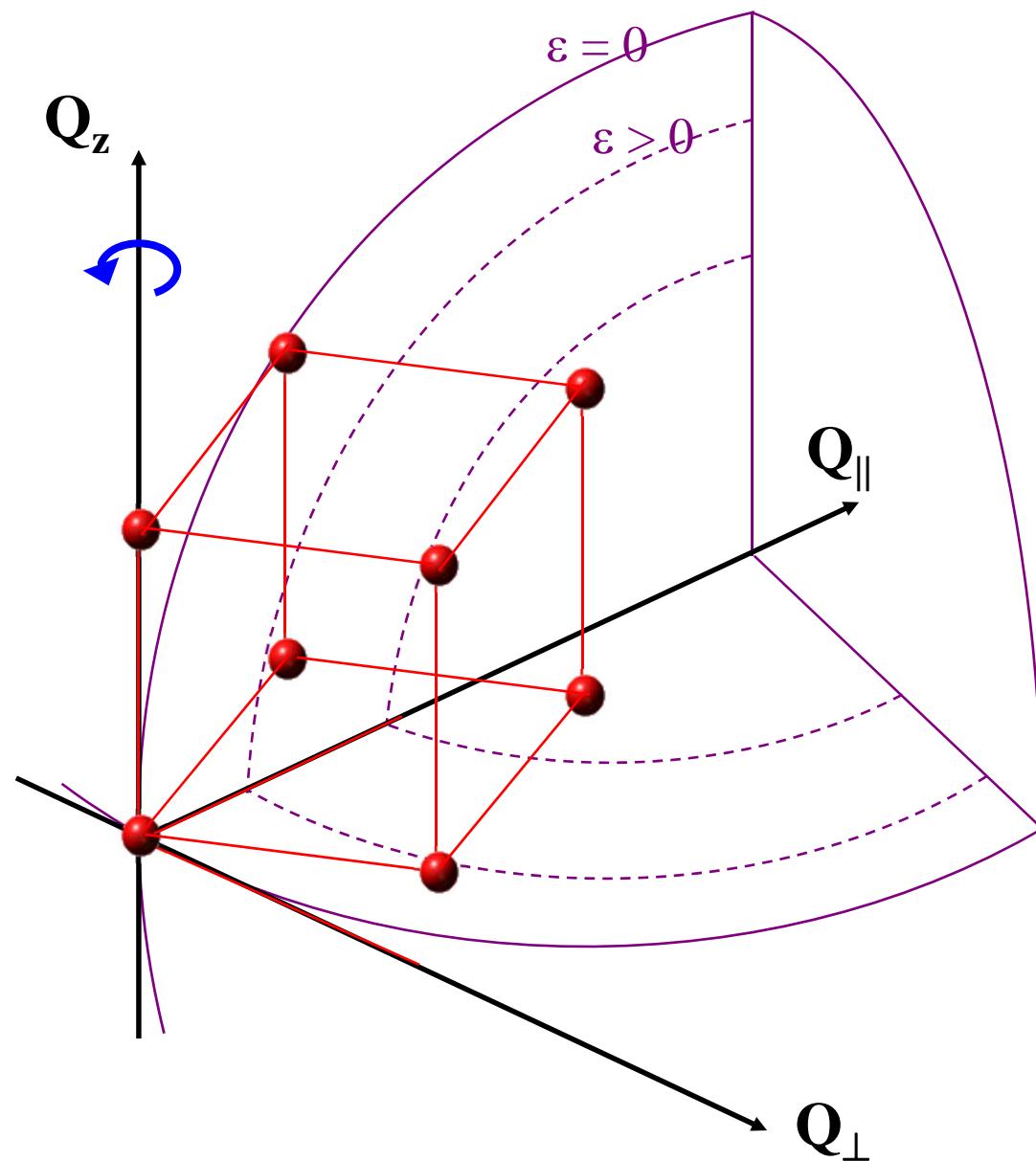
credit: T Perring

## 2-D (area) of detectors : rotate sample



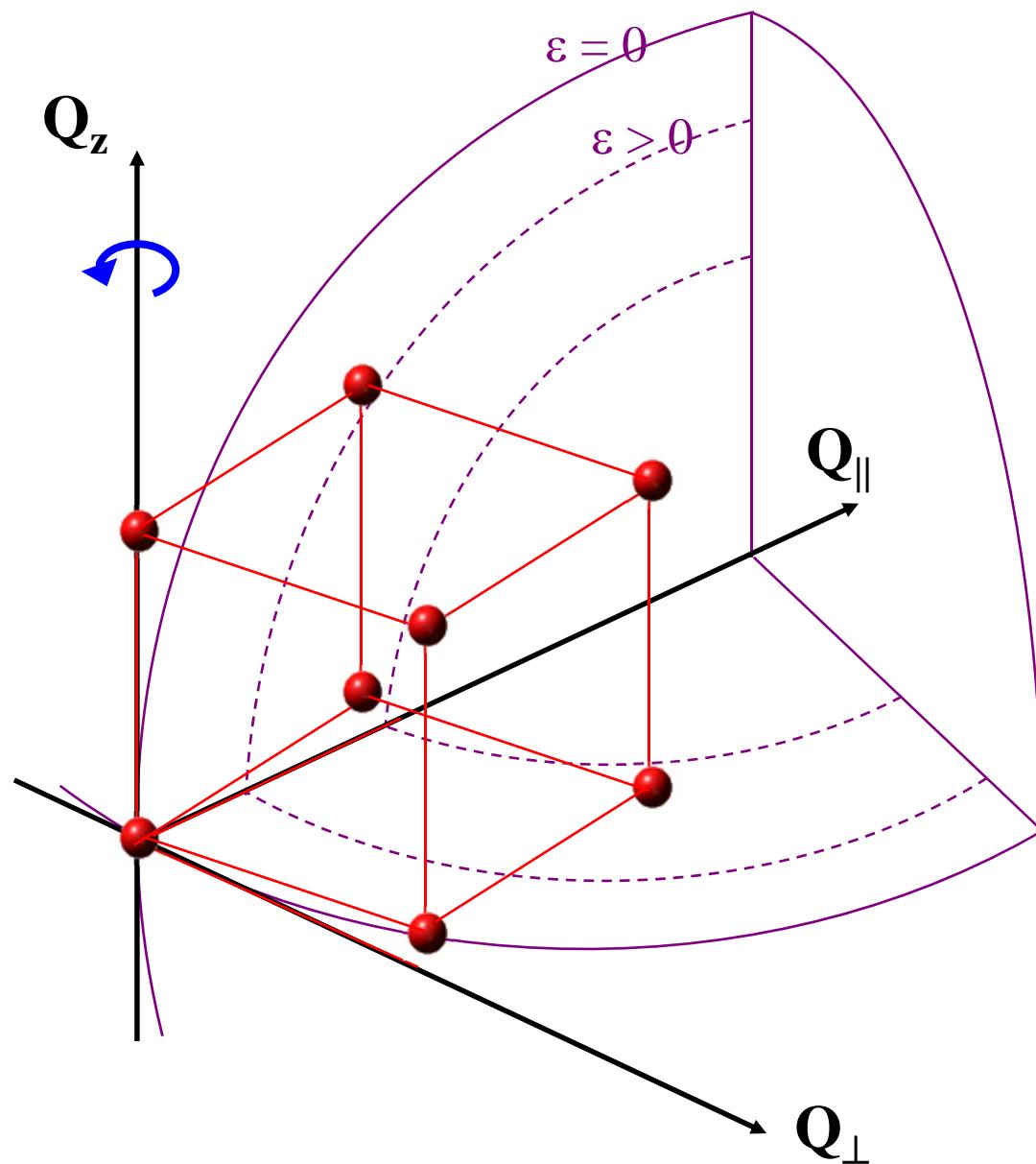
credit: T Perring

## 2-D (area) of detectors : rotate sample



credit: T Perring

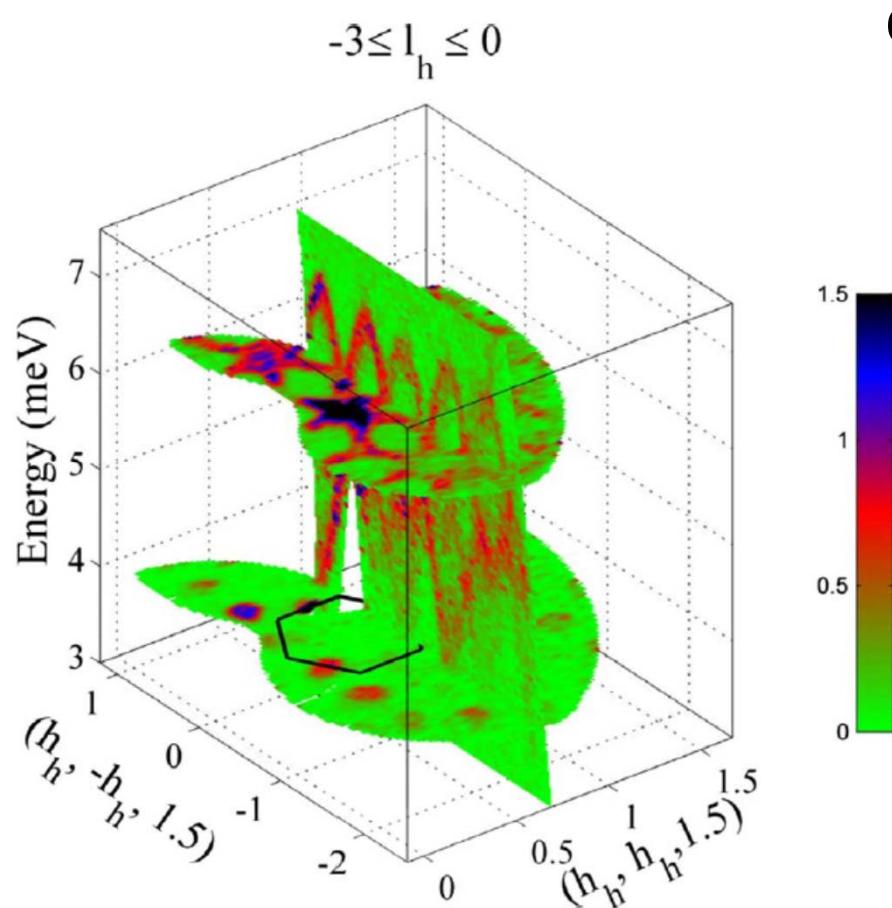
## 2-D (area) of detectors : rotate sample



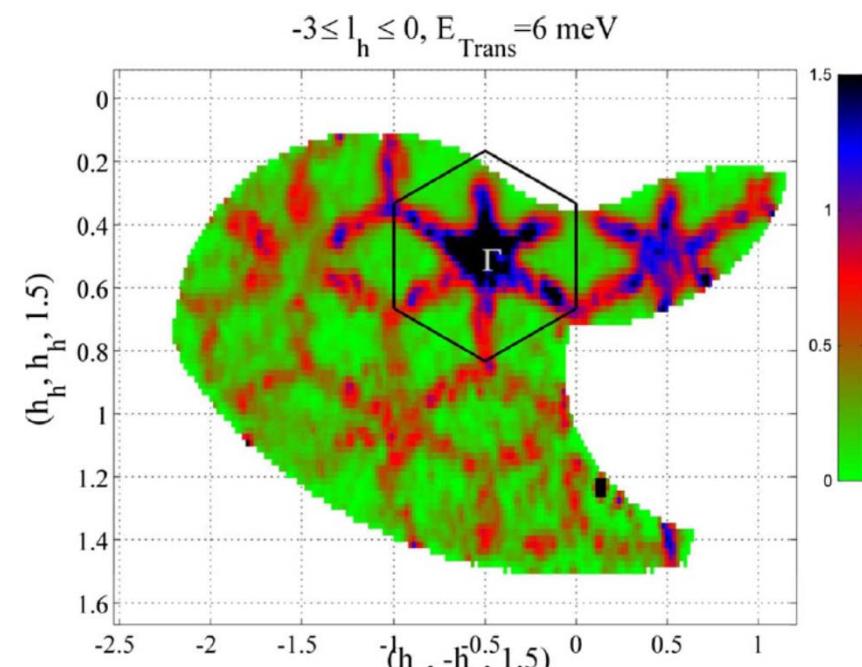
credit: T Perring

# Example of ‘Horaced’ data

Spin dimer antiferromagnet  $\text{Sr}_3\text{Cr}_2\text{O}_8$  MERLIN



cut through 4-D space of data

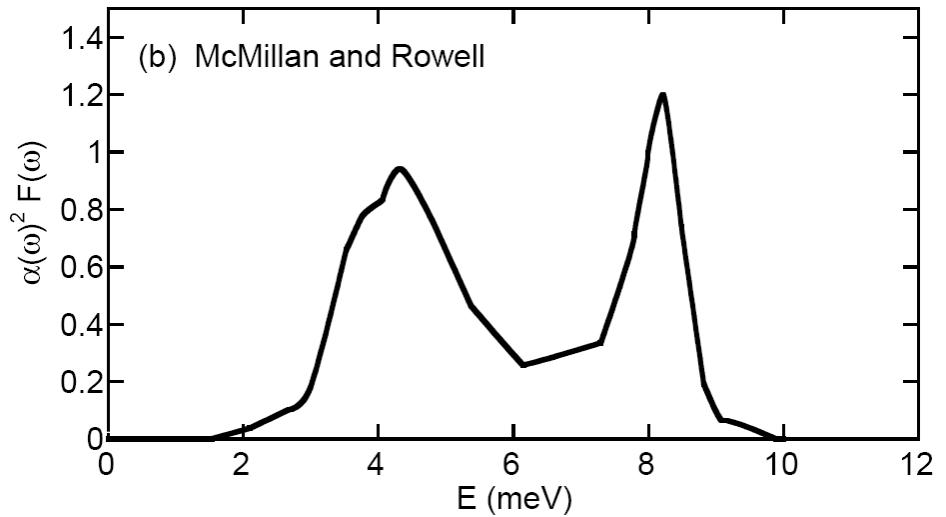
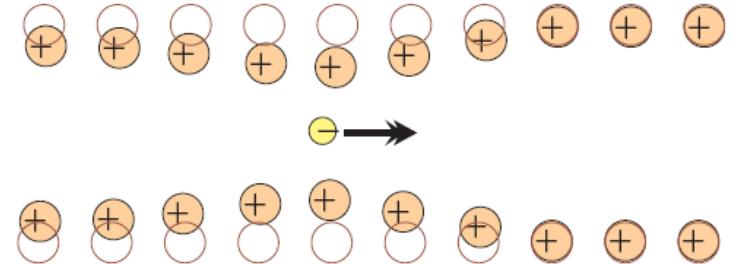
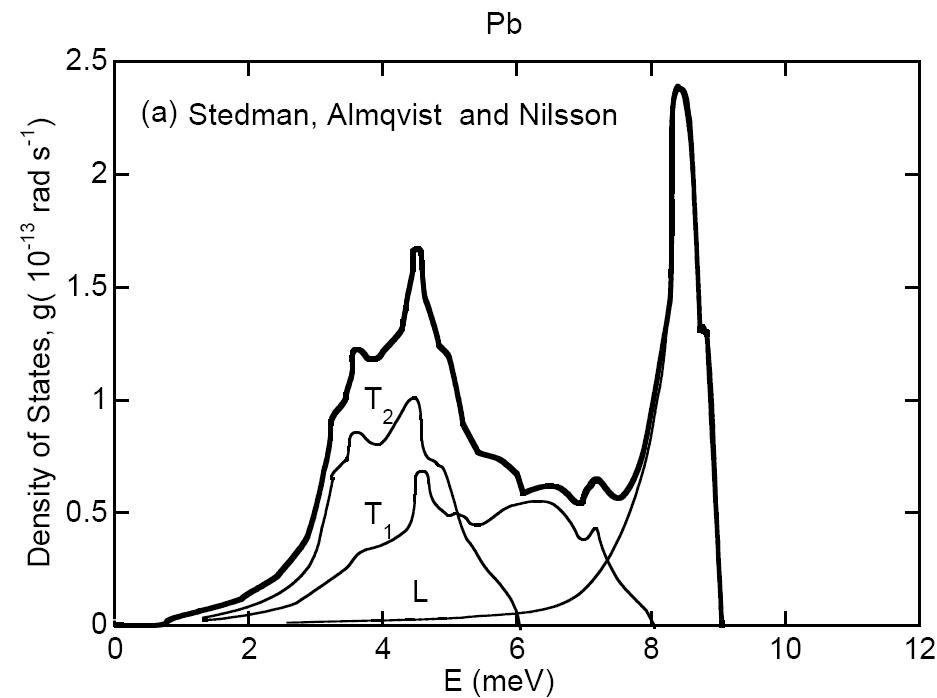


- I. Thinking about resolution (Three axis)
- II. Direct geometry time-of-time Spectroscopy
- III. Magnetic excitations in cuprates
- IV. Magnetic excitations in  $\text{Sr}_3\text{Ru}_2\text{O}_7$

# Magnetic excitations in cuprates

- I. The pairing glue
- II. Neutron scattering in  $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$
- III. Parent antiferromagnet
- IV. Optimal doping
- V. Overdoped
- VI. Summary

# "Phonon glue" in lead



- phonon DOS,  $F(\omega)$
- $\alpha^2(\omega) F(\omega)$  from SIS tunnelling
- SIS Tunnelling in Conventional Superconductors

# Magnetically mediated pairing

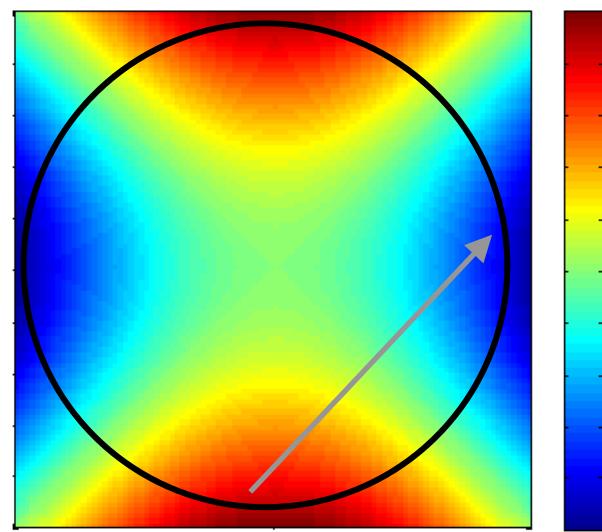
$$V_{kk'} = g^2 \chi(q) \quad \text{Pairing Interaction}$$

$$\Delta_k = - \sum_{k'} V_{kk'} \frac{\Delta_{k'}}{2\sqrt{\varepsilon_{k'}^2 + \Delta_{k'}^2}} \quad \text{Gap equation}$$

favours states where  $\Delta_k$  changes sign over Fermi surface

# Gap functions

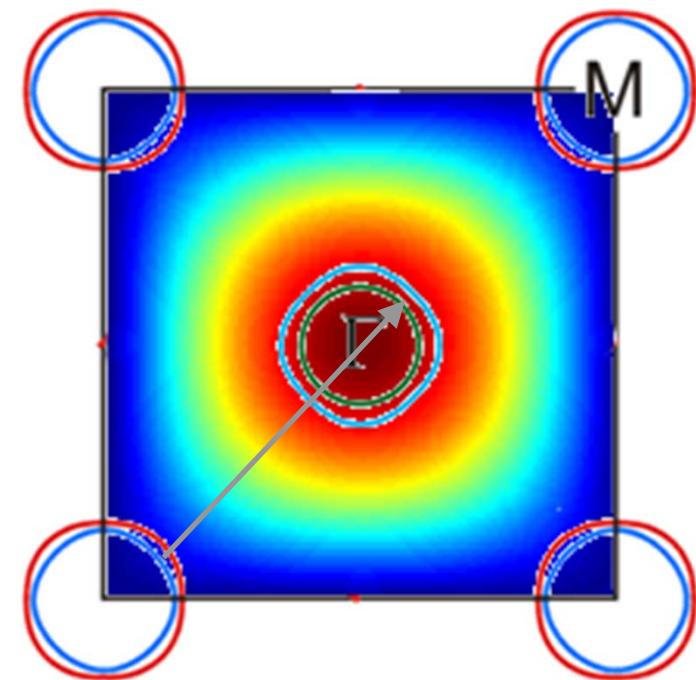
cuprates



$$d_{x^2-y^2}$$

$$\Delta_k = \Delta_0 (\cos k_x - \cos k_y)$$

pnicides



$$s_{\pm}$$

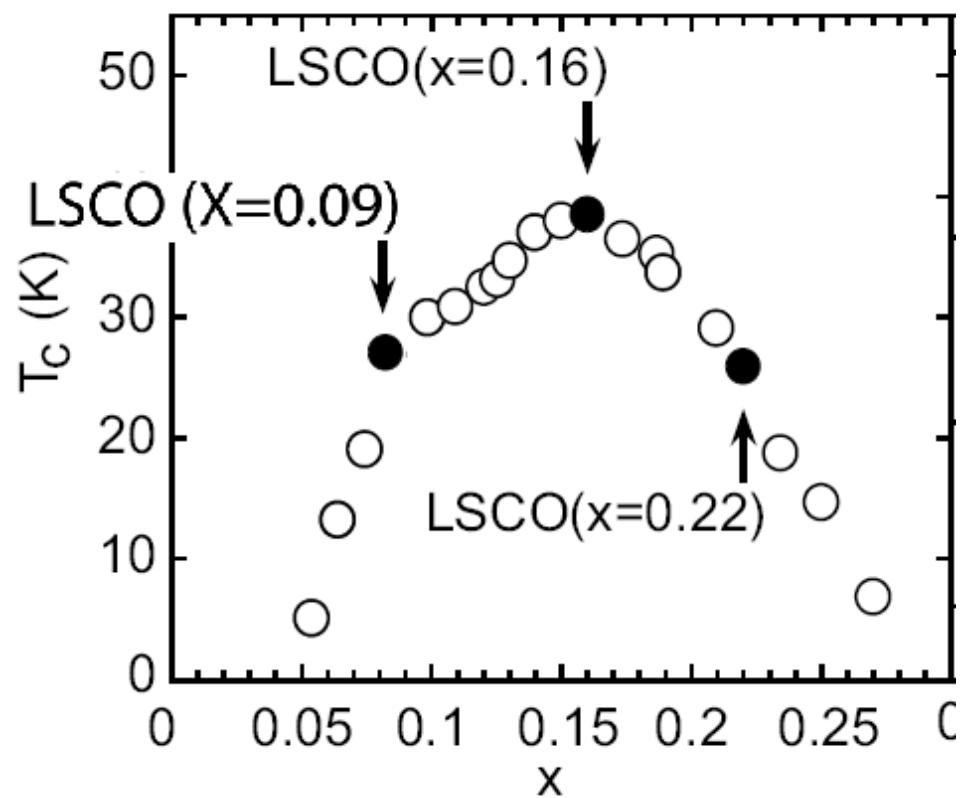
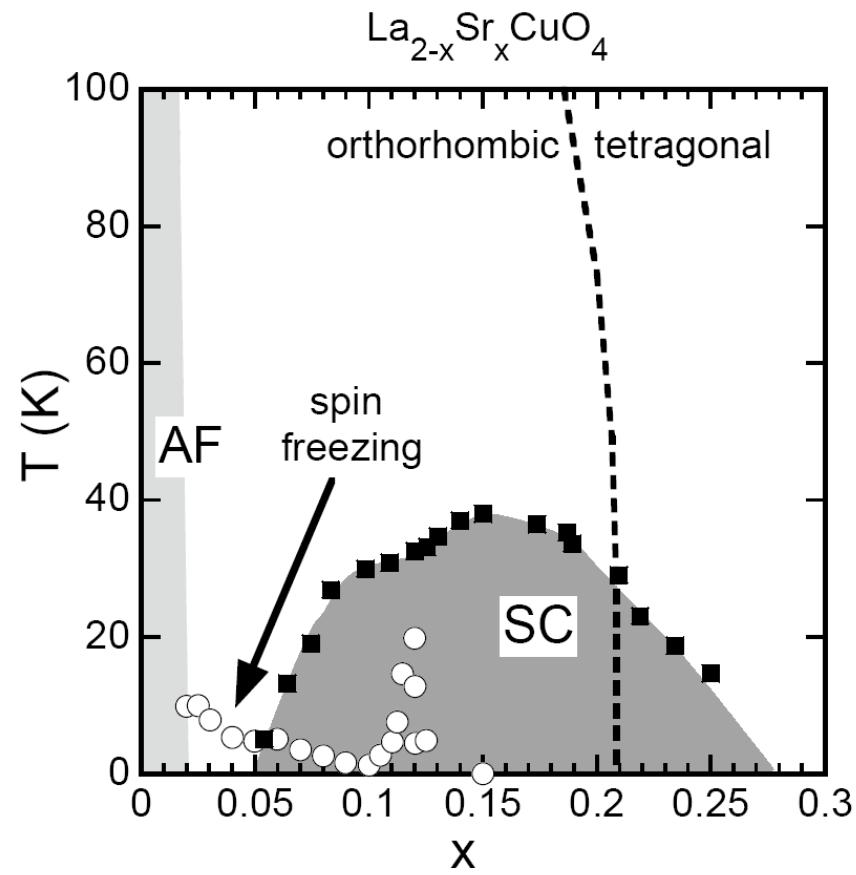
$$\Delta_k = \Delta_0 \cos k_x \cos k_y$$

Mazin et al, PRL 08

# Magnetic excitations in cuprates

- I. The pairing glue
- II. Neutron scattering in  $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$
- III. Parent antiferromagnet
- IV. Optimal doping
- V. Overdoped
- VI. Summary

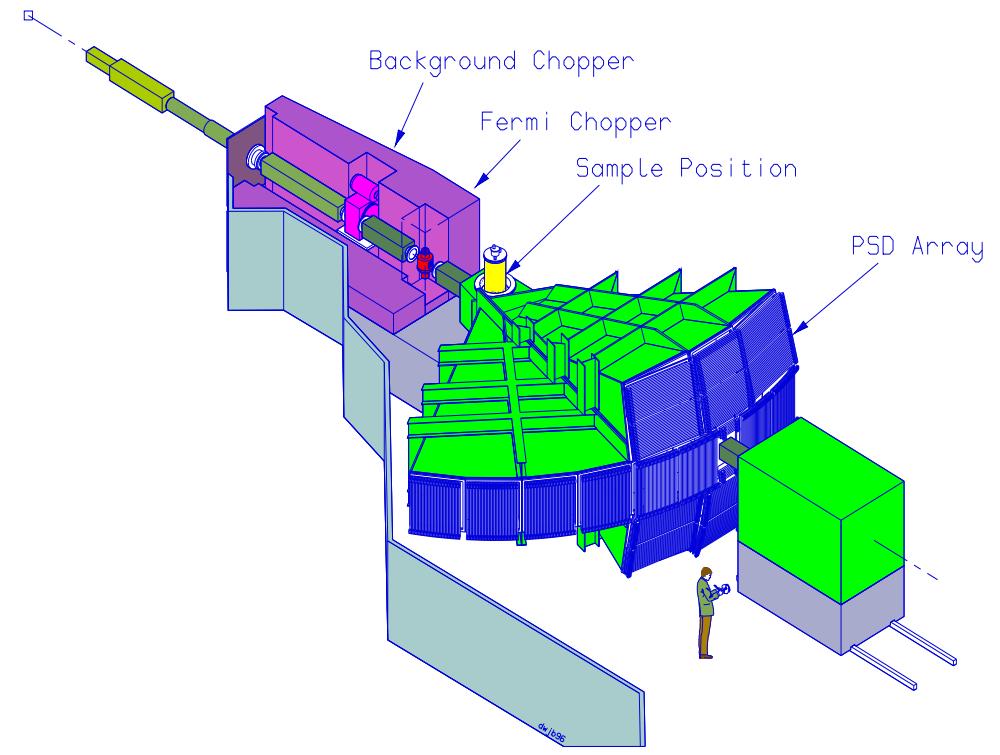
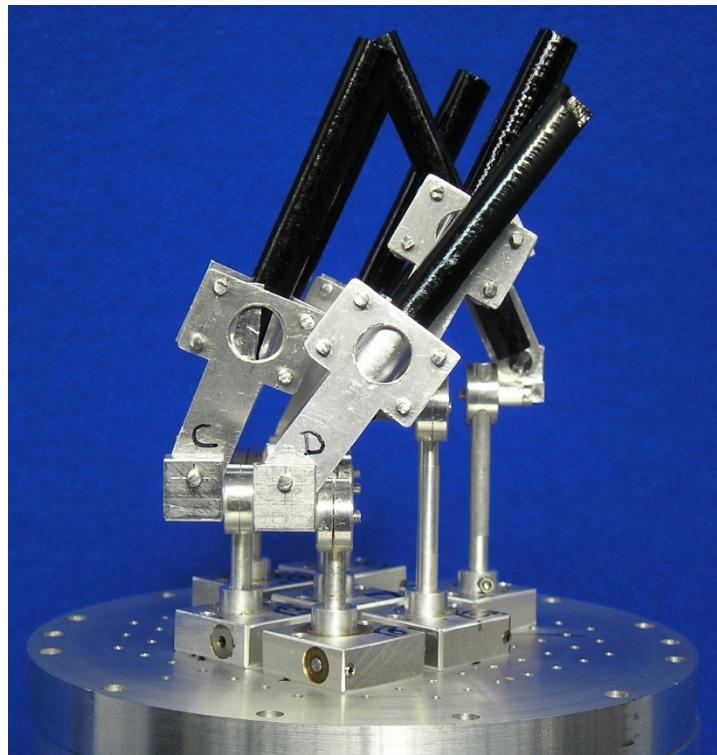
# Spin fluctuations in $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$



# Inelastic neutron scattering

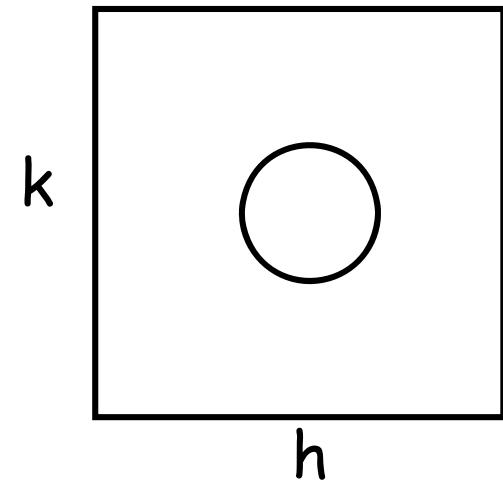
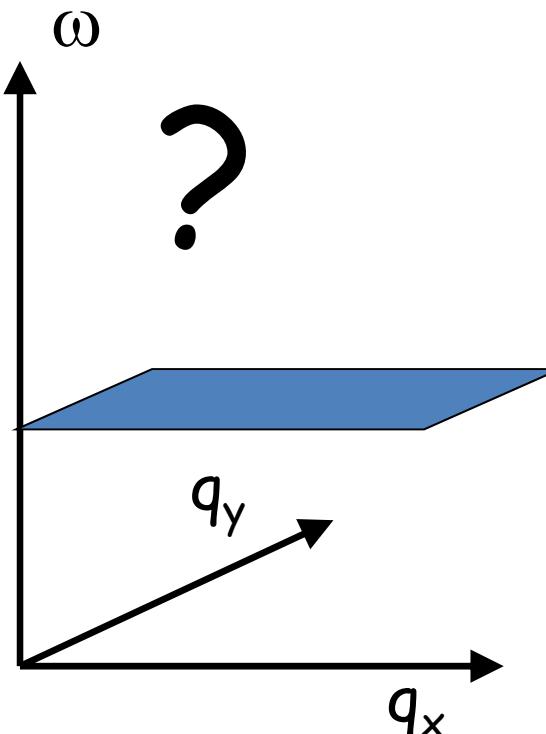
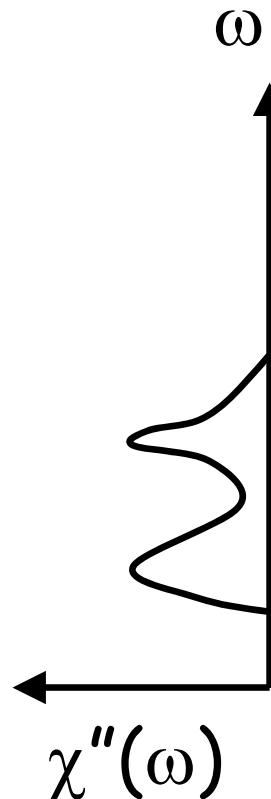
$$\frac{d^2\sigma}{d\Omega \, dE} = (\gamma r_e)^2 \frac{k_f}{k_i} |F(\mathbf{Q})|^2 \left( \frac{2/\pi g^2 \mu_B^2}{1 - \exp(-\hbar\omega/kT)} \right) \chi''(\mathbf{Q}, \omega).$$

$$\chi''_{\text{local}}(\omega) = \frac{\int_{BZ} \chi''(\mathbf{Q}, \omega) \, d\mathbf{Q}}{\int_{BZ} d\mathbf{Q}}$$



# The response function $\chi''(q,\omega)$

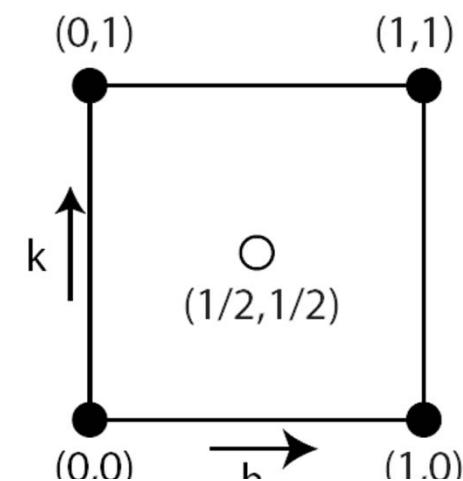
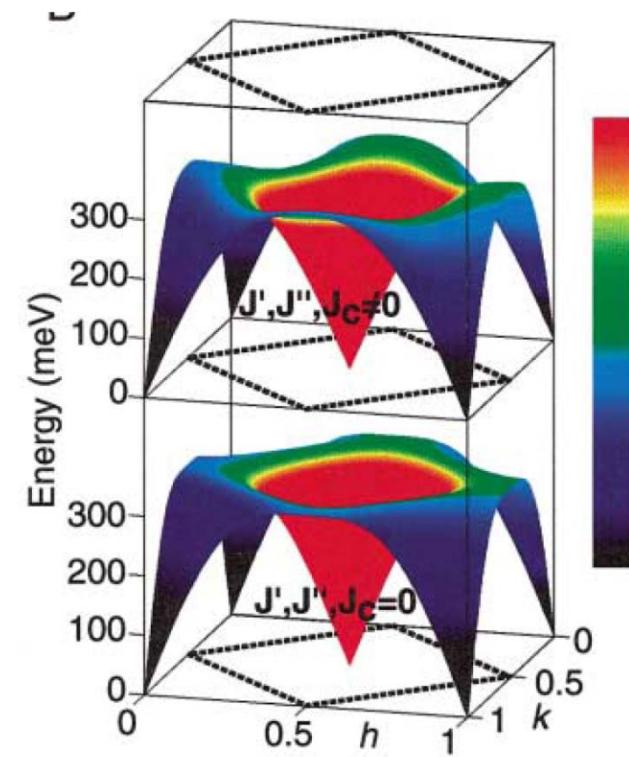
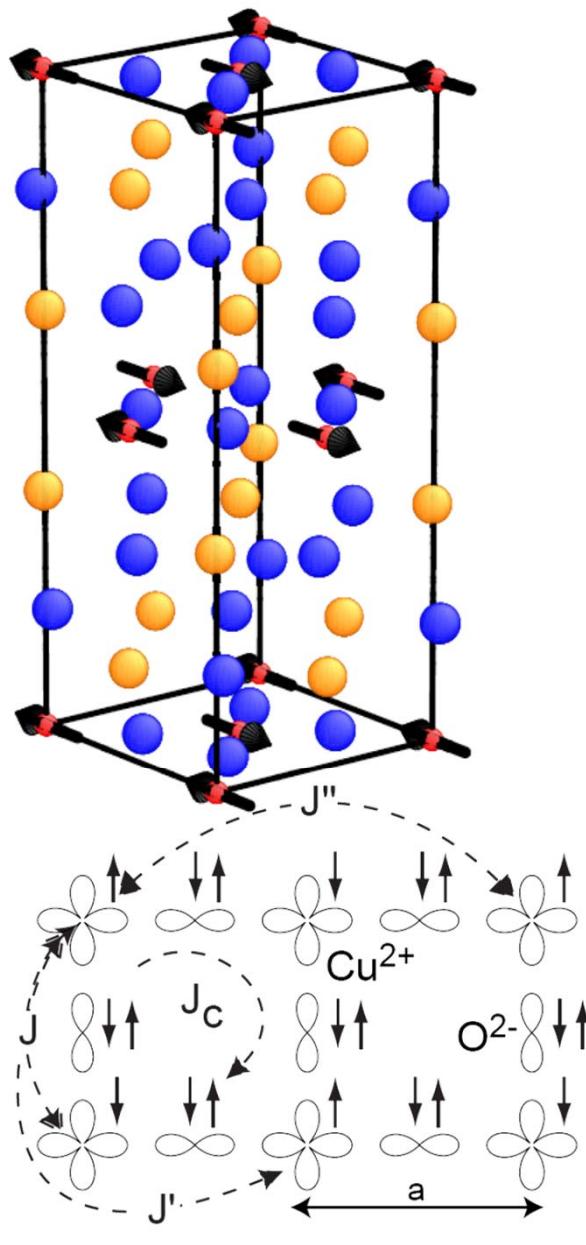
things we would like to know .....



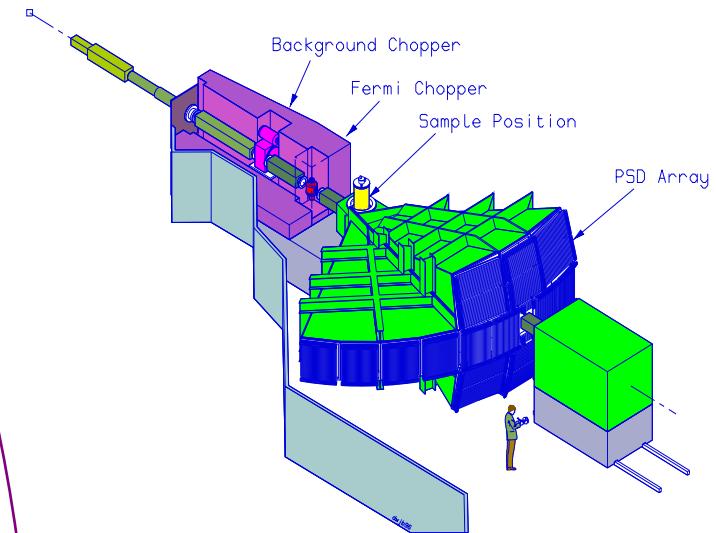
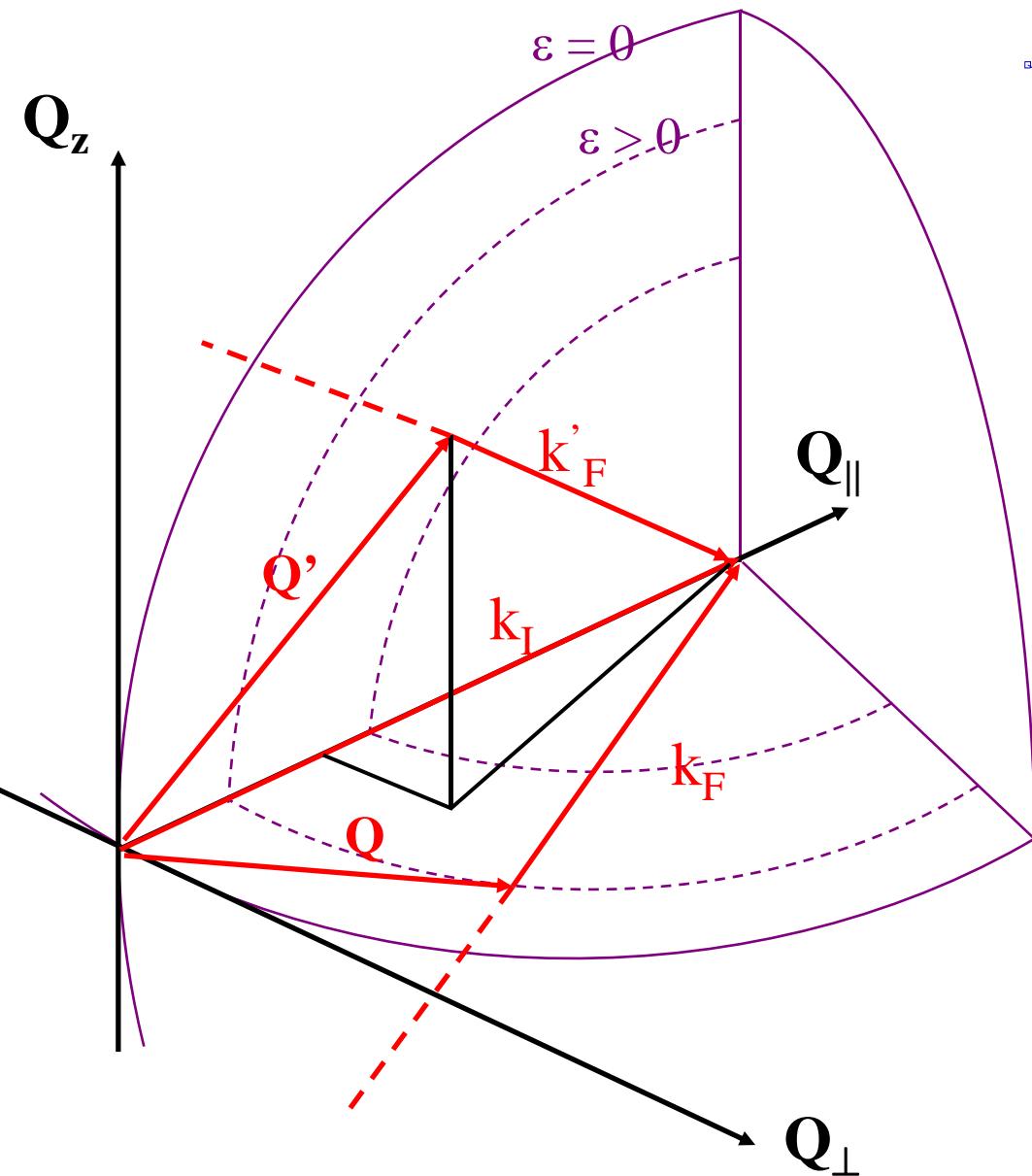
# Magnetic excitations in cuprates

- I. The pairing glue
- II. Neutron scattering in  $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$
- III. Parent antiferromagnet
- IV. Optimal doping
- V. Overdoped
- VI. Summary

# Magnetic Excitations in $\text{La}_2\text{CuO}_4$

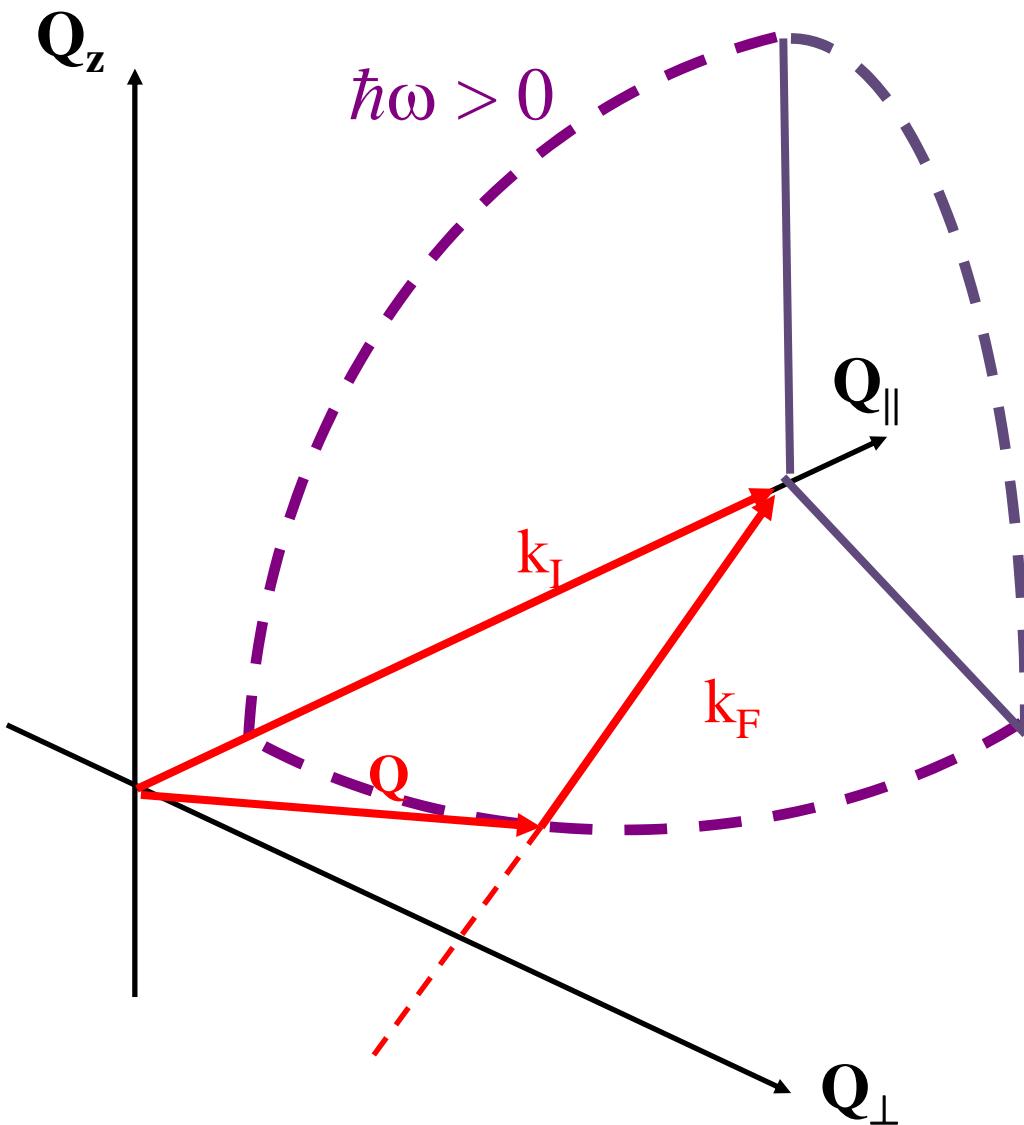


# MAPS with 2-D system



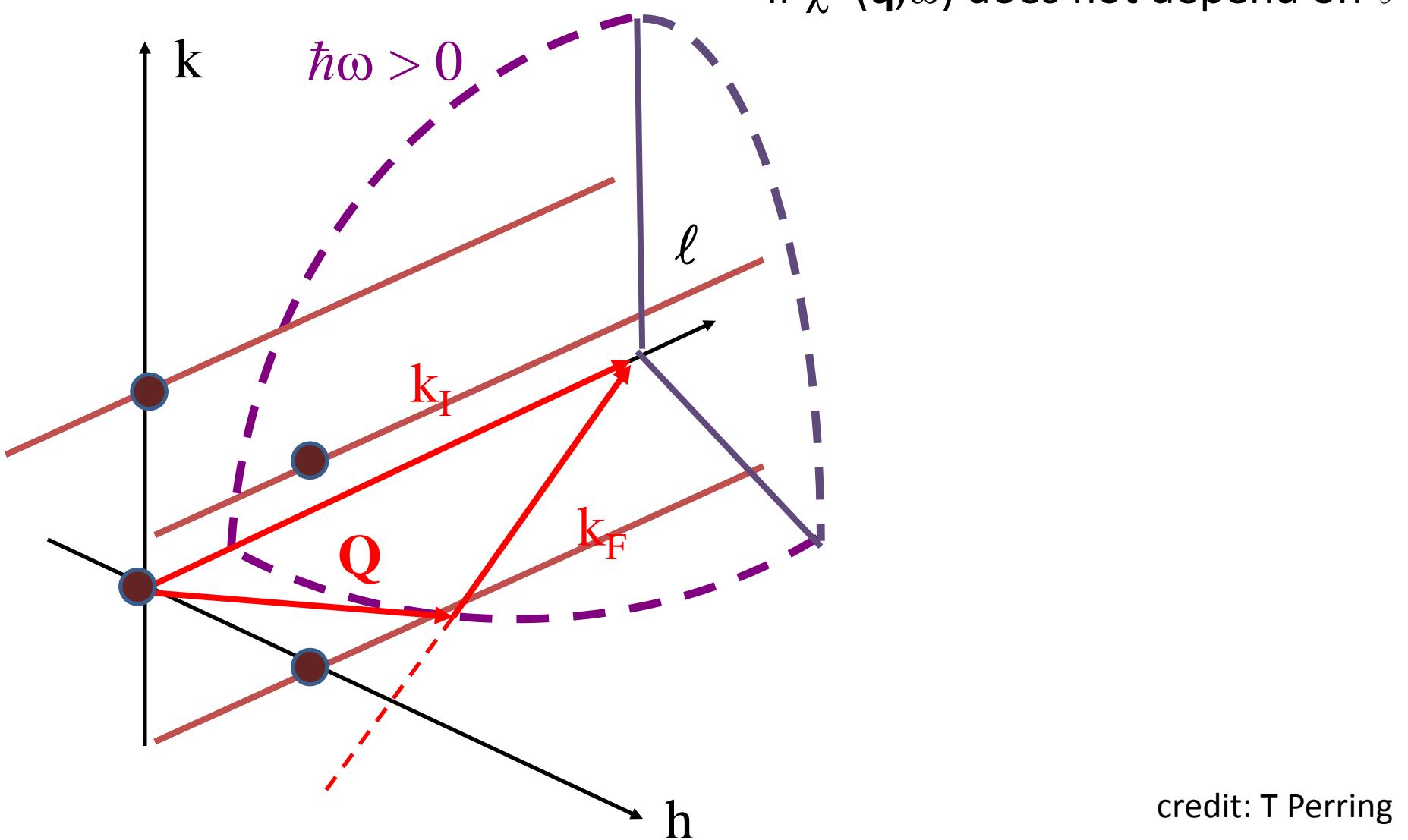
credit: T Perring

# MAPS with 2-D system

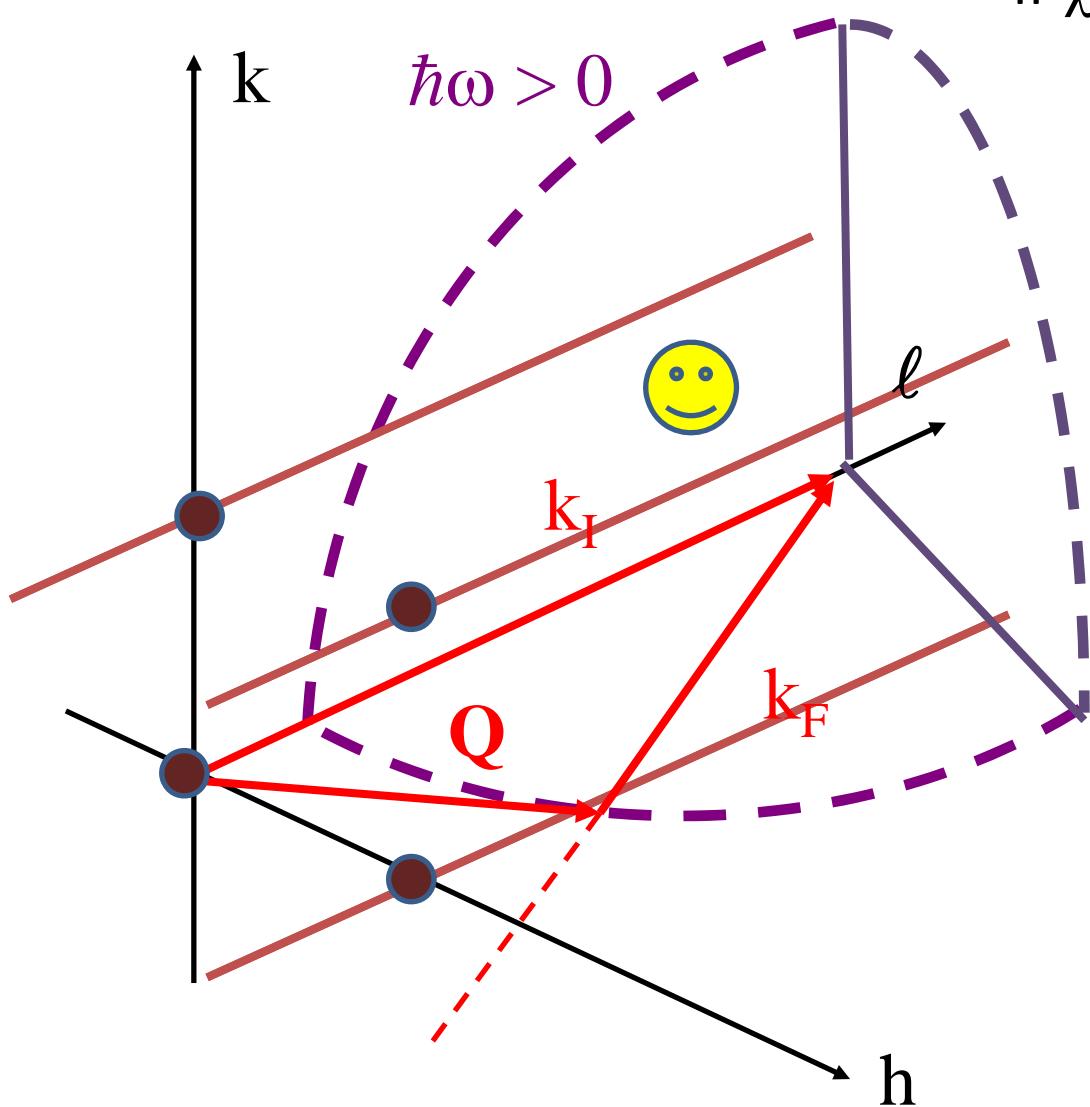


credit: T Perring

# MAPS with 2-D system

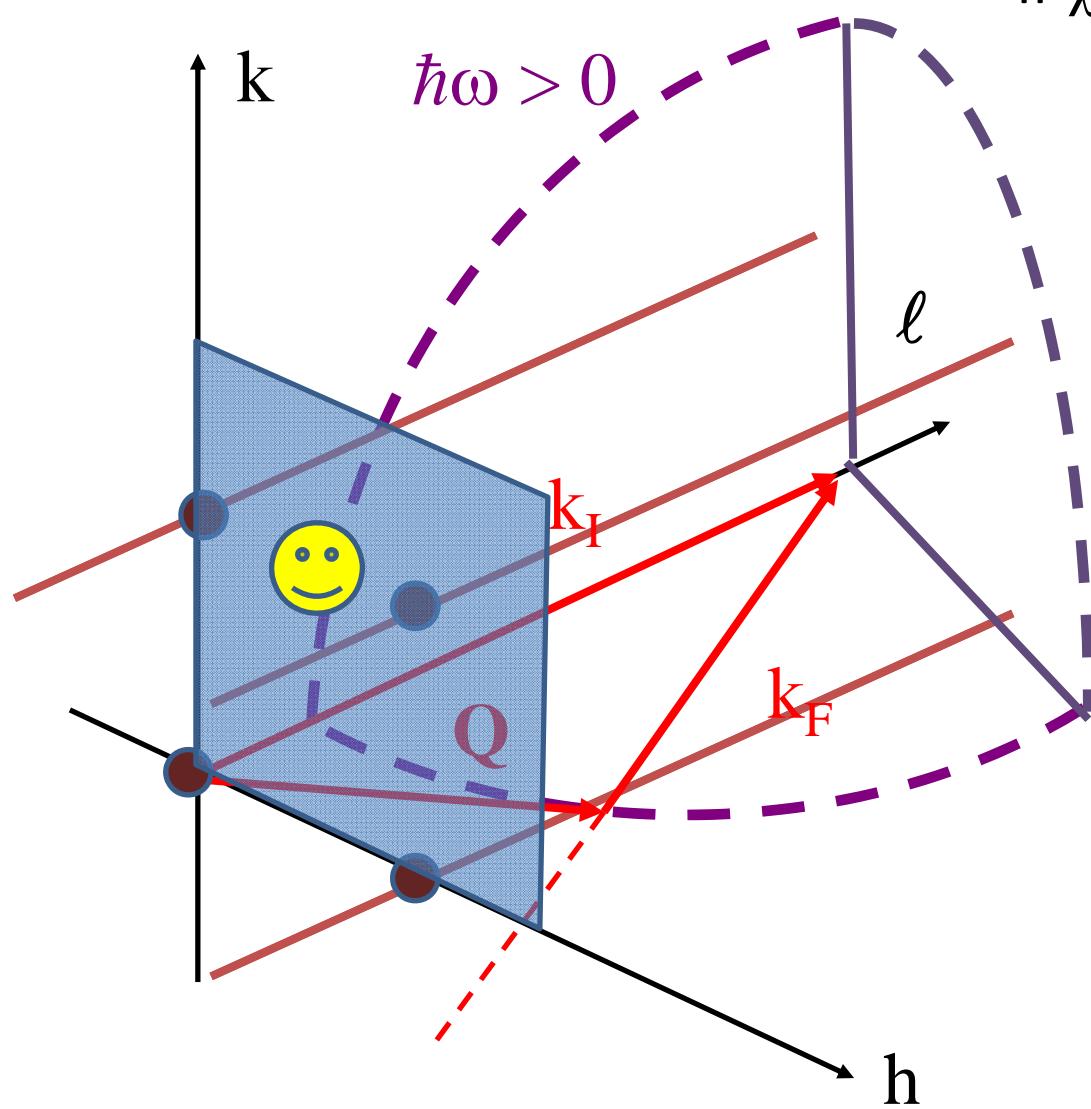


# MAPS with 2-D system

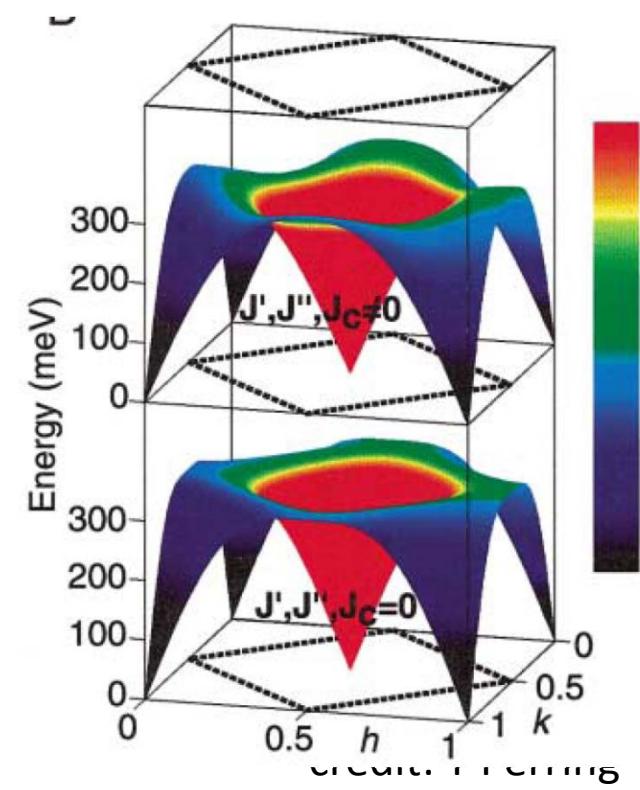


credit: T Perring

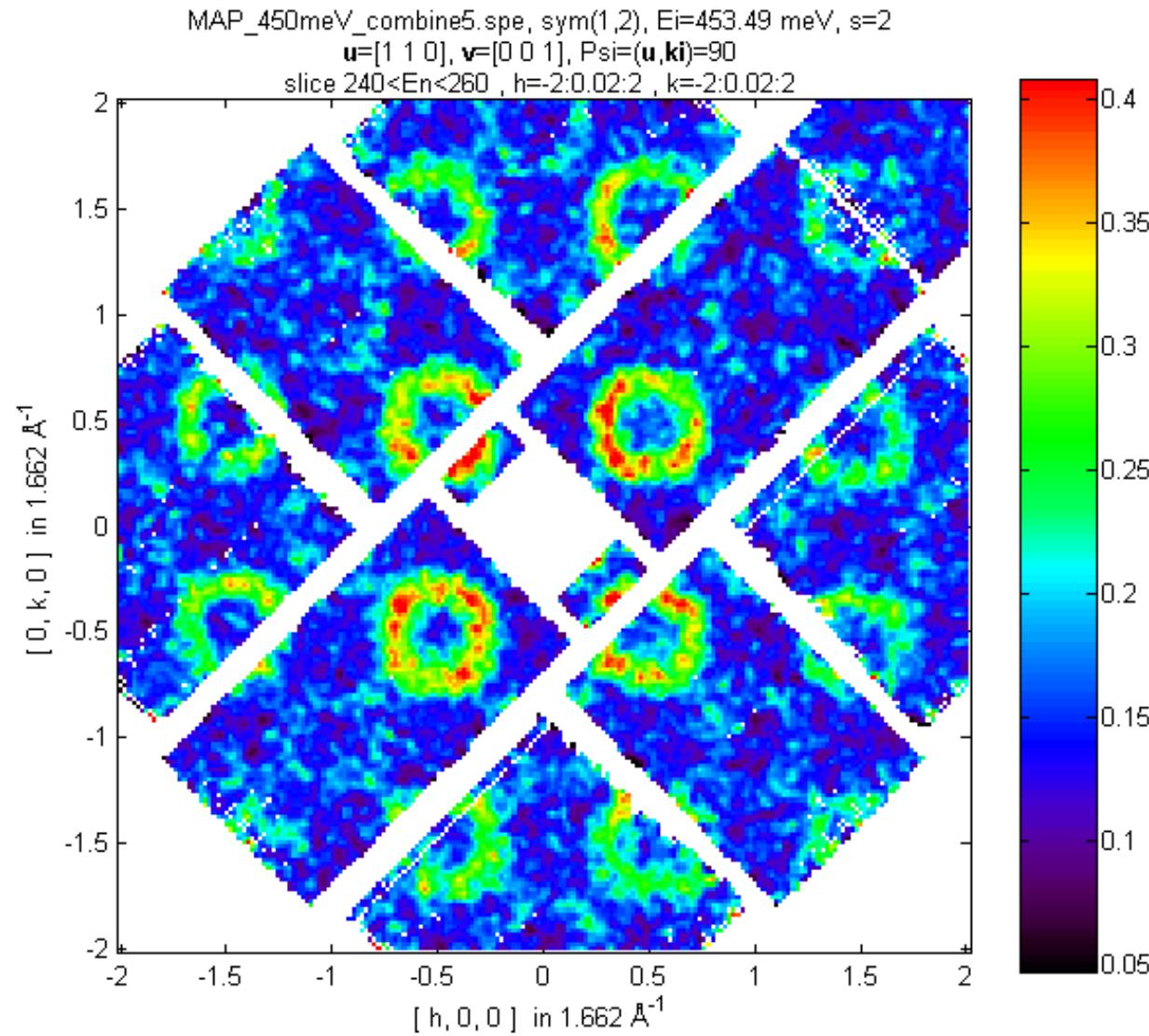
# MAPS with 2-D system



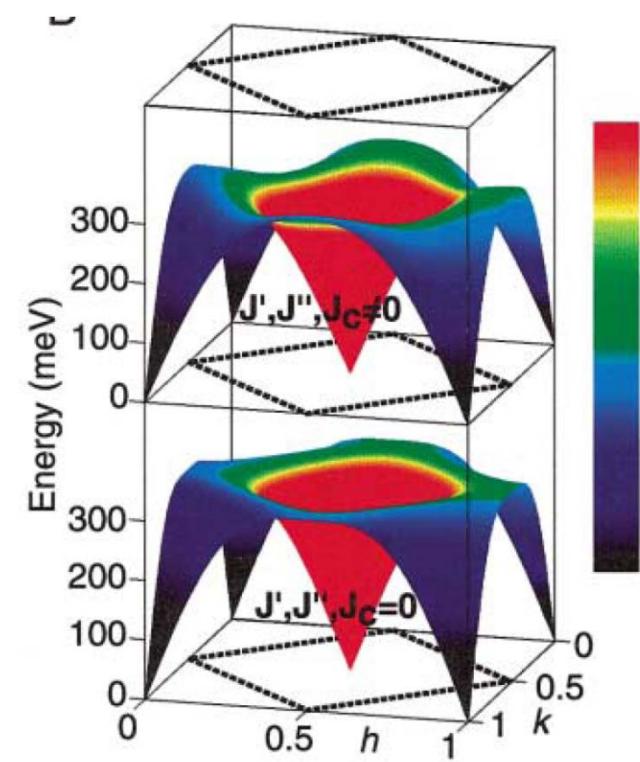
if  $\chi''(\mathbf{q}, \omega)$  does not depend on  $\ell$



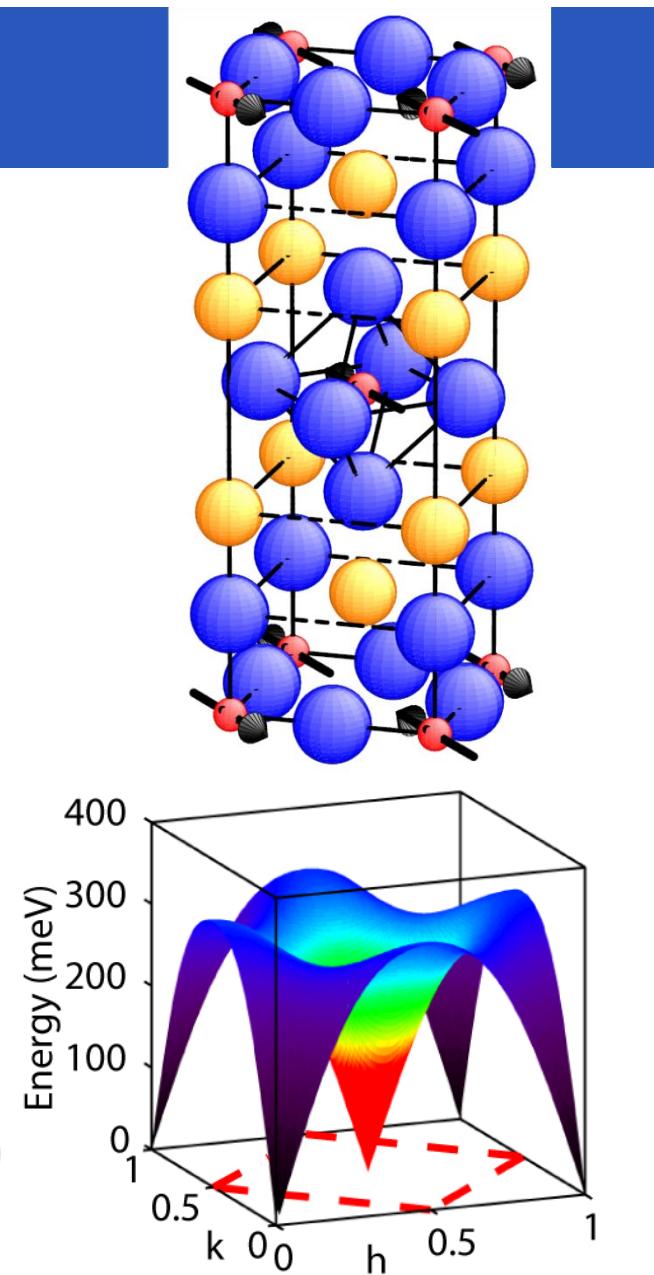
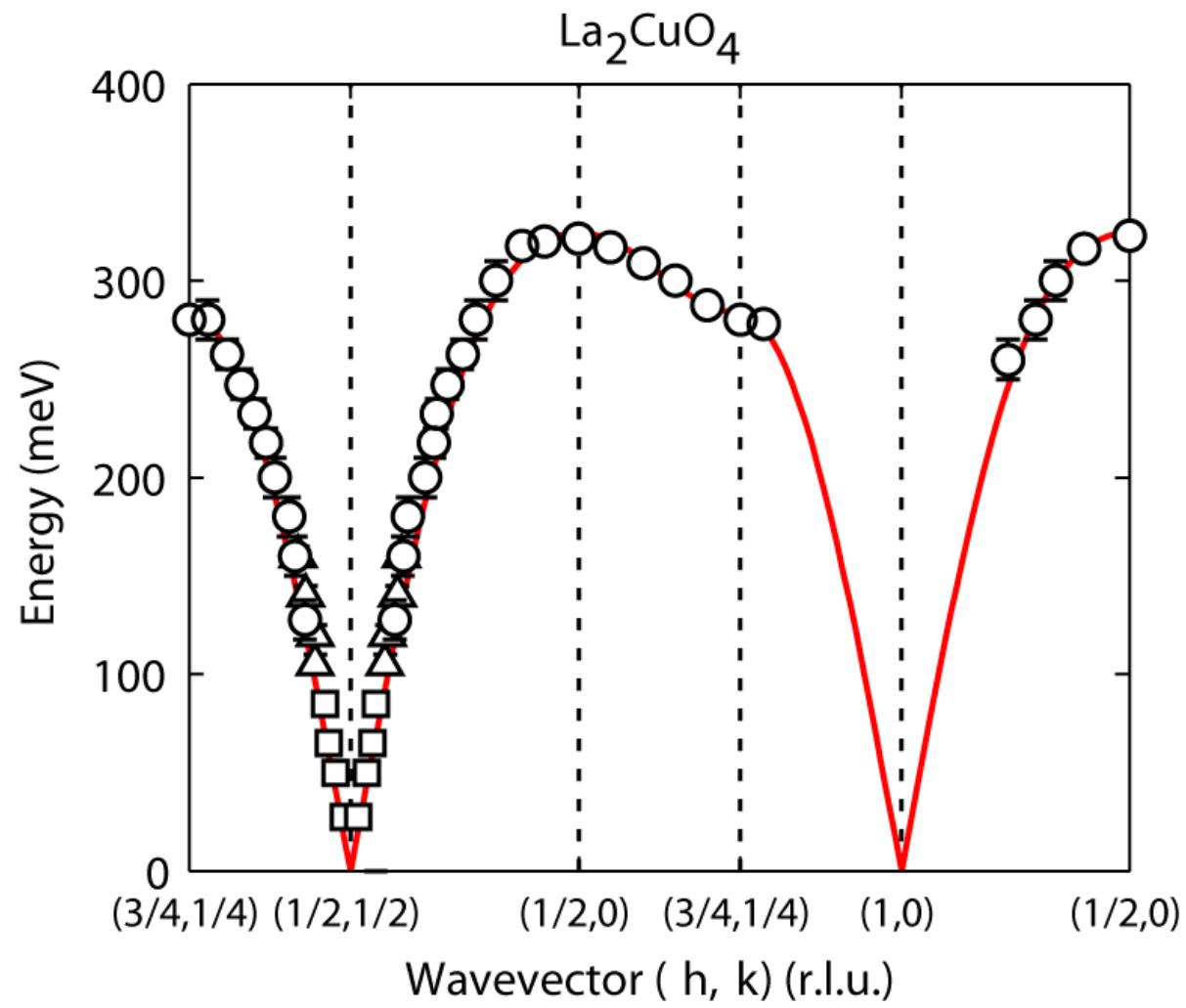
# Spin waves in $\text{La}_2\text{CuO}_4$ :MAPS



E=250 meV



# Spin waves in $\text{La}_2\text{CuO}_4$



Coldea, Hayden, Aeppli, Perring, Frost, Mason, Cheong, Fisk, PRL **86** 5377 (2001).

Headings, Hayden, Coldea, Perring, Phys. Rev. Lett. **105**, 247001 (2010)

# Spin wave dispersion

SW dispersion along BZ (3/4,1/4) to (1/2,0) subject to two competing effects):

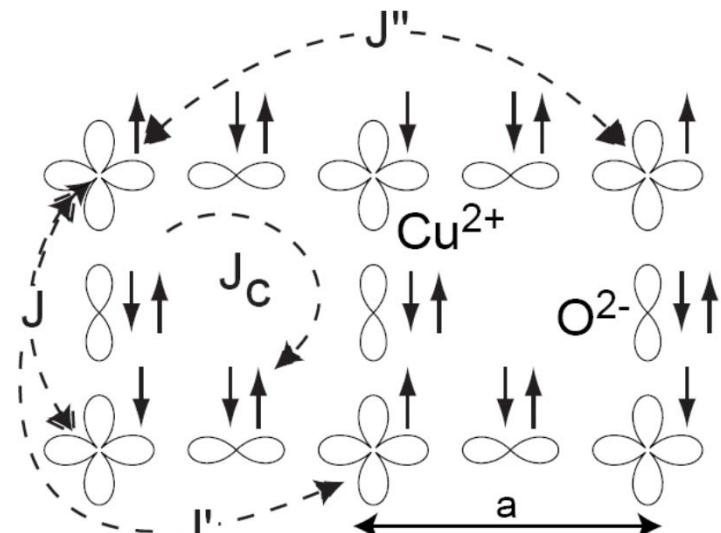
- 1) Quantum fluctuations (beyond SWT) raise (3/4,1/4) wrt (1/2,0)
- 2) for large  $t/U$  the opposite happens due to higher order exchange terms.

$$J = 4t^2/U - 24t^4/U^3$$

$$J_c = 80 t^4/U^3$$

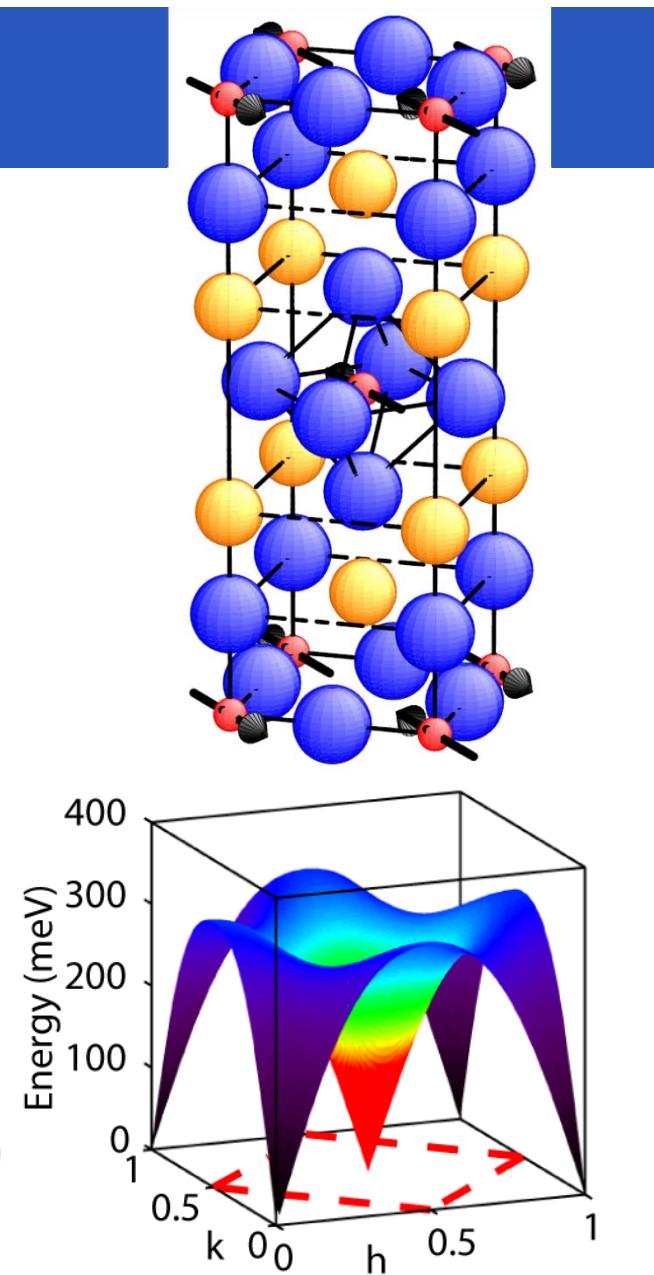
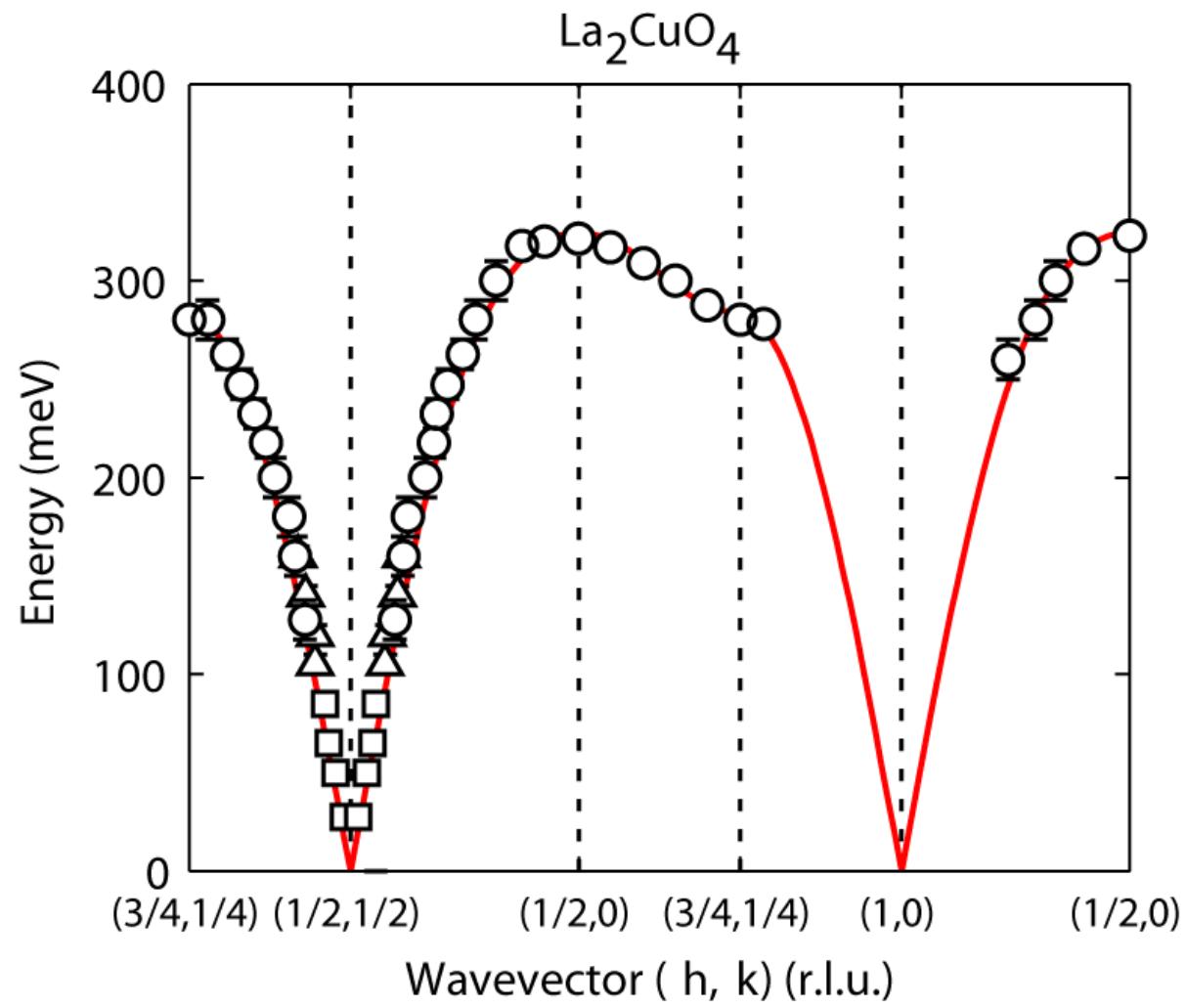
$$J' = J'' = 4t^4/U^3$$

$$t = 0.30 \pm 0.02 \text{ eV}, U = 2.2 \pm 0.4 \text{ eV}, \\ J = 146 \pm 4 \text{ meV}, \text{ and } J_c = 61 \pm 8 \text{ meV}.$$



A. H. MacDonald PRB (1990)

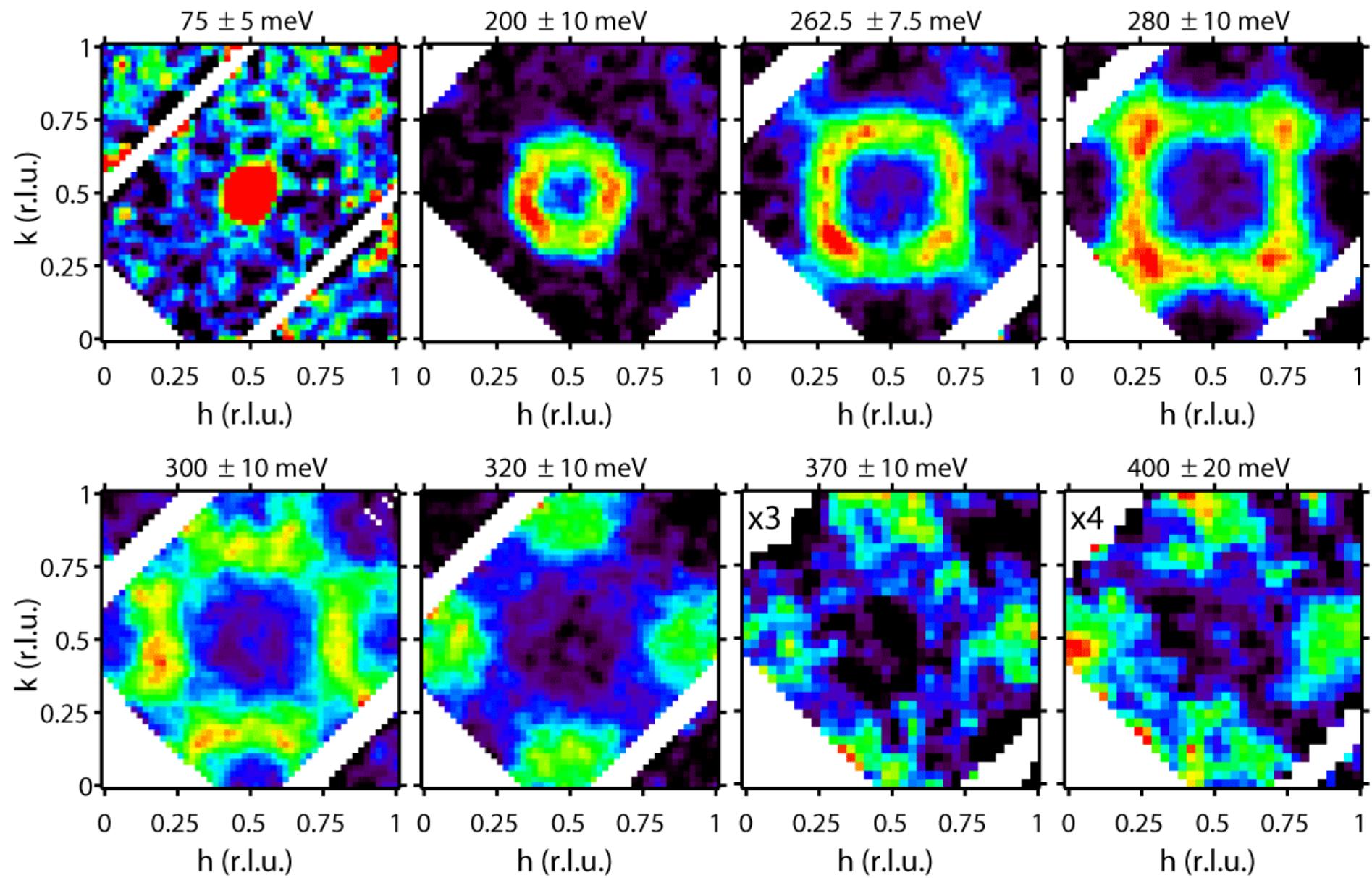
# Spin waves in $\text{La}_2\text{CuO}_4$



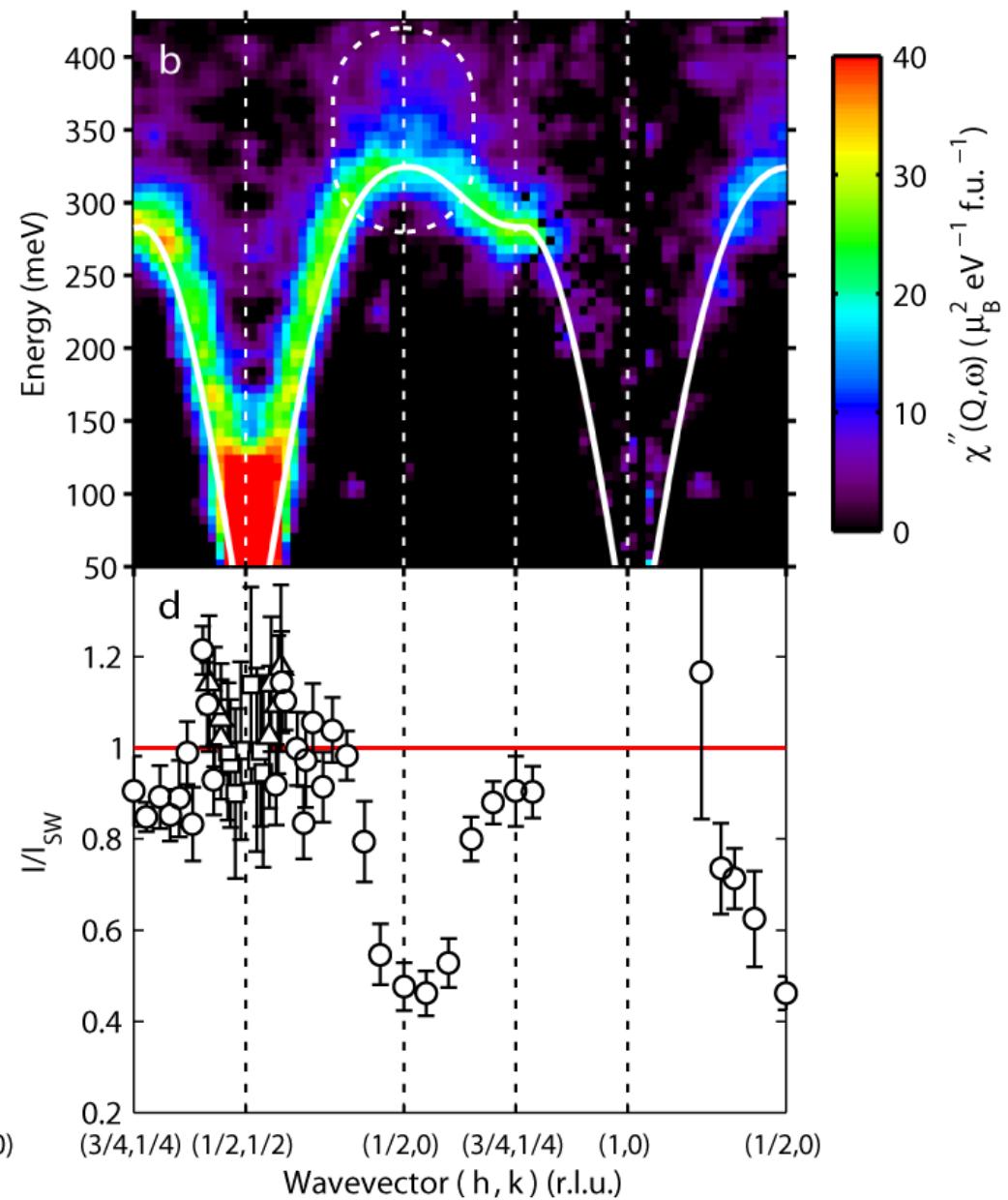
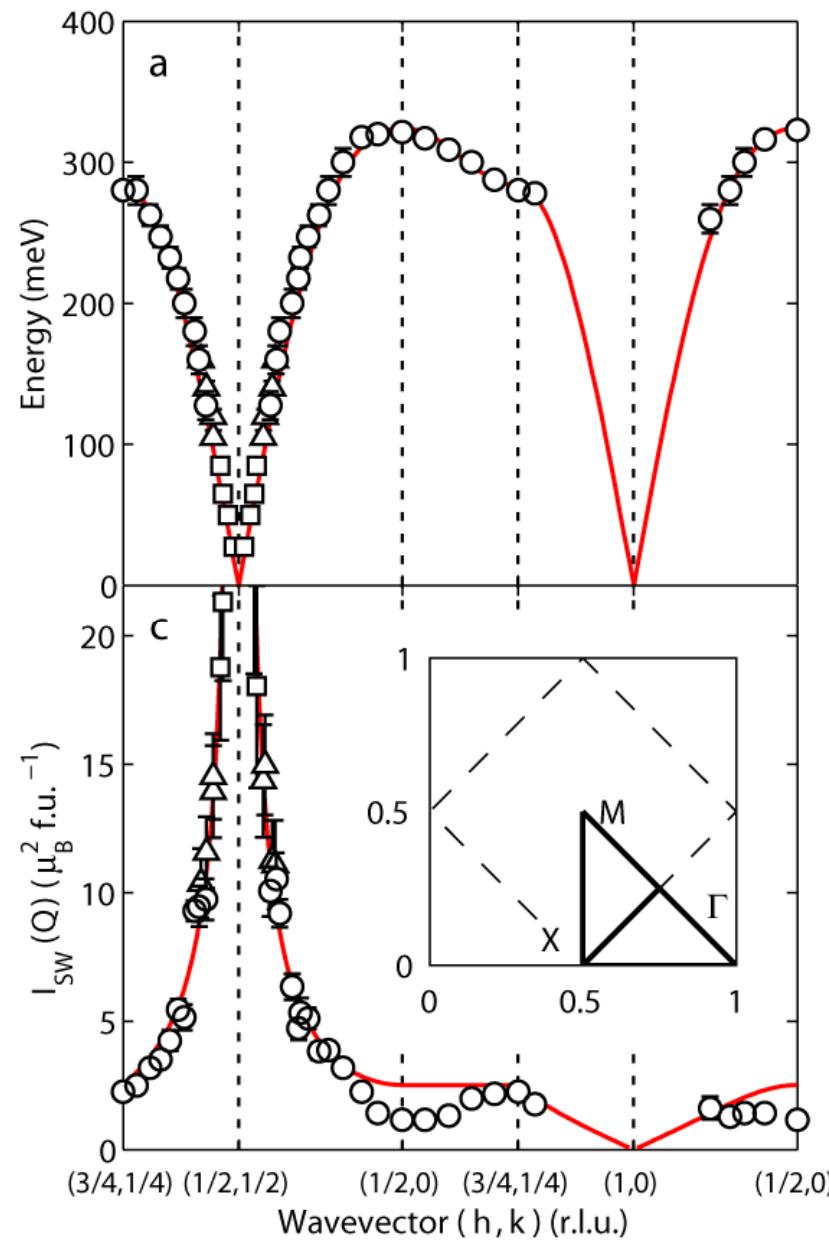
Coldea, Hayden, Aeppli, Perring, Frost, Mason, Cheong, Fisk, PRL **86** 5377 (2001).

Headings, Hayden, Coldea, Perring, Phys. Rev. Lett. **105**, 247001 (2010)

# $\text{La}_2\text{CuO}_4$ : $\mathbf{q}$ -dependent slices



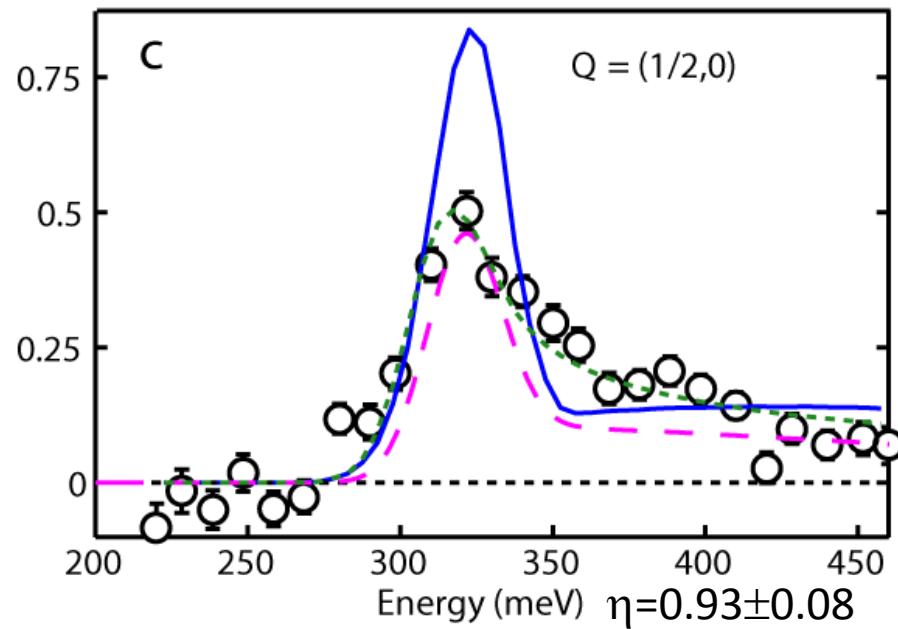
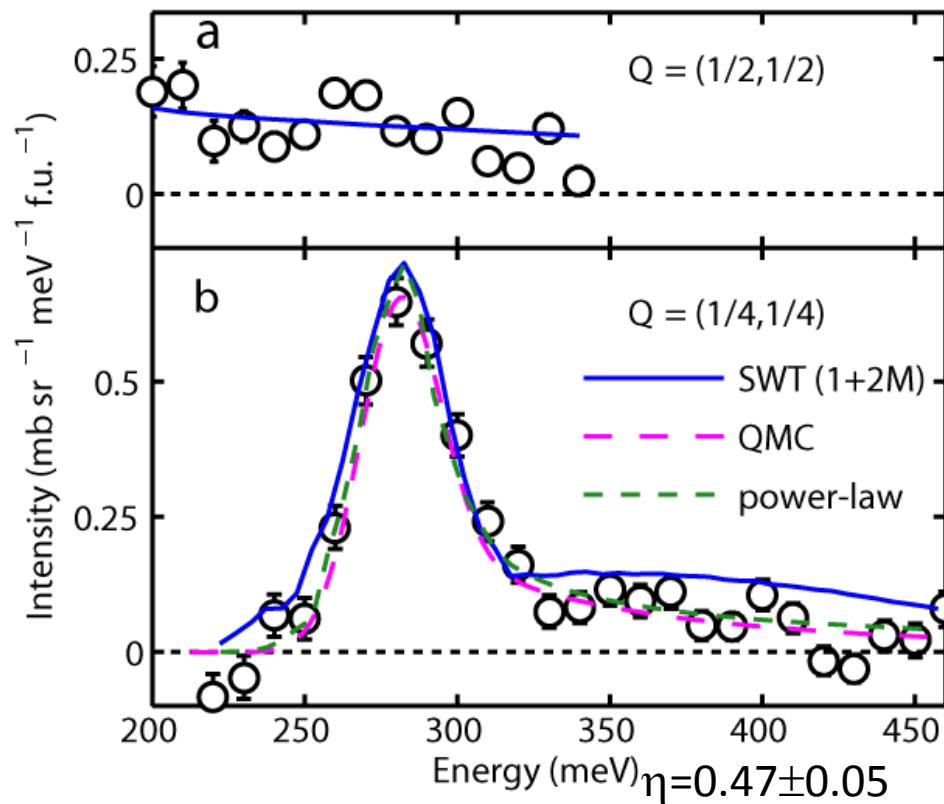
# Magnetic Excitations in $\text{La}_2\text{CuO}_4$



# “Spinon” power law

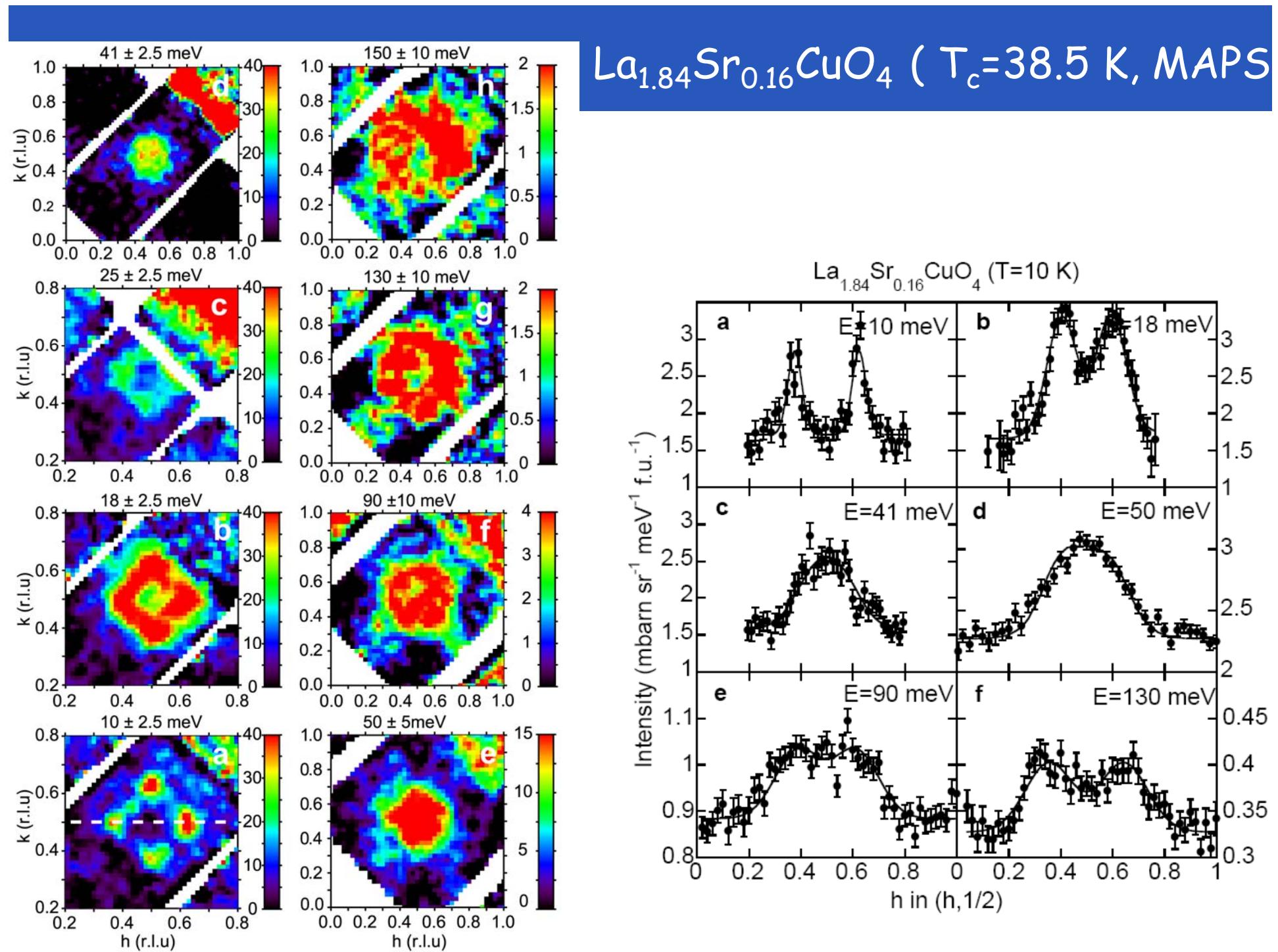
Generalization of continuum scattering lineshape of the 1D Heisenberg AF chain ( $\eta=1$ )

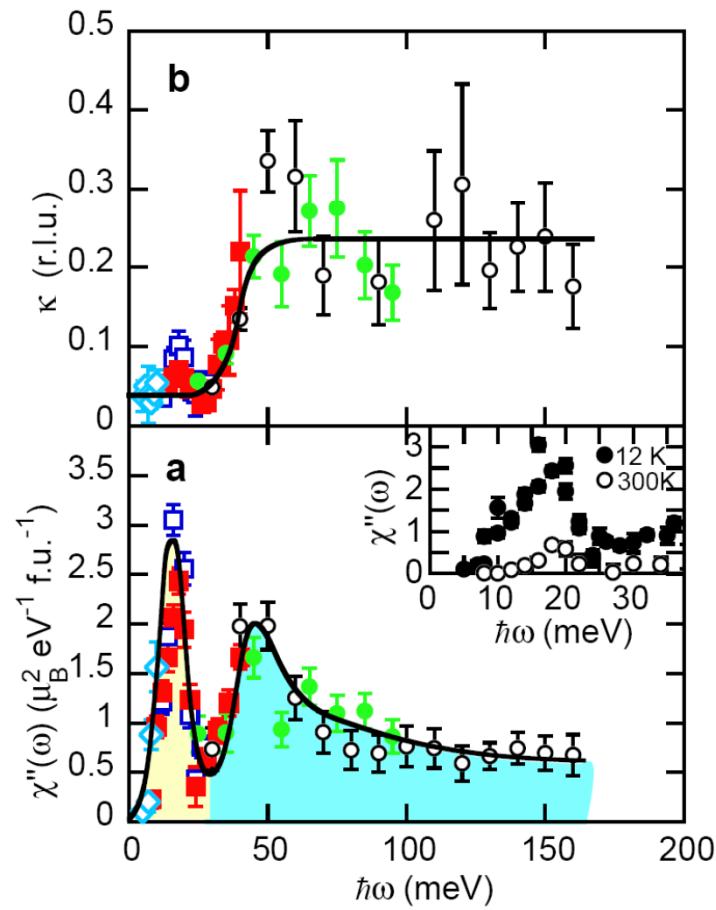
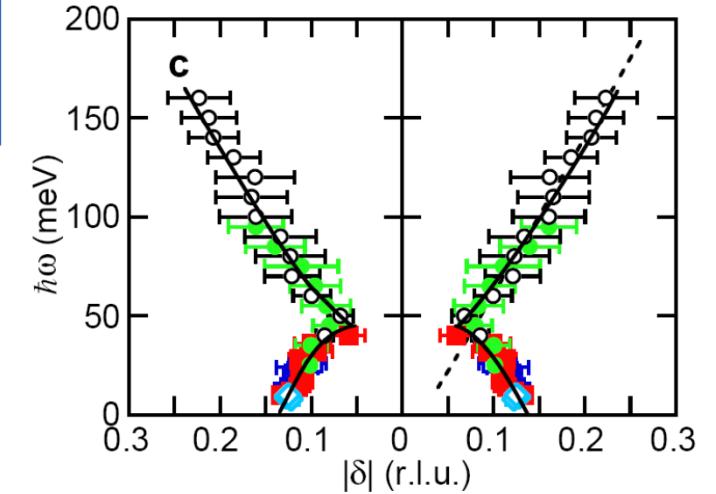
$$\chi''(\mathbf{q}, \omega) = A_{\mathbf{q}} \frac{\theta(\omega - \omega_{\mathbf{q}})}{(\omega^2 - \omega_{\mathbf{q}}^2)^{1-\eta/2}}$$



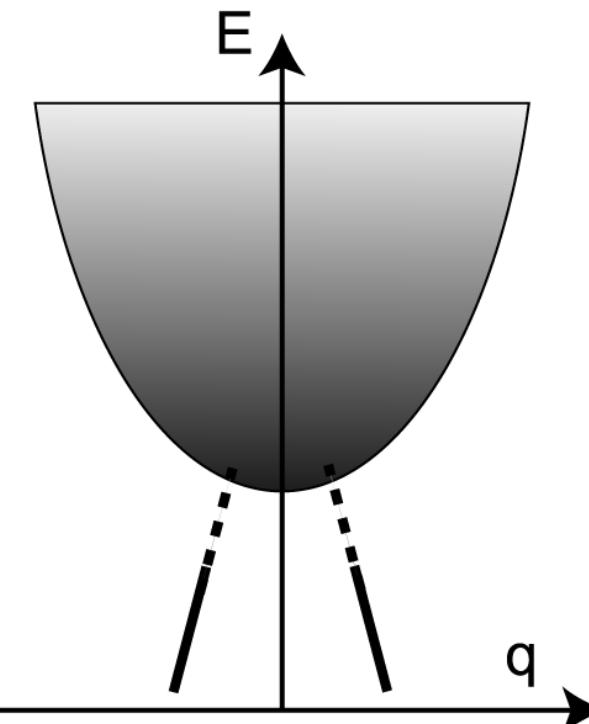
# Magnetic excitations in cuprates

- I. The pairing glue
- II. Neutron scattering in  $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$
- III. Parent antiferromagnet
- IV. Optimal doping
- V. Overdoped
- VI. Summary



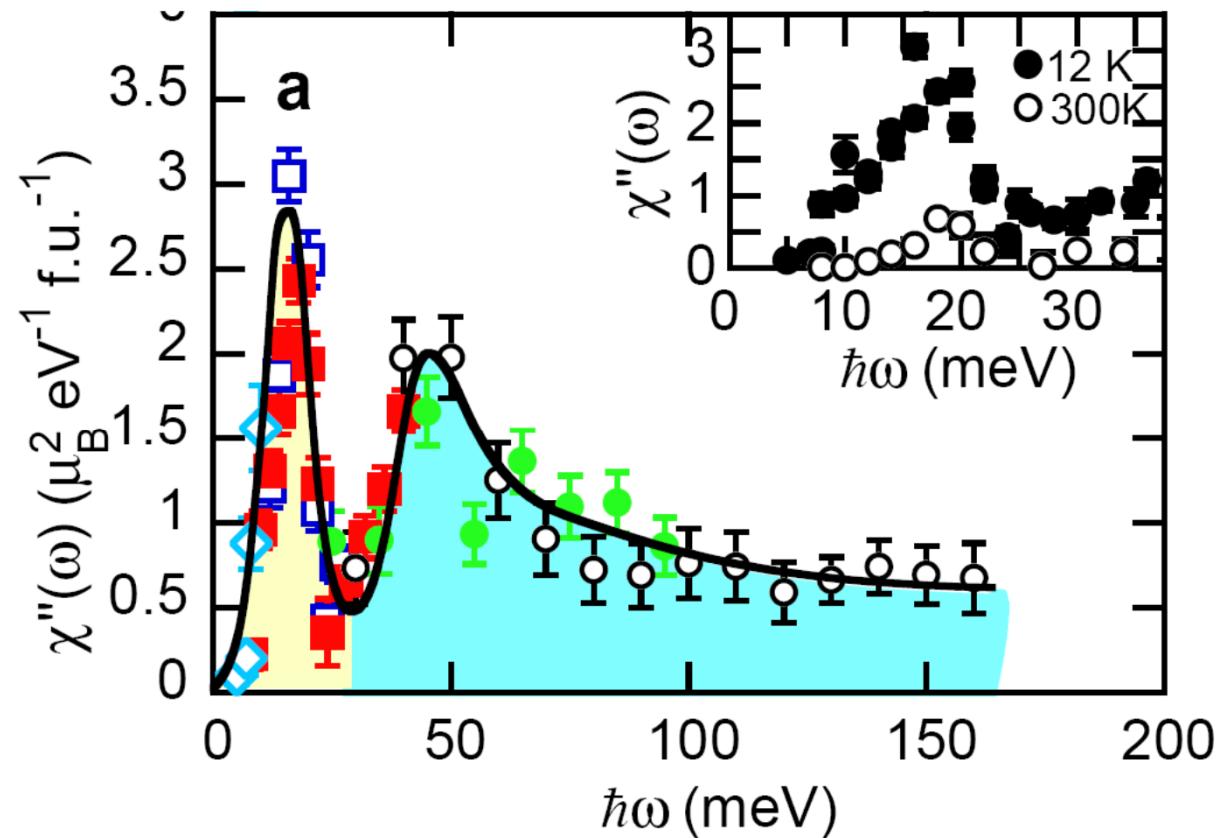


fitted parameters ( $x=0.16$ )



$$\chi''(\omega) = \int \chi''(\mathbf{Q}, \omega) d^3 Q / \int d^3 Q$$

# $\text{La}_{1.84}\text{Sr}_{0.16}\text{CuO}_4$



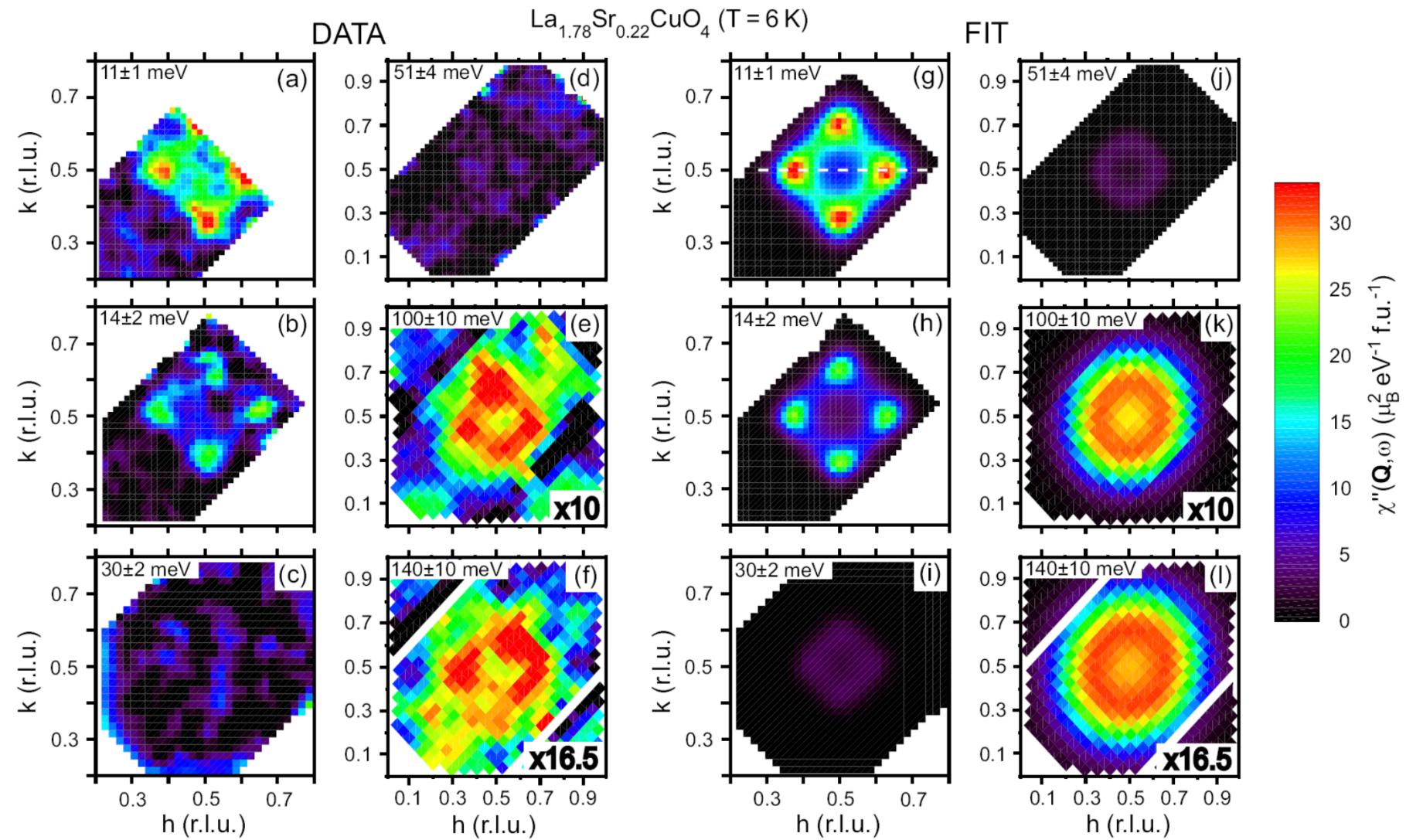
## key features

- double-peaked structure (two energy scales 18,50 meV)
- rapid broadening of pattern near 40 meV
- possible “rotation of pattern” at high energies

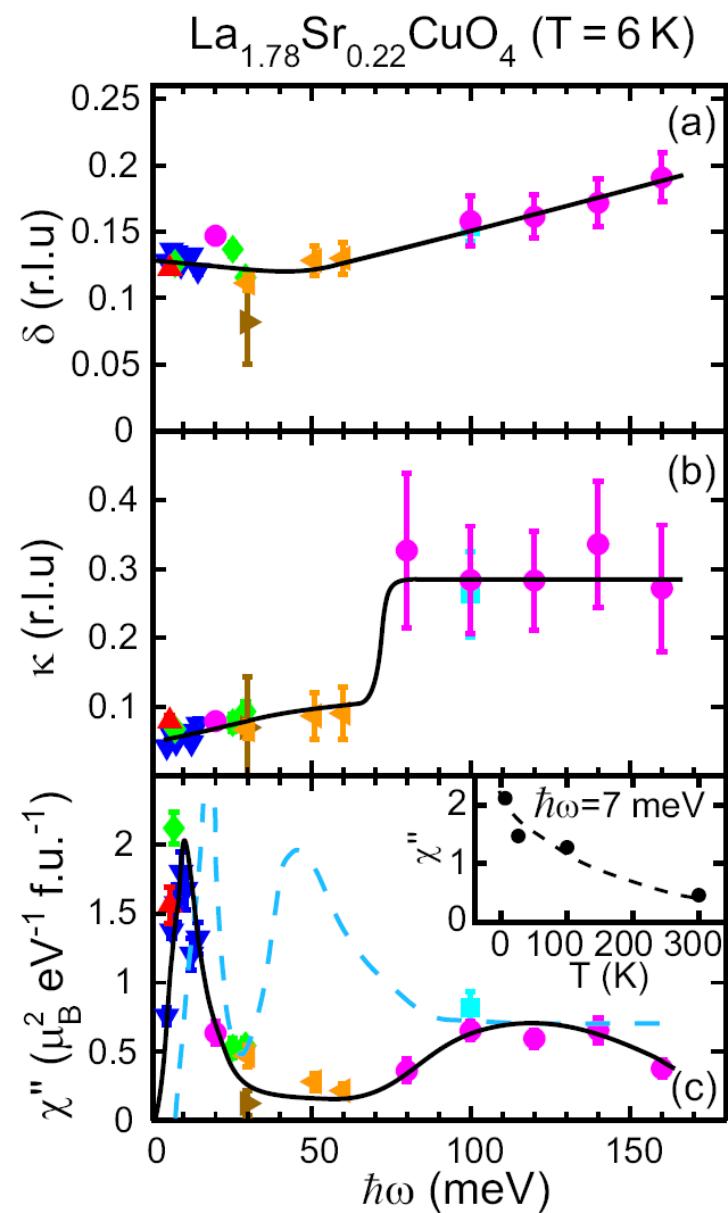
# Magnetic excitations in cuprates

- I. The pairing glue
- II. Neutron scattering in  $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$
- III. Parent antiferromagnet
- IV. Optimal doping
- V. Overdoped
- VI. Summary

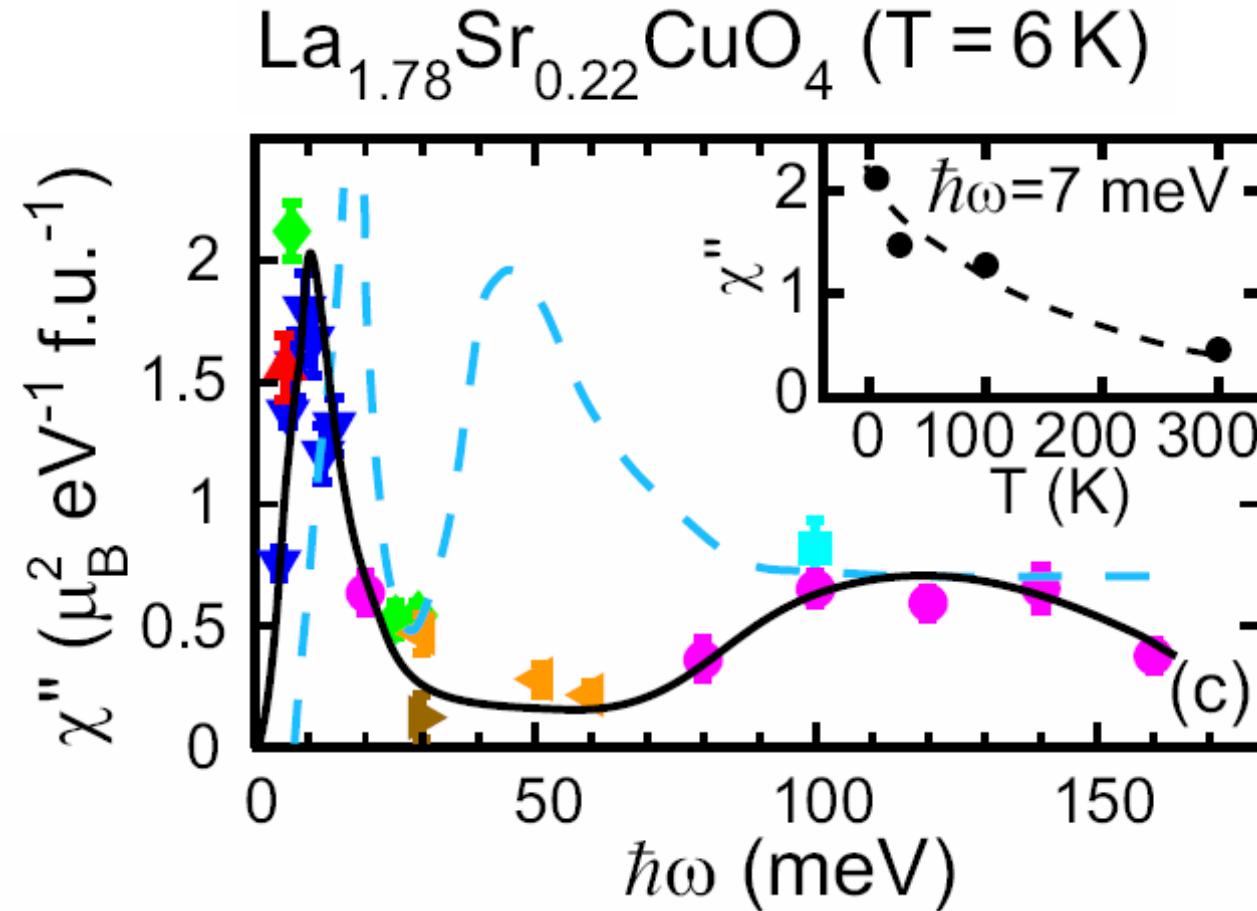
# Overdoped Sample $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$ ( $x=0.22, T_c=26\text{K}$ )



# Overdoped Sample $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$ ( $x=0.22, T_c=26\text{K}$ )



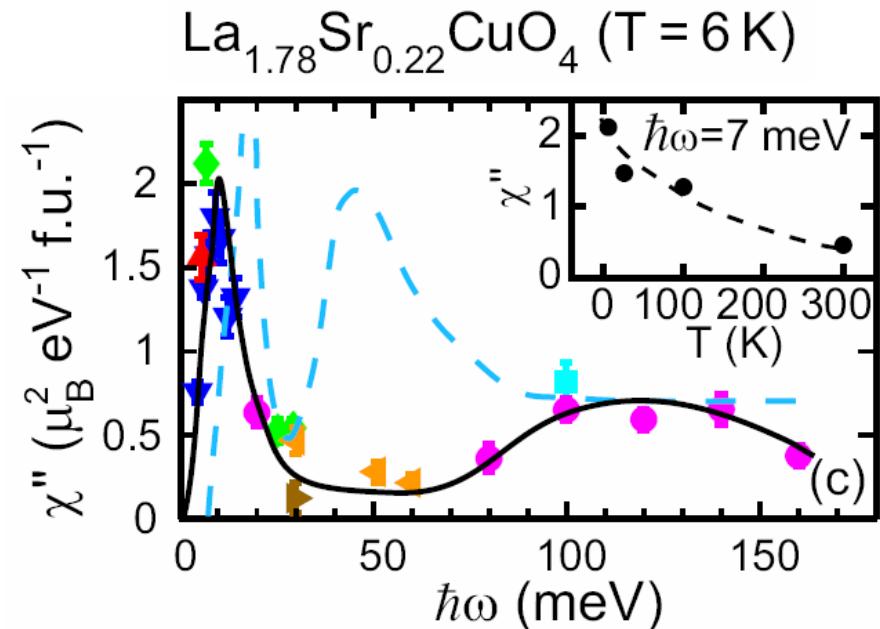
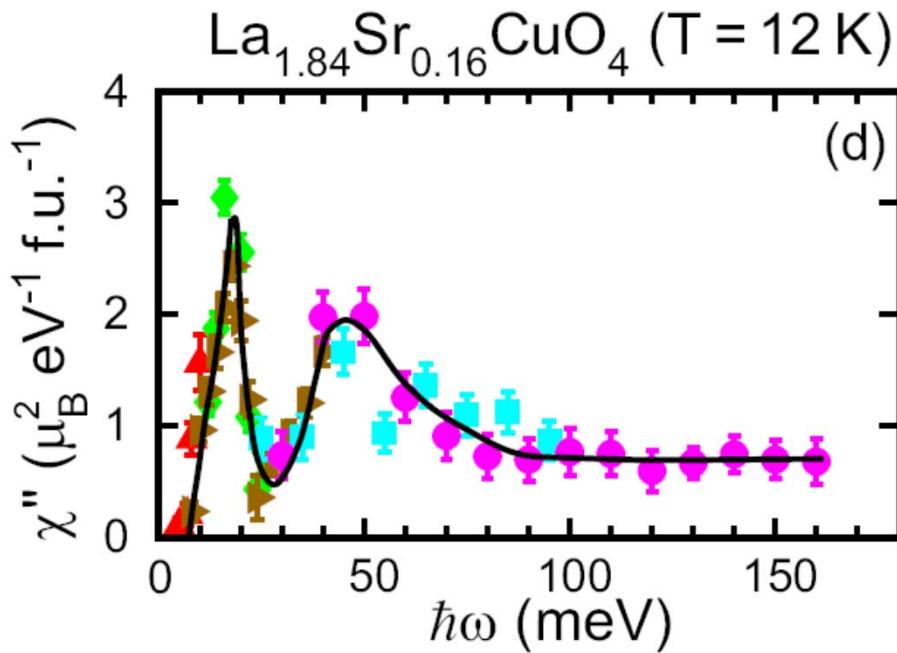
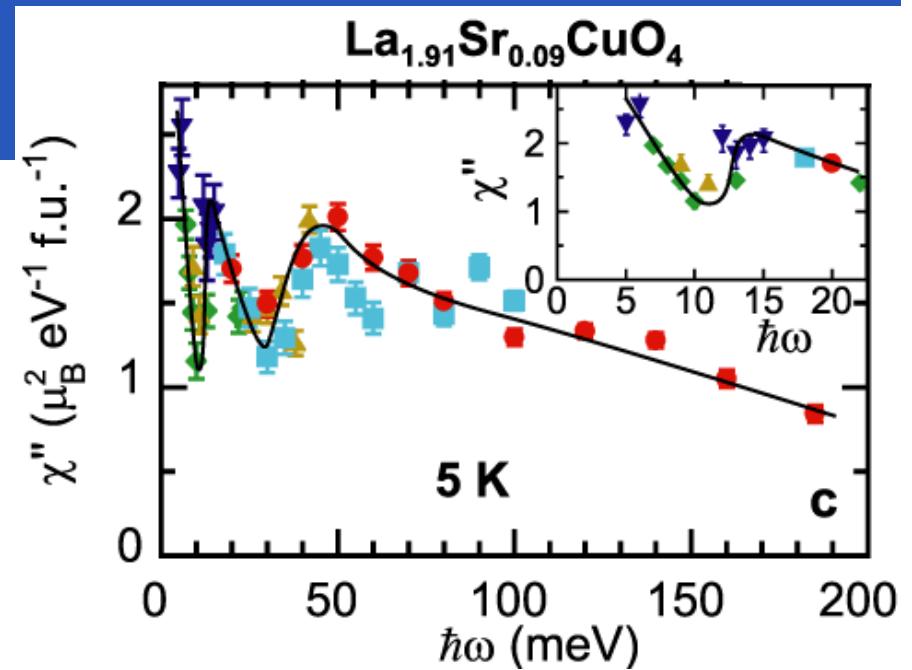
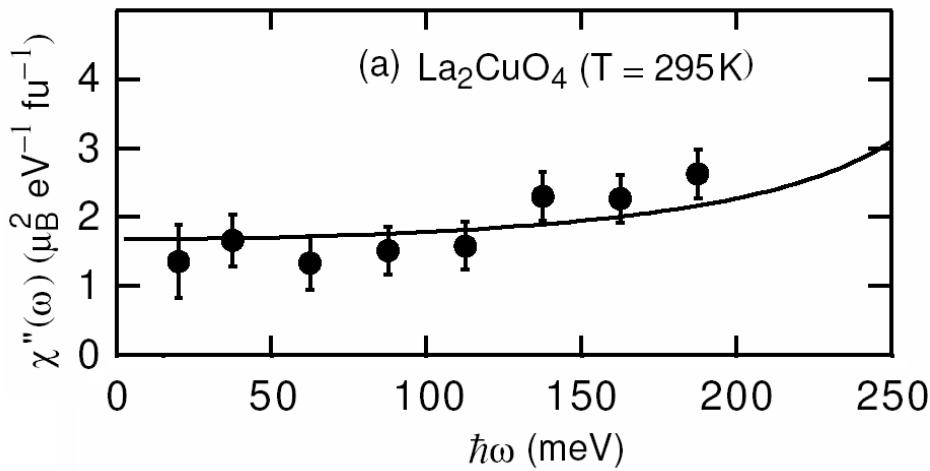
# Collapse of intermediate frequencies



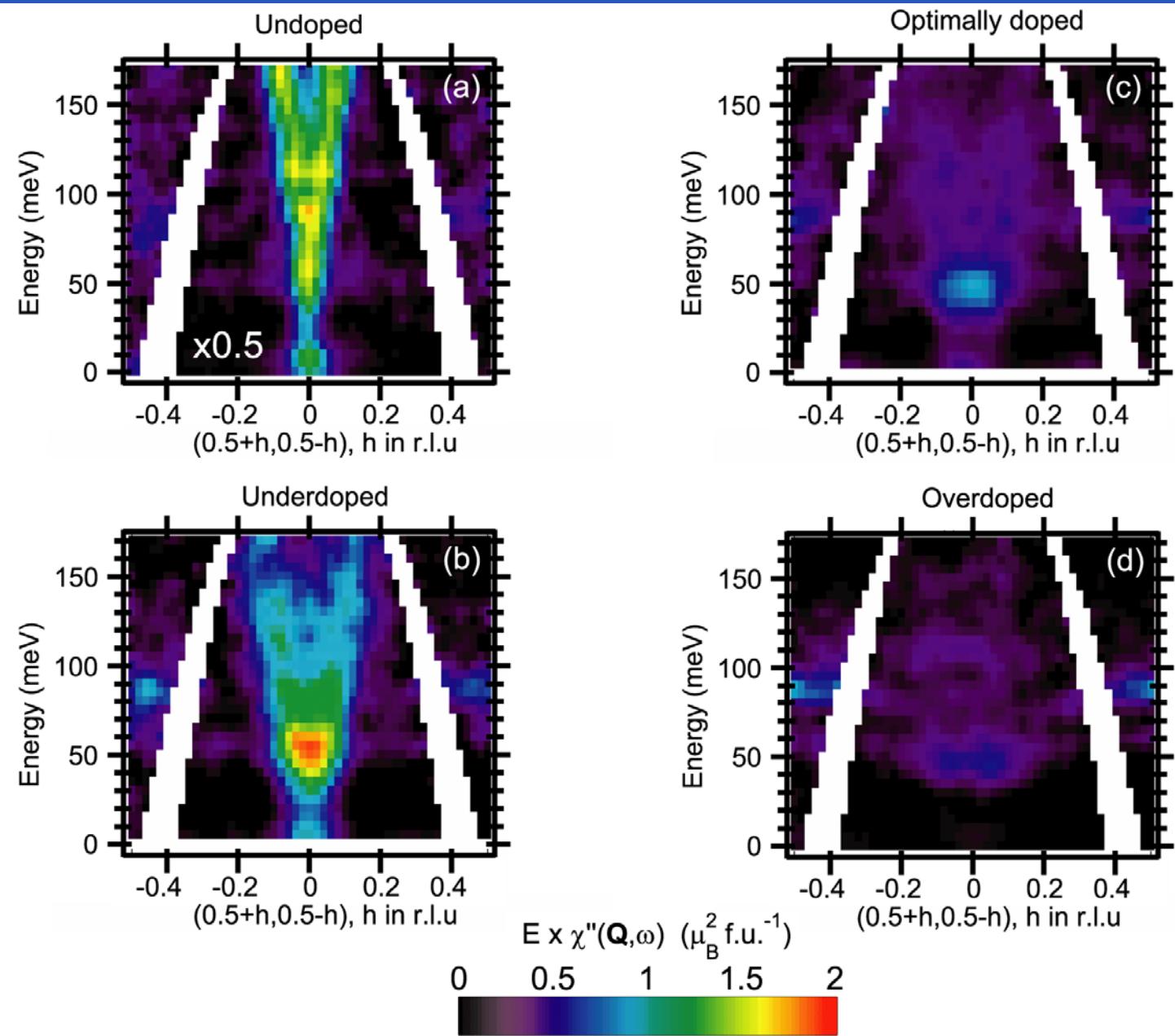
# Magnetic excitations in cuprates

- I. The pairing glue
- II. Neutron scattering in  $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$
- III. Parent antiferromagnet
- IV. Optimal doping
- V. Overdoped
- VI. Summary

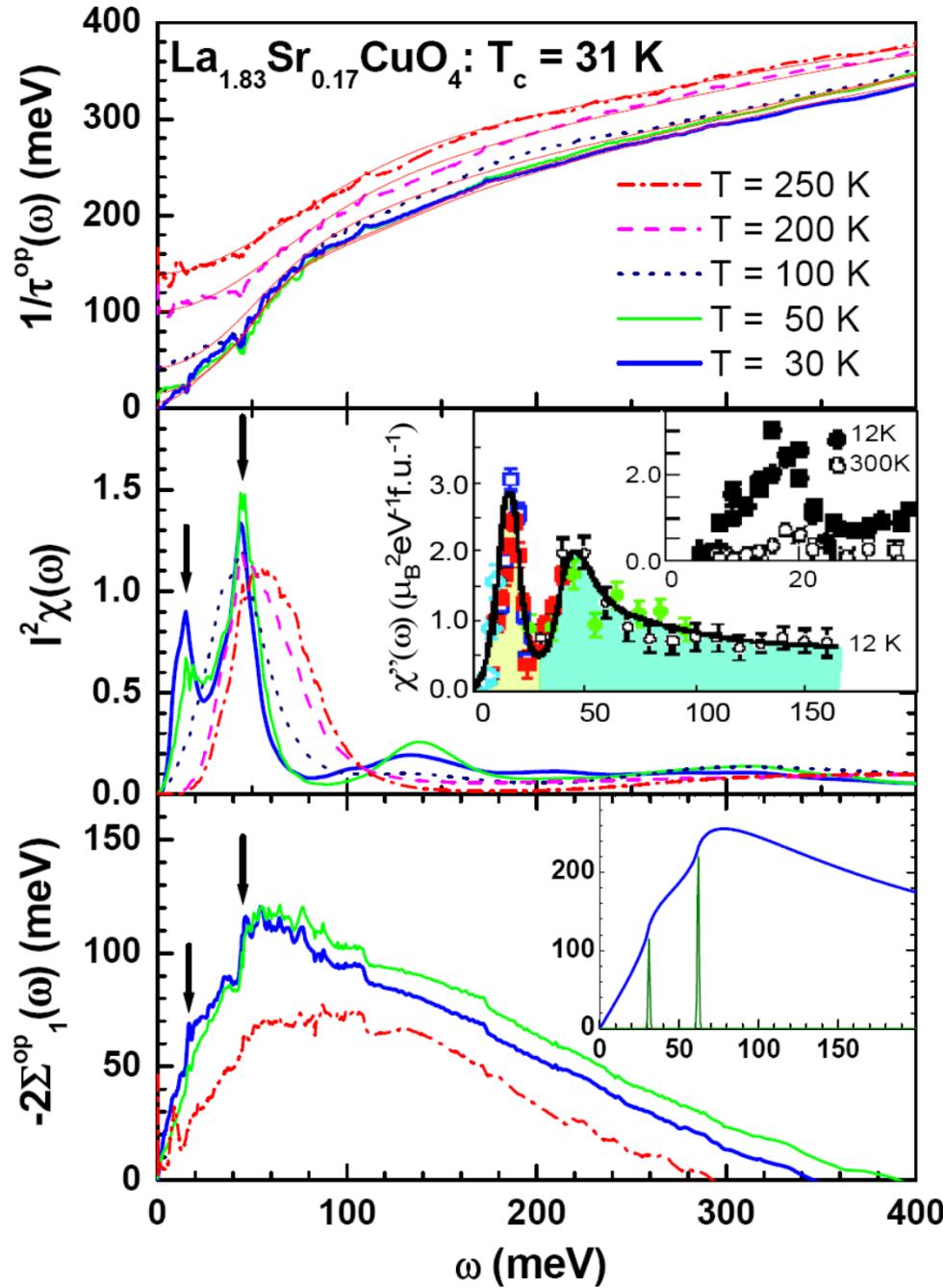
# Evolution with doping



# The evolution of $\chi''(q,\omega)$ in $\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$ with doping



# Optical Conductivity



J. Hwang et al.,  
PRL 100, 137005 (2008)

$$\begin{aligned} \sigma(\omega) &= i \frac{\omega_p^2}{4\pi} \frac{1}{\omega[1 + \lambda^{op}(\omega)] + i/\tau^{op}(\omega)} \\ &= i \frac{\omega_p^2}{4\pi} \frac{1}{\omega - 2\Sigma^{op}(\omega)} \end{aligned}$$

$$\frac{1}{\tau^{op}} \simeq \frac{2\pi}{\omega} \int_0^\infty d\Omega \alpha^2 F(\Omega) \int_0^{\omega-\Omega} d\omega' \tilde{N}(\omega')$$

spin fluctuations probably  
the glue...



*fin*