





### Lattice studies for MUonE analysis

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Radiative corrections and Monte Carlo tools for low-energy hadronic cross sections in  $e^+e^-$  collisions Jun 7 – 9, 2023





Contribution	value $\times 10^{11}$	
QED	116584718.931(104)	
Electroweak	153.6(1.0)	
HVP	6845(40)	
HLbL	92(18)	
Total SM value	116591810(43)	
Difference: $a_{\mu}^{\exp} - a_{\mu}^{SM}$	251(59)	
[arXiv:{2104.03281, 2006.04822, 2203.15810}]		

- HVP is the main contributor to the total theory uncertainty
- Lots of activities in lattice QCD to improve the precision
- • •
- Lattice QCD, strong isospin breaking & QED:
  - direct effects in HVP of order 1%, yet need to be measured precisely...
  - indirect effects through scale setting quantities
- MUonE experiment as an alternative method





#### 2 Analysis of a mock data for HVP using Padé based fit functions

### Dispersive method: a time-like approach to evaluate $a_{\mu}^{\rm HVP}$

• Relation between the  $\mathcal{R}e \Pi(Q^2)$  and  $\mathcal{I}m \Pi(Q^2)$ :

$$\Pi(Q^2) - \Pi(0) = \frac{Q^2}{\pi} \int_0^\infty ds \frac{\mathcal{I} m \, \Pi(s)}{s(s-Q^2)}$$

• Imaginary part of  $\Pi(s)$  is related to the experimental total cross-section in e+e- annihilation:

$$\mathcal{I}m \Pi(s) = \frac{\alpha}{3}R(s)$$

- Important contributions :  $ho, \omega, \phi, J/\psi$
- O(1000) channels
- Model calculations had to be used for some channels
- [Keshavarzi, Nomura, Teubner, Phys.Rev. D97 (2018) no.11]



### Time-like vs space-like evaluation of $a_{\mu}^{\text{HVP}}$



- $\Pi_{had}$  is the hadronic part of the photon vacuum polarization
- $Im\Pi_{had}(s)$  is related to the experimental total cross-section in  $e^+e^-$  annihilation

- (1) Formulate QCD on a (finite-size) lattice in Euclidean time
- (2) Generate ensembles of field configuration with MC simulations
- (3) Compute correlation function of fields as a function of time/momentum
- (4) Average over configurations
- (5) Extrapolate to continuum, infinite volume, and physical quark masses



# Space-like evaluation of $a_{\mu}^{\rm HVP}:$ lattice QCD

- $\bullet\,$  In WP20, lattice results (< Mar/2020) were averaged; uncertainty 2.6%
- BMW20 reported first lattice result with sub-percent uncertainty:
  - reduced tension with experiment:  $\sim 1.5\sigma$  ,
  - some tension with the R-ratio method (WP20);  $\sim 2.1\sigma$

	value $\times 10^{10}$	error $\%$
$a_{\mu}^{\text{HVP, LO}}$ (R-ratio, WP20)	693.1(4.0)	0.6%
$a_{\mu}^{\text{HVP, LO}}$ (lattice, WP20)	711.6(18.4)	2.6%
$a_{\mu}^{\text{HVP, LO}}$ (lattice, BMW20)	707.5(5.5)	0.8%

[WP20: arXiv:2006.04822, BMW20: arXiv:2002.12347]



• Intermediate Euclidean-time window values of  $a_{\mu}$  are introduced for further investigations

## Space-like evaluation of $a_{\mu}^{\text{HVP}}$ : lattice QCD

Dominant sources of error:

1) Determination of signal at small  $Q^2$ 

• The integrand peaks at  $Q^2$  about  $m_{\mu}^2/4$ 

$$\left(\frac{\alpha}{\pi}\right)^2 \int_0^\infty dQ^2 f(Q^2) \hat{\Pi}_{\rm had}(Q^2)$$

- These very low momenta cannot be directly accessed on current lattices ( $L\approx 10~{\rm fm}$  required)
- Analytic functions (like Padé) in combination of the method of time moments have been suggested & used to describe  $\hat{\Pi}(Q^2)$  over small values of  $Q^2$ ; increase of statistical error at higher moments
- In alternative, time-momentum representation the problem with small  $Q^2$  shows itself as exponential growth of the relative statistical error at large time
- 2) Continuum extrapolation, scale setting errors, finite volume effects, disconnected diagrams, ···

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- In alternative, time-momentum representation the problem with small  $Q^2$  shows itself as exponential growth of the relative statistical error at large time
- A hybrid method (with MUonE) can be used to circumvent this problem
- 2) Continuum extrapolation, scale setting errors, finite volume effects, disconnected diagrams, ···

## Space-like evaluation of $a_{\mu}^{\text{HVP}}$ : MUonE

• HVP contributes to the running of QED fine structure coupling

$$\alpha(Q^2) = \frac{\alpha_0}{1 - \hat{\Pi}(Q^2)}$$

- Comparing experimental data & perturbative calculations yields HVP through its contribution to  $\alpha(Q^2)$
- MUonE extracts  $\Delta \alpha_{had}(Q^2)$  from the shape of the differential  $\mu e$ scattering cross section



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# Space-like evaluation of $a_{\mu}^{\rm HVP}:$ hybrid method

Divide & Conquer:

$$\begin{split} a^{\rm HVP}_{\mu} &= I_0 + I_1 + I_2 \\ I_0 &= \left(\frac{\alpha}{\pi}\right)^2 \int_0^{0.14} dQ^2 f(Q^2) \hat{\Pi}_{\rm had}(Q^2) \\ I_1 &= \left(\frac{\alpha}{\pi}\right)^2 \int_{0.14}^{Q^2_{\rm max}} dQ^2 f(Q^2) \hat{\Pi}_{\rm had}(Q^2) \\ I_2 &= \left(\frac{\alpha}{\pi}\right)^2 \int_{Q^2_{\rm max}}^{\infty} dQ^2 f(Q^2) \hat{\Pi}_{\rm had}(Q^2) \end{split}$$

- $I_0:$  contains  $\sim$  84% of the  $a_\mu^{\rm had,\ LO}$  & can be calculated precisely with the MUonE experiment
- $I_1$ : use lattice QCD or R-ratio
- $I_2$ : use perturbation theory

#### MUonE: template fit to $\mu$ -e scattering data

• A proposed template fit inspired by contribution of a lepton-pair to the space-like photon vacuum polarization looks very good on test data

$$\Delta \alpha_{\rm had}(t) = k \left\{ -\frac{5}{9} - \frac{4M}{3t} + \left(\frac{4M^2}{3t^2} + \frac{M}{3t} - \frac{1}{6}\right) \frac{2}{\sqrt{1 - \frac{4M}{t}}} \log \left| \frac{1 - \sqrt{1 - \frac{4M}{t}}}{1 + \sqrt{1 - \frac{4M}{t}}} \right| \right\}$$

$$\begin{bmatrix} \frac{1}{9} & \cos \theta \\ \cos \theta$$

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$$\left| \frac{y}{s} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq 0}^{s \neq 0} \left| \frac{y}{s \neq 0} \left| \frac{y}{s \neq 0} \right|_{s \neq$$

- The above template fit can be potentially problematic with highly precise data
- Alternative methods such as using Padé based fits are suggested and explored



#### 2 Analysis of a mock data for HVP using Padé based fit functions

### HVP, Stieltjes functions, Padé approximants

In the context of lattice QCD, the use of Padé approximants for low  $Q^2$  regions  $\hat{\Pi}(Q^2)$  was suggestion by [Aubin, Blum, Golterman, Peris (2012); Golterman, Maltman, Peris (2013)]

- $\bullet\,$  It was introduced to deal with low signal at small  $Q^2$  in lattice QCD
- They investigated the Padé approximants using mock data from hadronic  $\tau$  decay with I = 1 non-strange spectral function:

$$\hat{\Pi}^{I=1}(Q^2) = Q^2 \int_{4m_\pi^2}^{\infty} ds \frac{\rho^{I=1}(s)}{s(s+Q^2)}$$

- They chose 40 equally-spaced points in  $Q^2$  (from 0.01 to 0.4  ${\rm GeV^2}$ ) from  $\hat{\Pi}^{I=1}(Q^2)$ , exploited an exaggerated lattice-based covariance matrix, and created a mock data set
- After performing their analysis, they compared their results with the "exact" value obtained directly from the original data set:

$$\tilde{a}_{\mu, \rm HLO} \Big|_{Q^2 \le 1 {\rm GeV}^2} = 1.204(27) \times 10^{-7}$$

#### Sample covariance matrix from lattice QCD

- We continue their study using a correlation matrix calculated on a lattice ensemble of size  $64^3\times128$  at  $a\approx0.066~{\rm fm}$
- We calculate current-current correlator and integrate over spacial dimensions

$$G(x_0) = -\frac{1}{3} \sum_{k=1}^{3} \int dx^3 \langle J_k(x) J_k(0) \rangle$$

After subtracting associated contact terms, we take FT to obtain Q<sup>2</sup>Π(Q<sup>2</sup>)
 We extract a correlation matrix and use it with chosen 40 equally-spaced points of Π<sup>I=1</sup>(Q<sup>2</sup>) (curtesy of Kim Maltman)



#### Padé-based fit functions

• Writing the HVP in terms of a Stieltjes function warrants an existence of the converging sequence of order [N-1, N] and [N, N] Padé approximants (PAs), defined as

$$\Delta \alpha_{\rm had}(Q^2) = c_0 + Q^2 \left( a_0 + \sum_{i=1}^N \frac{a_i}{b_i + Q^2} \right)$$

where  $Q^2 = -t$  and  $a_0 = 0$  in [N - 1, N] PAs



### Some fit results



Mock data: 20 correlated data are included in the fits

Fit func: Padé-based fits and a lepton-like ansatz

Challenge: Small eigenvalues of the covariance matrices make the fit more challenging; manifested in large  $\chi^2/{\rm ndof}$ 

Message: Better fit functions are needed with very precise covariance matrices

### Summary & concluding remarks

- $\bullet\,$  We performed a test study using a mock data from dispersive  $\tau\text{-based}$  data
- $\bullet$  We used Padé based functions and a lepton-like ansatz to investigate  $Q^2$  dependence of the mock data
- Our investigation showed, the lepton-like ansatz works as good as higher order Pade-based fits if we ignore correlations
- With correlated data, however, one has to exploit other functions such as the Padé approximants
- Padé approximants (and similar, analytic functions) can be used for extrapolating both lattice and MUonE data, and comparing them if needed

# Thanks for Your Attention!