

Muonic atom spectroscopy with radioactive targets

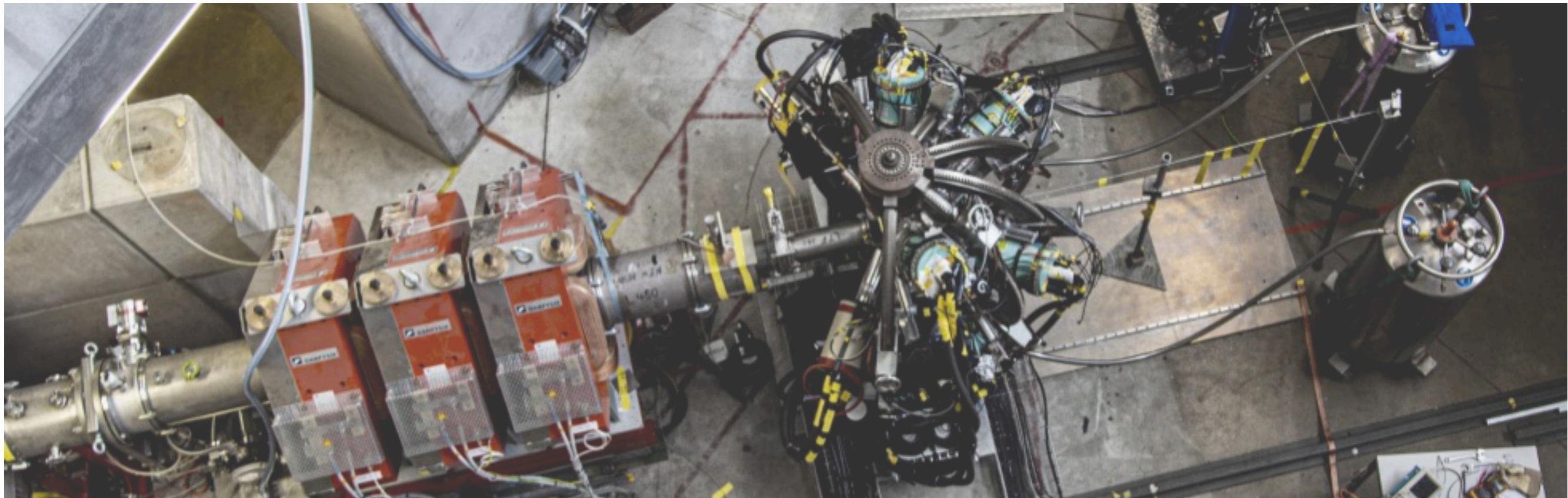
Stella Vogiatzi

On behalf of the muX collaboration*

17.10.2022

PSI2022: Physics of fundamental symmetries and interactions

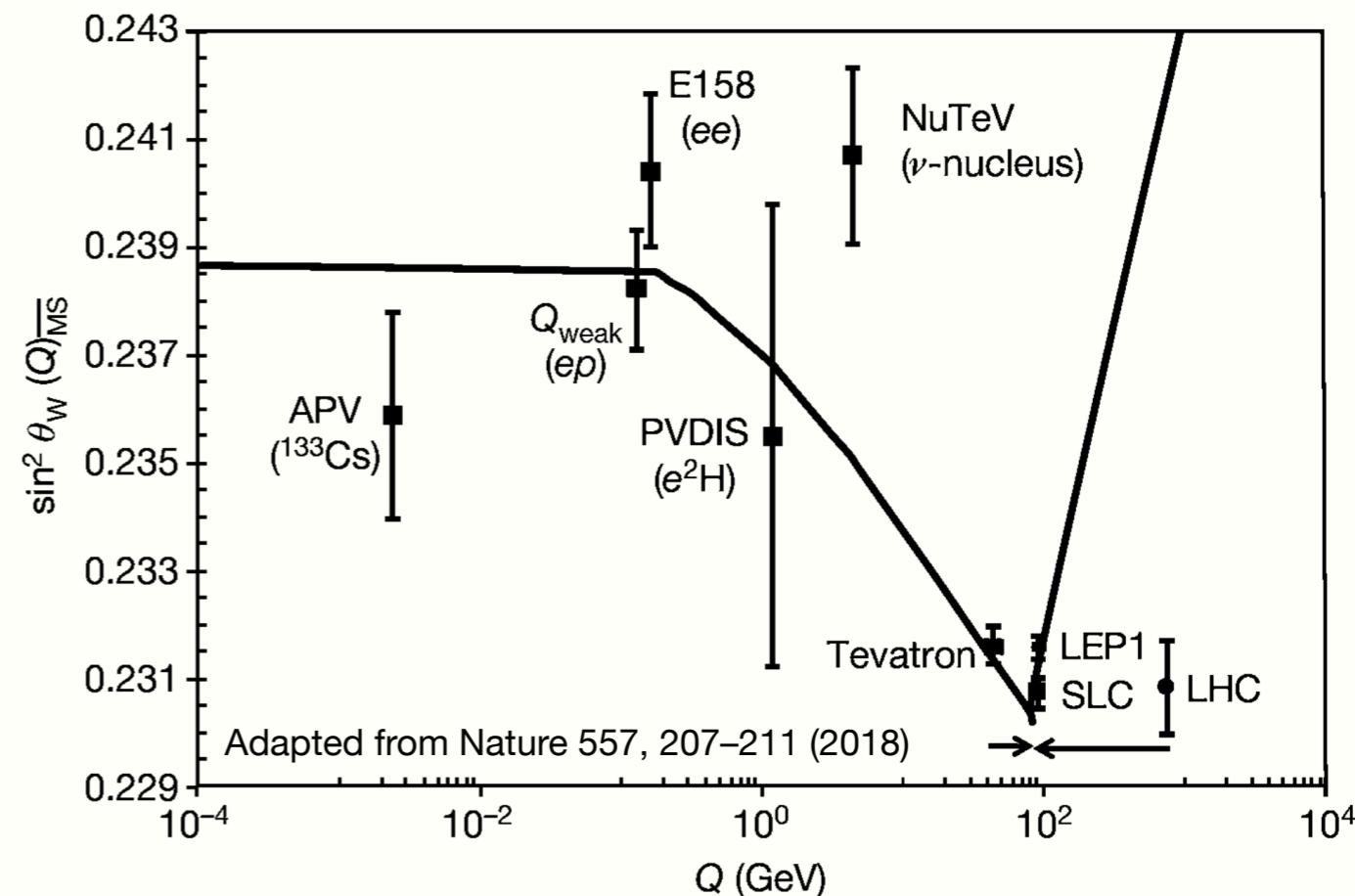
* Institute of Nuclear Physics, Polish Academy of Sciences, Krakow, Poland | Paul Scherrer Institut, Villigen, Switzerland | ETH Zürich, Switzerland | Johannes Gutenberg University Mainz, Germany | KU Leuven, Belgium | GSI Helmholtzzentrum für Schwerionenforschung Darmstadt, Germany | Helmholtz Institute Mainz, Germany | Institut für Kernphysik, Universität zu Köln, Germany | LKB Paris, France | University of Groningen, The Netherlands | University of Pisa and INFN, Pisa, Italy | University of Victoria, Canada | Perimeter Institute, Waterloo, Canada | CSNSM, Université Paris Sud, CNRS/IN2P3, Université Paris Saclay, Orsay Campus, France



Motivation: atomic parity violation in radium

- Atomic Parity Violation (APV) experiments in atoms probe the low transfer momentum Q region in the running of the $\sin^2(\theta_W)$ plot
- Weak interactions of nucleus and leptons enable APV transitions in atoms
- APV is magnified proportionally to $\gtrsim Z^3$
 - ⇒ heavy atoms are good candidates
 - ⇒ experimental efforts to trap / laser cool $^{226}\text{Ra}^+$

Hyperfine Interact. 199, 9 (2011)
Phys. Rev. Lett. 122, 223001 (2019)



$$\text{APV transitions} \propto k \cdot Q_W$$

weak charge (related to
Weinberg angle)

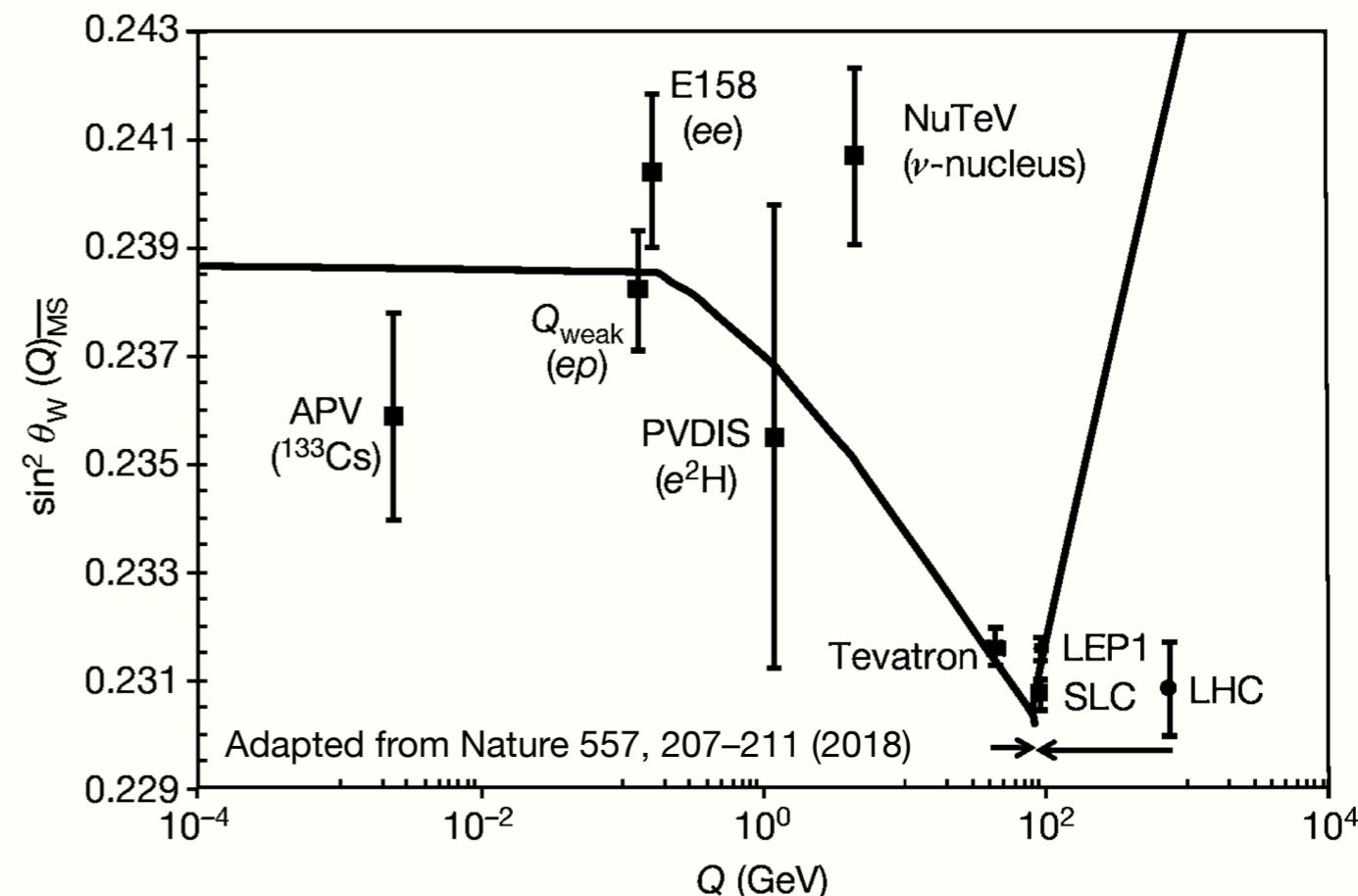
precision atomic
calculations

The absolute nuclear charge radius of ^{226}Ra at the level of 0.2% is needed
→ muX aims to measure it using muonic atoms

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What about other
radioactive elements?

^{248}Cm

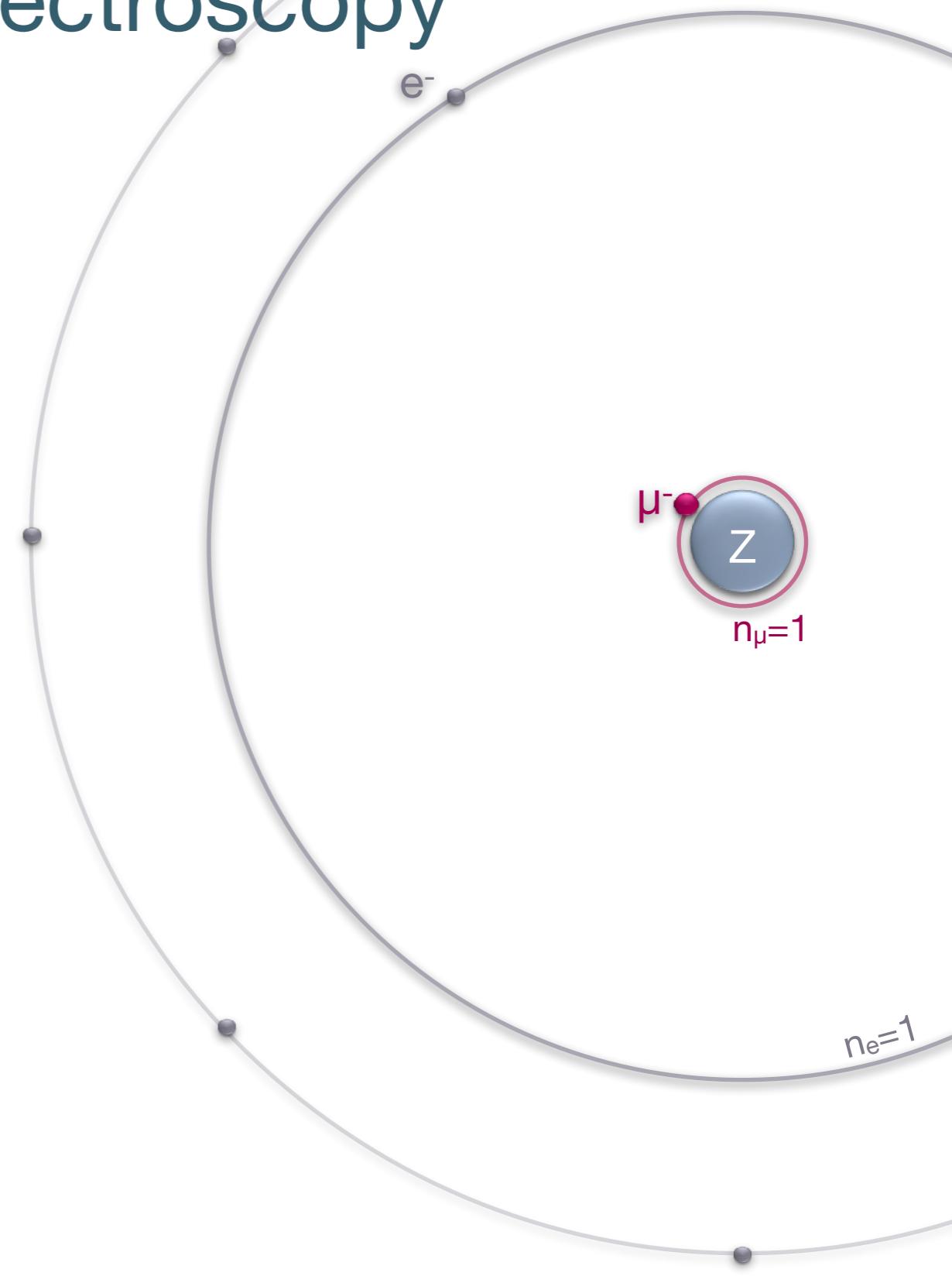
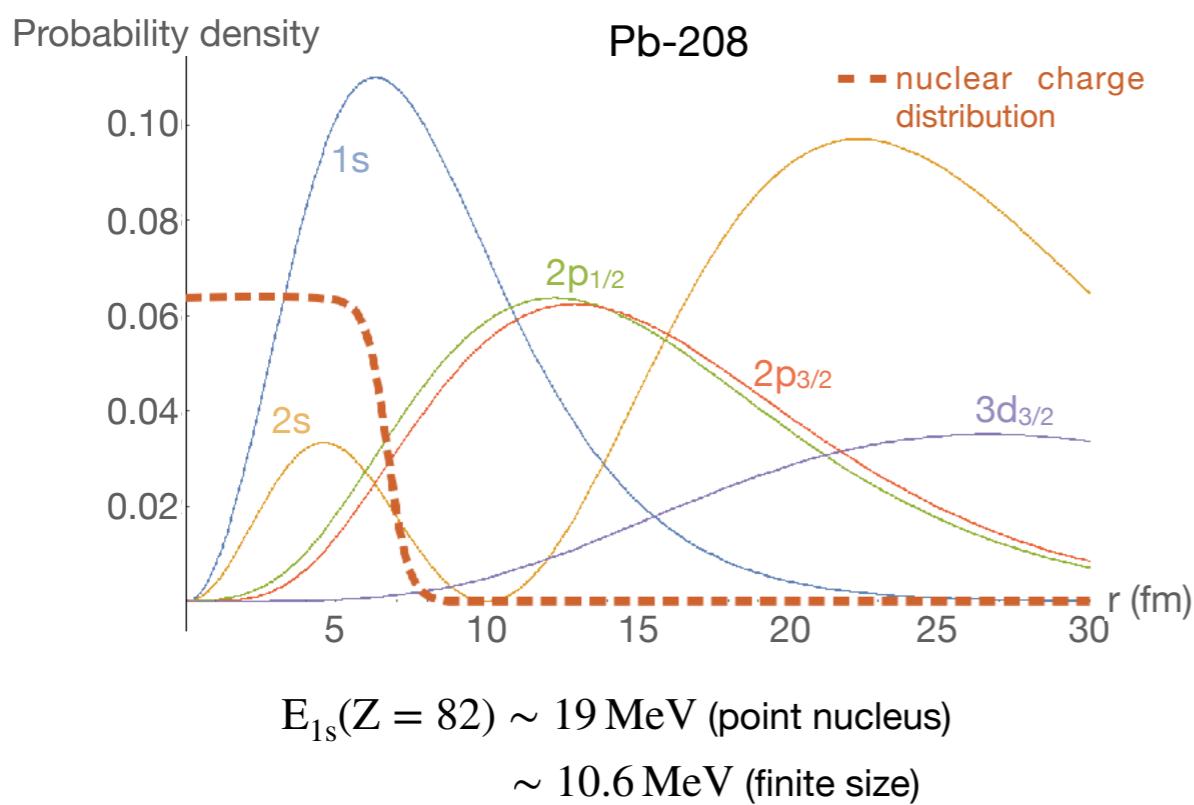
Muonic atom spectroscopy

- Due to the higher muon mass, there is large overlap of the low-lying muonic states with the nuclear charge distribution
⇒ The energy of these levels is highly affected by the nuclear structure details
- The measurement of the muonic energy levels allows to extract properties of the nucleus such as the charge radius r_c :

$$r_c^2 = \langle r^2 \rangle = \frac{1}{Ze} \int d^3\vec{r} \rho(\vec{r}) r^2$$

2p-1s transition most sensitive to r_c

nuclear shape? ——————→



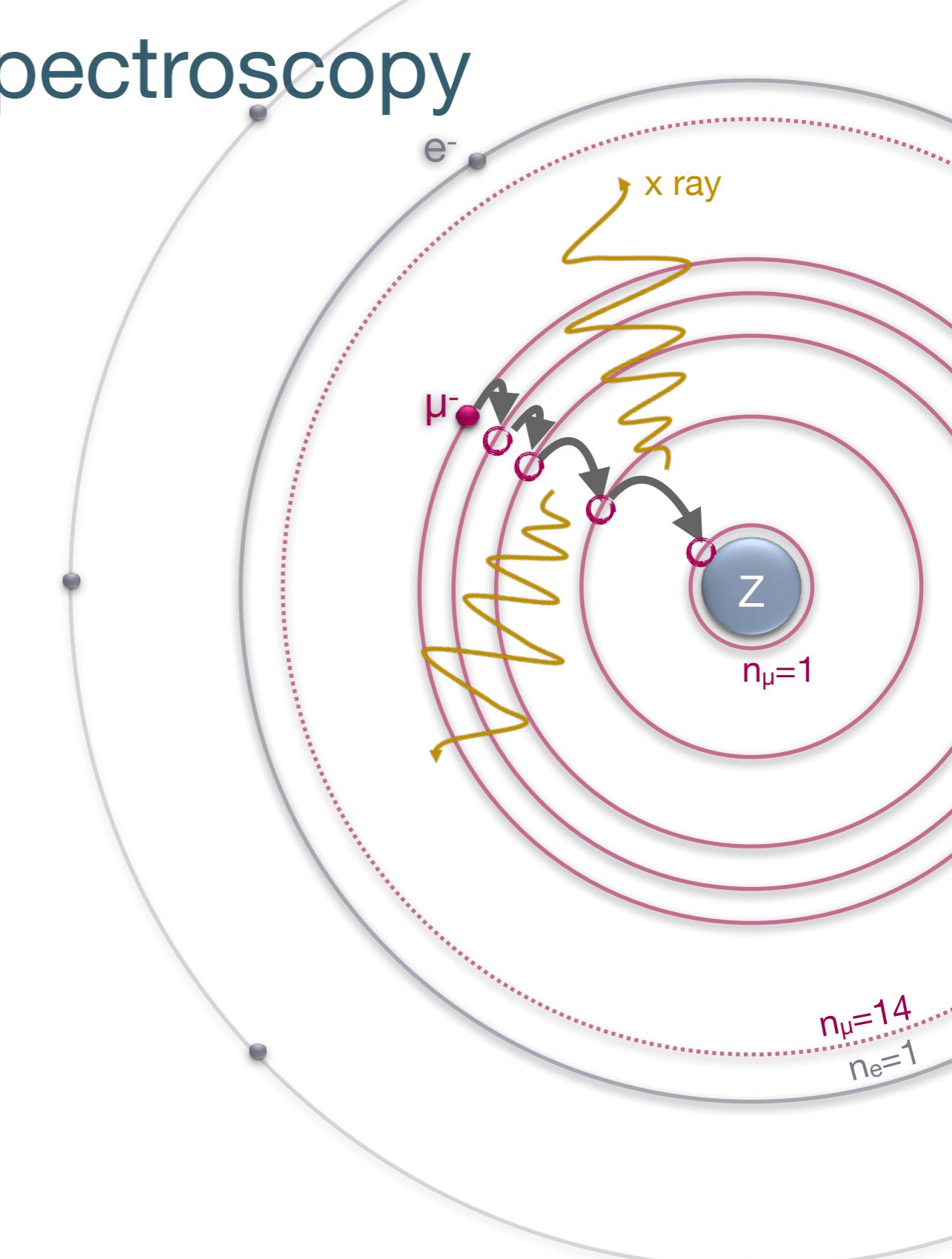
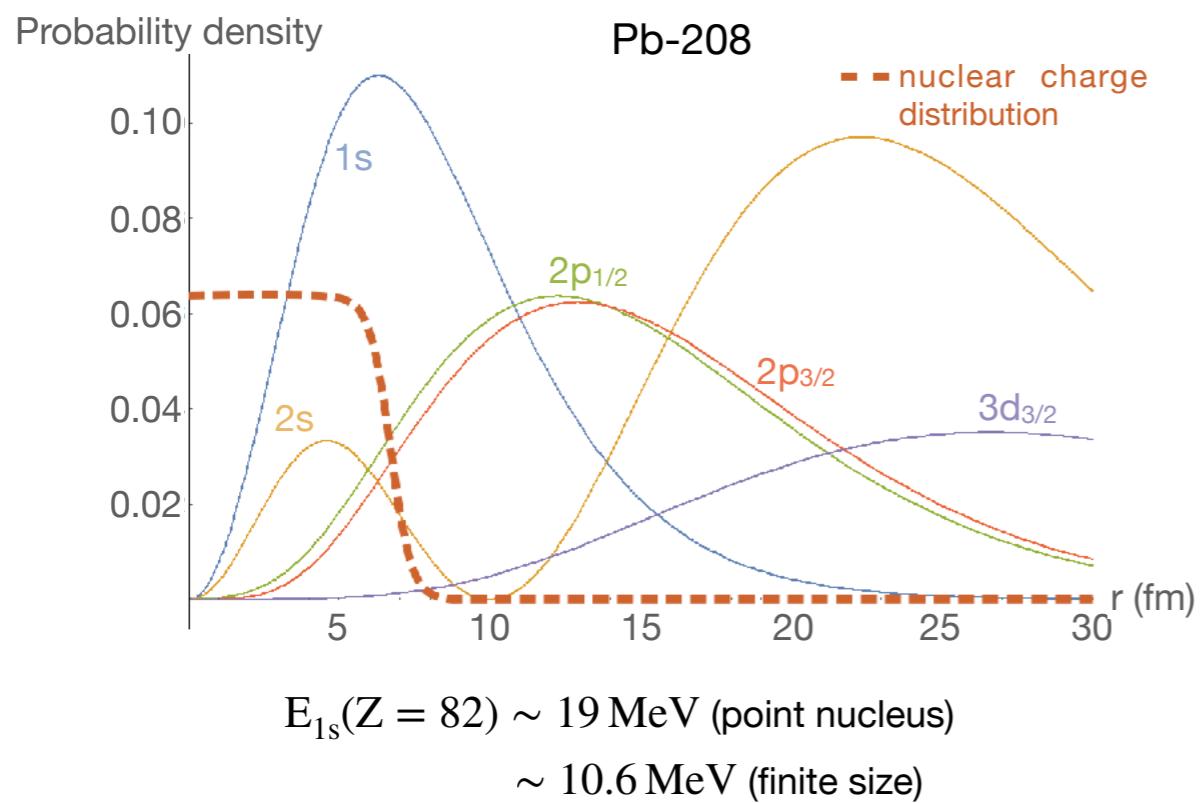
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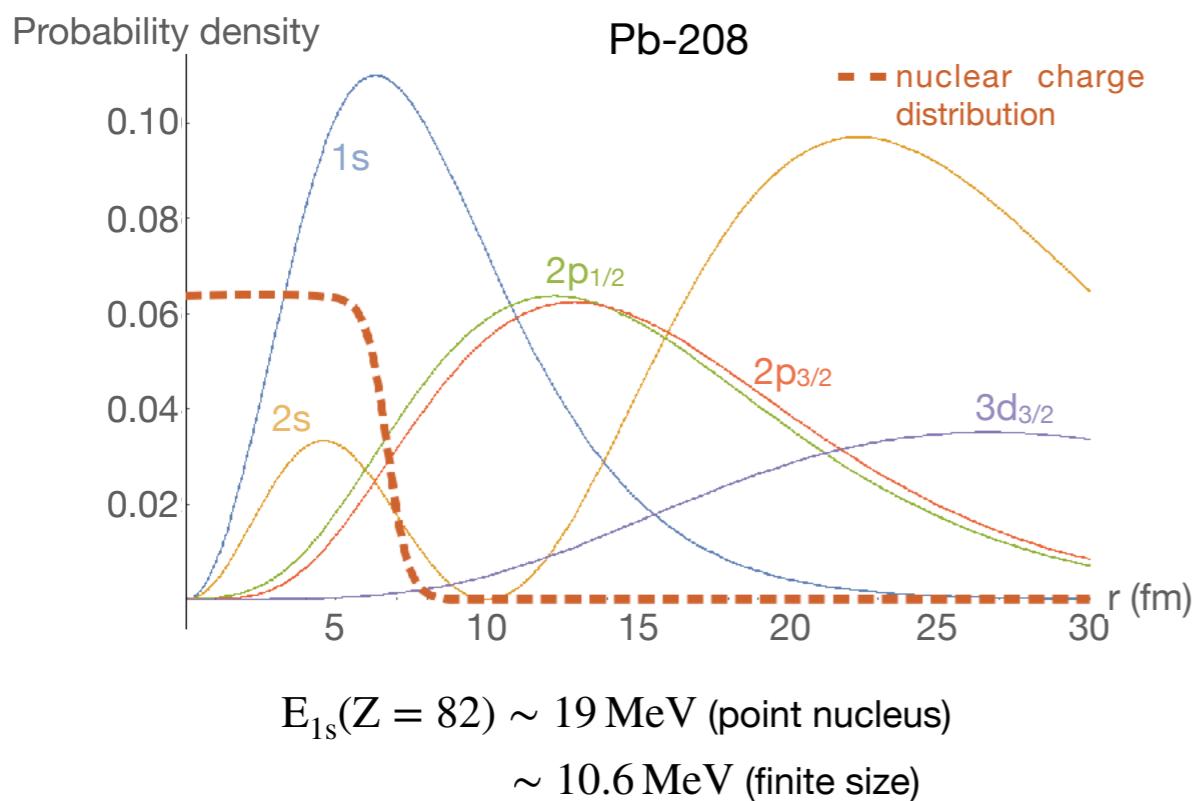
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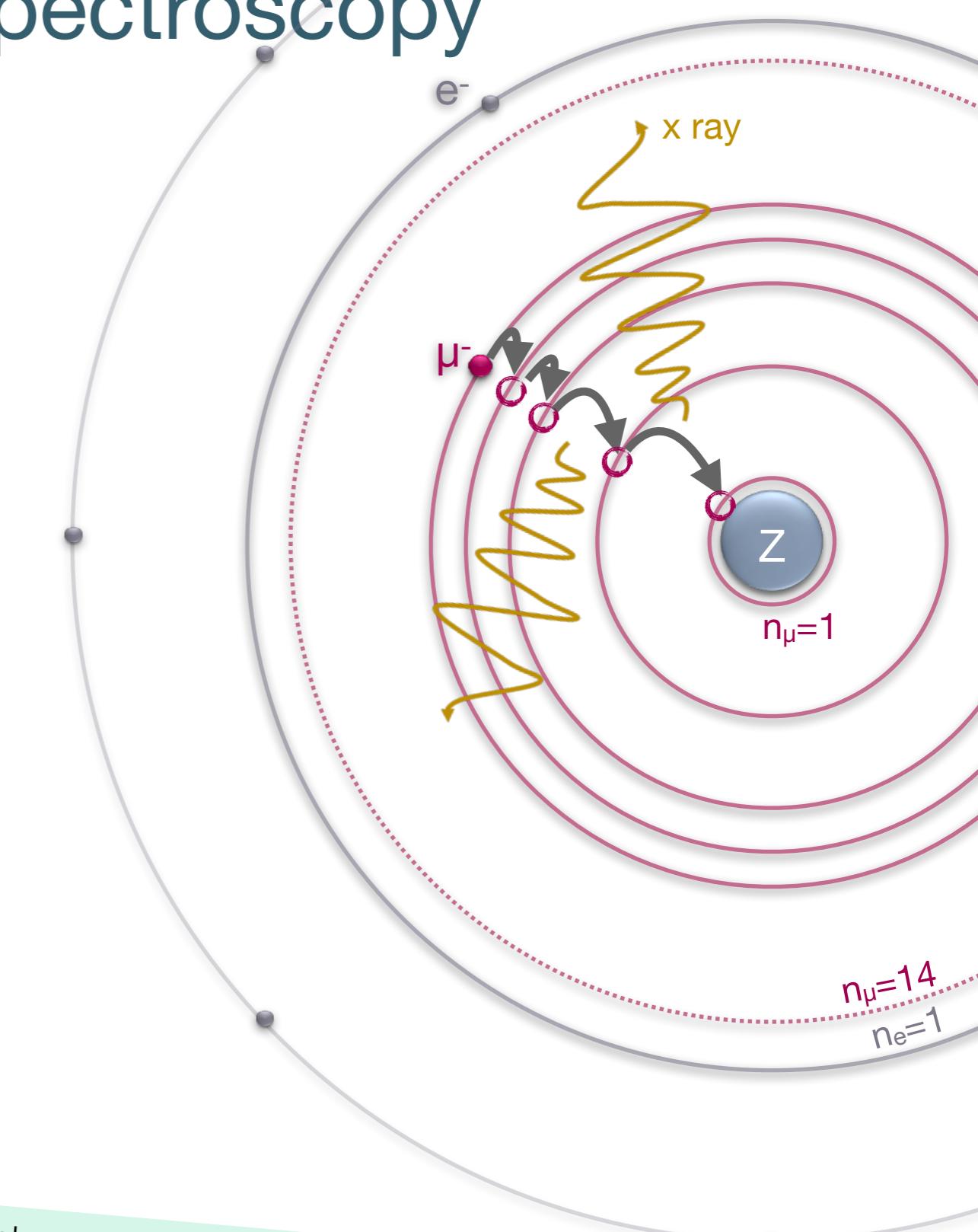
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... only possible with $\gtrsim 100 \text{ mg}$ targets... radioactive targets are typically restricted to $\sim \mu\text{g}$ quantities in the experimental areas



Where does this happen?



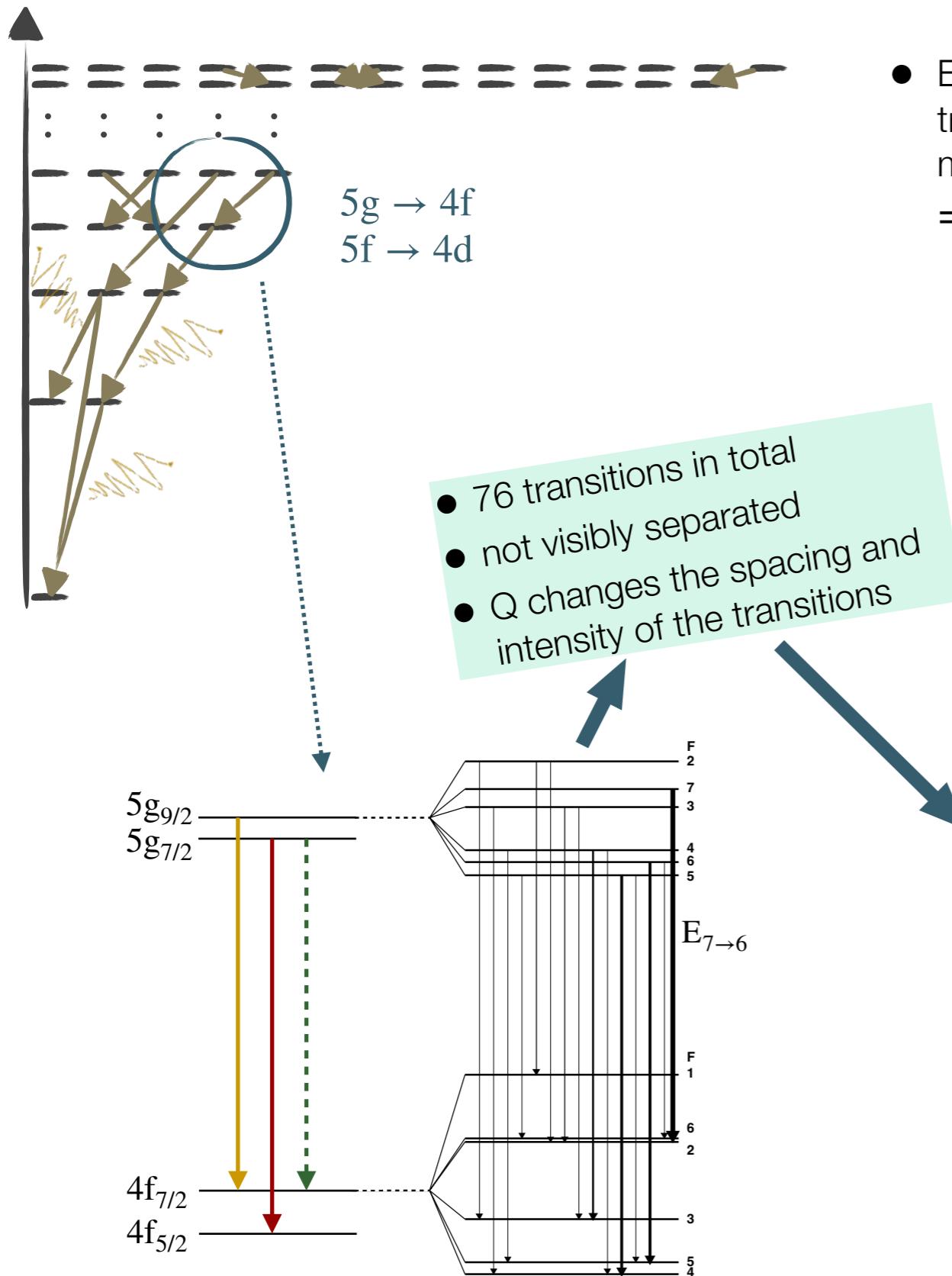
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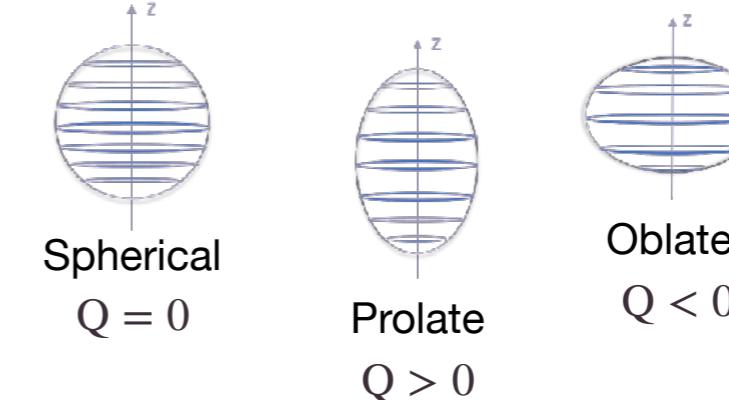
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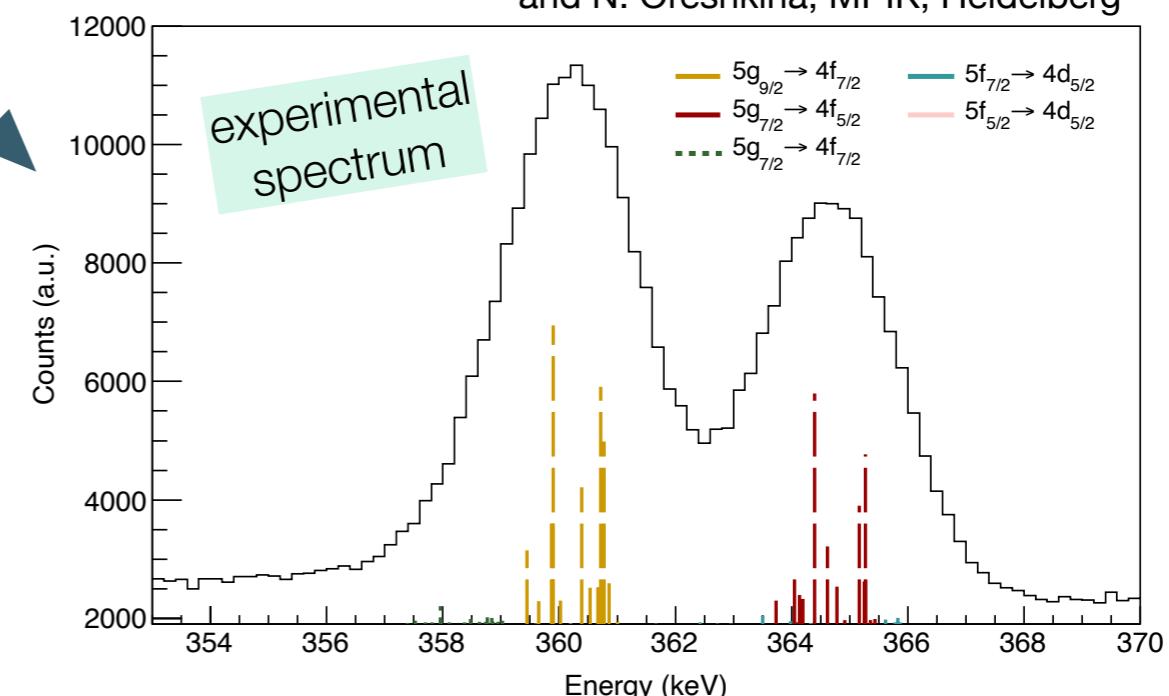
Spectroscopic quadrupole moment in $^{185,187}\text{Re}$



- Being sufficiently far away from the nucleus, the $5 \rightarrow 4$ transition is not sensitive to the details of the nuclear moments distribution
⇒ we can extract the spectroscopic quadrupole moment Q

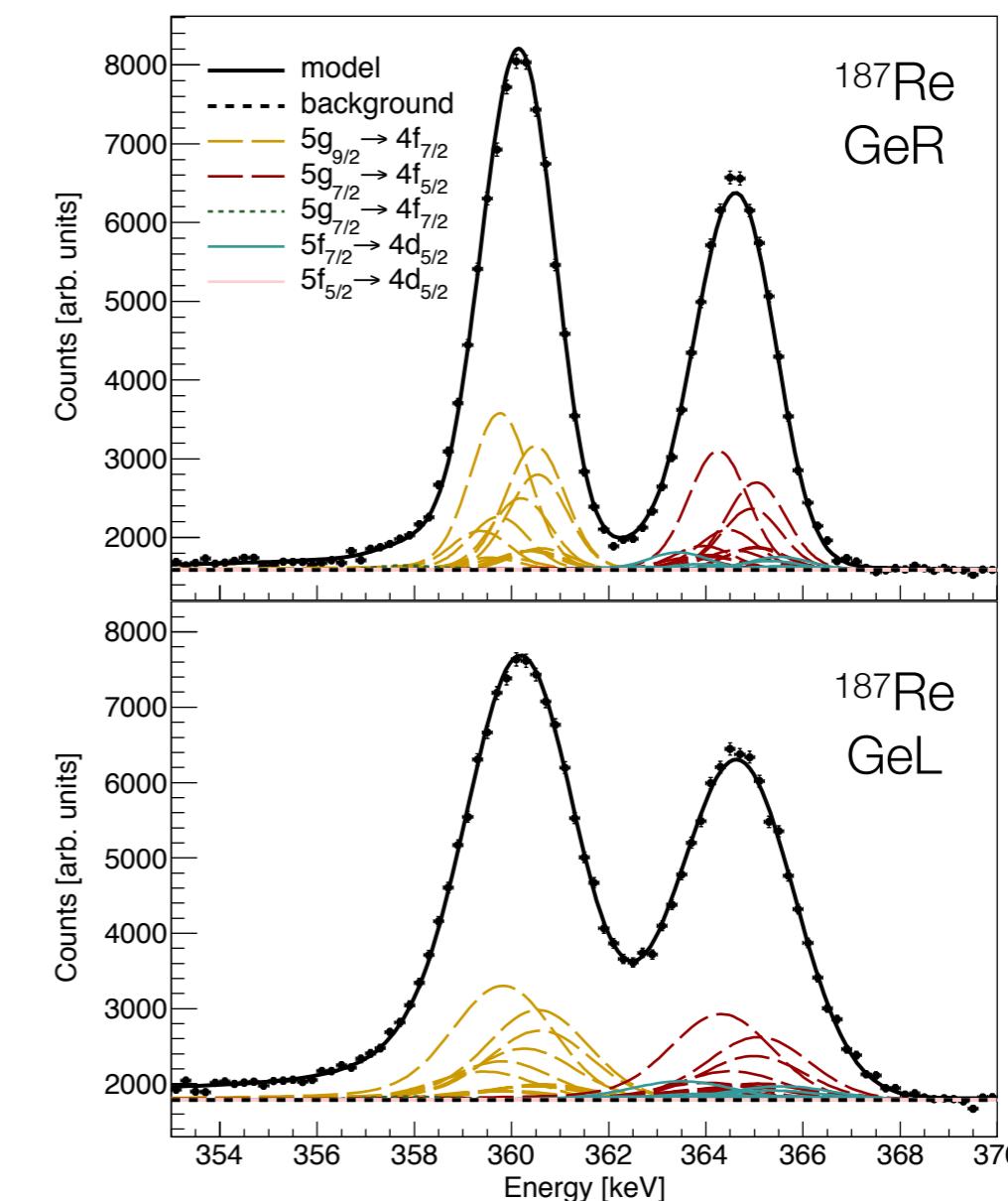
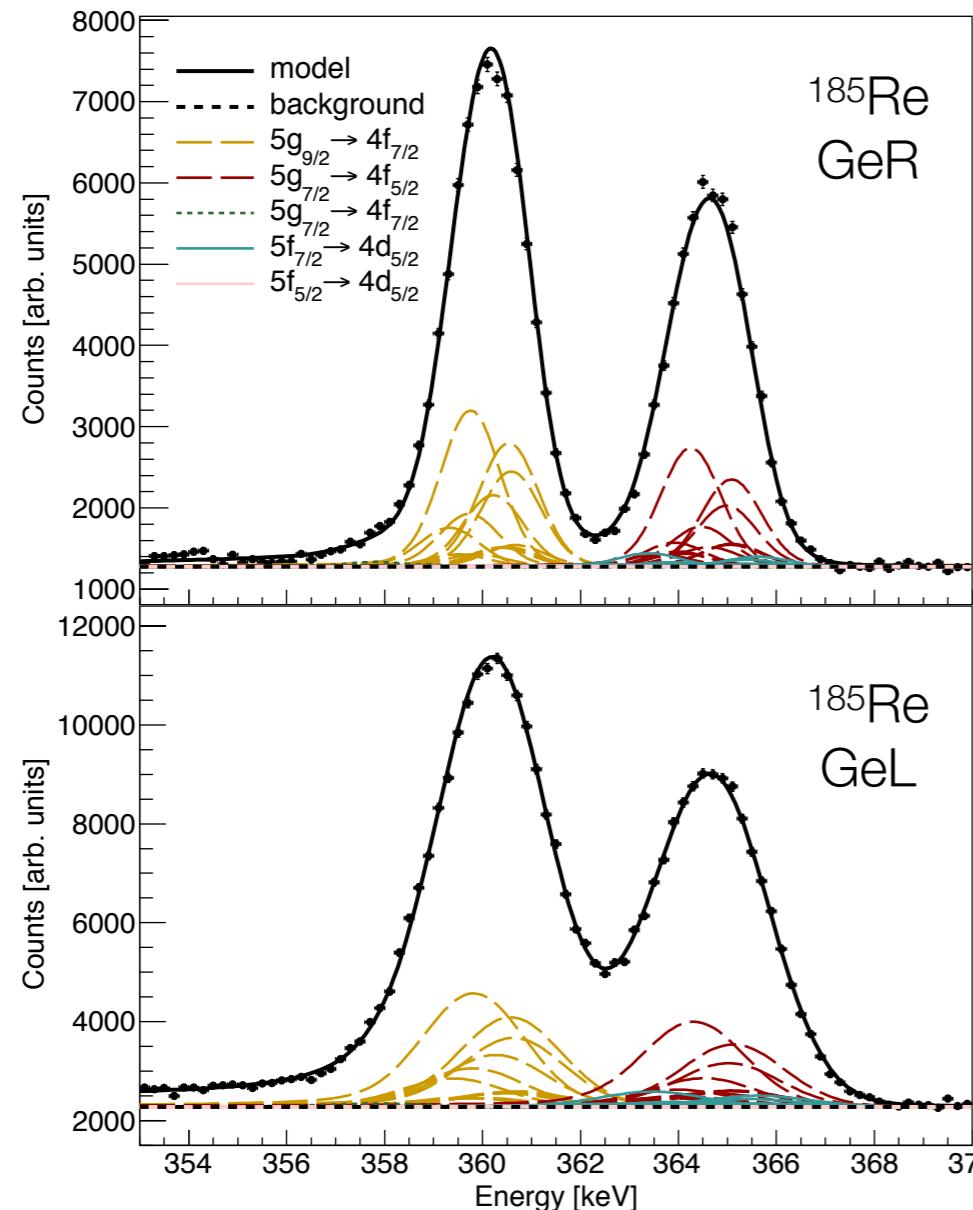
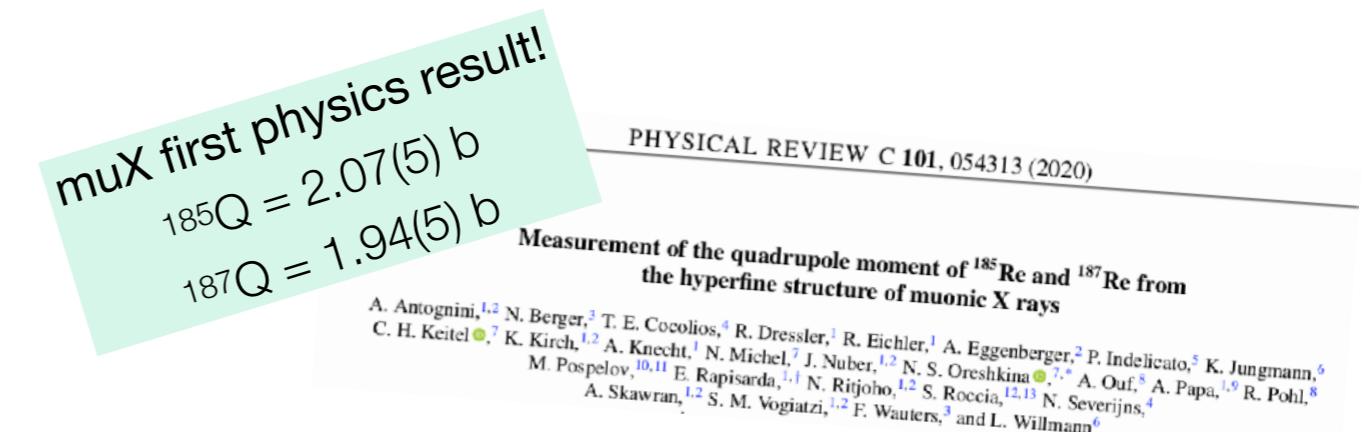


Theoretical calculations by N. Michel and N. Oreshkina, MPIK, Heidelberg



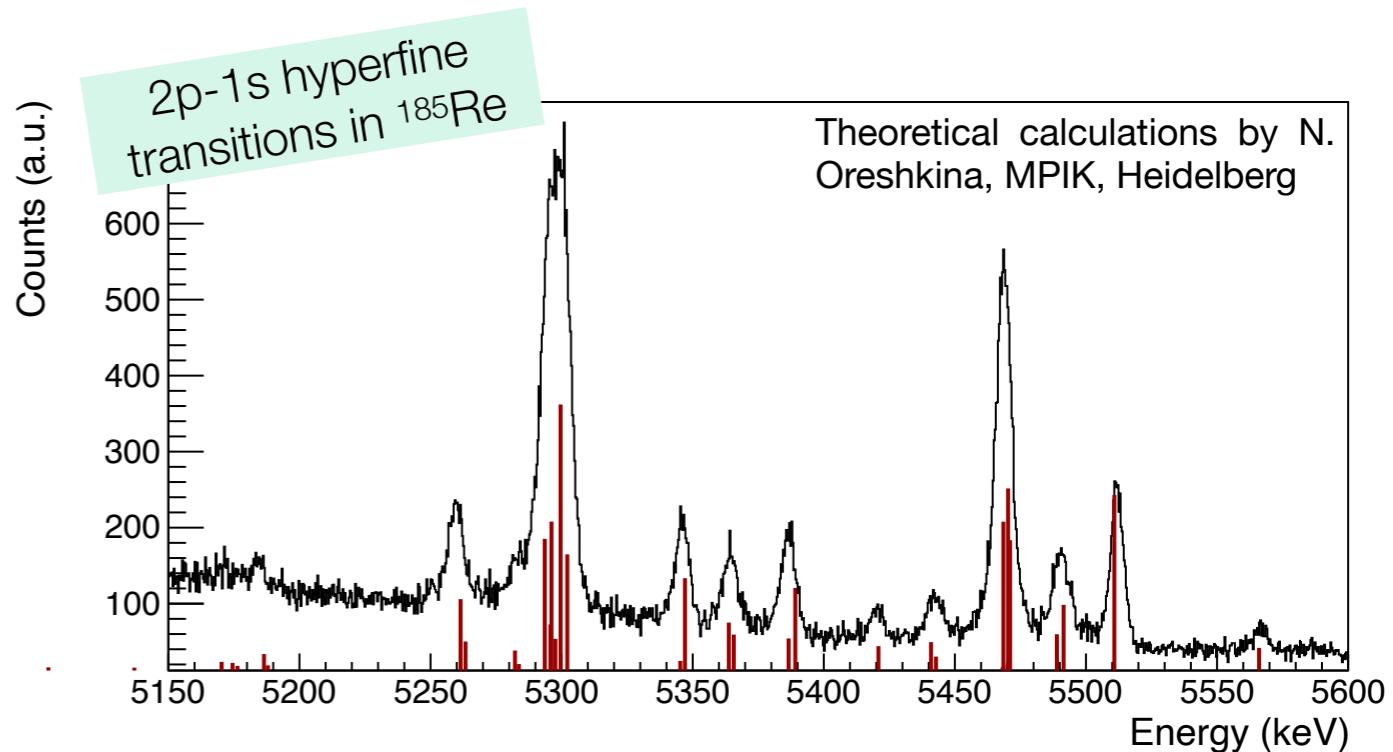
Spectroscopic quadrupole moment in $^{185,187}\text{Re}$

- Fitting of the theoretical predictions in the experimental spectrum with the quadrupole moment as a free parameter
- Cross checks in two HPGe datasets



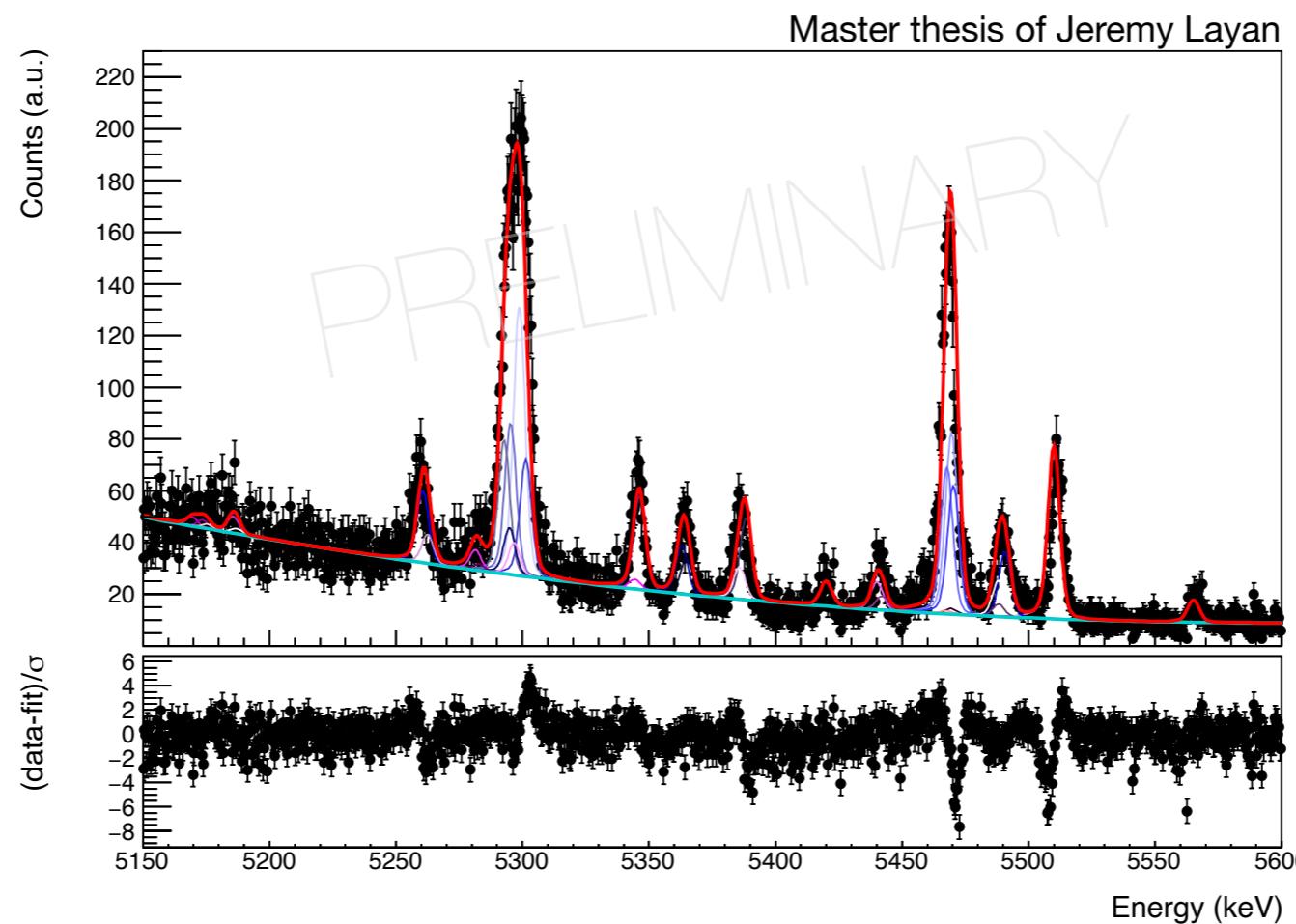
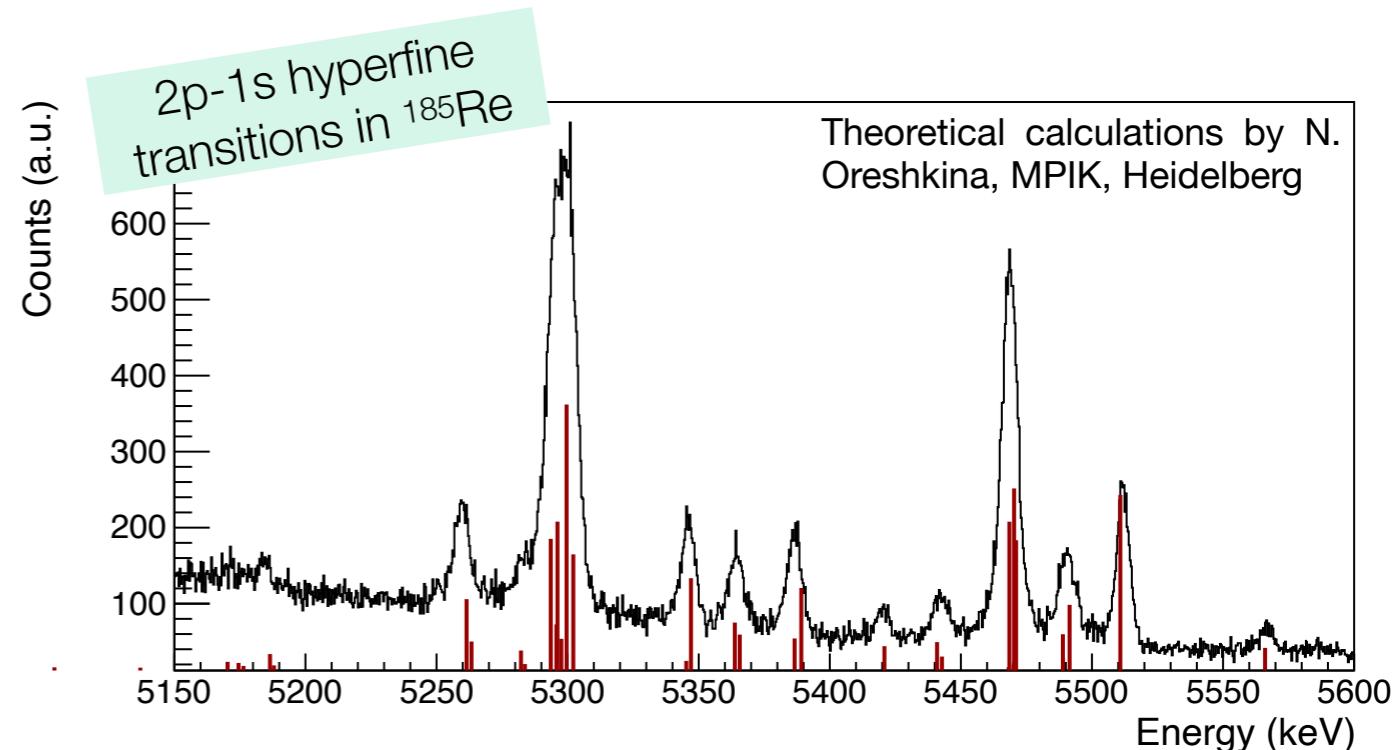
Outlook: charge radius of $^{185,187}\text{Re}$

- Currently ongoing analysis of the 2p-1s hyperfine transitions in $^{185,187}\text{Re}$ for the extraction of its absolute nuclear charge radius
⇒ has not been measured before
- Complicated hyperfine structure due to the dynamic effect

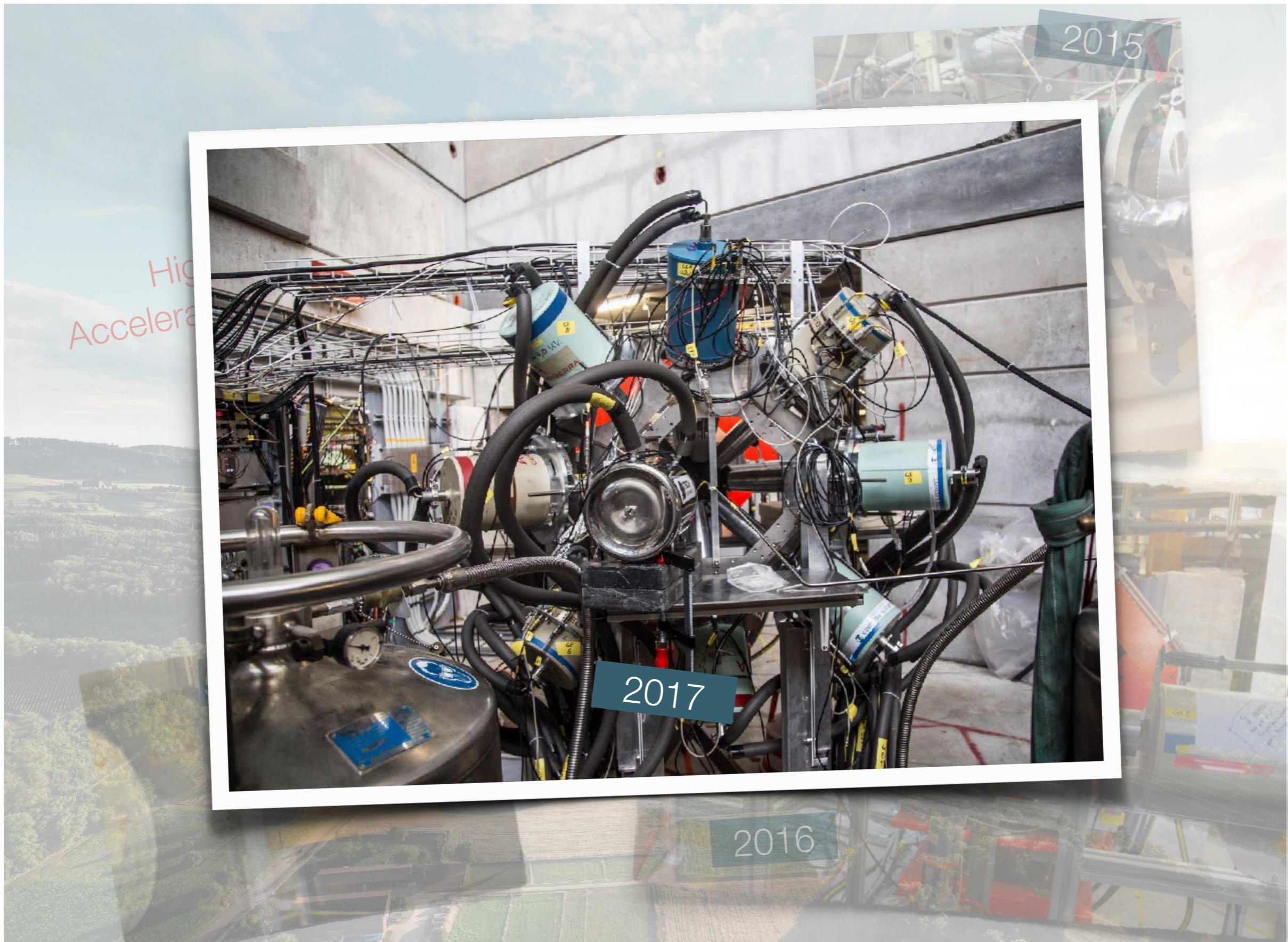


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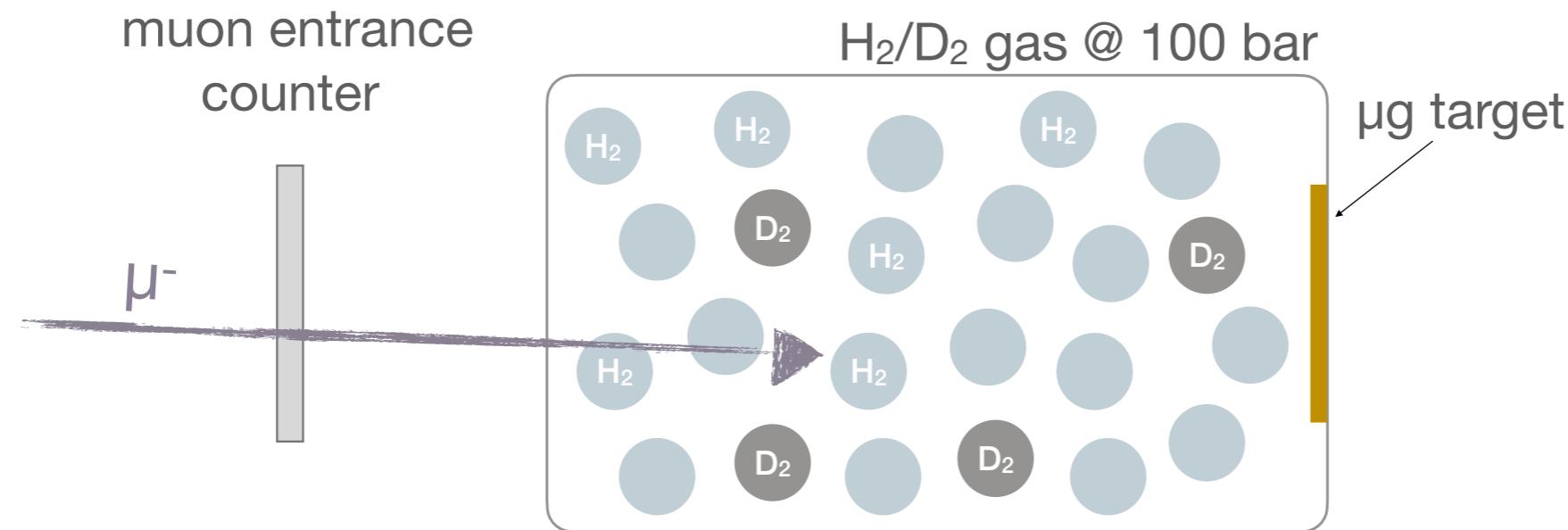


Towards the measurement of μ g targets



Muon transfer to microgram targets

Inspired by the work of Strasser *et al.* and Kraiman *et al.*

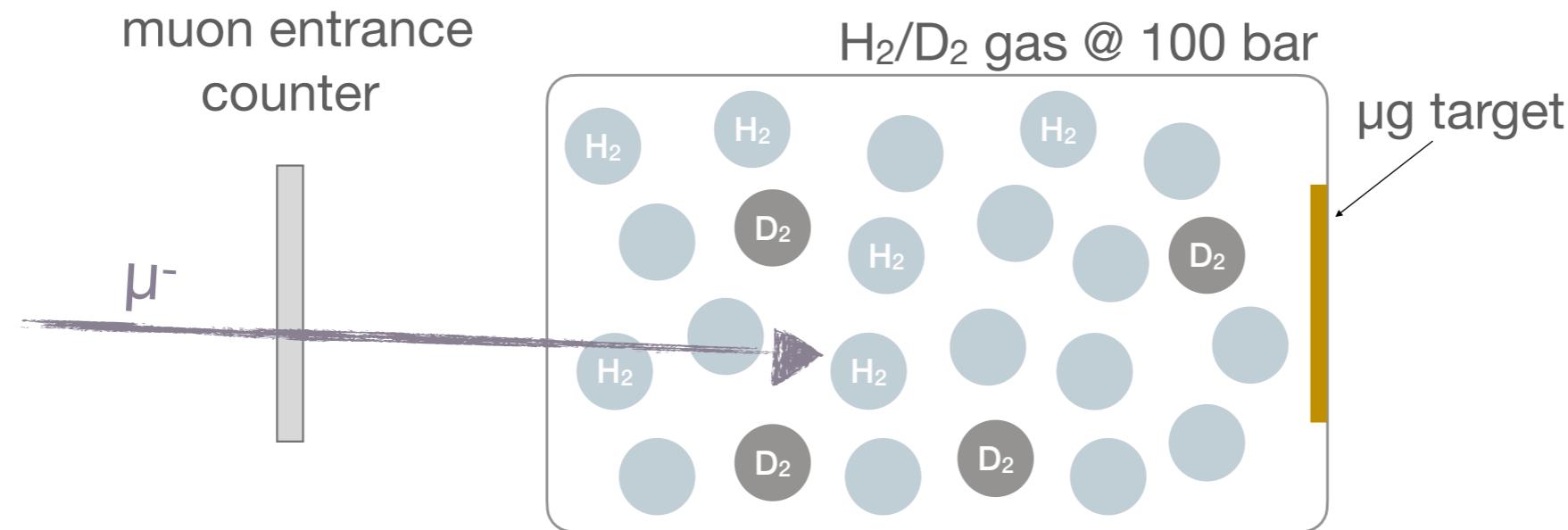


1. μ^- stops in 100 bar of $H_2 + 0.1\text{-}1.5\% D_2$ & forms muonic hydrogen μp
2. transfer to deuterium $\mu p \rightarrow \mu d$
3. μd moves almost freely in the H_2 gas (Ramsauer-Townsend effect¹)
4. transfer to high-Z element $\mu d \rightarrow \mu Z$ when hitting target & emission of x rays during the atomic cascade

¹ Physical Review A 73, 034501 (2006)

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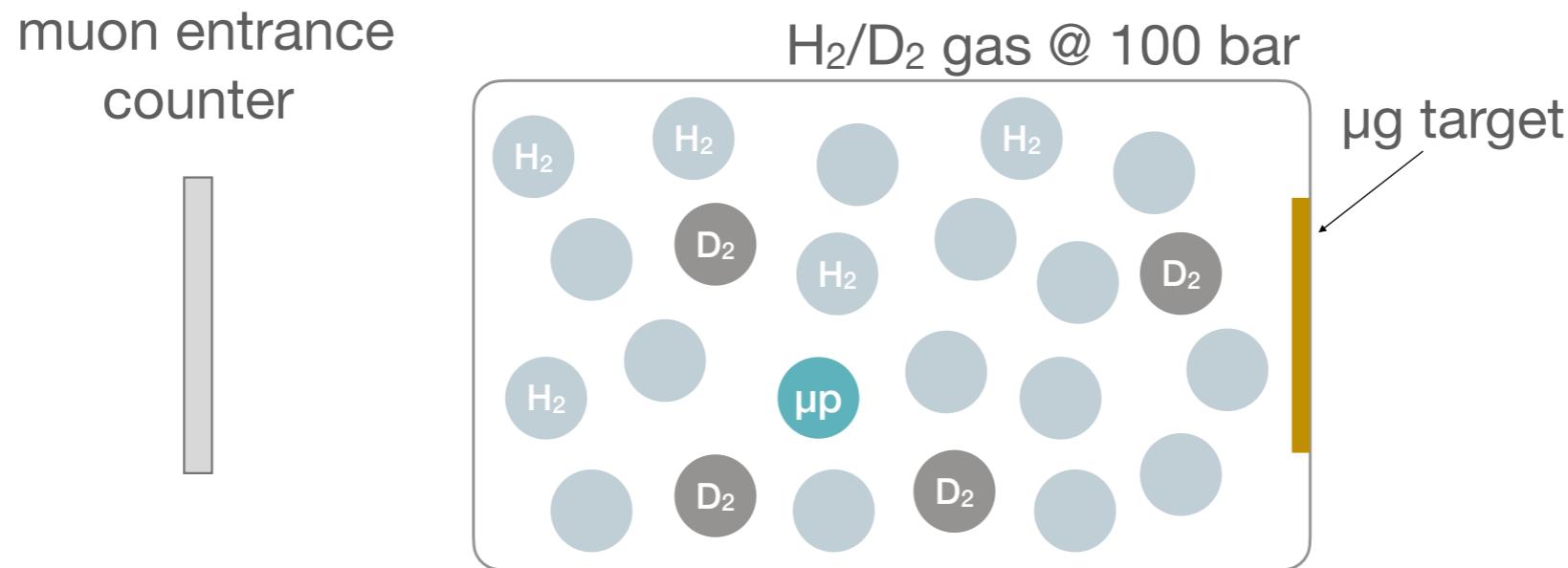


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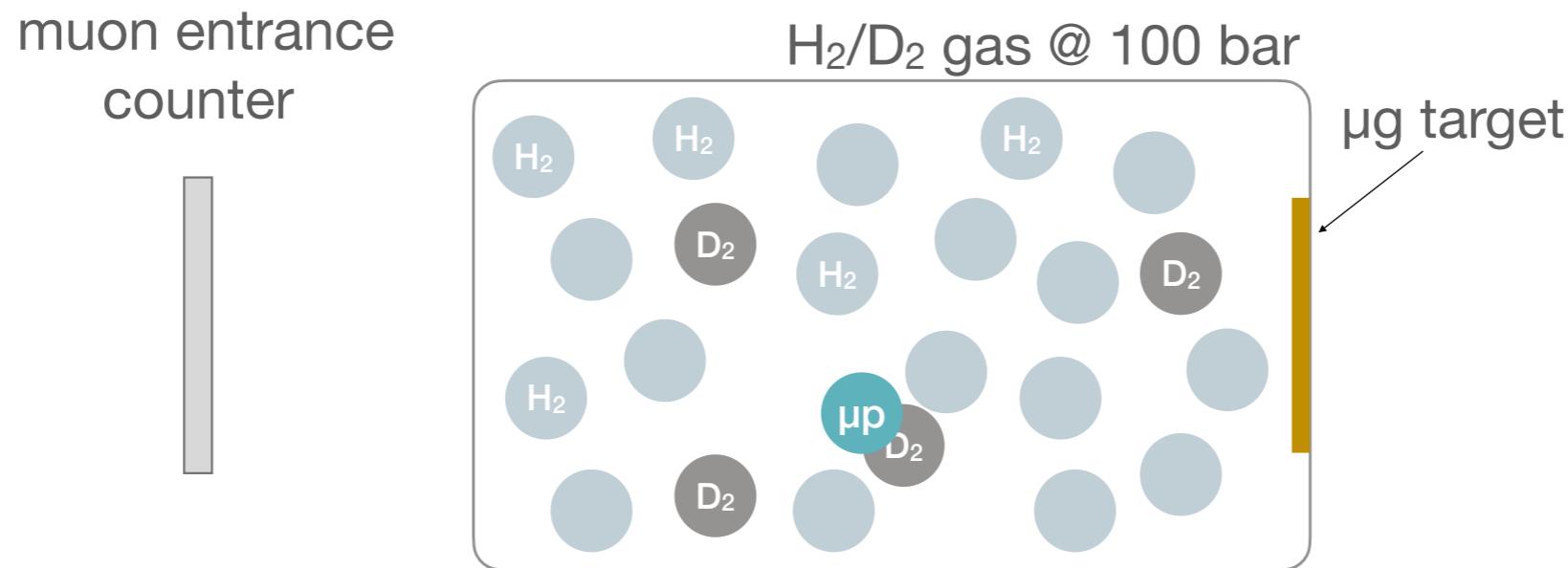


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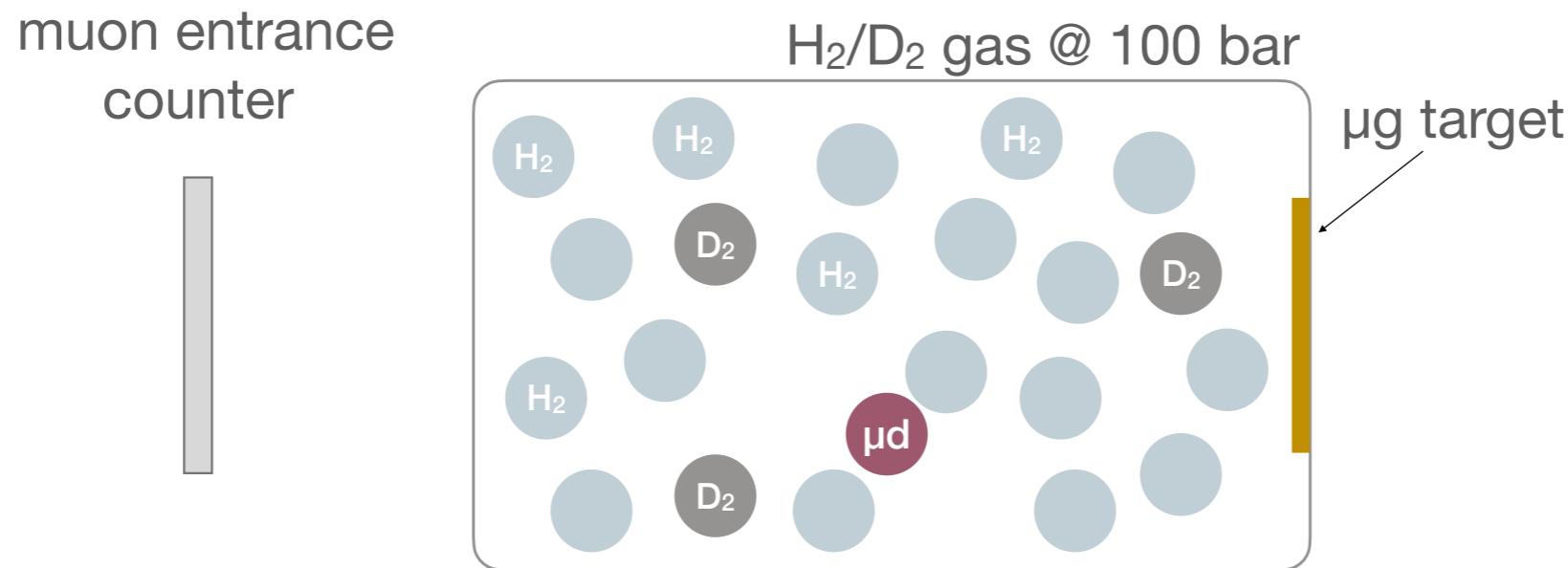


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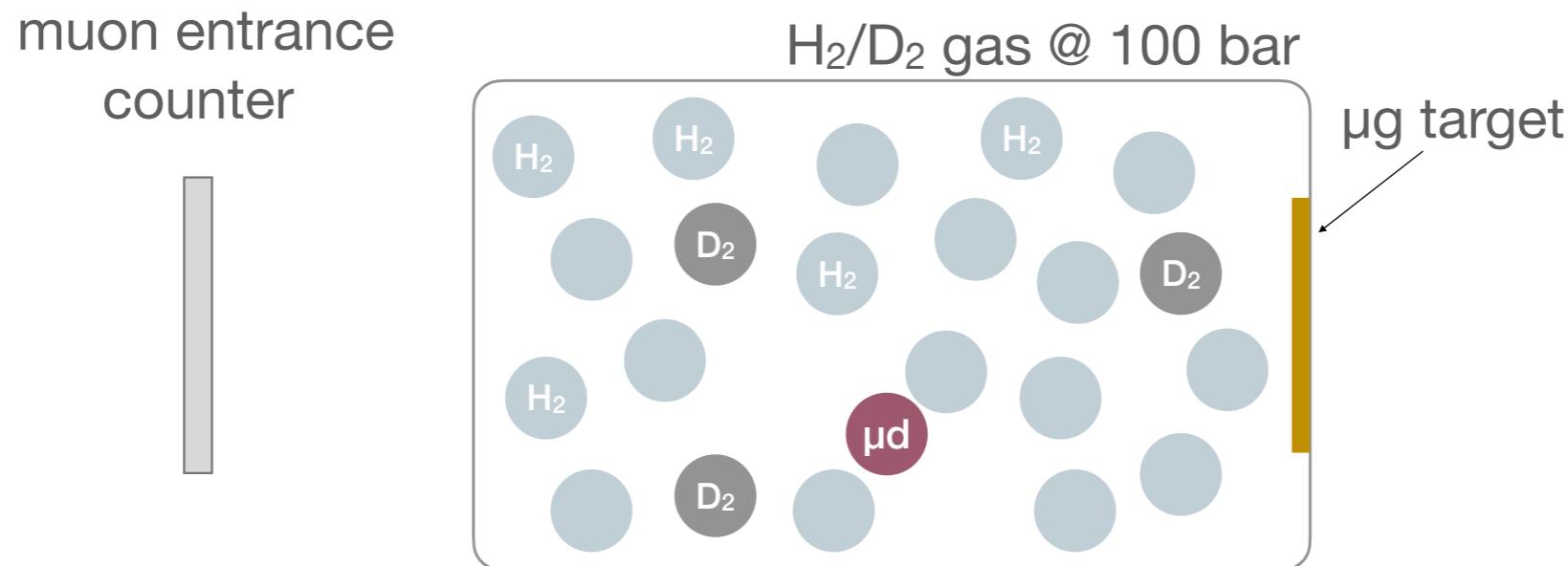


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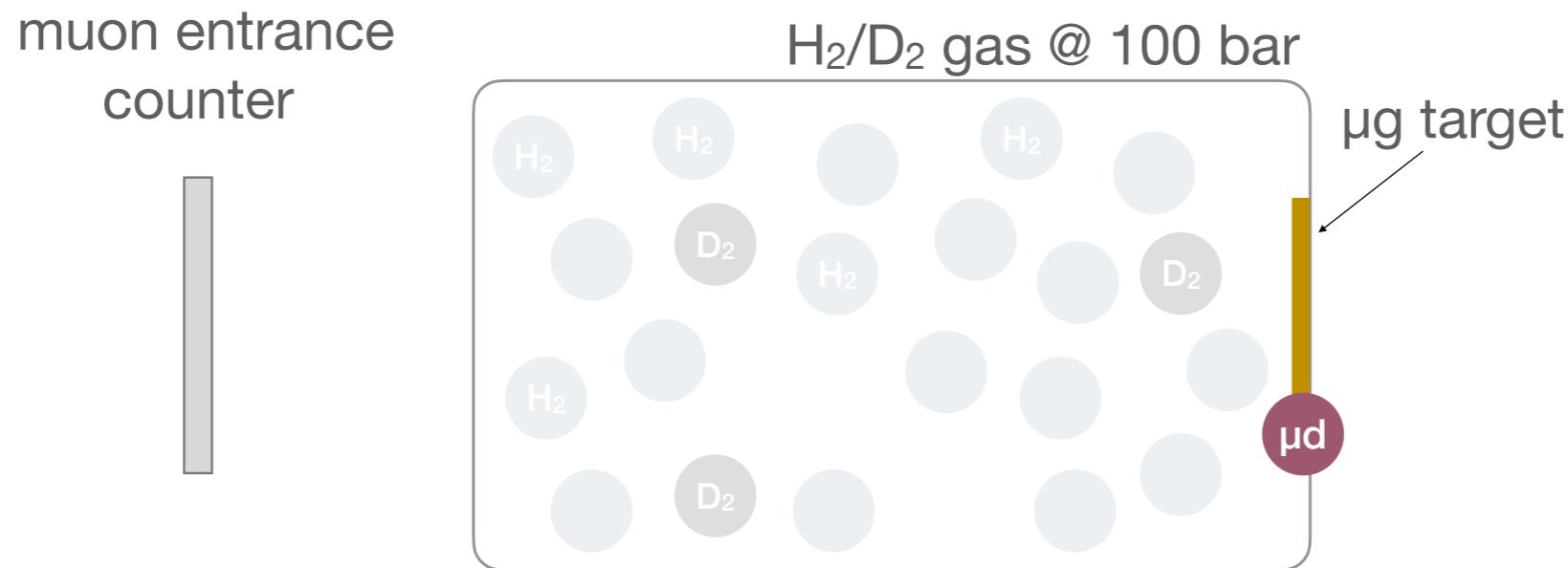


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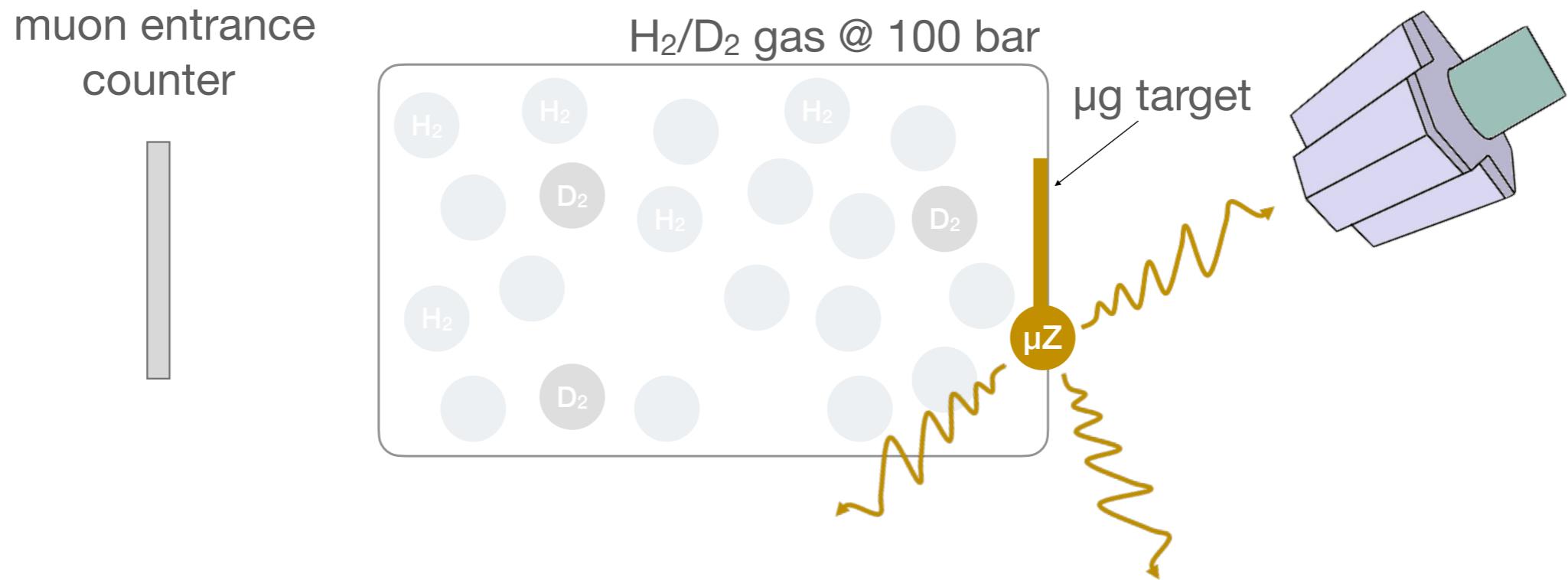


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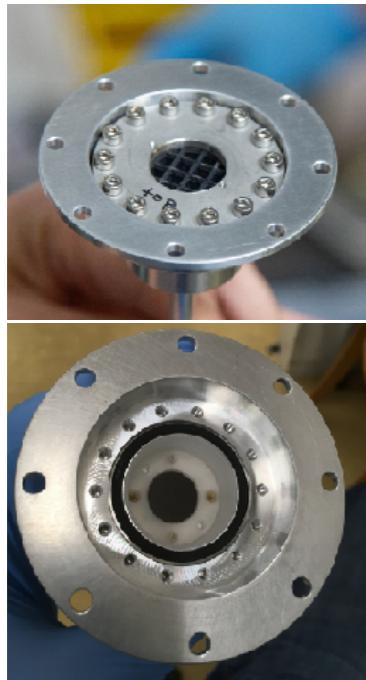
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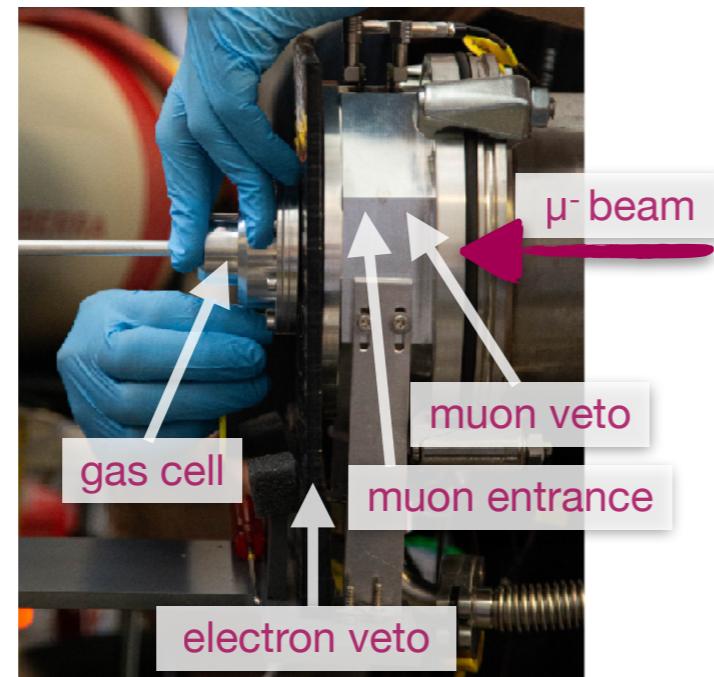
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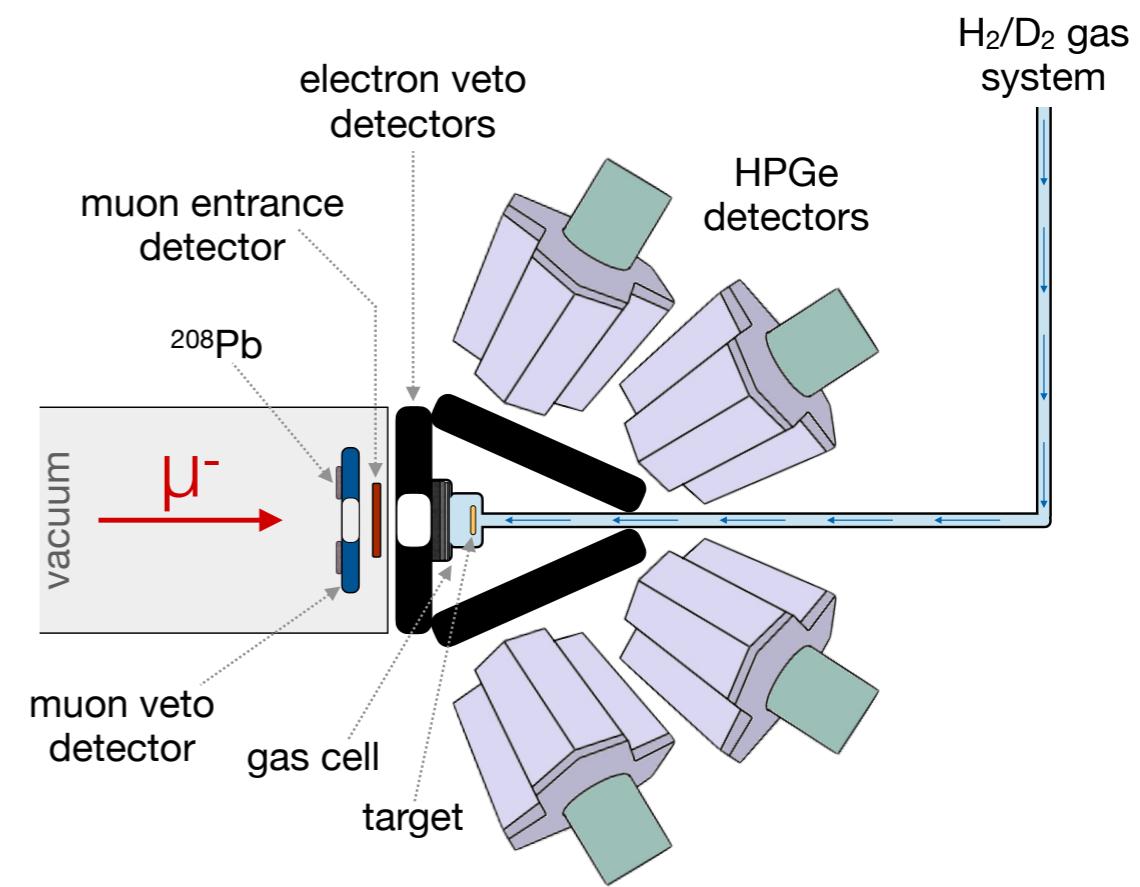
muX detectors



The gas cell

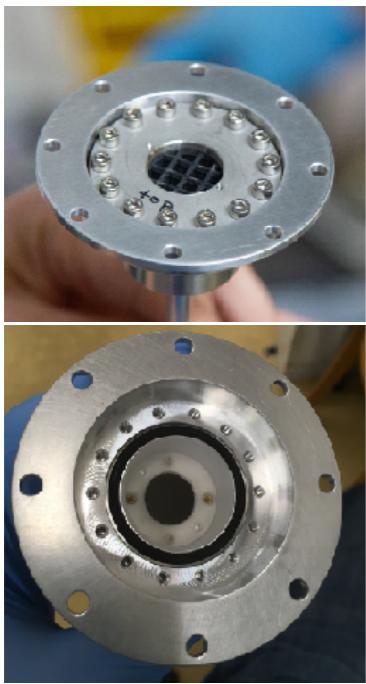


The muon and electron counters

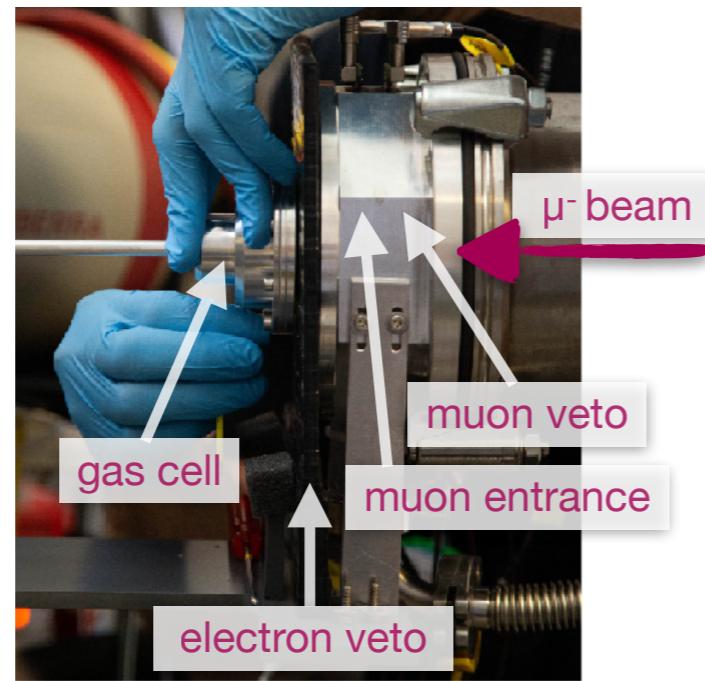


Schematics of detectors setup

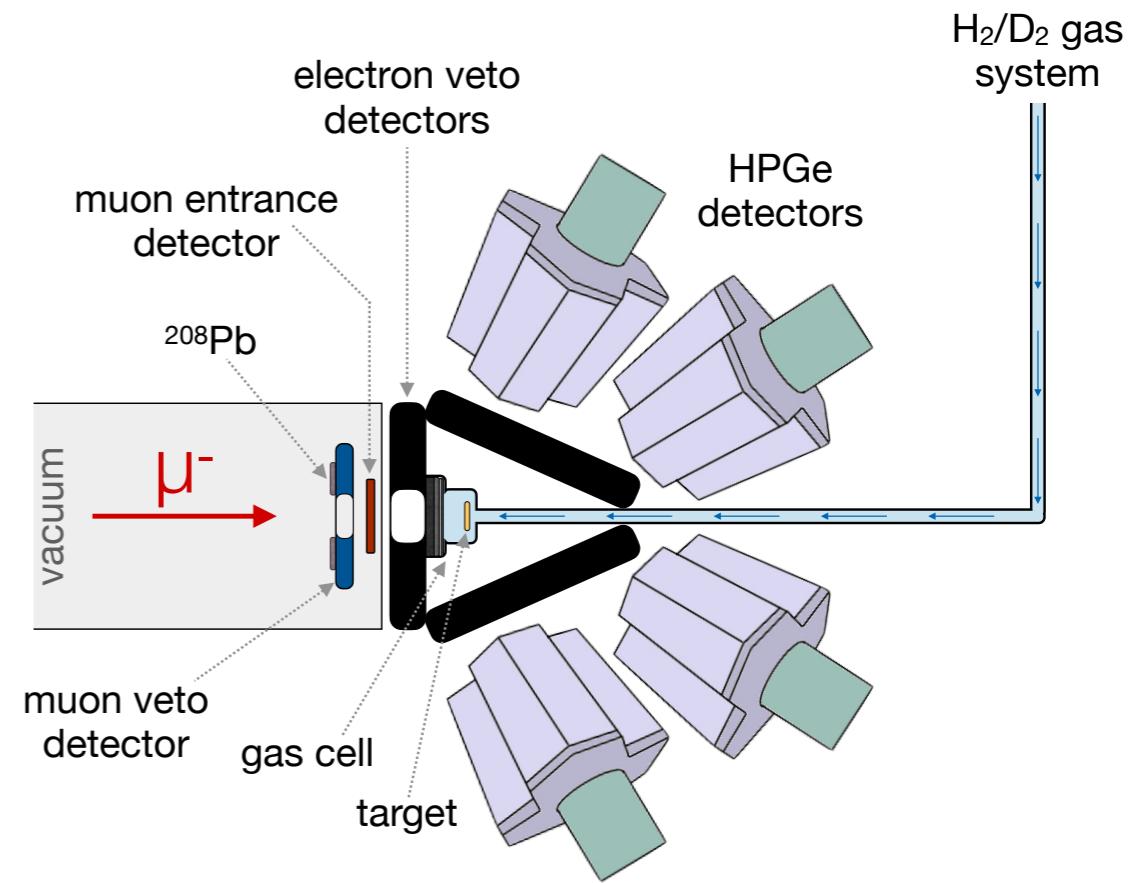
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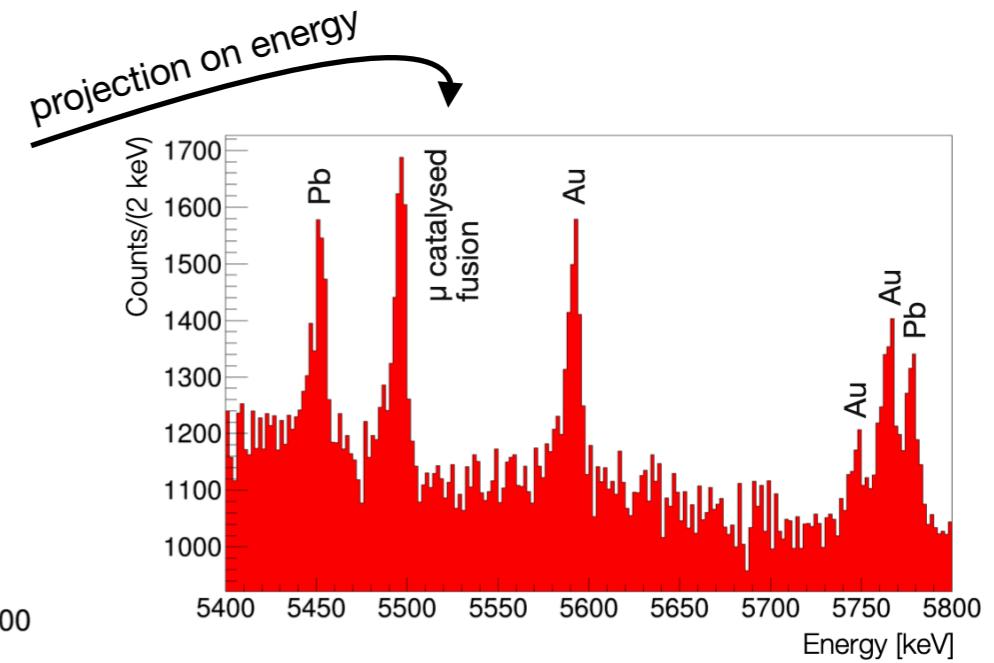
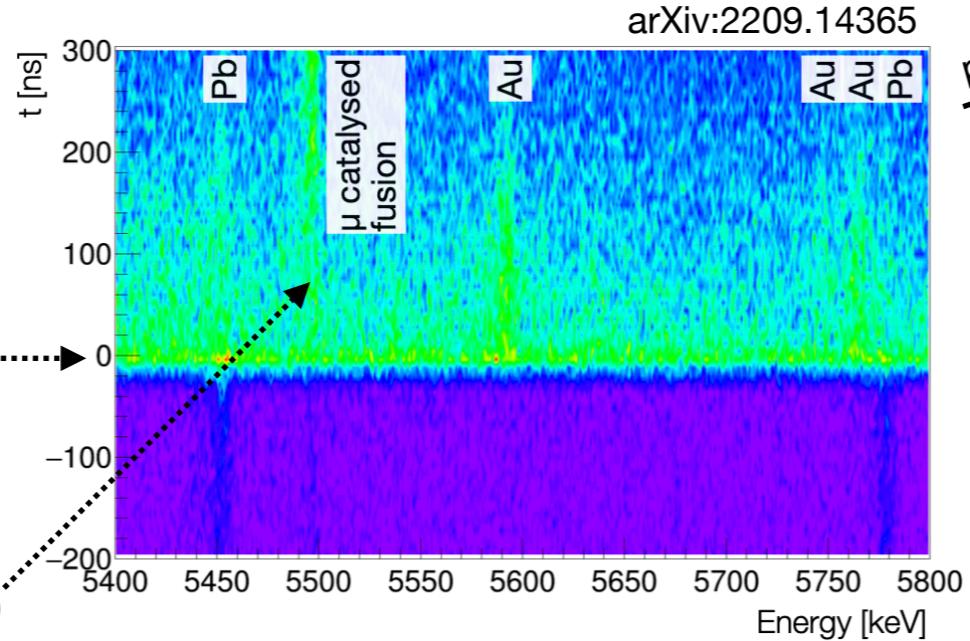
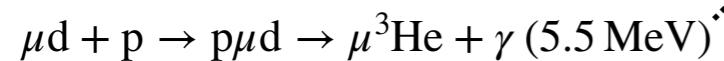
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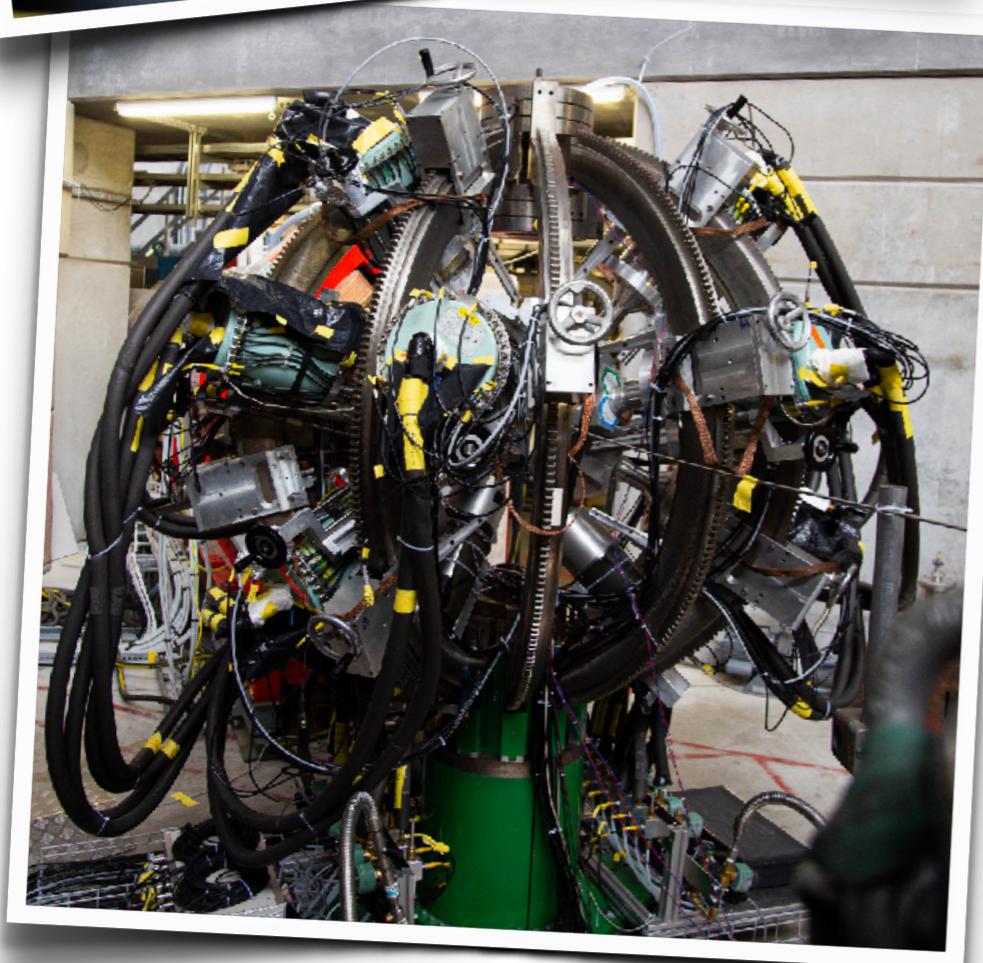
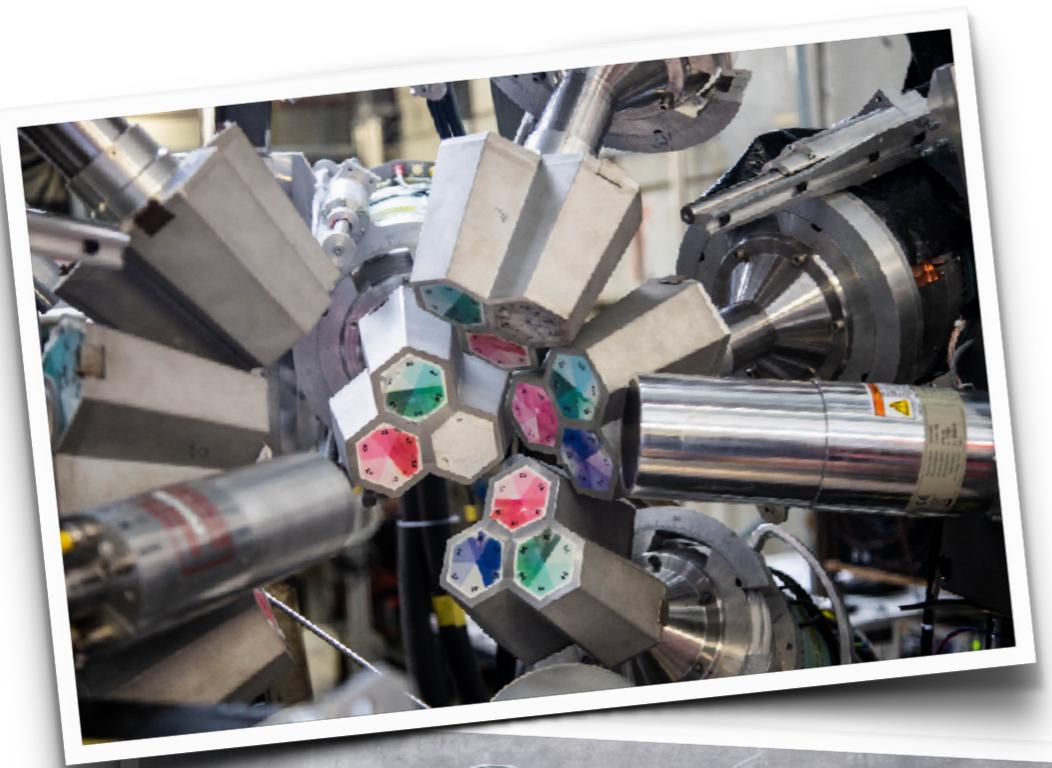
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Demonstration of principle in a 5 µg gold target in 2017, 18.5 h of measurement

the time the muon enters the gas cell

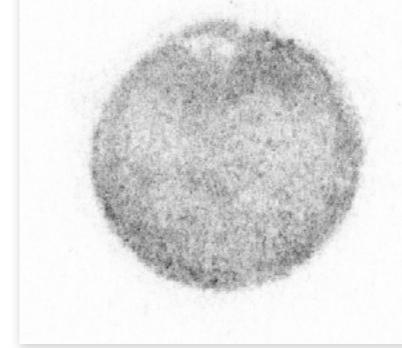
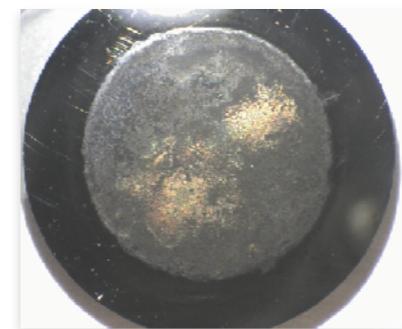


2019 measurement of ^{248}Cm and ^{226}Ra



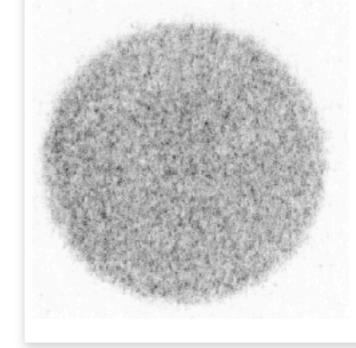
- 8 Miniball¹ germanium clusters and 2 standalone germanium detectors making a total of 26 HPGe crystals were operating ¹Eur. Phys. J. A 49, 40 (2013)
- Radiation protection restrictions at PSI allow for 16 µg of ^{248}Cm and 5.5 µg of ^{226}Ra
- Targets were produced by the radiochemistry group of the University of Mainz

^{248}Cm target



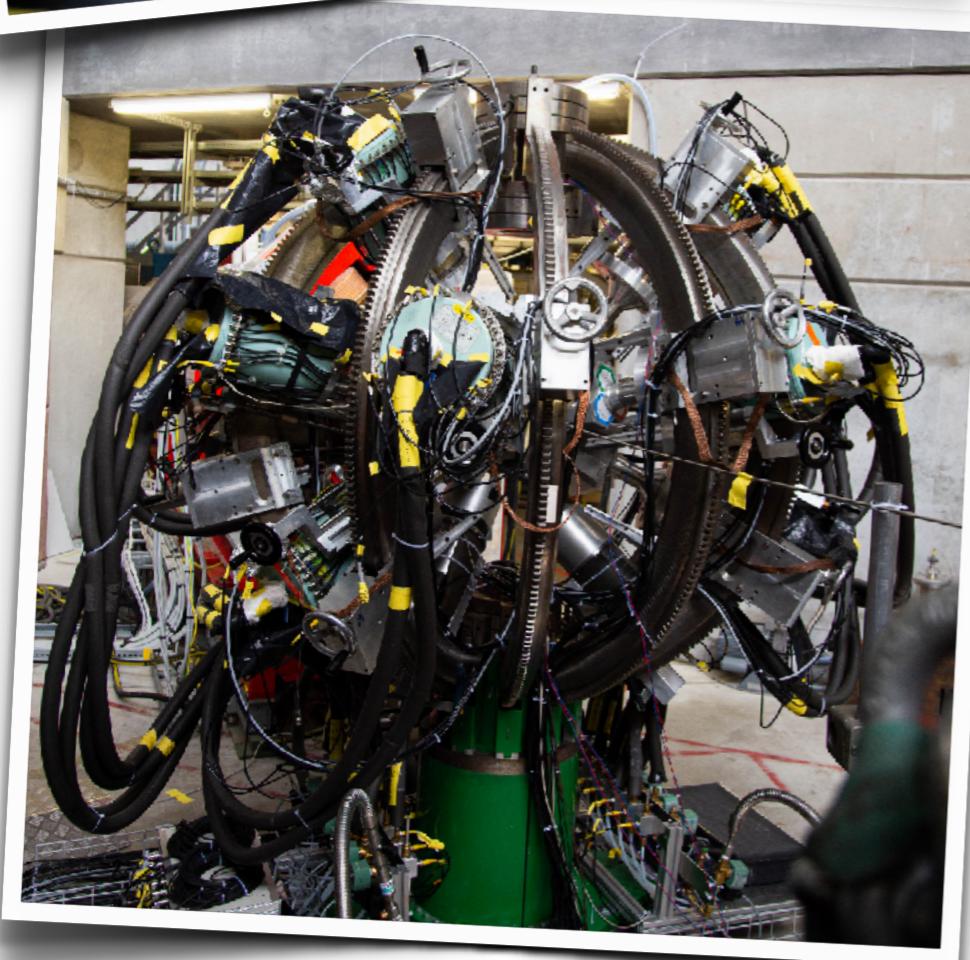
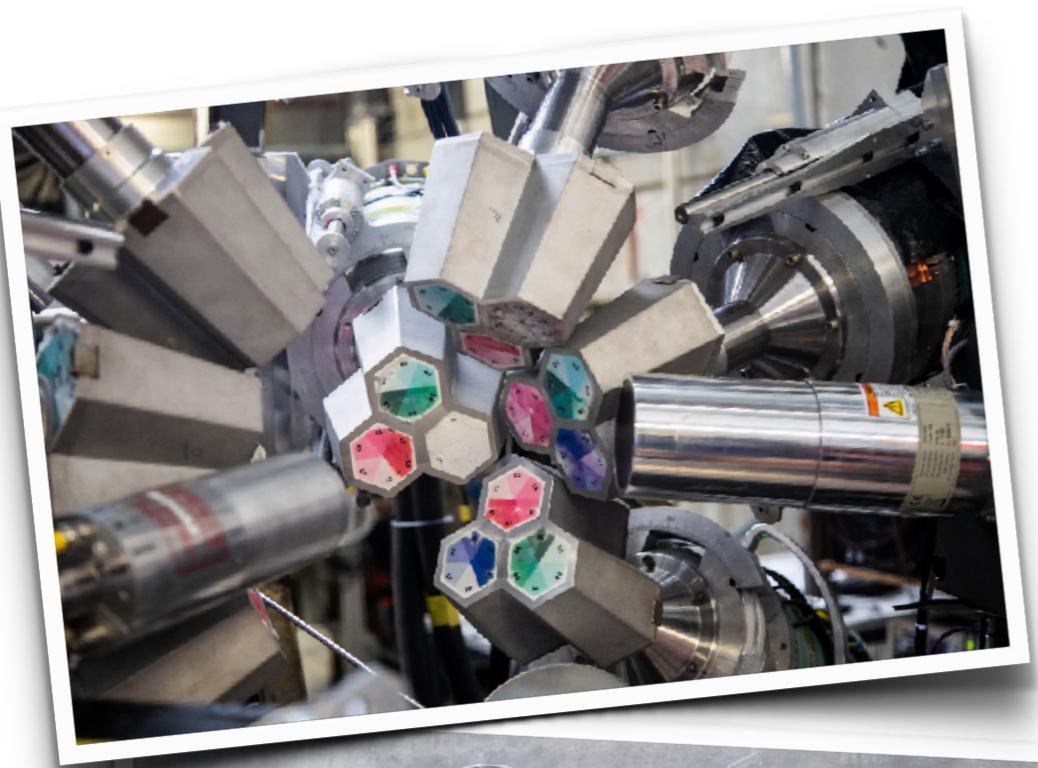
15.46 µg, uniformly distributed

^{226}Ra targets



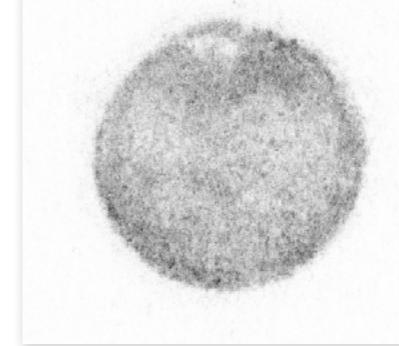
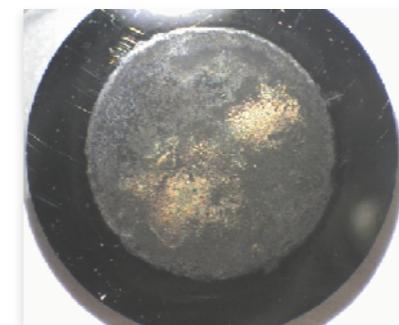
1.35 µg, uniformly distributed 4.37 µg, ring structure

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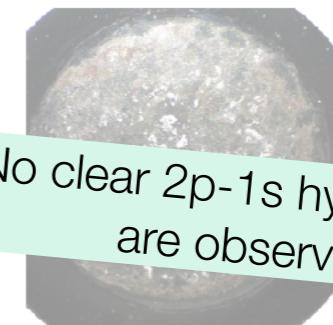
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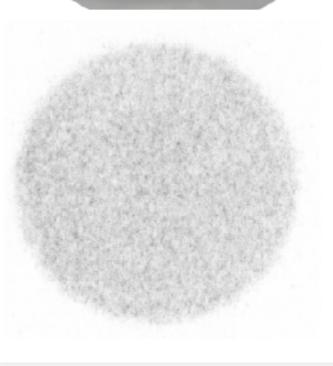


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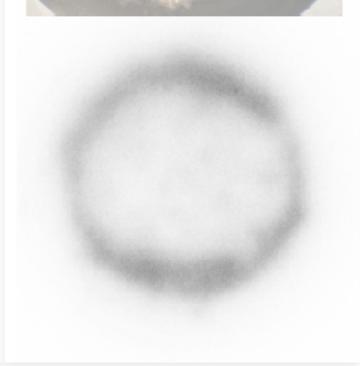
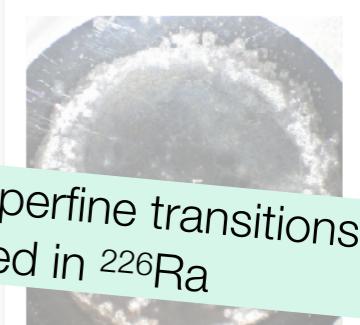
^{226}Ra targets



No clear 2p-1s hyperfine transitions are observed in ^{226}Ra



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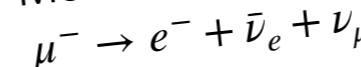


How to get to good S/B ratio at high energies?

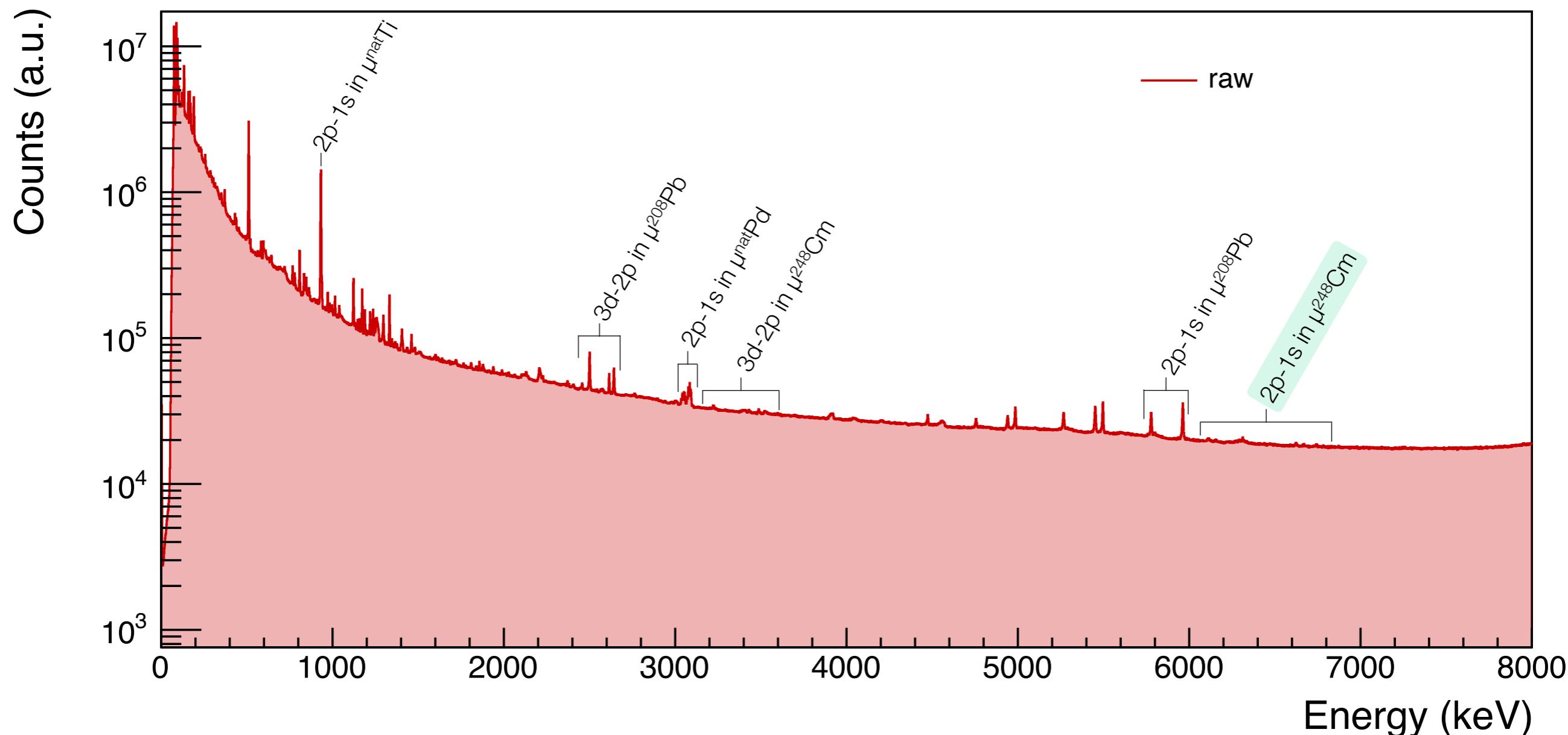
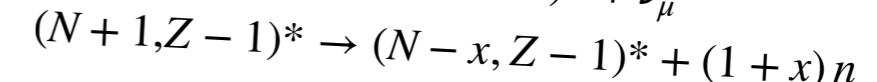
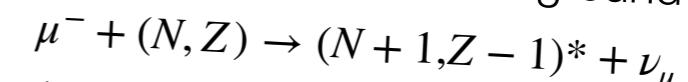
Cuts included:

- Electron veto anti-coincidences
- Clustering of events using the Miniball geometry
- Gain stabilisation
- Baseline correction

Most electron background is removed



Remaining neutron background

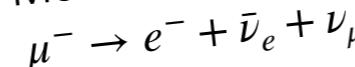


How to get to good S/B ratio at high energies?

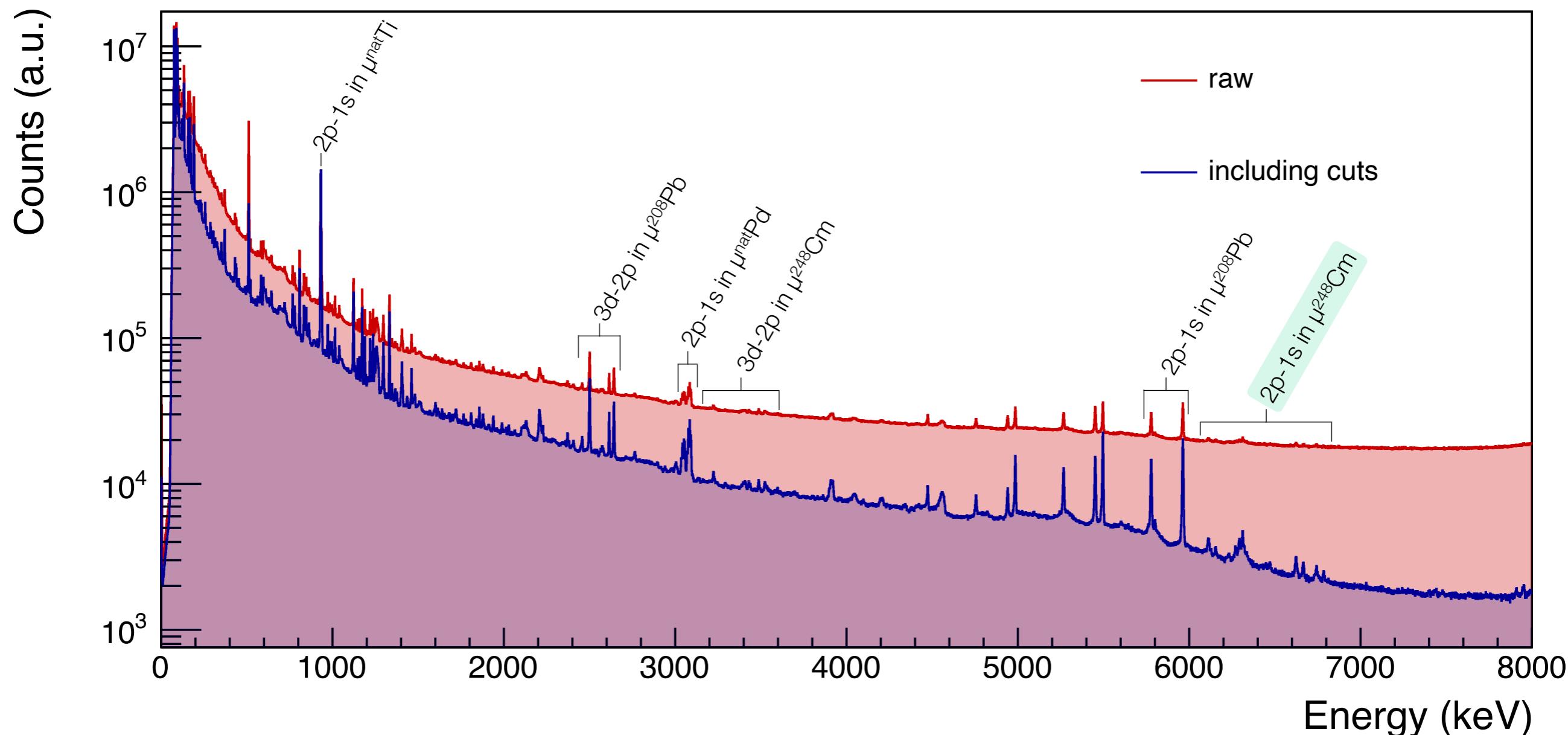
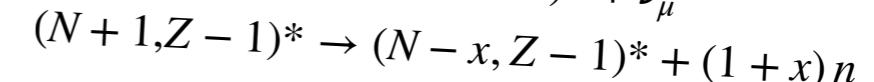
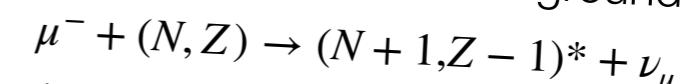
Cuts included:

- Electron veto anti-coincidences
- Clustering of events using the Miniball geometry
- Gain stabilisation
- Baseline correction

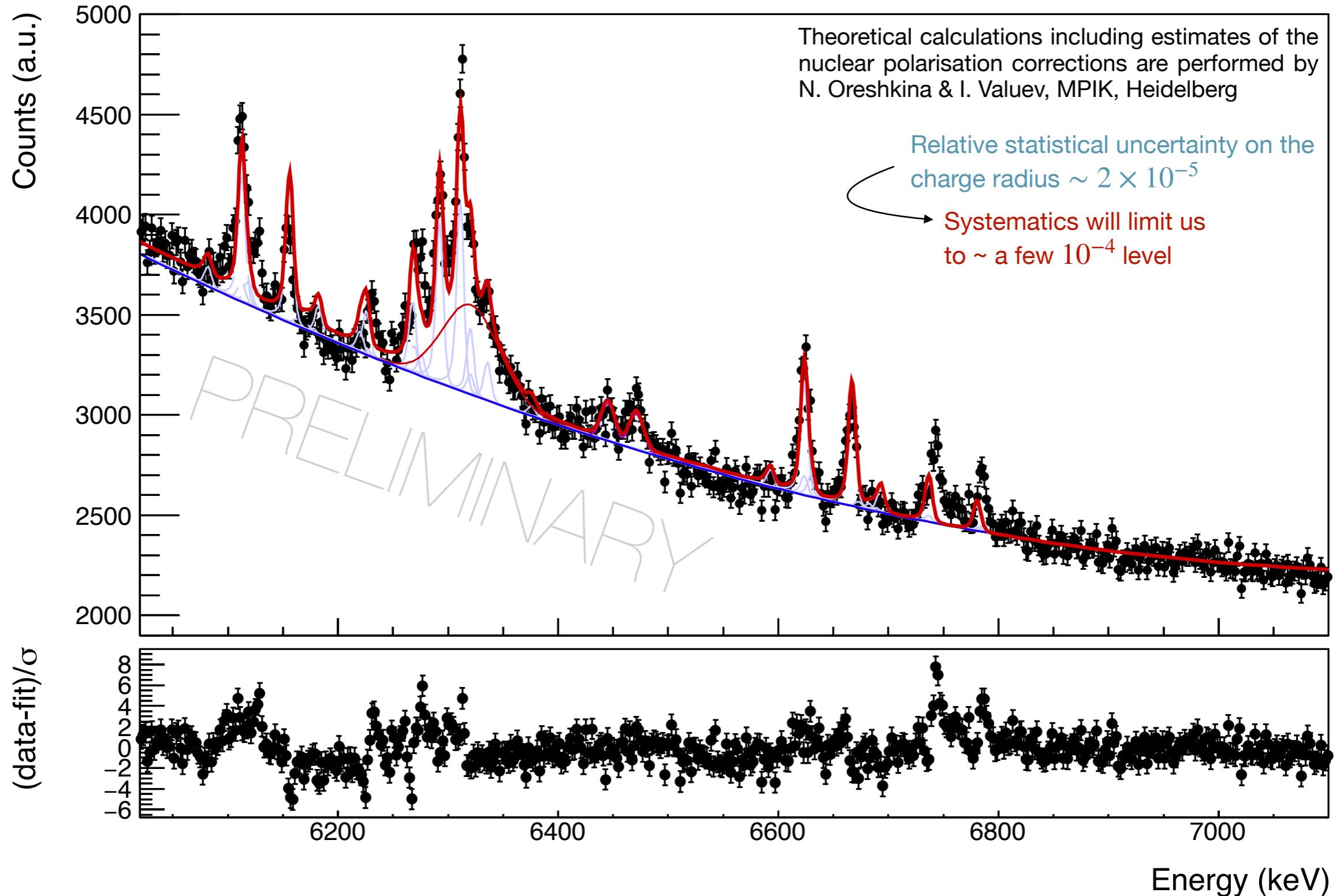
Most electron background is removed



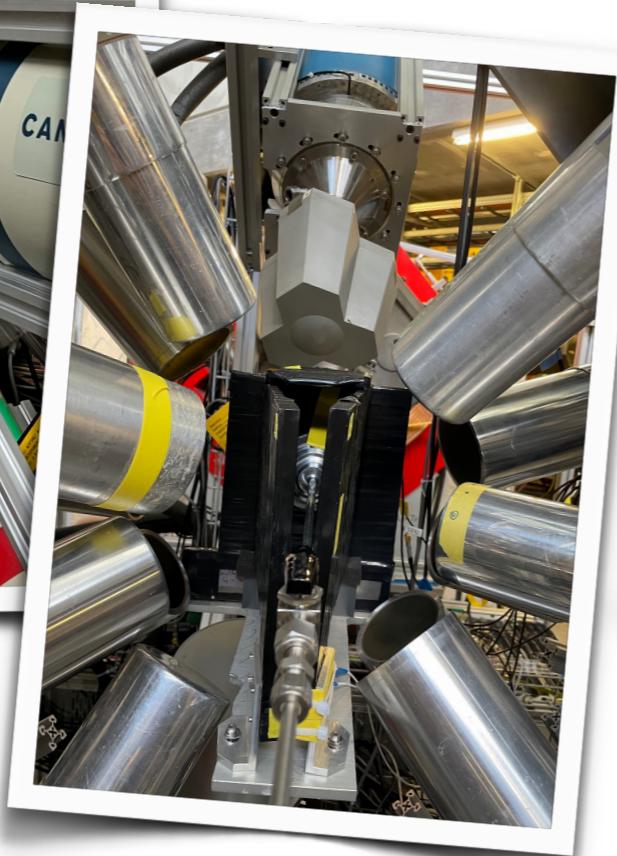
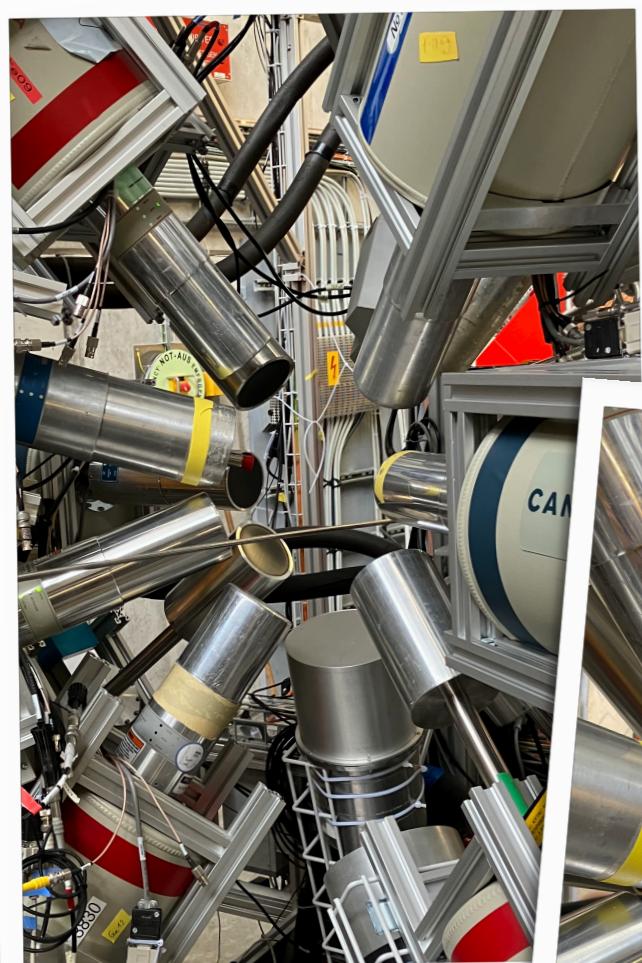
Remaining neutron background



$2p \rightarrow 1s$ hyperfine transitions in ^{248}Cm



September 2022 measurement



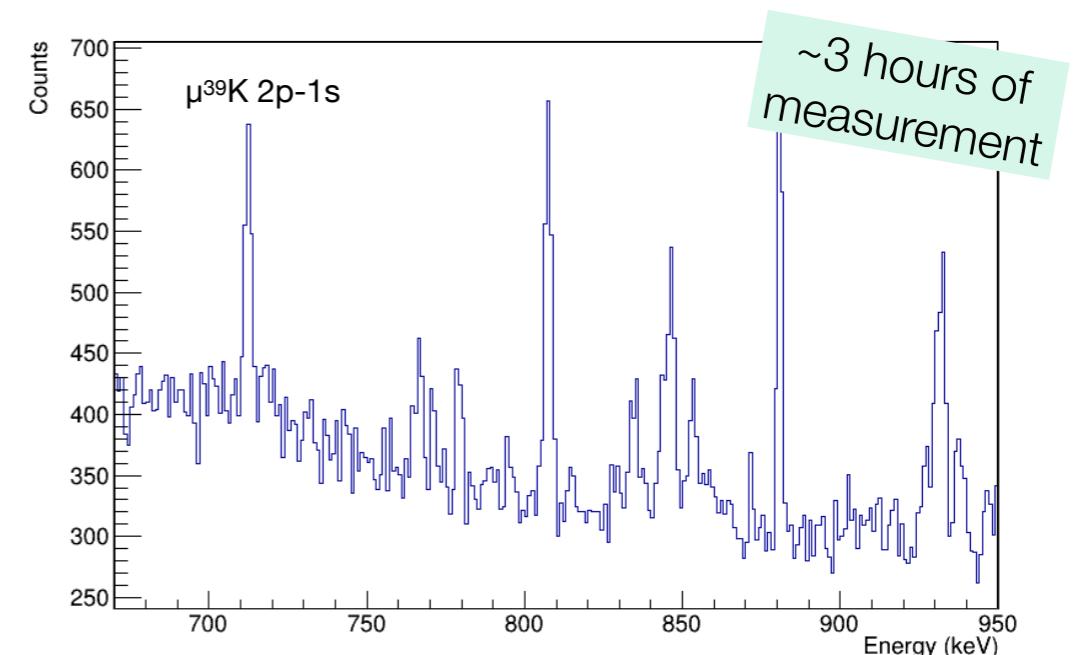
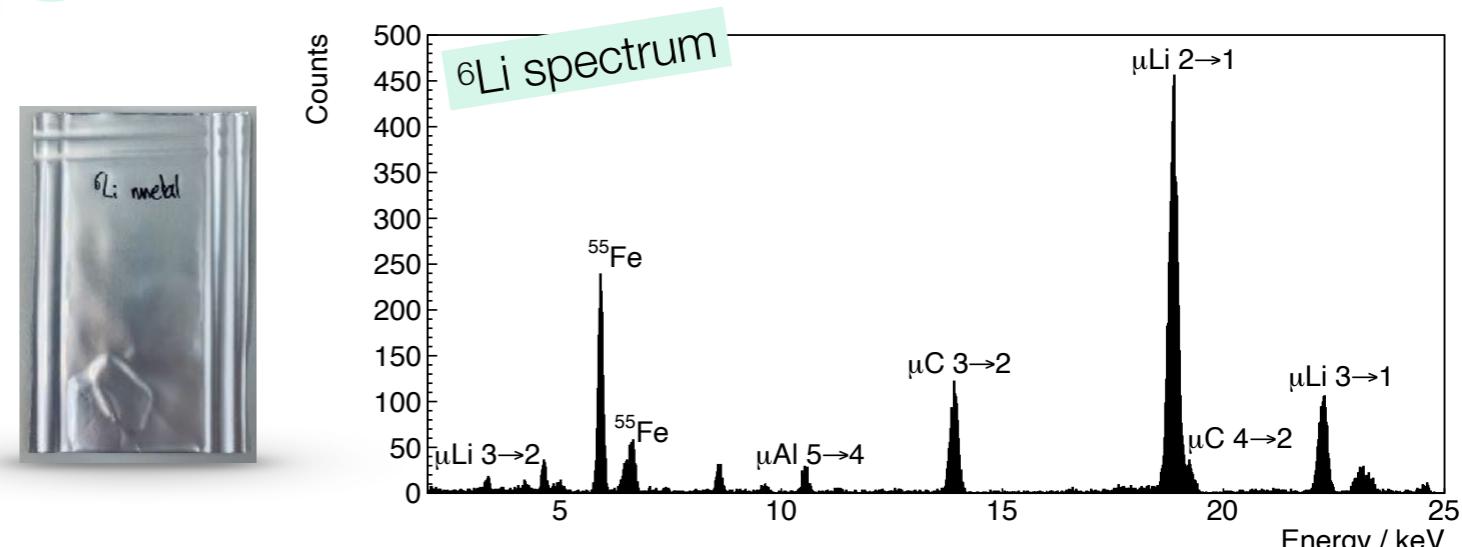
A

Towards the measurement of ^{226}Ra :

- measured implanted gold targets
- measured barium deposition targets (drop-on-demand, molecular plating)

B

First **low-Z** elements tests with enriched ^6Li and $^{\text{nat}}\text{Li}$ samples using a silicon drift detector



C

Expand towards **medium-Z** elements:

- measured ^{39}K (5.7 μg) using the muX method
- future goal: measure triplet ^{39}K , ^{40}K , ^{41}K
- improve laser spectroscopy results

What's next?

- Extraction of the nuclear charge radius of $^{185,187}\text{Re}$
- Determination of the ^{248}Cm nuclear charge radius
- Remeasure ^{226}Ra

New experimental proposals:

- Measure ^{39}K , ^{40}K , ^{41}K to improve laser spectroscopy results
- Measure low-Z elements (up to neon) using high resolution x-ray detectors to improve the charge radius results

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see more on
Tuesday 18/10 at
4-5.30 pm

Microcalorimetric high-resolution spectroscopy of muonic lithium
WHGA / Foyer and Tent, PSI

Katharina von Schoeler
16:16 - 16:17

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The 2s-1s transition in muonic atoms and atomic parity violation

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Frederik Wauters



16:06 - 16:07

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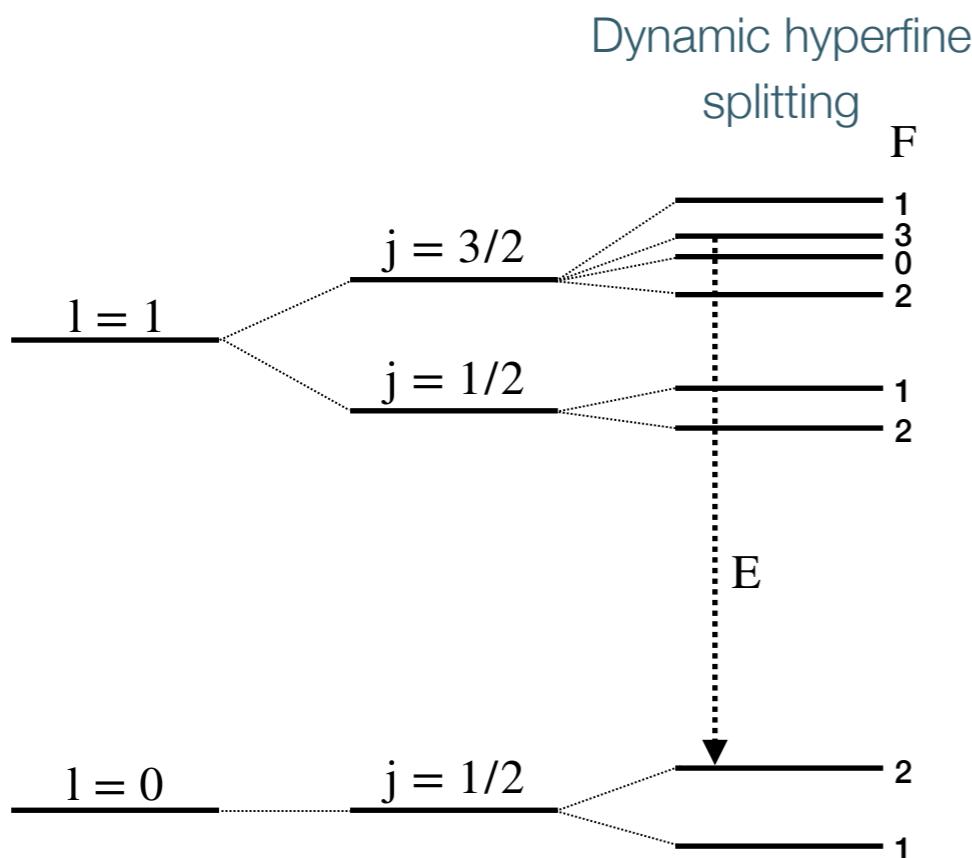
Frederik Wauters

16:06 - 16:07

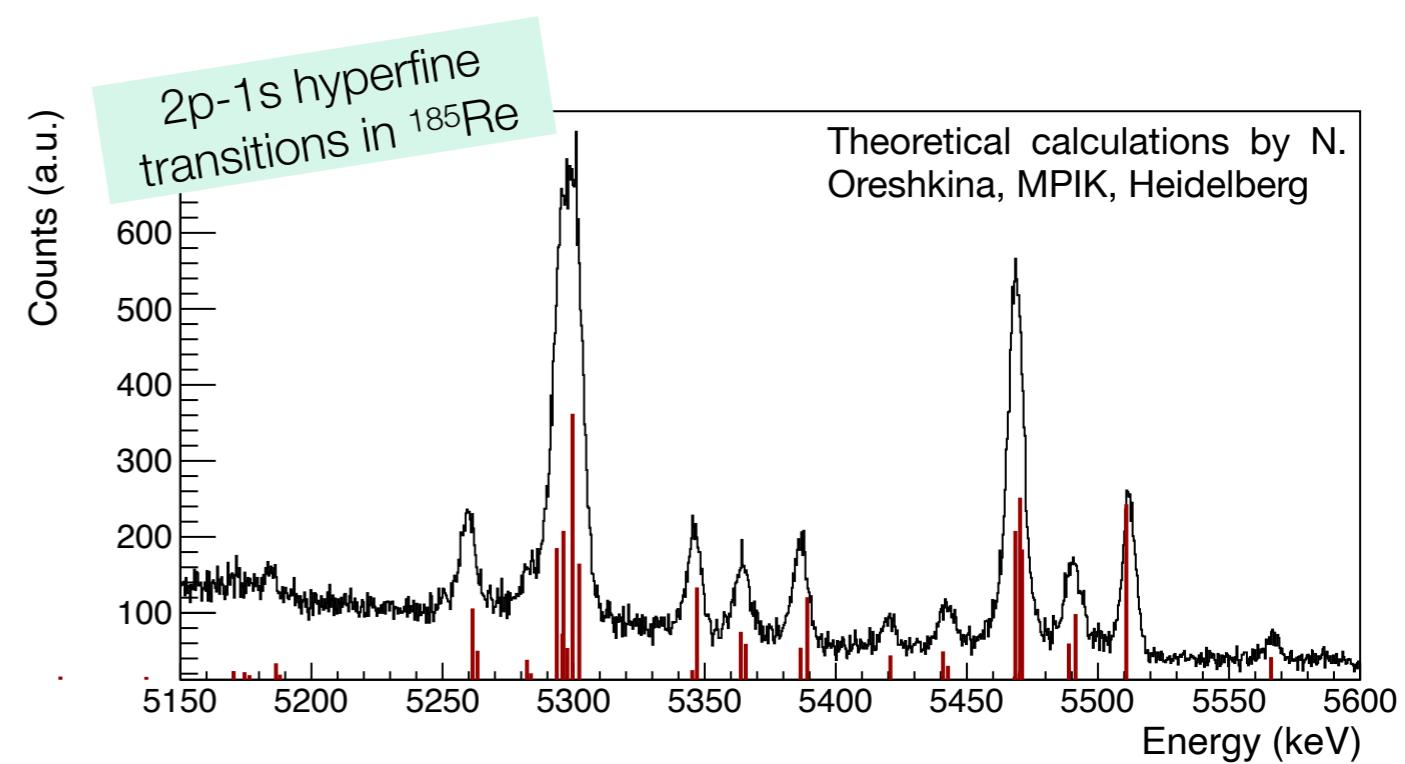
Thank you!

Backup Slides

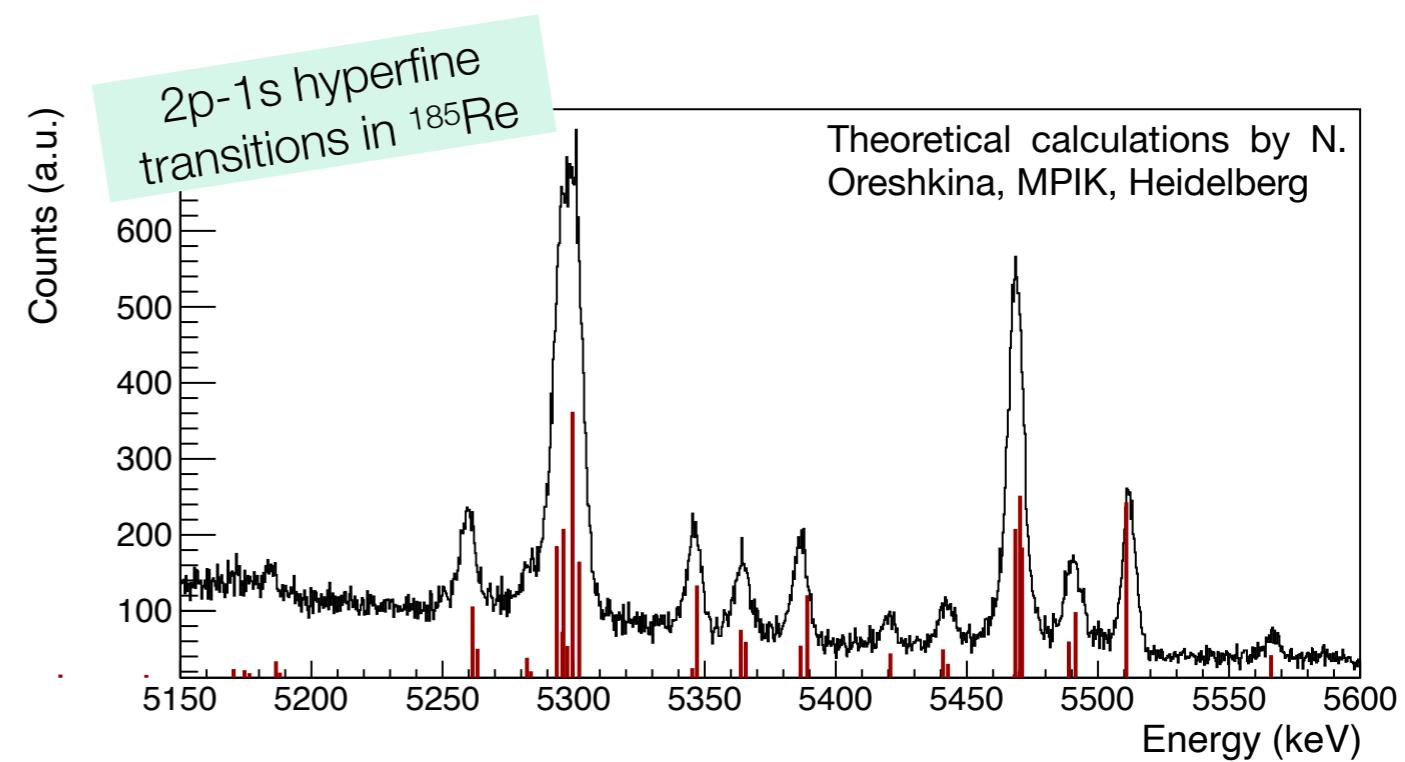
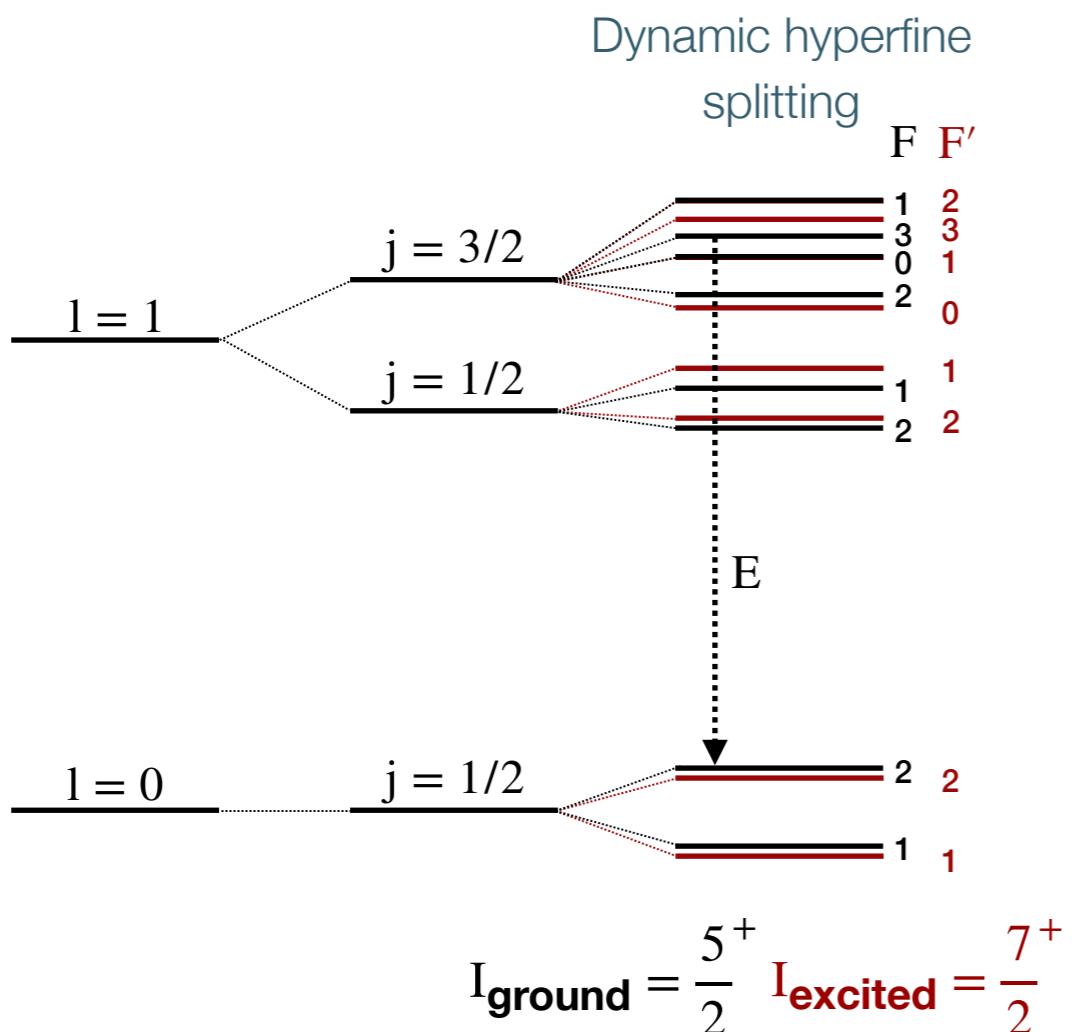
Dynamic effect in ^{185}Re



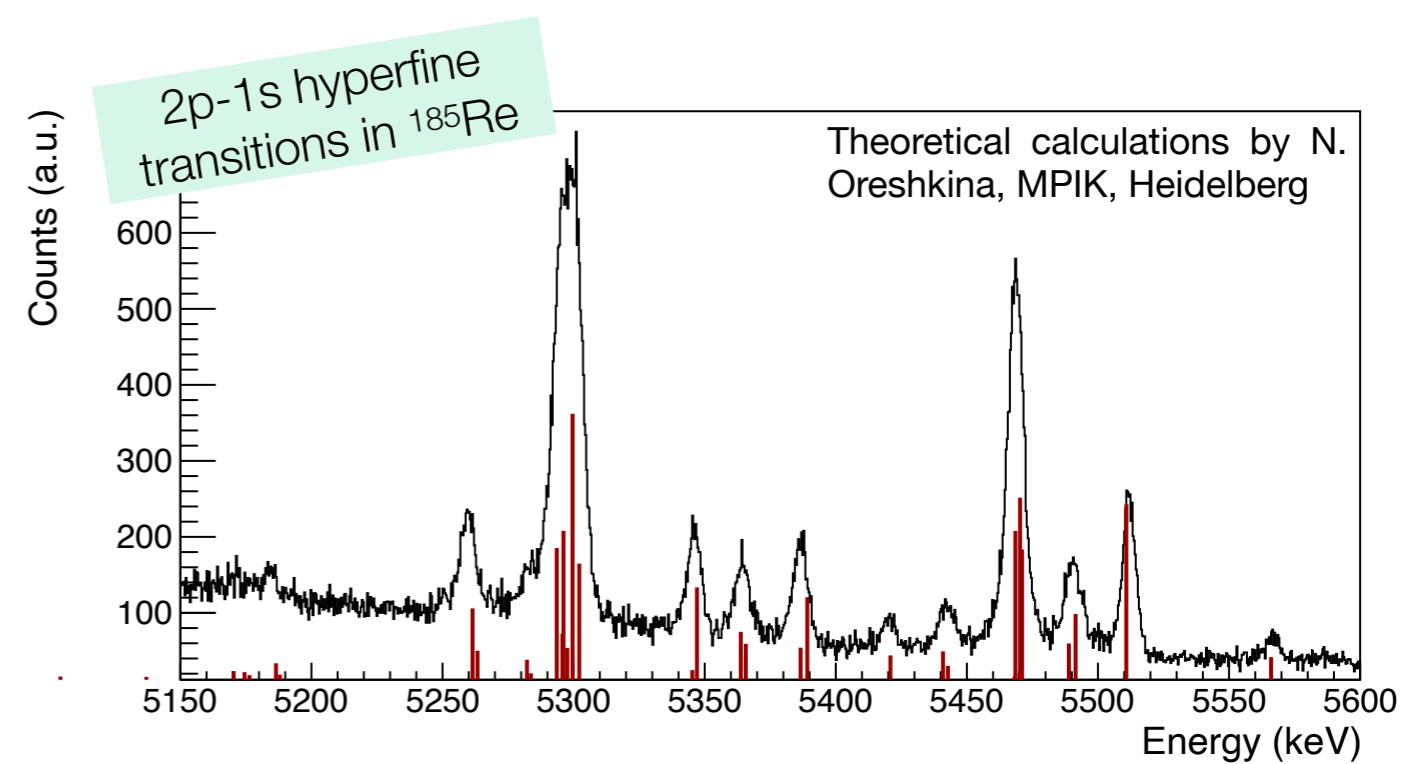
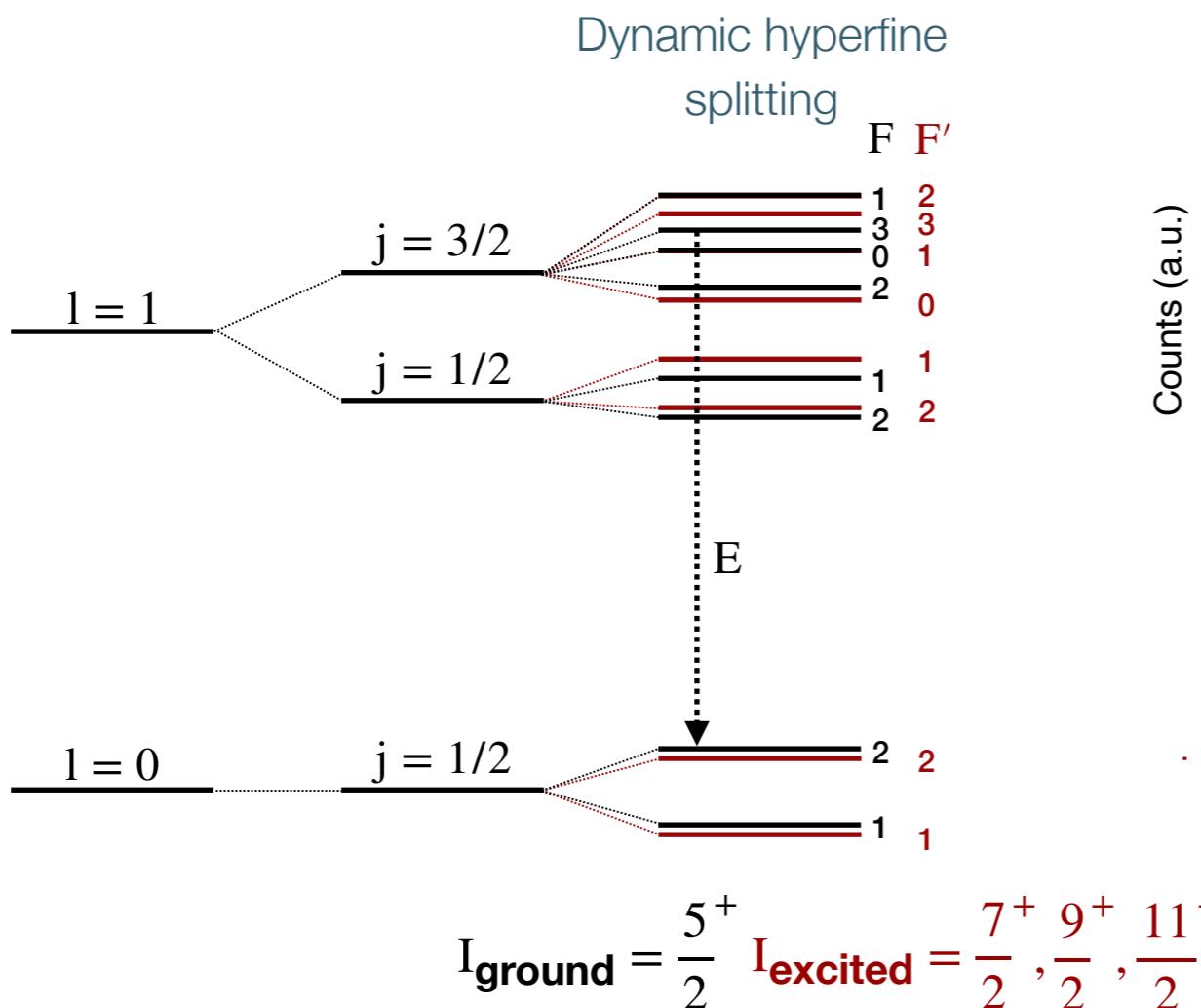
$$I_{\text{ground}} = \frac{5}{2}^+$$



Dynamic effect in ^{185}Re



Dynamic effect in ^{185}Re



Theory limitations

- Charge radius can be extracted with excellent relative precision
For ^{208}Pb : $r_C = \langle r^2 \rangle^{\frac{1}{2}} = 5.5031(11) \text{ fm}$
 $\Rightarrow 2 \cdot 10^{-4}$ relative precision

- Experimental accuracy at the level of $\sim 0.1 \text{ keV}$

Phys. Rev. C 37, 2821 (1988)

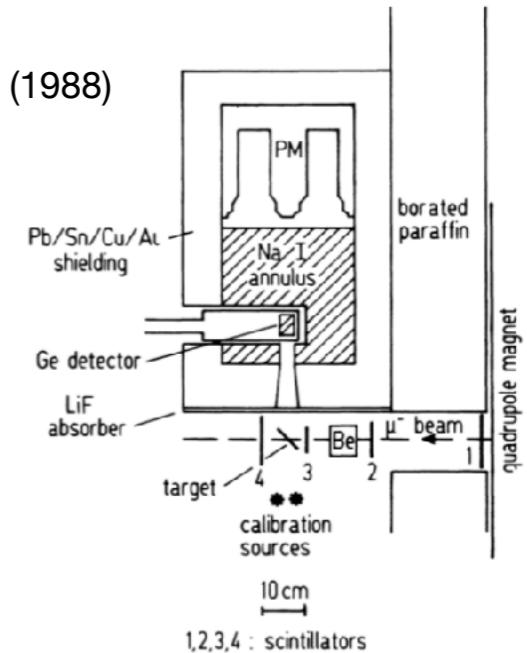


TABLE V. Experimental muonic transition energies (keV) in ^{208}Pb (recoil corrected).

Transition	Kessler (Ref. 9)	Hoehn (Ref. 27)	This experiment
$2p_{3/2}-1s_{1/2}$	5 962.770(420)		5 962.854(90)
$2p_{1/2}-1s_{1/2}$	5 777.910(400)		5 778.058(100)
$3d_{3/2}-2p_{1/2}$	2 642.110(60)	2642.292(23)	2 642.332(30)
$3d_{5/2}-2p_{3/2}$	2 500.330(60)	2500.580(28)	2 500.590(30)
$3d_{3/2}-2p_{3/2}$	2 457.200(200)		2 457.569(70)

- Limitation from theory \Rightarrow nuclear polarisation effect
- Due to the electrostatic interaction of the muon and nucleus the system is excited to virtual excited states resulting in the increase of the binding energy of the levels
- Limited knowledge of the highly excited nuclear states \Rightarrow nuclear polarisation determines the charge radius accuracy

TABLE II. Theoretical nuclear polarization corrections in ^{208}Pb .

Energy (MeV)	I^π	$B(E\lambda)\uparrow$ ($e^2 b^{2\lambda}$)	$1s_{1/2}$ (eV)	$2s_{1/2}$ (eV)	$2p_{1/2}$ (eV)	$2p_{3/2}$ (eV)	$3p_{1/2}$ (eV)	$3p_{3/2}$ (eV)	$3d_{3/2}$ (eV)	$3d_{5/2}$ (eV)
2.615	3^-	0.612	135	12	90	84	26	26	111	-63
4.085	2^+	0.318	198	20	182	180	76	84	6	4
4.324	4^+	0.155	14	1	8	7	2	2	1	1
4.842	1^-	0.001 56	7	1	-9	-8	0	0	1	1
5.240	3^-	0.130	27	2	16	15	5	5	2	2
5.293	1^-	0.002 04	9	2	-27	-19	0	-1	1	1
5.512	1^-	0.003 80	16	3	-90	-53	-1	-1	1	1
5.946	1^-	0.000 07	0	0	3	-30	0	0	0	0
6.193	2^+	0.050 5	29	3	22	21	7	7	0	0
6.262	1^-	0.000 24	1	0	3	5	0	0	0	0
6.312	1^-	0.000 22	1	0	3	4	0	0	0	0
6.363	1^-	0.000 14	1	0	2	2	0	0	0	0
6.721	1^-	0.000 75	3	1	6	7	0	-1	0	0
7.064	1^-	0.001 56	6	1	9	11	-1	-1	0	0
7.083	1^-	0.000 75	3	1	4	5	-1	-1	0	0
7.332	1^-	0.002 04	8	1	10	11	-2	-2	0	0
Total low-lying states			458	48	233	242	111	117	123	-53
13.5	0^-	0.047 872	906	315	64	38	24	15	1	0
22.8	0^+	0.043 658	546	147	43	26	15	10	0	0
13.7	1^-	0.537 672	1454	221	786	738	255	258	66	54
10.6	2^+	0.761 038	375	37	237	222	67	68	33	30
21.9	2^+	0.566 709	207	21	108	99	29	29	8	7
18.6	3^-	0.497 596	77	7	40	36	11	11	3	2
33.1	3^-	0.429 112	53	5	25	23	7	7	2	1
> 3 ^a			176	15	80	71	21	21	4	4
Total high-lying states			3794	768	1383	1253	429	419	117	98
Total			4252	816	1616	1495	540	536	240	45

^aValues from Ref. 7. Positive NP values mean that the respective binding energies are increased.

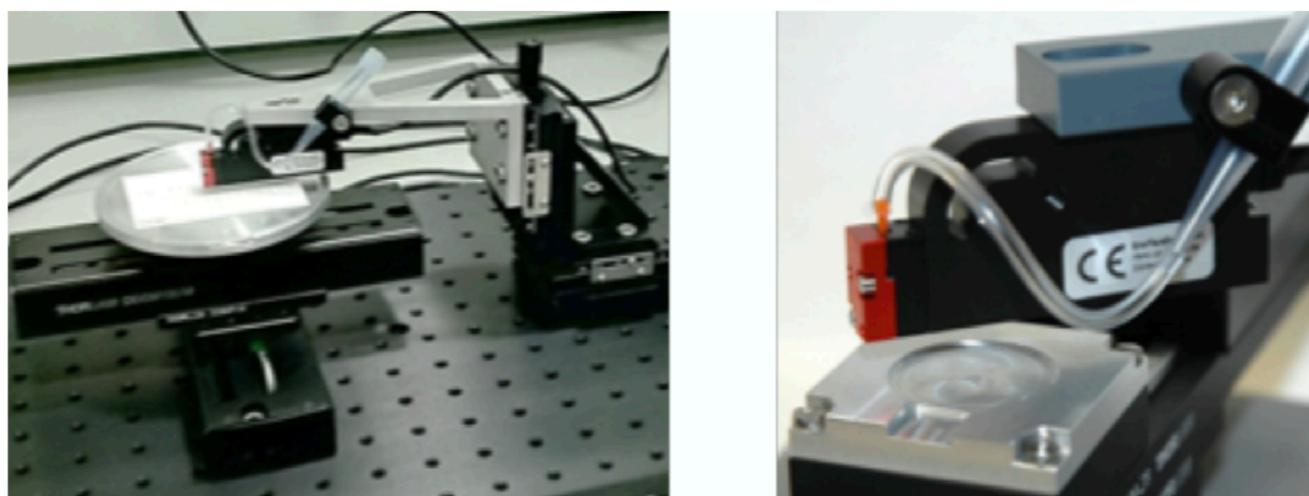
If nuclear polarisation calculations are improved,
more precise charge radius results can be extracted

Target production in 2019

Table 1: Summary of all targets produced for the muX beamtime 2019. The given production method ED refers to electrodeposition and DoD refers to drop on demand printing. The maximum allowed quantities in the experimental halle of PSI are 16 μg for ^{248}Cm (given the presence of ^{246}Cm) and 5.5 μg for ^{226}Ra . See text for more details.

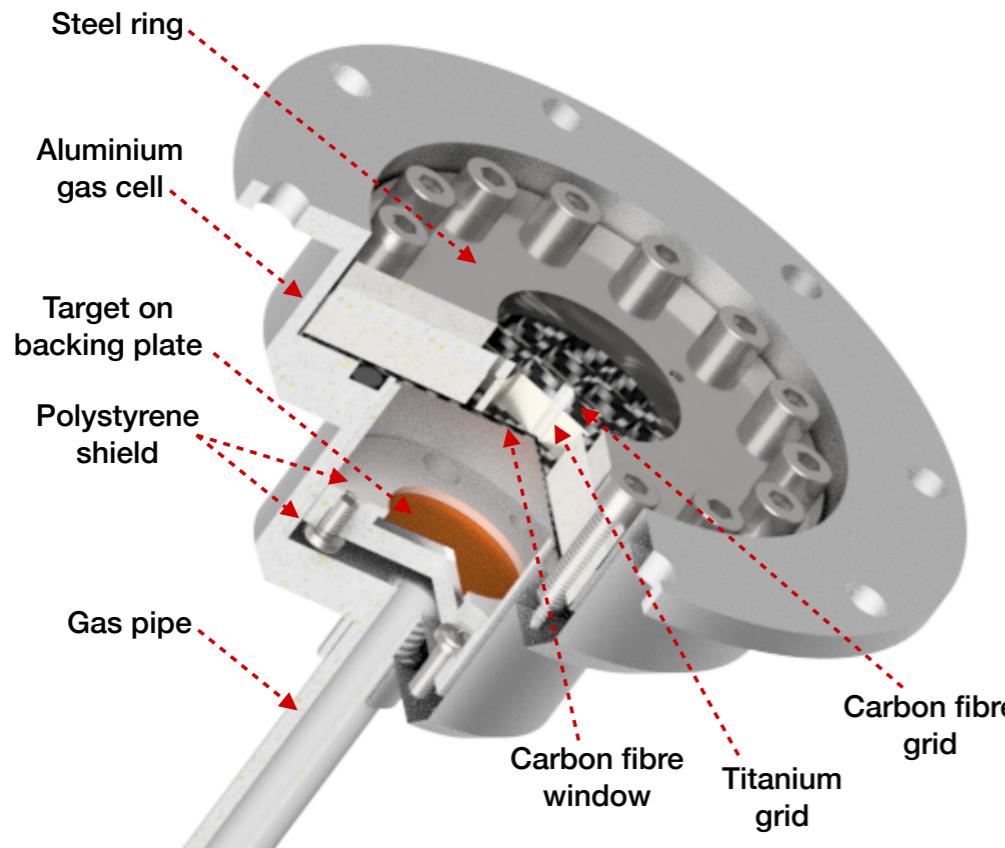
Target #	Nuclide	m / μg	A / kBq	Method
Cm-MP2-DoD	Cm-248	15.91	2.50	ED+DoD
Cm-MP3-DoD	Cm-248	15.46	2.43	ED+DoD
Cm-Aceton1	Cm-248	13.72	2.15	DoD (acetone)
Cm-Aceton2	Cm-248	13.94	2.19	DoD (acetone)
Ra-MP1	Ra-226	1.35	49.5	ED+DoD
Ra-MP2	Ra-226	2.50	91.6	ED
Ra-MP3	Ra-226	4.37	160.1	ED+DoD

muX Progress Report 2020 (2019)



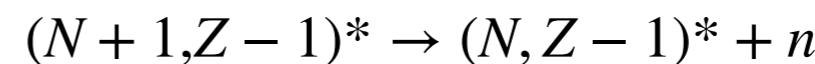
Drop-on-demand printing device at the Johannes Gutenberg-Universitat Mainz

Neutron background



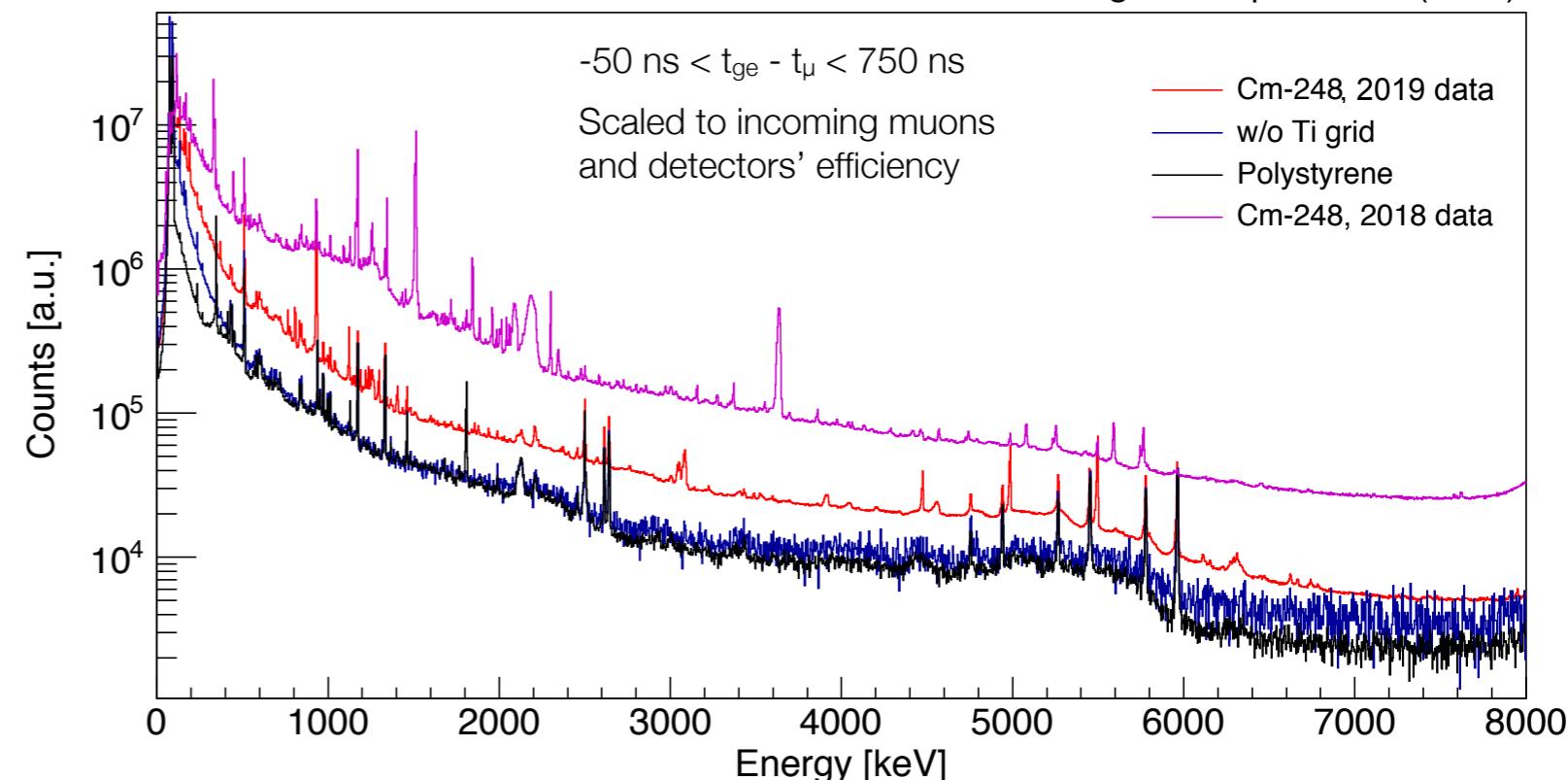
Beam time	2018	2019
Target backing material	copper	glassy carbon
Polystyrene shielding	No	yes

- Another source of background at ~6 MeV: neutrons emitted during the nuclear capture of the muon e.g.



- Those neutrons interact with the Ge crystal and create a continuous background of negative slope
- Effort for only low-Z material in the gas cell

muX Progress Report 2020 (2019)



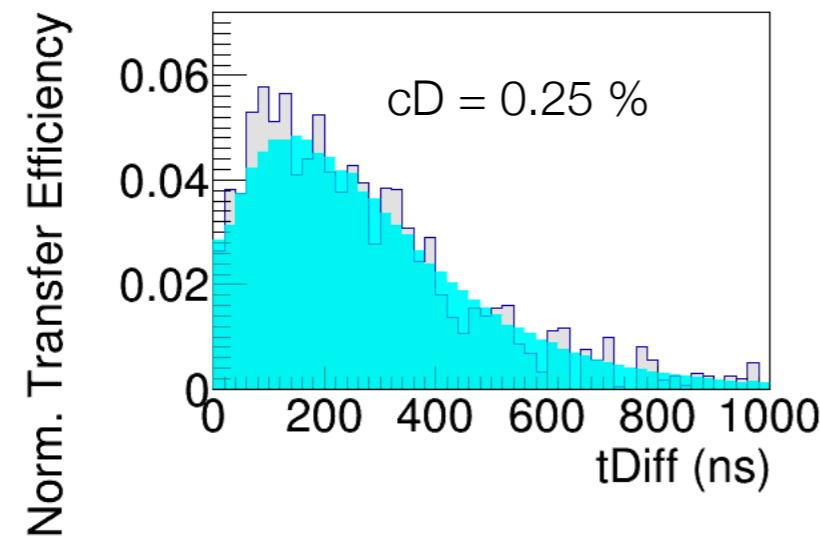
Optimisation of the x-ray yield using a gold target

- A 0.2 mg Au target was mounted inside the gas cell
- The amount of the 2p-1s μ Au x rays was measured by scanning the:
 - ▶ cD \rightarrow D₂ admixture in H₂ gas (cD)
 - ▶ p \rightarrow stopping position of the muon beam

↳

- Maximum yield for 27.25 MeV/c & 0.25% D₂
- Measurements and simulations are in good agreement

- The time distribution of the observed 2p_{1/2}-1s_{1/2} gold x ray after transfer for different D₂ concentrations was simulated



PhD thesis of A. Skawran
Simulations by J. Nuber

