

# Search for axion-like particles in muon decays

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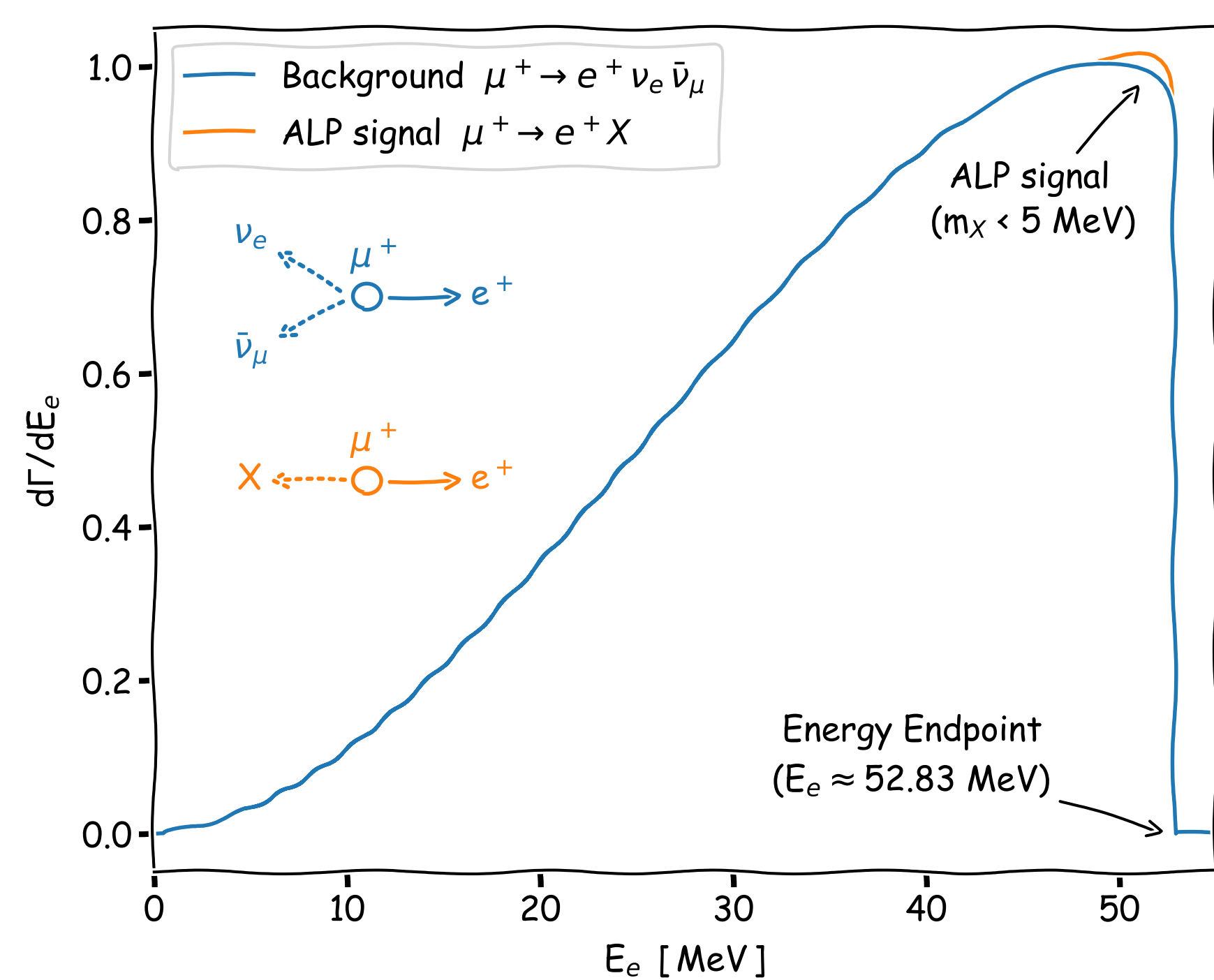


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## Motivations

- The search for **charged Lepton Flavour Violation (cLFV)** in rare **muon decays** is a key tool to test the Standard Model (SM).
- The **MEG II** ( $\mu^+ \rightarrow e^+ \gamma$ ) and **Mu3e** ( $\mu^+ \rightarrow e^+ e^- e^+$ ) experiments at PSI are competitive in searching for decays involving a light neutral boson  $X$ , which remains invisible.
- This particle can be an **Axion-Like Particle (ALP)** arising from the spontaneous breaking of a global  $U(1)$  symmetry.
- A possible process is the two-body decay  $\mu^+ \rightarrow e^+ X$ .
- Its signature is a monochromatic peak close to the endpoint of the positron spectrum of the  $\mu^+ \rightarrow e^+ \nu_e \bar{\nu}_\mu$  background.
- The theoretical uncertainty at the endpoint is enhanced by the emission of **soft photons**, reducing the signal sensitivity.
- The hunt for such an elusive signal requires extremely accurate theoretical predictions for simulations and data analysis.



## The MCMULE framework

- MCMULE: Monte Carlo for **MUons** and other **LEptons**.
- A numerical framework for the fully differential computation of higher-order **QED corrections** for decay and scattering processes involving leptons, mainly at low energy.
- The precision goal is Next-to-Next-to-Leading Order (**NNLO**).
- All divergences are treated with dimensional regularisation, while renormalisation is performed in OS scheme.
- Soft singularities are subtracted by using the **FKS<sup>2</sup> scheme**.
- Collinear singularities are eliminated by keeping all fermion masses at their physical value ( $m \neq 0$ ).
- Phase space is integrated with the adaptive VEGAS algorithm.
- For a process implemented in the code, the user can obtain any differential distribution with any cut, for example to reproduce detector acceptances or analysis selections.
- Both signal  $\mu^+ \rightarrow e^+ X$  and background  $\mu^+ \rightarrow e^+ \nu_e \bar{\nu}_\mu$  are implemented in MCMULE for ALP searches.

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## Theoretical predictions

- The positrons produced by **polarised** muon decays are fully characterised by the inclusive distribution

$$\frac{m_\mu}{2} \frac{d^2\Gamma}{dE_e d\cos\theta_e} = \frac{G_F^2 m_\mu^5}{192 \pi^3} \left[ F(E_e) + P_\mu \cos\theta_e G(E_e) \right]$$

$m_\mu$ : Muon mass       $G_F$ : Fermi constant  
 $E_e$ : Positron energy       $P_\mu$ : Muon polarisation  
 $\theta_e$ : Angle between  $e^+$  momentum and  $\mu^+$  spin  
 $F$ : Isotropic function  $\rightarrow$  Energy spectrum  
 $G$ : Anisotropic function  $\rightarrow$  Polarisation effect

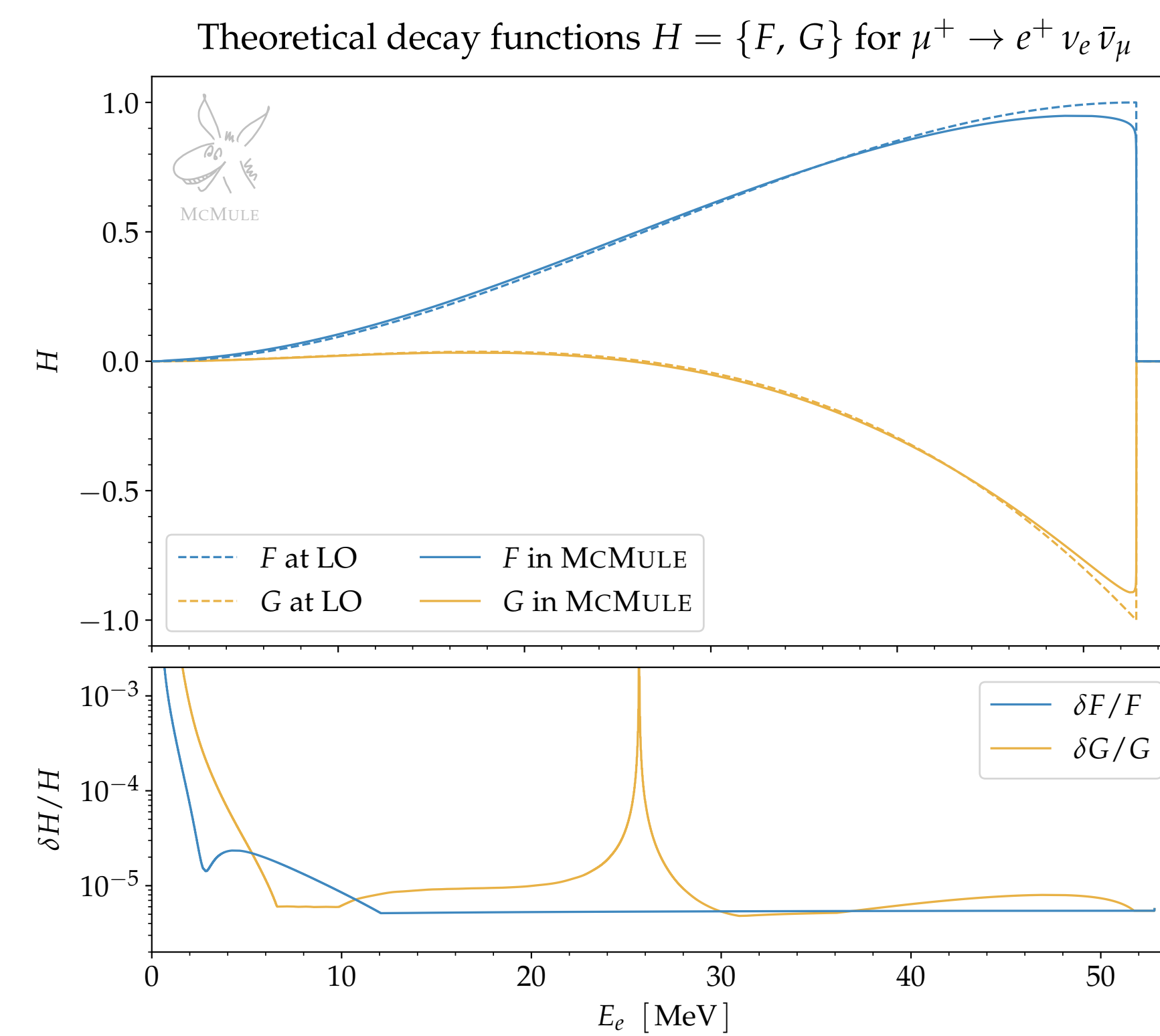
- The positron dynamic is therefore determined by the two dimensionless functions  $F(E_e)$  and  $G(E_e)$ .
- These functions have been computed with MCMULE for both  $\mu^+ \rightarrow e^+ X$  and  $\mu^+ \rightarrow e^+ \nu_e \bar{\nu}_\mu$  assuming **inclusive photons**.
- A generic muon polarisation is assumed and the positron mass is not neglected. The centre-of-mass frame is always used.

## Background $\mu^+ \rightarrow e^+ \nu_e \bar{\nu}_\mu$

- The background  $\mu^+ \rightarrow e^+ \nu_e \bar{\nu}_\mu$  is computed in Fermi theory

$$\mathcal{L}_F = -\frac{G_F}{\sqrt{2}} (\bar{\psi}_e \gamma^\rho (1 - \gamma^5) \psi_\mu) (\bar{\psi}_{\nu_e} \gamma_\rho (1 - \gamma^5) \psi_{\nu_\mu}) + \mathcal{L}_{\text{QED}}$$

- Exact QED corrections are added up to **NNLO**, including the (hadronic) vacuum polarisation effects.
- The **collinear logarithmic** terms  $\propto \log(m_e/m_\mu)$  are included up to N<sup>3</sup>LO with NLL accuracy, while the **soft logarithmic** terms  $\propto \log(1 + m_e^2/m_\mu^2 - 2E_e/m_\mu)$  are analytically resummed to all orders with NNLL accuracy.
- The leading **weak corrections** are considered as well.
- The resulting **theory error** for the positron energy spectrum is about  $5 \cdot 10^{-6}$ , the smallest achieved so far!

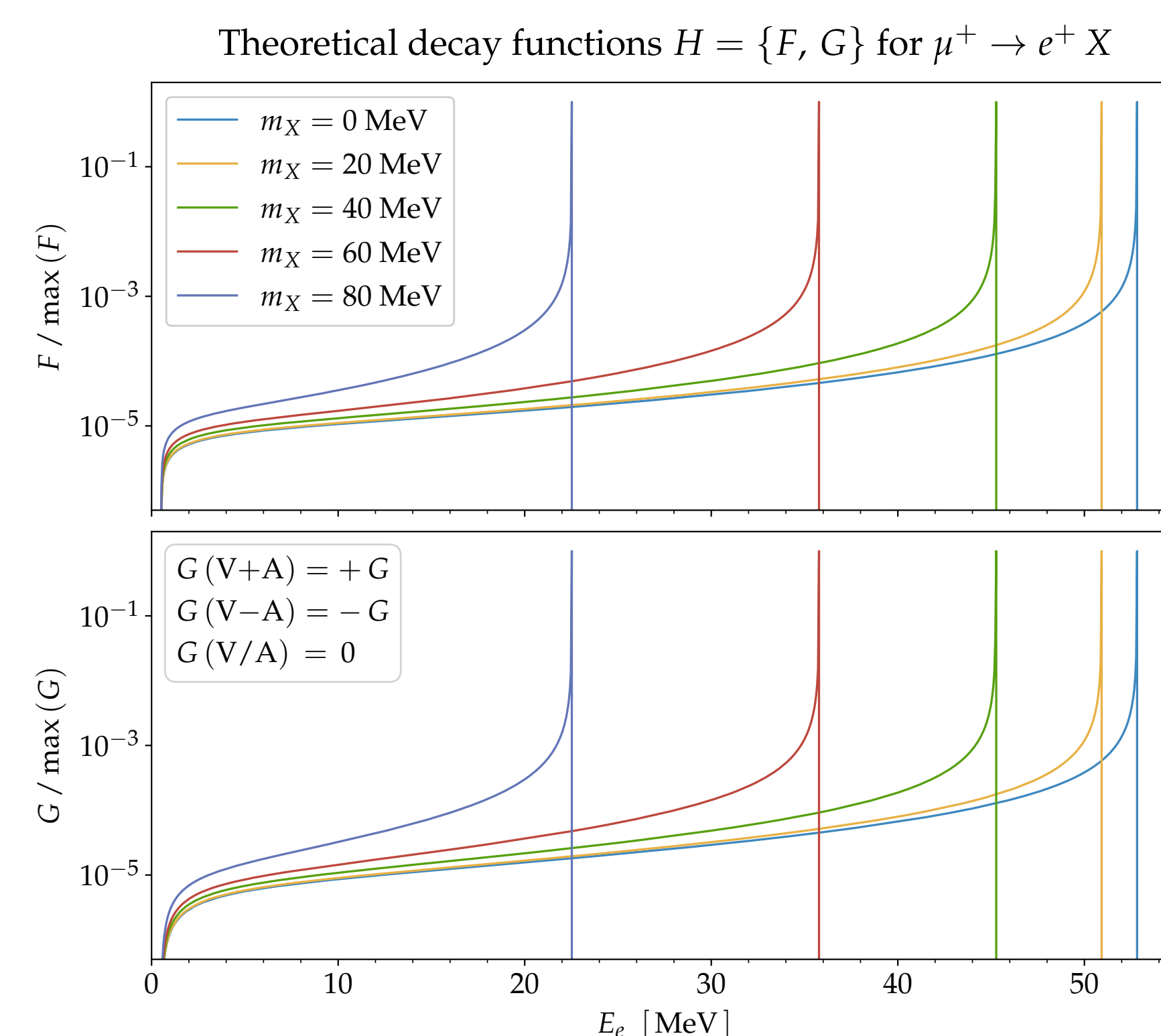


## Signal $\mu^+ \rightarrow e^+ X$

- The signal  $\mu^+ \rightarrow e^+ X$  is computed using an **effective model**, which accounts for different ALP masses and couplings

$$\mathcal{L}_X = \frac{1}{\Lambda} (\partial_\rho X) \bar{\psi}_e (\gamma^\rho g_V + \gamma^\rho \gamma^5 g_A) \psi_\mu + \mathcal{L}_{\text{QED}}$$

- The contribution is suppressed by a **large energy scale  $\Lambda$** .
- The coupling constants  $g_V$  and  $g_A$  can be chosen in order to obtain typical chiral structures, such as left-handed (**V-A**), right-handed (**V+A**), vector-like (**V**) or pseudovector-like (**A**).
- The QED corrections at NLO are included with the effect of introducing a **radiative tail** to the positron energy spectrum.



## Toy analysis

- The new predictions can be used to estimate the experimental sensitivity and evaluate the impact of the theory error.
- To this end a **simplified model** of the MEG II and Mu3e positron spectrometers has been defined

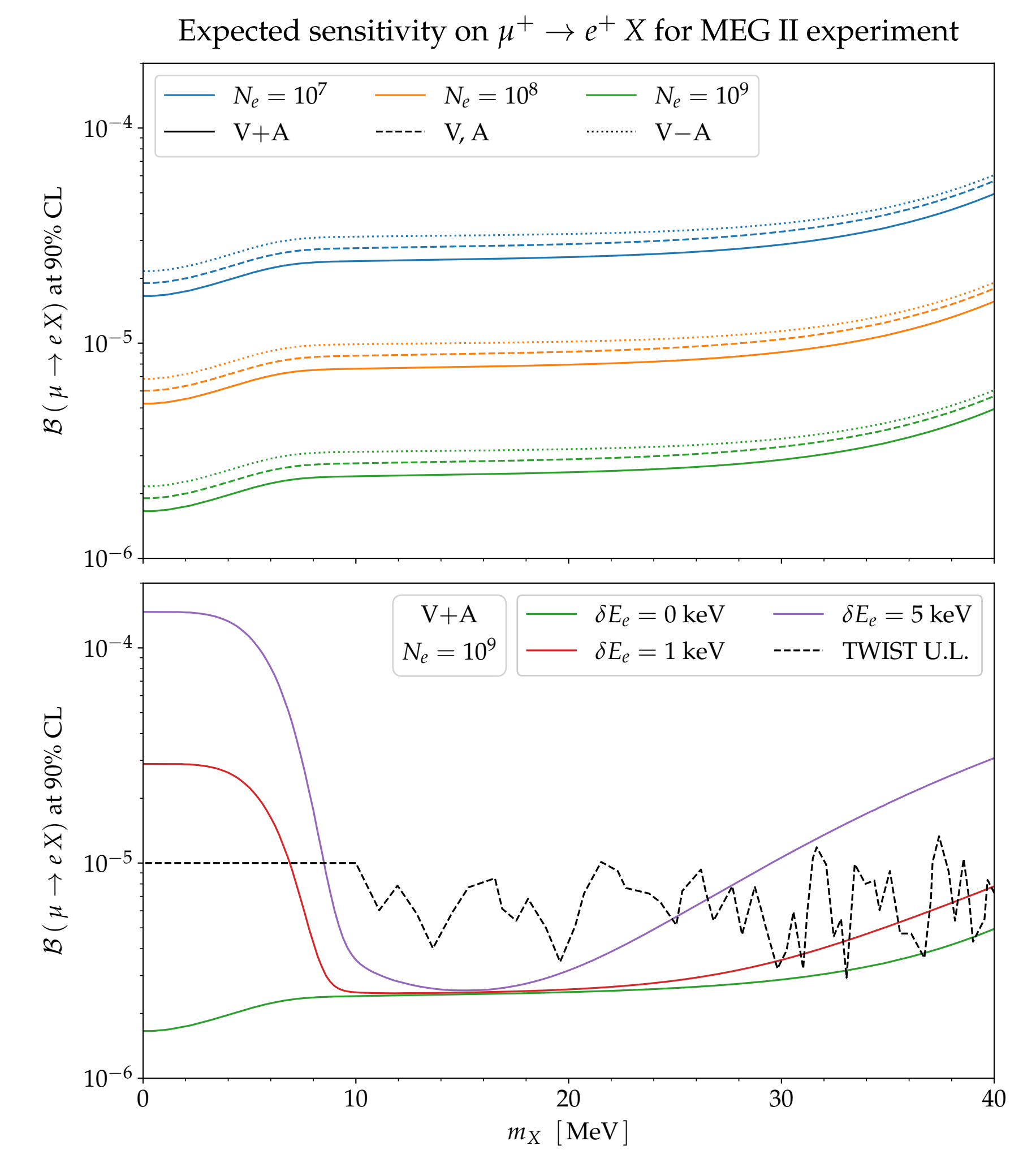
$$\mathcal{P}_e(E_e) = \int dE'_e \left[ \mathcal{H}_e(E'_e) \times \mathcal{A}_e(E'_e) \times \mathcal{S}_e(E_e - E'_e) \right]$$

$\mathcal{P}_e$ : **Expected** positron energy spectrum  
 $\mathcal{H}_e$ : **Theoretical** positron energy spectrum  
 $\mathcal{A}_e$ : Positron energy **acceptance** function  
 $\mathcal{S}_e$ : Positron energy **resolution** function

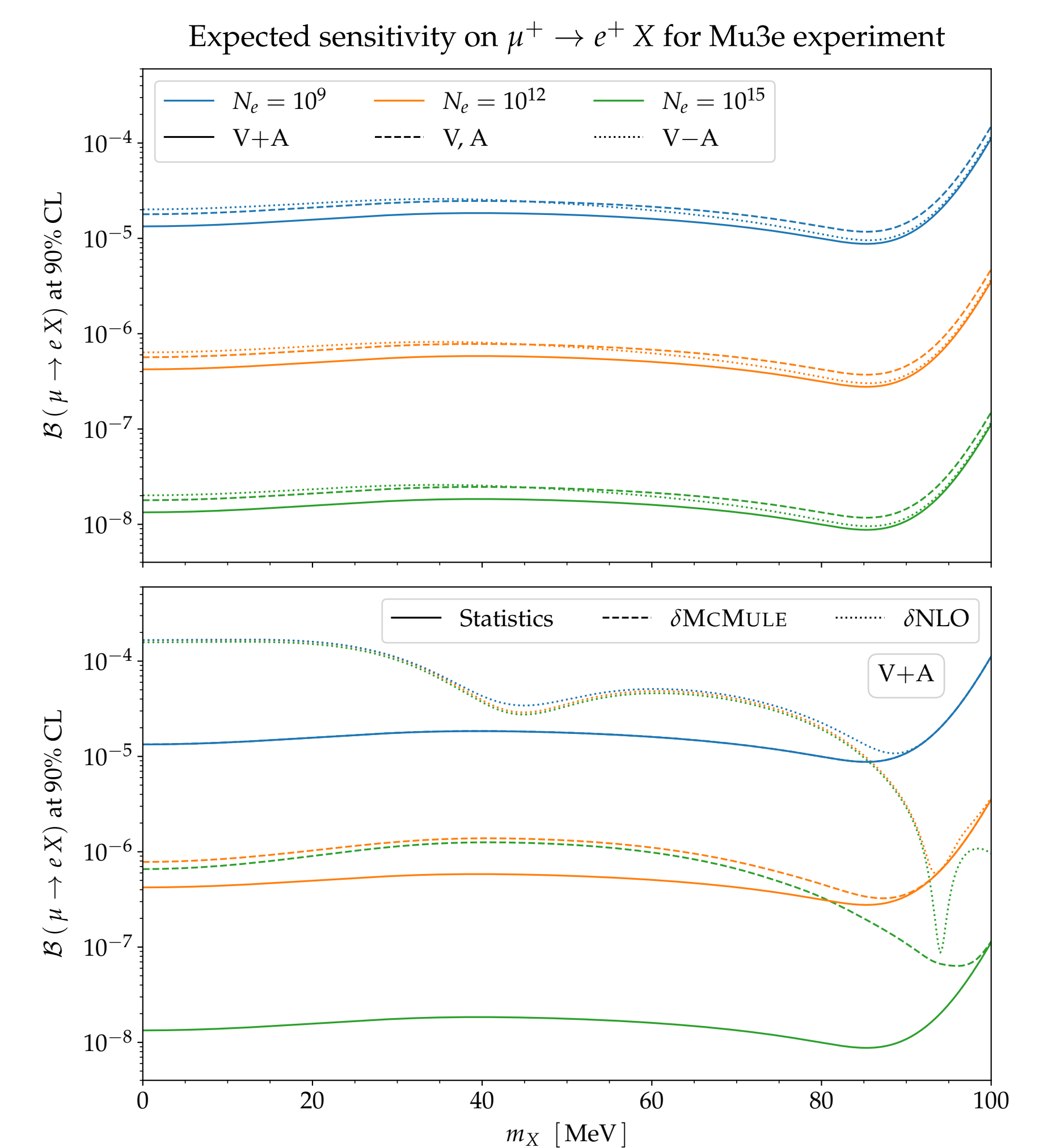
- The input functions have been parametrised according to the **nominal geometry** and the **expected performance** of the involved detectors.

## Expected sensitivity

- The sensitivity on  $\mathcal{B}(\mu^+ \rightarrow e^+ X)$  at 90% CL for MEG II and Mu3e has been estimated by using a **cut-and-count** procedure, accounting for statistical, theoretical and systematic errors.
- Different positron events  $N_e$  and ALP masses  $m_X$  are assumed.
- In the V+A case, the signal and background positrons tend to be emitted in opposite directions, giving a better sensitivity.
- The impact of a systematic error  $\delta E_e$  in the positron energy reconstruction is reported for MEG II. Since an energy offset in the positron spectrum has the same shape of a signal at the endpoint, the effect is enhanced for small ALP masses.



- The impact of the theoretical uncertainty is reported for Mu3e. The new MCMULE predictions make it possible to reach a sensitivity of  $\mathcal{B} \sim 10^{-6}$ , while a simple NLO computation would have limited it to  $\mathcal{B} \sim 10^{-4}$ .



## Conclusions

- The experimental search for small ALP masses is limited by the systematic error on the positron energy. The development of dedicated **calibration tools** is essential to avoid signal biases.
- The reduced theoretical error turned out to be indispensable to achieve competitive sensitivities and even more accurate predictions may be needed for Mu3e.
- The sensitivity for the V+A case can be further improved by designing a dedicated **forward detector**, placed in the opposite direction to the muon beam polarisation, where the SM background is minimal. Preliminary studies suggest a sensitivity of  $\mathcal{B} \sim 10^{-8}$ , including the theory error effect.
- The new MCMULE predictions are ready to be implemented in experimental software for more detailed studies.
- The search for flavour-violating ALPs in muon decays is an excellent opportunity for MEG II and Mu3e to extend their physics programme beyond  $\mu^+ \rightarrow e^+ \gamma$  and  $\mu^+ \rightarrow e^+ e^- e^+$ .