

Deterministic Dynamics of initially random systems

Vincent Pasquier

IPhT, Saclay

Entanglement Scaling and Criticality with Tensor Networks

- Ziga Krajnick, Enej Ilievski, Tomaz Prosen,
Ljubljana
- Johannes Schmidt,
T.U. Berlin
- Help of Jean-Marc Luck
IPhT

- Motivation
- I – Motivation, XXZ, Landau Lifshitz
- II – Large deviations in simple models
- III – Landau Lifshitz
- Conclusion

I - Motivation

- Transport out of equilibrium
- Spin transport in XXZ spin chain
- intégrable \Rightarrow analytic results? difference with nonintegrable?

interesting quantity:

Large deviation of current.

[Marco Znidaric **PRL**, 2011]

Transport of spin current in XXZ chain

Infinite temperature transport:

- Ballistic for $\Delta < 1$
- Anomalous for $\Delta = 1$
- Diffusive for $\Delta > 1$

[Marco Znidaric **PRL**, 2011]

Transport of spin current in XXZ chain

Infinite temperature transport:

- Ballistic for $\Delta < 1$
- Anomalous for $\Delta = 1$
- Diffusive for $\Delta > 1$

Infinite temperature transport of magnetization in Landau Lifshitz: the same.

[Ziga Krajnick, Tomaz Prosen **PRL**, 2019]

I - The coin problem, central limit theorem

A player tosses a coin t times, gains $1/2$ for each tail, $-1/2$ for each head. S_t his gain.

I - The coin problem, central limit theorem

A player tosses a coin t times, gains $1/2$ for each tail, $-1/2$ for each head. S_t his gain.

- The probability of Gain:

$$P(S_t = X) = \frac{1}{2^t} \binom{t}{\frac{t}{2} + X}$$

I - The coin problem, central limit theorem

A player tosses a coin t times, gains $1/2$ for each tail, $-1/2$ for each head. S_t his gain.

- The probability of Gain:

$$P(S_t = X) = \frac{1}{2^t} \binom{t}{\frac{t}{2} + X}$$

- His expected gain will be zero.

I - The coin problem, central limit theorem

A player tosses a coin t times, gains $1/2$ for each tail, $-1/2$ for each head. S_t his gain.

- The probability of Gain:

$$P(S_t = X) = \frac{1}{2^t} \binom{t}{\frac{t}{2} + X}$$

- His expected gain will be zero.
- The fluctuations will be of order \sqrt{T}

I - The coin problem, central limit theorem

A player tosses a coin t times, gains $1/2$ for each tail, $-1/2$ for each head. S_t his gain.

- The probability of Gain:

$$P(S_t = X) = \frac{1}{2^t} \binom{t}{\frac{t}{2} + X}$$

- His expected gain will be zero.
- The fluctuations will be of order \sqrt{T}
- For $X = x\sqrt{t}$ with $x = O(1)$, we have in the large t limit
The probability of Gain S_T at typical scale is the normal law:

$$P(S_t/\sqrt{t} \in [x, x + dx]) = \sqrt{\frac{2}{\pi}} \exp(-2x^2) dx$$

I - The coin problem, central limit theorem

A player tosses a coin t times, gains $1/2$ for each tail, $-1/2$ for each head. S_t his gain.

- The probability of Gain:

$$P(S_t = X) = \frac{1}{2^t} \binom{t}{\frac{t}{2} + X}$$

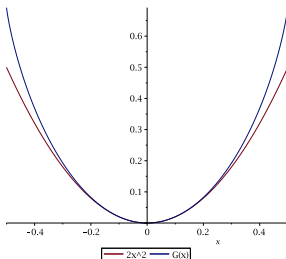
- His expected gain will be zero.
- The fluctuations will be of order \sqrt{T}
- For $X = x\sqrt{t}$ with $x = O(1)$, we have in the large t limit
The probability of Gain S_T at typical scale is the normal law:

$$P(S_t/\sqrt{t} \in [x, x + dx]) = \sqrt{\frac{2}{\pi}} \exp(-2x^2) dx$$

- **Central limit theorem: Moivre**, the probability is Gaussian.

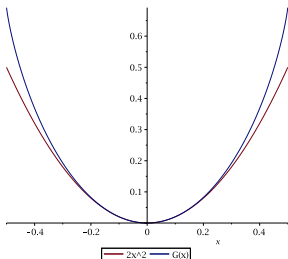
I - The coin problem, large deviations

- Ask for the probability of a gain of order t :
 - $P(S_t/t \in [x, x + dx]) = \exp(-tI(x))dx$



I - The coin problem, large deviations

- Ask for the probability of a gain of order t :
 - $P(S_t/t \in [x, x + dx]) = \exp(-tI(x))dx$



Rate function: $I(x)$. Its explicit expression follows from Stirling formula.

I - The coin problem, Cumulants

- **Large deviation principle:** The probability decays as e^{-TG}
- **Cumulant Generating Function (CGF)** behaves similarly:

$$\langle e^{\lambda S_t} \rangle = \int dx e^{-t(I(x) - \lambda x)} = e^{-tF(\lambda)}$$

- F =Legendre transform of rate function I :

$$F(\lambda) = \min_x (I(x) - \lambda x)$$

- Gartner Ellis Thm: It can apply with **non independant variables** under some conditions on MGF: $F(\lambda)$ differentiable? **complex analytic?**
- Consequence, all cumulants scale as t .

Comparison with thermodynamics

Similar to Thermodynamics:

Comparison with thermodynamics

Similar to Thermodynamics:

- Cumulant generating function:

$$F(\lambda) \leftrightarrow \text{free energy } F(\beta)$$

Comparison with thermodynamics

Similar to Thermodynamics:

- Cumulant generating function:

$$F(\lambda) \leftrightarrow \text{free energy } F(\beta)$$

- rate function:

$$I(x) \leftrightarrow \text{entropy } S(U)$$

Comparison with thermodynamics

Similar to Thermodynamics:

- Cumulant generating function:

$$F(\lambda) \leftrightarrow \text{free energy } F(\beta)$$

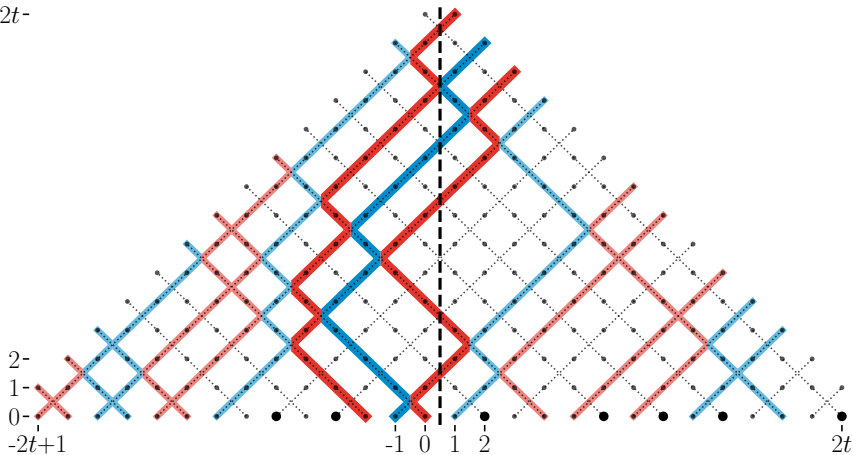
- rate function:

$$I(x) \leftrightarrow \text{entropy } S(U)$$

- Some discrepancy with the standard large deviation theory should be expected in the presence of dynamical phase transitions as free energy is not always extensive in the vicinity of thermodynamical phase transitions.

A simple model

$2t-$



Rate function

- **uncharged model** Two players R and L toss a coin t times independently, G_L and G_R their respective gains. The number of particles which have passed through the origin at large time coincides with the difference of their respective gains:

$$J_P = G_L - G_R$$

Rate function

- **uncharged model** Two players R and L toss a coin t times independently, G_L and G_R their respective gains. The number of particles which have passed through the origin at large time coincides with the difference of their respective gains:

$$J_P = G_L - G_R$$

- With probability

$$P(J_P) = \frac{1}{2^{2N}} \sum_{G_L - G_R = J_P} \binom{t}{G_R} \binom{t}{G_L}$$

Rate function

- **uncharged model** Two players R and L toss a coin t times independently, G_L and G_R their respective gains. The number of particles which have passed through the origin at large time coincides with the difference of their respective gains:

$$J_P = G_L - G_R$$

- With probability

$$P(J_P) = \frac{1}{2^{2N}} \sum_{G_L - G_R = J_P} \binom{t}{G_R} \binom{t}{G_L}$$

- **charged model** Color independently blue (plus) or red (minus) each particle trajectory and ask the amount of charge J_c that has passed through the origin.

single file property

- The particles obey **single file property**: They keep their ordering. Expect subdiffusive behaviour.

single file property

- The particles obey **single file property**: They keep their ordering. Expect subdiffusive behaviour.
- probability of charge transmission:

$$P(J_c) = \sum_P P(J_P) P(J_c | J_P)$$

single file property

- The particles obey **single file property**: They keep their ordering. Expect subdiffusive behaviour.
- probability of charge transmission:

$$P(J_c) = \sum_P P(J_P) P(J_c | J_P)$$

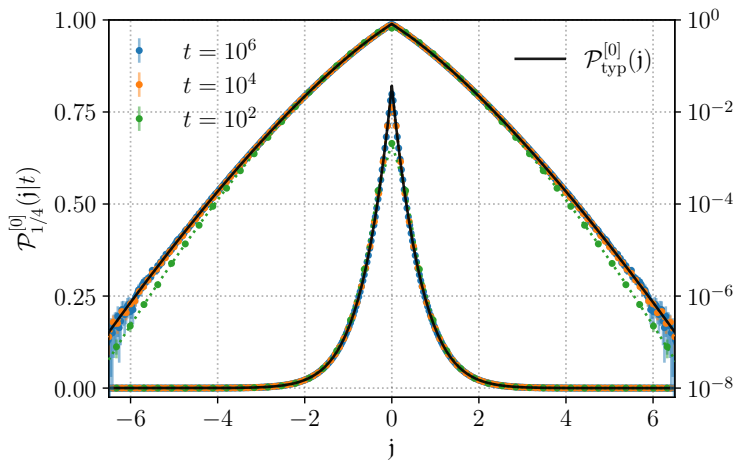
- explicitly:

$$P(J_c) = \frac{1}{2^{2N}} \sum_{G_L - G_R = J_P} \binom{t}{G_R} \binom{t}{G_L} \binom{|J_P|}{|J_P|/2 + J_c}$$

- Gaussian approximation $J_P = t^{1/2}x$, $J_c = t^{1/4}j$

$$P(j \in dj) = dj \frac{2}{\pi} \int_0^\infty \frac{dx}{\sqrt{x}} e^{-2(x^2 + \frac{j^2}{x})}$$

Non Gaussian profile



Cumulant generating function CGF

- Scaling of lambda follows from that of $J_c = O(t^{1/4})$

$$\lambda t^{1/4} = w$$

Cumulant generating function CGF

- Scaling of lambda follows from that of $J_c = O(t^{1/4})$

$$\lambda t^{1/4} = w$$

- CGF at scale $t^{1/4}$

$$F(w) = \frac{w^4}{32} + \ln(1 + \operatorname{erf}(\frac{w^2}{4}))$$

Cumulant generating function CGF

- Scaling of lambda follows from that of $J_c = O(t^{1/4})$

$$\lambda t^{1/4} = w$$

- CGF at scale $t^{1/4}$

$$F(w) = \frac{w^4}{32} + \ln(1 + \operatorname{erf}(\frac{w^2}{4}))$$

- Cumulants scale as:

$$c_{2n} \sim t^{\frac{n}{2}}$$

Cumulant generating function CGF

- Scaling of lambda follows from that of $J_c = O(t^{1/4})$

$$\lambda t^{1/4} = w$$

- CGF at scale $t^{1/4}$

$$F(w) = \frac{w^4}{32} + \ln(1 + \operatorname{erf}(\frac{w^2}{4}))$$

- Cumulants scale as:

$$c_{2n} \sim t^{\frac{n}{2}}$$

- Do not all scale with the same power of t unlike large deviation predictions.

Cumulant generating function CGF

- Scaling of lambda follows from that of $J_c = O(t^{1/4})$

$$\lambda t^{1/4} = w$$

- CGF at scale $t^{1/4}$

$$F(w) = \frac{w^4}{32} + \ln(1 + \operatorname{erf}(\frac{w^2}{4}))$$

- Cumulants scale as:

$$c_{2n} \sim t^{\frac{n}{2}}$$

- Do not all scale with the same power of t unlike large deviation predictions.
- expect non analyticity of CGF in the vicinity of $\lambda = 0$.

Phase transition

- CGF= partition function with t the size of the system:

$$Z(\lambda) = \frac{1}{2^{2N}} \sum_{G_L - G_R = J_P} \binom{t}{G_R} \binom{t}{G_L} \cosh(\lambda)^{|J_P|} e^{hJ_P}$$

Phase transition

- CGF= partition function with t the size of the system:

$$Z(\lambda) = \frac{1}{2^{2N}} \sum_{G_L - G_R = J_P} \binom{t}{G_R} \binom{t}{G_L} \cosh(\lambda)^{|J_P|} e^{hJ_P}$$

- Curie Weiss type model
 J_p magnetization, $\ln(\cosh \lambda)$ temperature.

Phase transition

- CGF= partition function with t the size of the system:

$$Z(\lambda) = \frac{1}{2^{2N}} \sum_{G_L - G_R = J_P} \binom{t}{G_R} \binom{t}{G_L} \cosh(\lambda)^{|J_P|} e^{hJ_P}$$

- Curie Weiss type model
 J_p magnetization, $\ln(\cosh \lambda)$ temperature.
- Free energy, define order parameter $m = J_p/t$,

$$f = -\frac{\ln(Z)}{t} \sim 2m^2 - \beta|m| + hm$$

Phase transition

- CGF= partition function with t the size of the system:

$$Z(\lambda) = \frac{1}{2^{2N}} \sum_{G_L - G_R = J_P} \binom{t}{G_R} \binom{t}{G_L} \cosh(\lambda)^{|J_P|} e^{hJ_P}$$

- Curie Weiss type model

J_p magnetization, $\ln(\cosh \lambda)$ temperature.

- Free energy, define order parameter $m = J_p/t$,

$$f = -\frac{\ln(Z)}{t} \sim 2m^2 - \beta|m| + hm$$

- h symmetry breaking field, $h = 0$ first order phase transition.

Phase transition

- CGF= partition function with t the size of the system:

$$Z(\lambda) = \frac{1}{2^{2N}} \sum_{G_L - G_R = J_P} \binom{t}{G_R} \binom{t}{G_L} \cosh(\lambda)^{|J_P|} e^{hJ_P}$$

- Curie Weiss type model

J_p magnetization, $\ln(\cosh \lambda)$ temperature.

- Free energy, define order parameter $m = J_p/t$,

$$f = -\frac{\ln(Z)}{t} \sim 2m^2 - \beta|m| + hm$$

- h symmetry breaking field, $h = 0$ first order phase transition.
- λ real double well,
 λ imaginary single well,
 $\lambda = 0$ Dynamical critical point signals breaking of large deviation principle.

zeros of partition function

- CGF using saddle point approximation:

$$Z(\lambda) = ((1 + \cosh \lambda)(1 + 1/\cosh \lambda))^t + O(1)$$

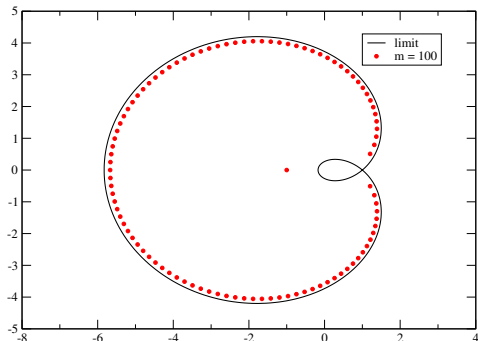
zeros of partition function

- CGF using saddle point approximation:

$$Z(\lambda) = ((1 + \cosh \lambda)(1 + 1/\cosh \lambda))^t + O(1)$$

- Natural boundary $|Z^{1/t}| = 1$

Zeros in complex c plane



Bernoulli curve !

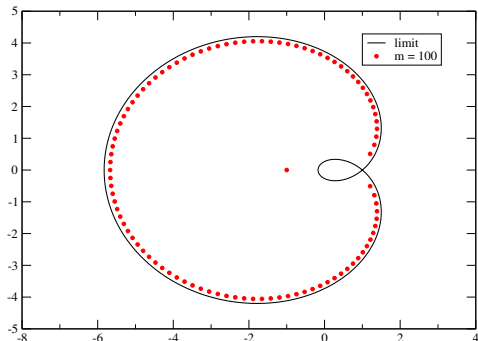
zeros of partition function

- CGF using saddle point approximation:

$$Z(\lambda) = ((1 + \cosh \lambda)(1 + 1/\cosh \lambda))^t + O(1)$$

- Natural boundary $|Z^{1/t}| = 1$

Zeros in complex c plane



Bernoulli curve !

- Classical limit of XXZ spin chain.

- Classical limit of XXZ spin chain.
- Landau Lifshitz evolution:

$$S_t = S \times S_{xx} + S \times JS$$

- Classical limit of XXZ spin chain.
- Landau Lifshitz evolution:

$$\dot{S}_t = S \times S_{xx} + S \times JS$$

- Why discretize? **Infinite temperature**
On each lattice site j , $S_j(t=0)$ is a random vector on the unit sphere.

- Classical limit of XXZ spin chain.
- Landau Lifshitz evolution:

$$\dot{S}_t = S \times S_{xx} + S \times JS$$

- Why discretize? **Infinite temperature**
On each lattice site j , $S_j(t=0)$ is a random vector on the unit sphere.
- exchange relation imply **discrete integrability**:

$$L(u, S_1)L(v, S_2) = L(v, S'_2)L(u, S'_1)$$

Discrete time evolution

- $J = 0$ dynamics:

$$S_{t+1} = \frac{1}{\sigma^2 + \tau^2} (\sigma^2 S + \tau^2 V + \tau S \times V)$$

$$V' = \frac{1}{\sigma^2 + \tau^2} (\sigma^2 V + \tau^2 S + \tau V \times S)$$

- The vertex

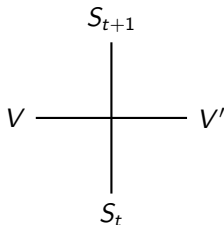


Figure: Time evolution of the spin

Discrete time evolution

- $J = 0$ dynamics:

$$S_{t+1} = \frac{1}{\sigma^2 + \tau^2} (\sigma^2 S + \tau^2 V + \tau S \times V)$$

$$V' = \frac{1}{\sigma^2 + \tau^2} (\sigma^2 V + \tau^2 S + \tau V \times S)$$

- The vertex

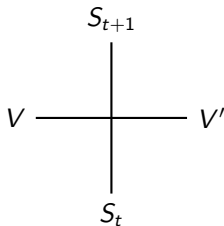
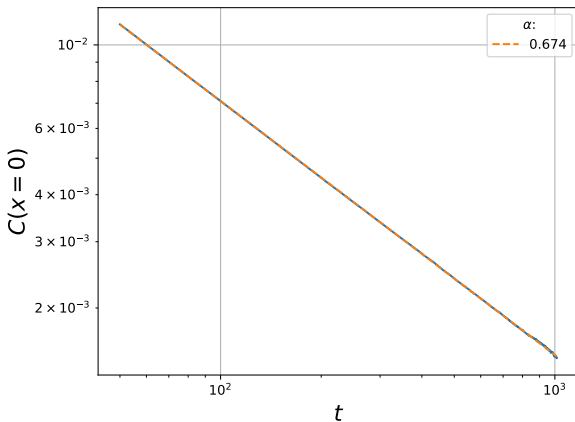


Figure: Time evolution of the spin

- Yang Baxter map

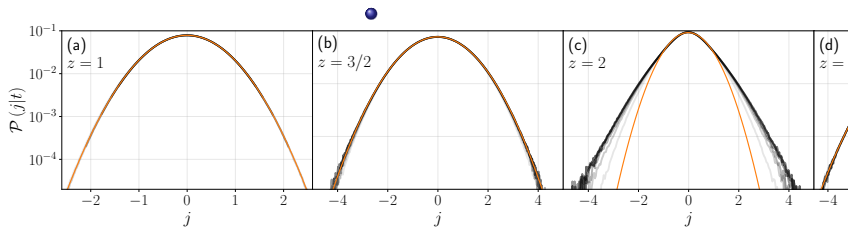
Dynamical exponent $3/2$

- Dynamical exponent



Large deviation of magnetization current

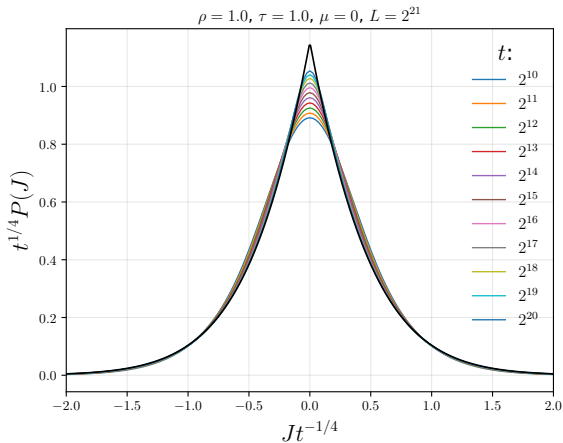
- Rate Function



- easy plane, isotropic, easy axis

Non Gaussianity

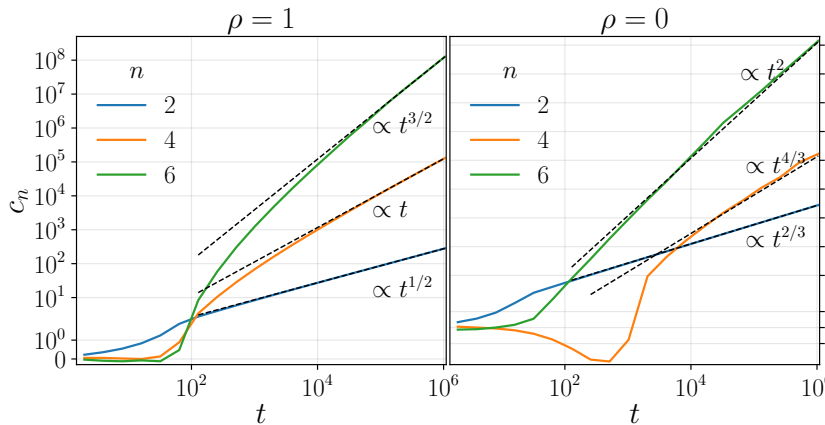
- Easy axis rate Function



Not Gaussian, again cusp

Cumulants

- cumulants



Conclusion

- High temperature transport in XXZ chain not so well understood.

Conclusion

- High temperature transport in XXZ chain not so well understood.
- The classical limit **Landau Lifshitz** has very analogous properties then the quantum model.

Conclusion

- High temperature transport in XXZ chain not so well understood.
- The classical limit **Landau Lifshitz** has very analogous properties then the quantum model.
- Our work suggests dynamical phase transition.

Conclusion

- High temperature transport in XXZ chain not so well understood.
- The classical limit **Landau Lifshitz** has very analogous properties then the quantum model.
- Our work suggests dynamical phase transition.
- Is the $3/2$ dynamical exponent related to **KPZ** universality class?