

# **muEDM in the landscape of storage ring EDM searches**

**– LTPhD Seminar (26 September 2023) –**

Tim Hume, Muon Physics Group

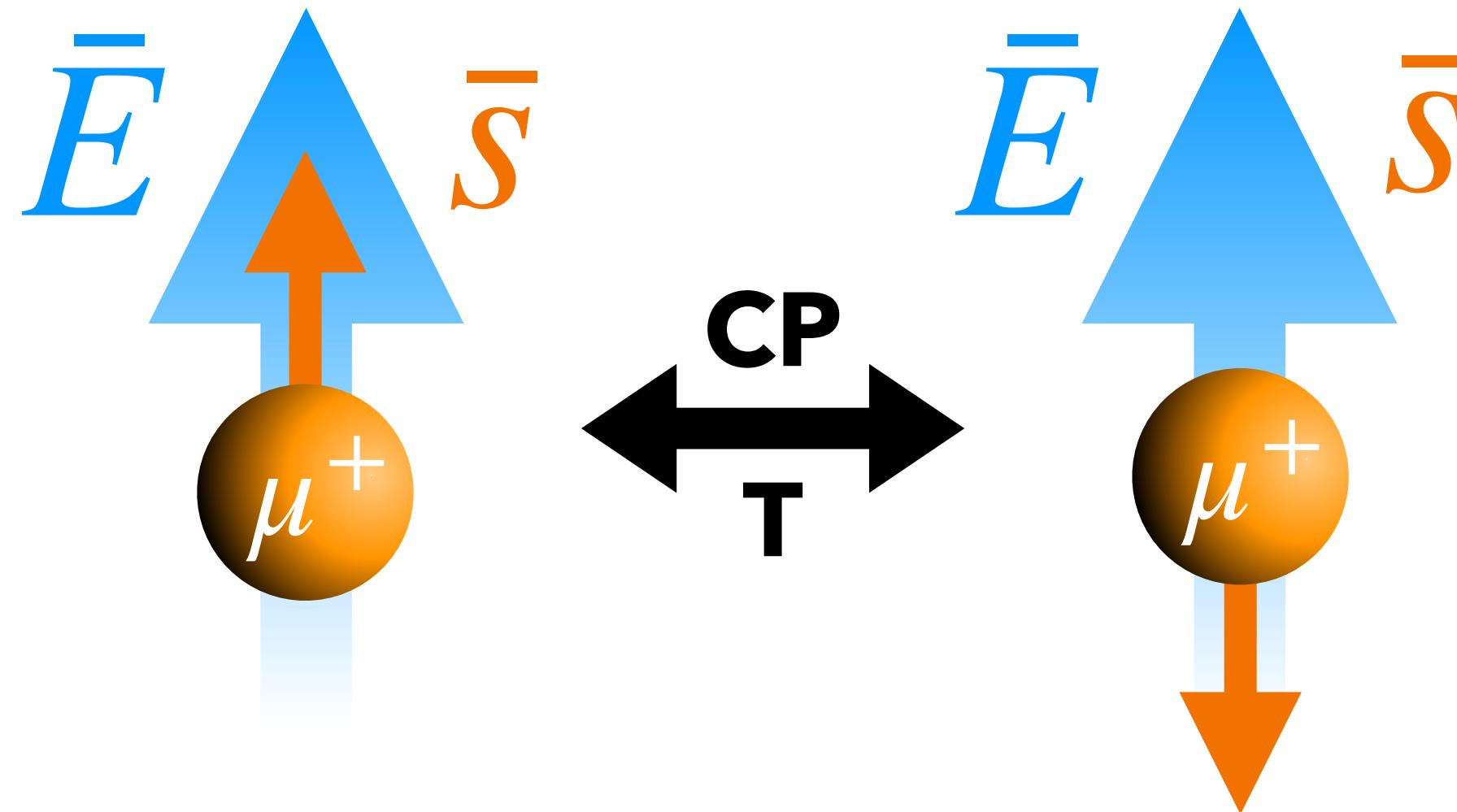
Supervised by Dr. Philipp Schmidt-Wellenburg



[timothy.hume@psi.ch](mailto:timothy.hume@psi.ch)

# Muon Electric Dipole Moment

A permanent EDM requires T violation,  
equivalently CP violation by the CPT Theorem.



$$H_{\mu}^{EDM} \xrightarrow{\beta \rightarrow 0} d_{\mu} \bar{\sigma} \cdot \bar{E}$$

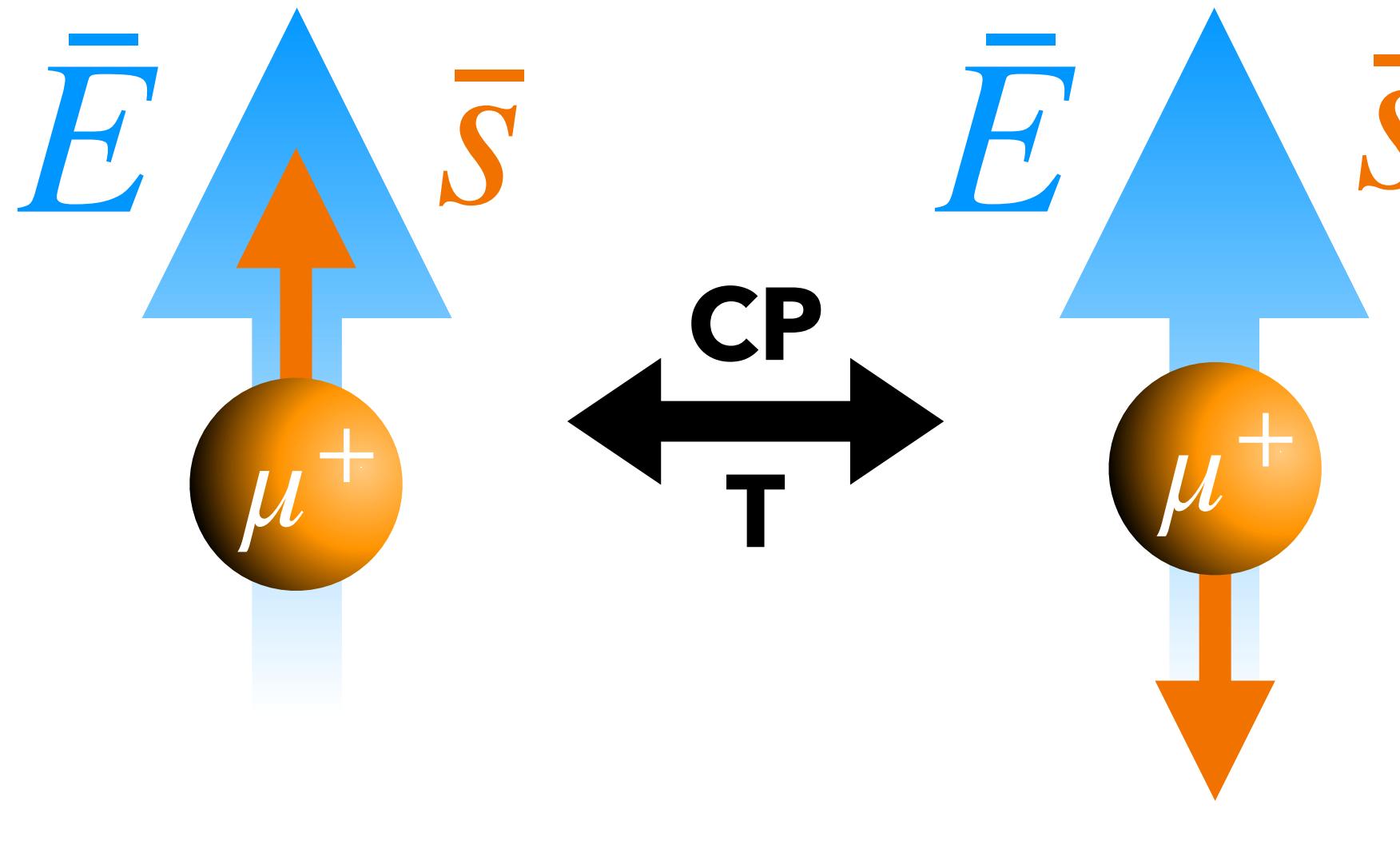
Hamiltonian EDM term is CP violating

SM Prediction:  $d_{\mu}^{\text{SM}} = 1.4 \times 10^{-38} e \cdot \text{cm}$

(Yamaguchi & Yamanaka, 2020)

# Muon Electric Dipole Moment

A permanent EDM requires T violation,  
equivalently CP violation by the CPT Theorem.

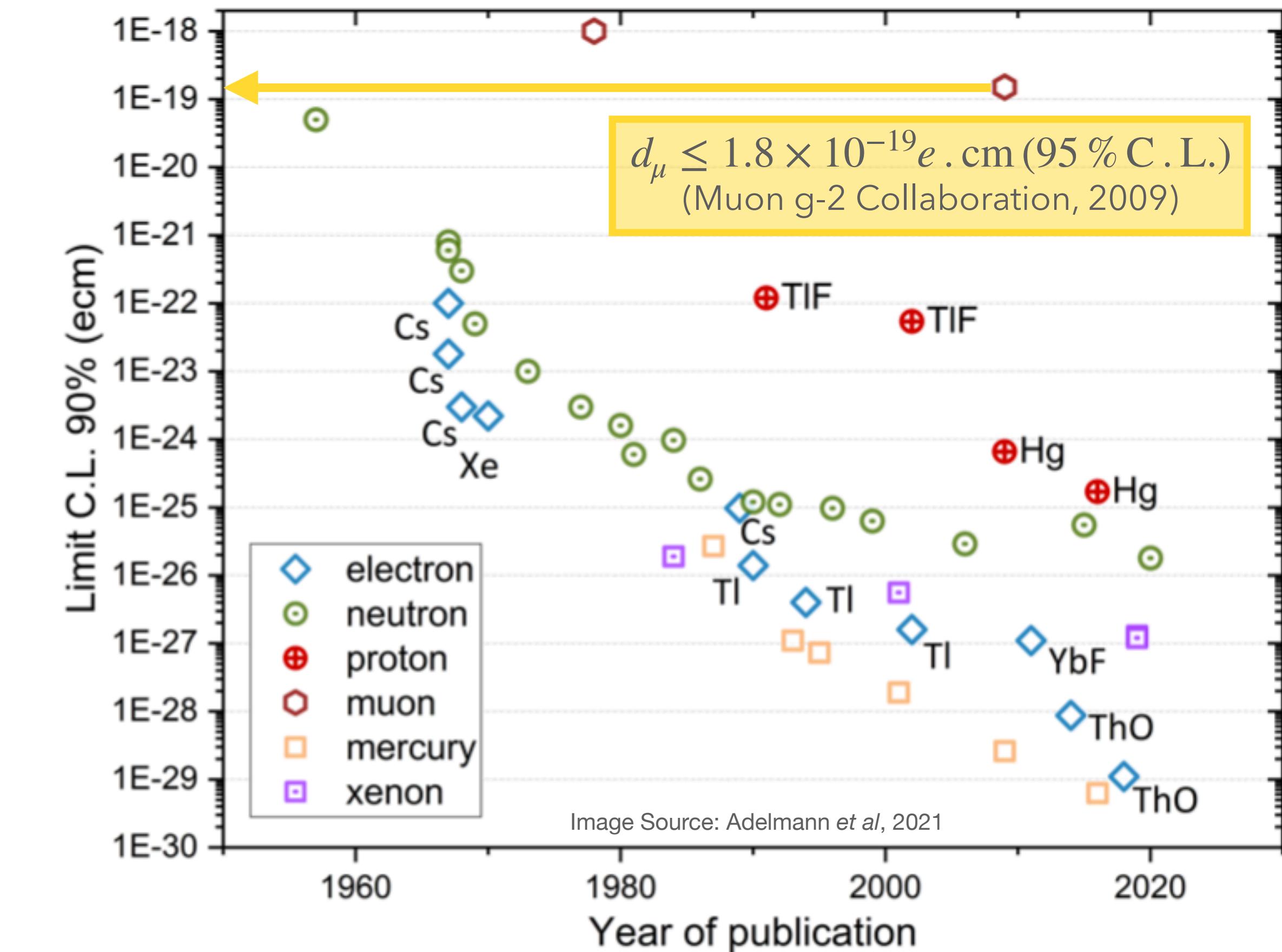


$$H_{\mu}^{EDM} \xrightarrow{\beta \rightarrow 0} d_{\mu} \bar{\sigma} \cdot \bar{E}$$

Hamiltonian EDM term is CP violating

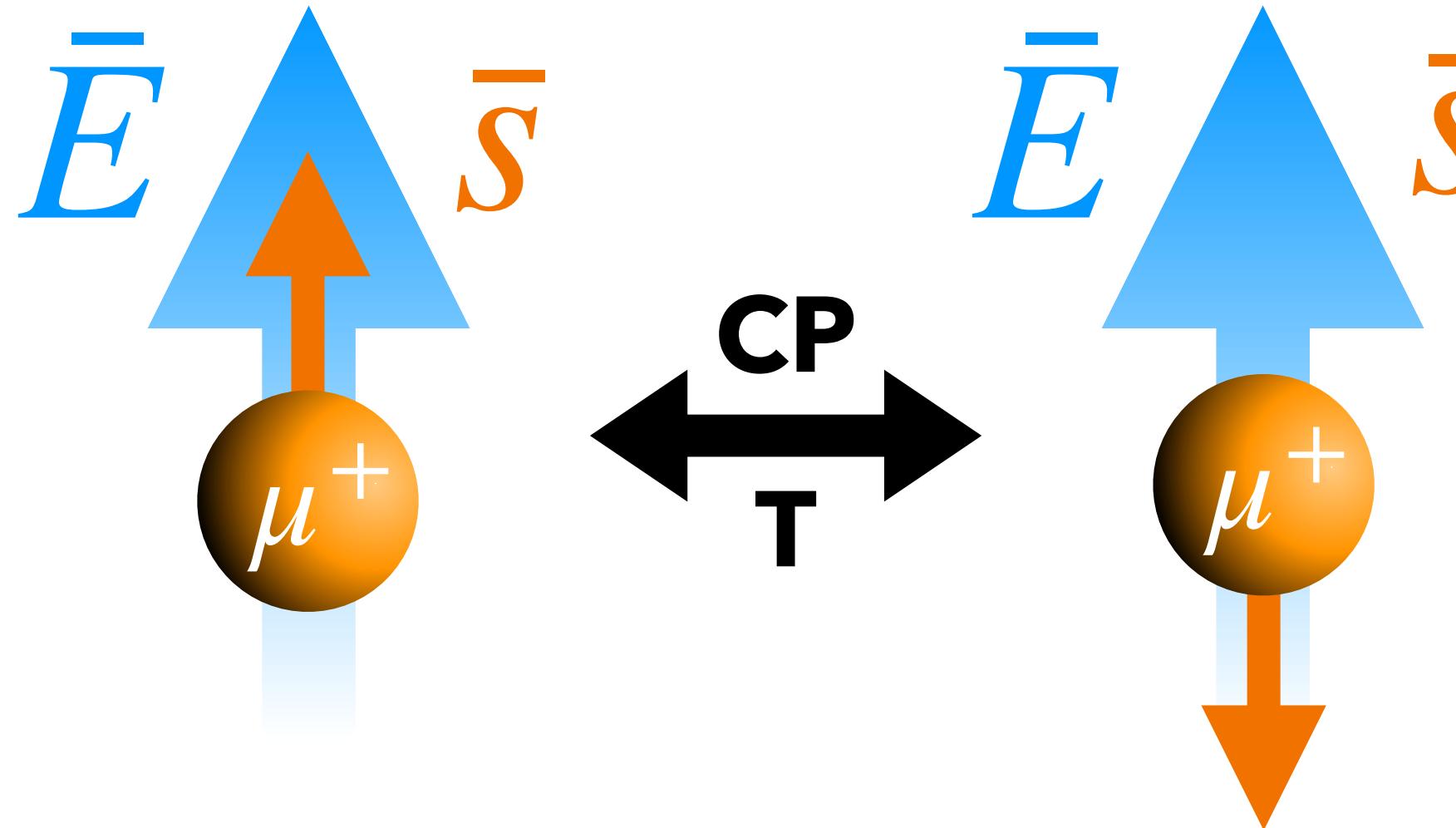
SM Prediction:  $d_{\mu}^{SM} = 1.4 \times 10^{-38} e \cdot \text{cm}$

(Yamaguchi & Yamanaka, 2020)



# Muon Electric Dipole Moment

A permanent EDM requires T violation,  
equivalently CP violation by the CPT Theorem.



$$H_{\mu}^{EDM} \xrightarrow{\beta \rightarrow 0} d_{\mu} \bar{\sigma} \cdot \bar{E}$$

Hamiltonian EDM term is CP violating

SM Prediction:  $d_{\mu}^{SM} = 1.4 \times 10^{-38} e \cdot \text{cm}$

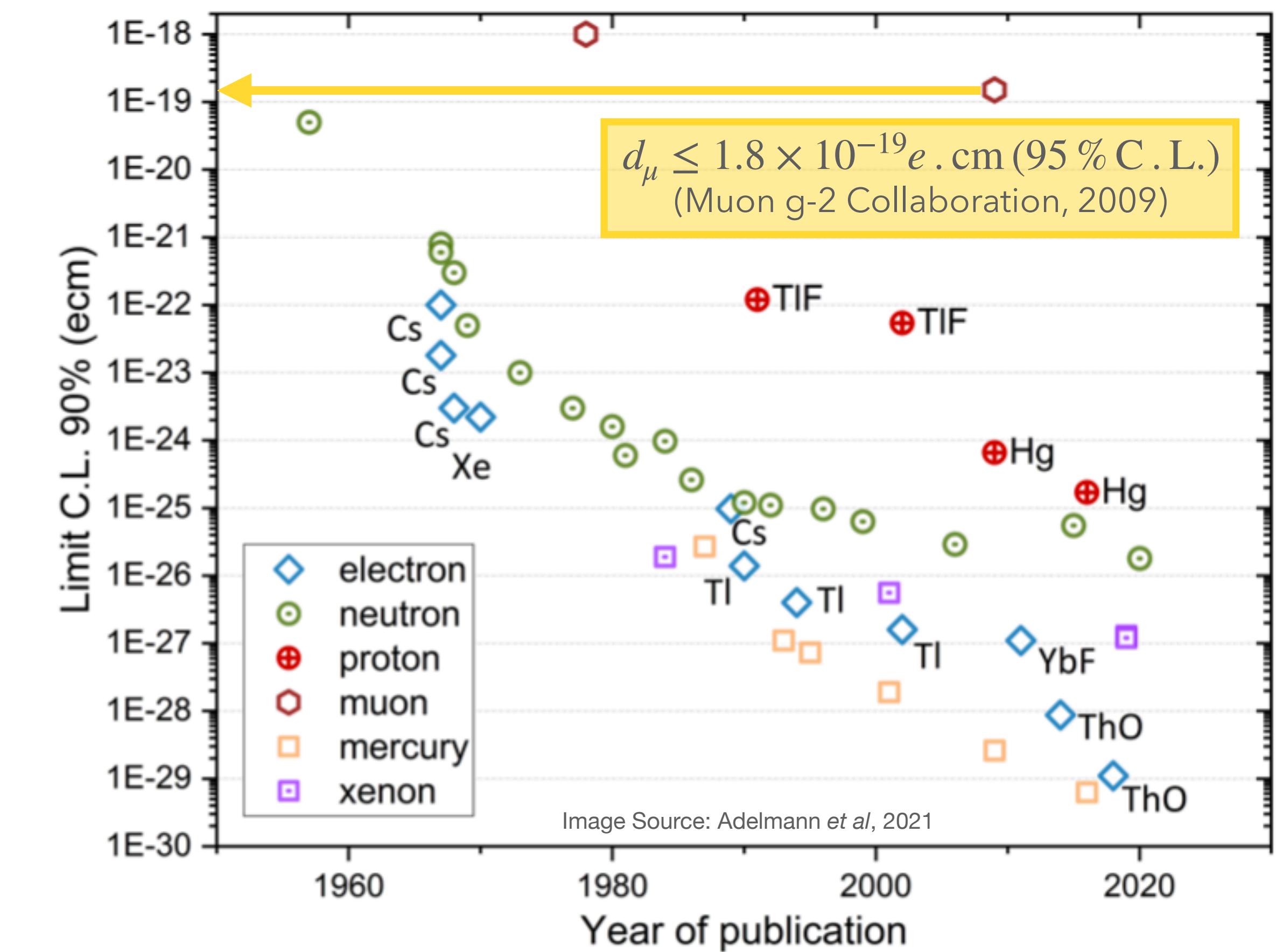
(Yamaguchi & Yamanaka, 2020)

$$d_e \leq 4.1 \times 10^{-30} e \cdot \text{cm}$$

(Roussy et al., 2023)

LFU  
⇒?

$$d_{\mu} \leq \frac{m_{\mu}}{m_e} d_e = 6.0 \times 10^{-28} e \cdot \text{cm}$$



# Standard Model Prediction

PHYSICAL REVIEW LETTERS 125, 241802 (2020)

## Large Long-Distance Contributions to the Electric Dipole Moments of Charged Leptons in the Standard Model

Yasuhiro Yamaguchi\*

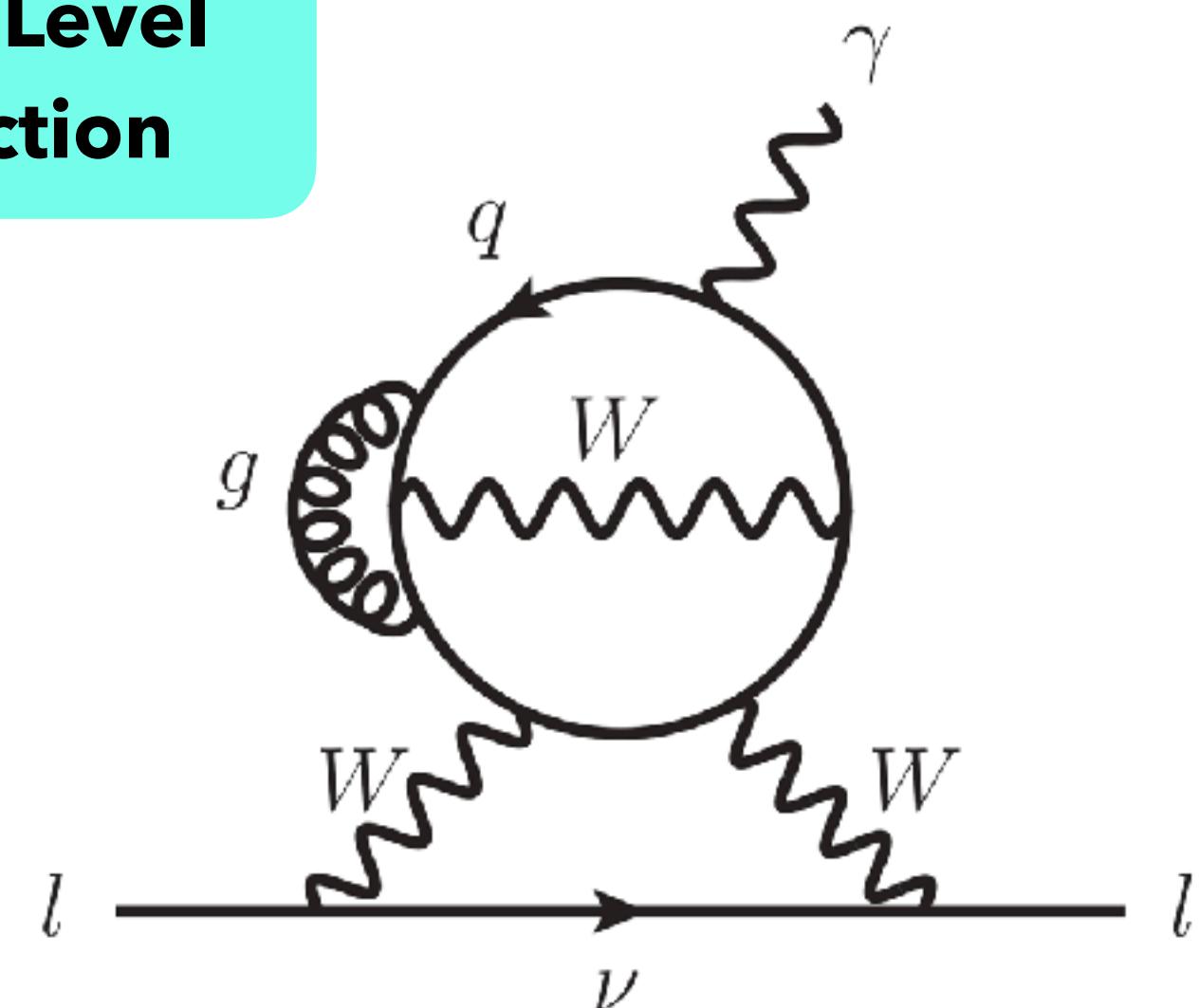
Advanced Science Research Center, Japan Atomic Energy Agency (JAEA), Tokai 319-1195, Japan  
and RIKEN Nishina Center, RIKEN, Wako, Saitama 351-0198, Japan

Nodoka Yamanaka<sup>†</sup>

Amherst Center for Fundamental Interactions, Department of Physics, University of Massachusetts,  
Amherst, Massachusetts 01003, USA  
and Yukawa Institute for Theoretical Physics, Kyoto University, Kitashirakawa-Oiwake, Kyoto 606-8502, Japan

(Received 4 June 2020; accepted 9 November 2020; published 10 December 2020)

### Quark Level Prediction

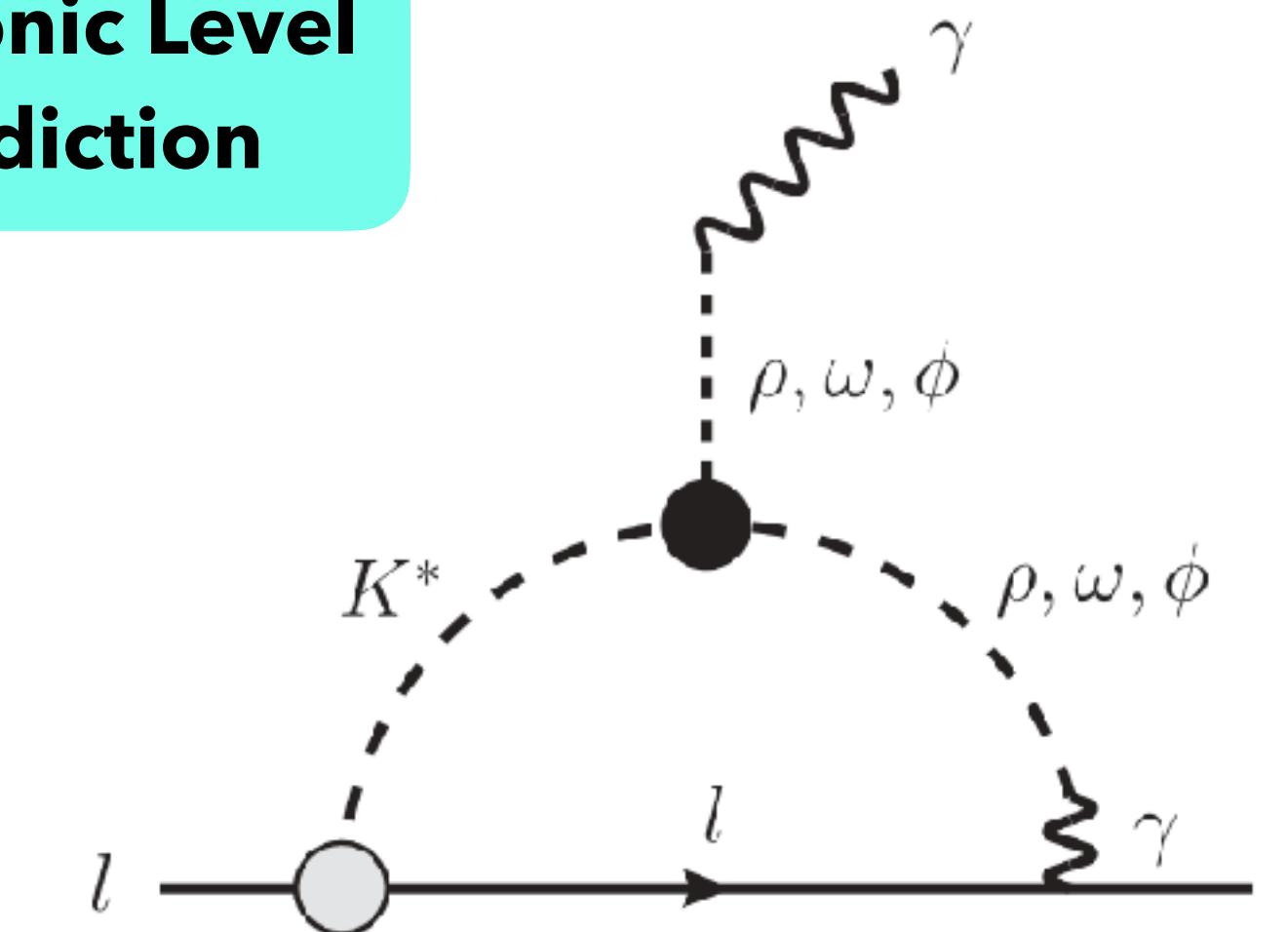


4-loop effect

Cancellation through GIM Mechanism

$$d_\mu \sim \mathcal{O}(10^{-48} \text{ ecm})$$

### Hadronic Level Prediction



1-loop effect

Less cancellation due to different momenta

$$d_\mu = 1.4 \times 10^{-38} \text{ ecm}$$

# Lepton Flavour Violation

PHYSICAL REVIEW D 98, 113002 (2018)

## Combined explanations of $(g - 2)_{\mu,e}$ and implications for a large muon EDM

Andreas Crivellin,<sup>1</sup> Martin Hoferichter,<sup>2</sup> and Philipp Schmidt-Wellenburg<sup>1</sup>

<sup>1</sup>*Paul Scherrer Institut, CH-5232 Villigen PSI, Switzerland*

<sup>2</sup>*Institute for Nuclear Theory, University of Washington, Seattle, Washington 98195-1550, USA*

(Received 3 August 2018; published 7 December 2018)

A BSM model assuming Minimal Flavour Violation (MFV) leads to a scaling of the lepton EDMS with mass, as expected with Lepton Flavour Universality (LFU):

$$d_e \leq 4.1 \times 10^{-30} e \cdot \text{cm} \xrightarrow{\text{MFV}} d_\mu \leq \frac{m_\mu}{m_e} d_e = 6.0 \times 10^{-28} e \cdot \text{cm}$$

(Roussy et al., 2023)

# Lepton Flavour Violation

PHYSICAL REVIEW D 98, 113002 (2018)

## Combined explanations of $(g - 2)_{\mu,e}$ and implications for a large muon EDM

Andreas Crivellin,<sup>1</sup> Martin Hoferichter,<sup>2</sup> and Philipp Schmidt-Wellenburg<sup>1</sup>

<sup>1</sup>Paul Scherrer Institut, CH-5232 Villigen PSI, Switzerland

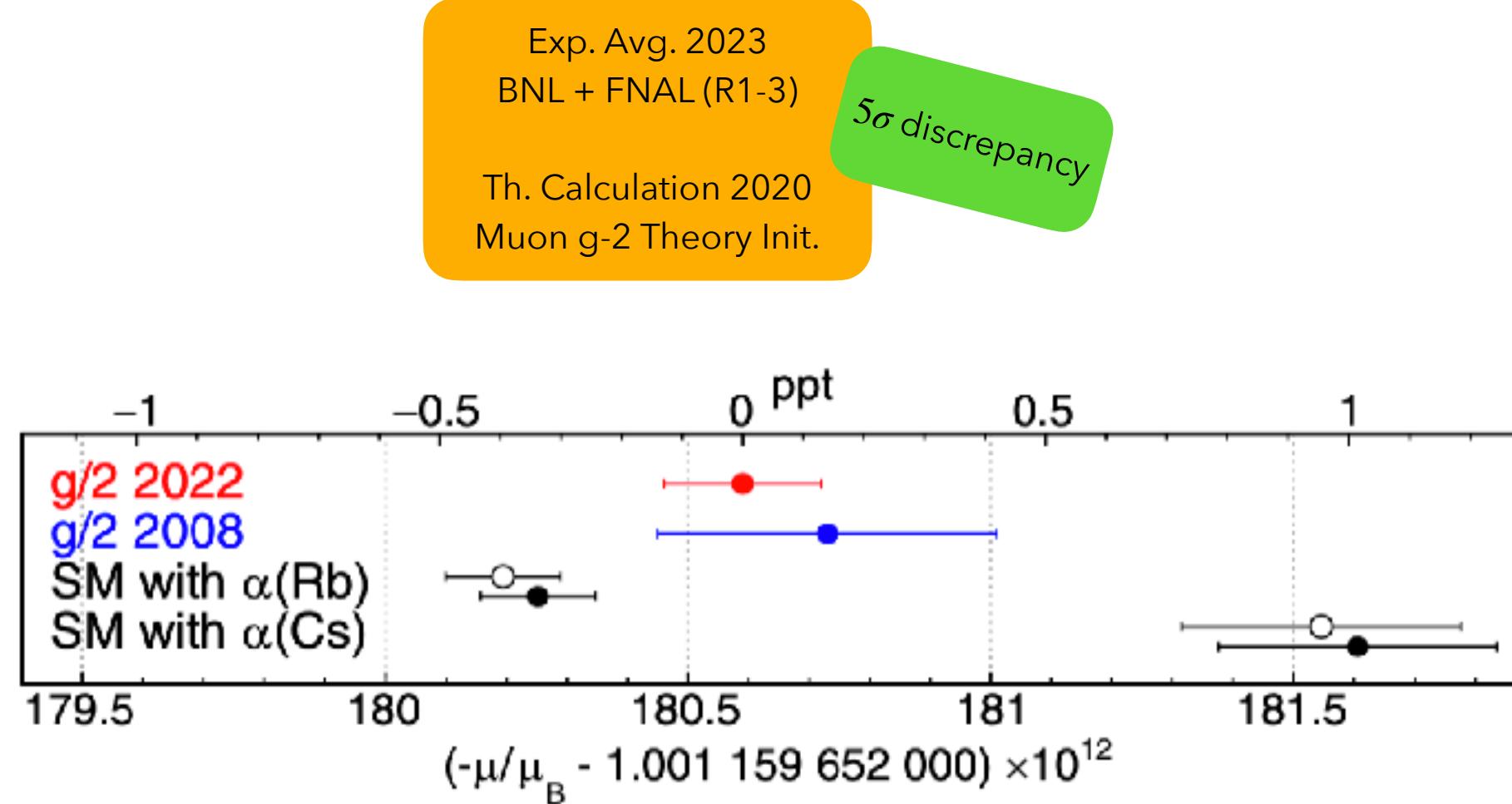
<sup>2</sup>Institute for Nuclear Theory, University of Washington, Seattle, Washington 98195-1550, USA

(Received 3 August 2018; published 7 December 2018)

A BSM model assuming Minimal Flavour Violation (MFV) leads to a scaling of the lepton EDMS with mass, as expected with Lepton Flavour Universality (LFU):

$$d_e \leq 4.1 \times 10^{-30} e \cdot \text{cm} \xrightarrow{\text{MFV}} d_\mu \leq \frac{m_\mu}{m_e} d_e = 6.0 \times 10^{-28} e \cdot \text{cm}$$

(Roussy et al., 2023)

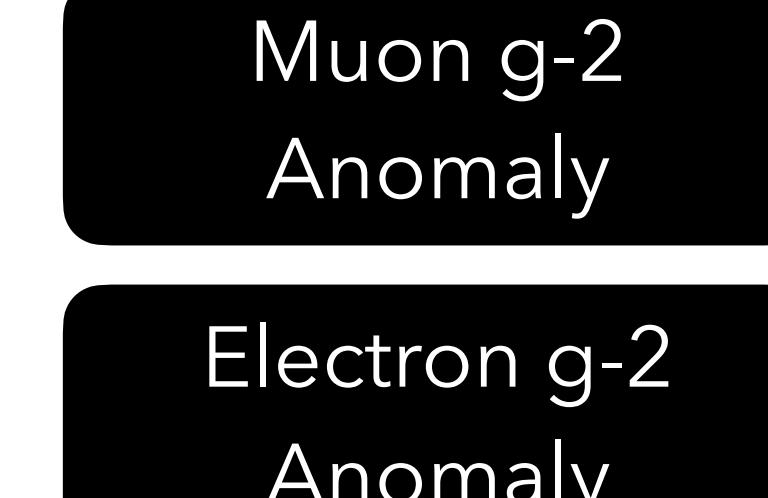


$$\text{Br}[\mu \rightarrow e\gamma] < 4.2 \times 10^{-13} \text{ (90 \% C.L.)}$$

(MEG Collab., 2016)

$$\Delta a_\mu = a_\mu^{\text{exp}} - a_\mu^{\text{SM}}$$

$$\Delta a_e = a_e^{\text{exp}} - a_e^{\text{SM}}$$



Given the constraints from MEG, and while the g-2 anomalies persist, a BSM theory addressing such anomalies should decouple  $e$  and  $\mu$  sectors, accommodating a large muon EDM.

# Lepton Flavour Violation

PHYSICAL REVIEW D 98, 113002 (2018)

Combined explanations of  $(g - 2)_{\mu,e}$  and implications for a large muon EDM

Andreas Crivellin,<sup>1</sup> Martin Hoferichter,<sup>2</sup> and Philipp Schmidt-Wellenburg<sup>1</sup>

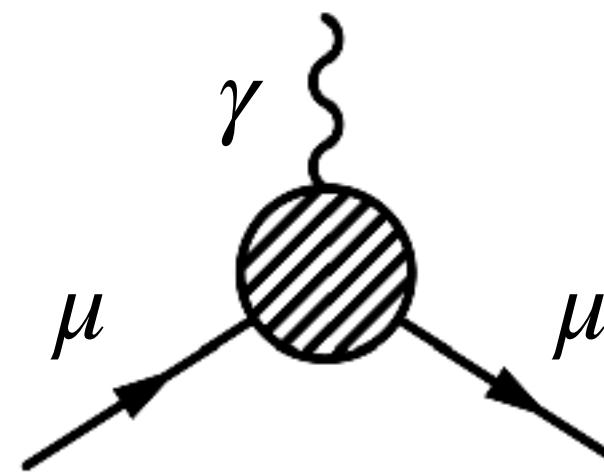
<sup>1</sup>Paul Scherrer Institut, CH-5232 Villigen PSI, Switzerland

<sup>2</sup>Institute for Nuclear Theory, University of Washington, Seattle, Washington 98195-1550, USA

(Received 3 August 2018; published 7 December 2018)

The EFT includes terms for the magnetic ( $\sigma_{\alpha\beta}$ ) and electric ( $\sigma_{\alpha\beta}\gamma^5$ ) dipole moments:

$$H_{\text{eff}} = c_R^{\ell_f \ell_i} \bar{\ell}_f \sigma_{\alpha\beta} P_R \ell_i F^{\alpha\beta} + \text{h.c.}$$



Expanding the terms ( $P_R = (1 + \gamma^5)/2$ ) and reducing to low energy limit (dipole form factors as  $q^2 \rightarrow 0$ ), gives:

$$a_\mu^{(\text{eff})} = \frac{4m_\mu}{e} \text{Re}(c_R^{\mu\mu}) \quad d_\mu^{(\text{eff})} = -2\text{Im}(c_R^{\mu\mu})$$

# Lepton Flavour Violation

PHYSICAL REVIEW D 98, 113002 (2018)

## Combined explanations of $(g - 2)_{\mu,e}$ and implications for a large muon EDM

Andreas Crivellin,<sup>1</sup> Martin Hoferichter,<sup>2</sup> and Philipp Schmidt-Wellenburg<sup>1</sup>

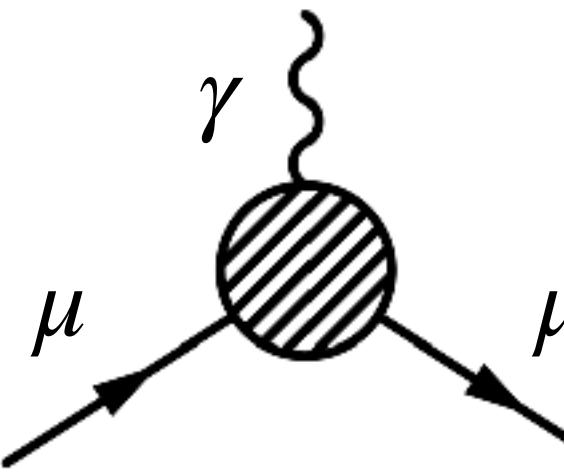
<sup>1</sup>Paul Scherrer Institut, CH-5232 Villigen PSI, Switzerland

<sup>2</sup>Institute for Nuclear Theory, University of Washington, Seattle, Washington 98195-1550, USA

(Received 3 August 2018; published 7 December 2018)

The EFT includes terms for the magnetic ( $\sigma_{\alpha\beta}$ ) and electric ( $\sigma_{\alpha\beta}\gamma^5$ ) dipole moments:

$$H_{\text{eff}} = c_R^{\ell_f \ell_i} \bar{\ell}_f \sigma_{\alpha\beta} P_R \ell_i F^{\alpha\beta} + \text{h.c.}$$

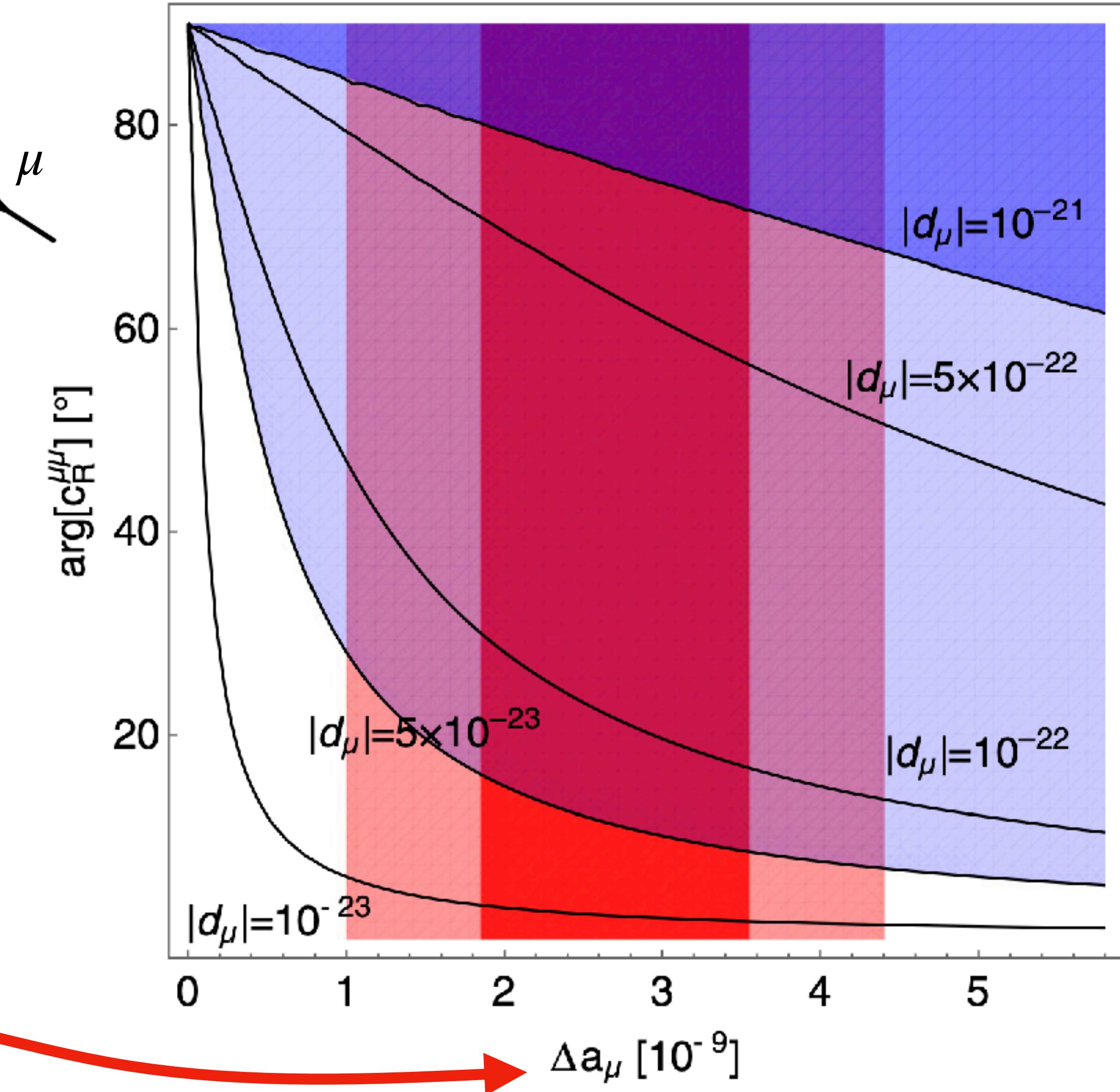


Expanding the terms ( $P_R = (1 + \gamma^5)/2$ ) and reducing to low energy limit (dipole form factors as  $q^2 \rightarrow 0$ ), gives:

$$a_\mu^{(\text{eff})} = \frac{4m_\mu}{e} \text{Re}(c_R^{\mu\mu}) \quad d_\mu^{(\text{eff})} = -2\text{Im}(c_R^{\mu\mu})$$

Precision eEDM searches have constrained  $\text{Im}(c_R^{ee})$ , but that of the muon  $\text{Im}(c_R^{\mu\mu})$  remains largely unconstrained:

$$\arg(c_R^{\mu\mu}) = \arctan \left( \frac{2m_\mu d_\mu^{(\text{eff})}}{e a_\mu^{(\text{eff})}} \right)$$



# Lepton Flavour Violation

PHYSICAL REVIEW D 98, 113002 (2018)

## Combined explanations of $(g - 2)_{\mu,e}$ and implications for a large muon EDM

Andreas Crivellin,<sup>1</sup> Martin Hoferichter,<sup>2</sup> and Philipp Schmidt-Wellenburg<sup>1</sup>

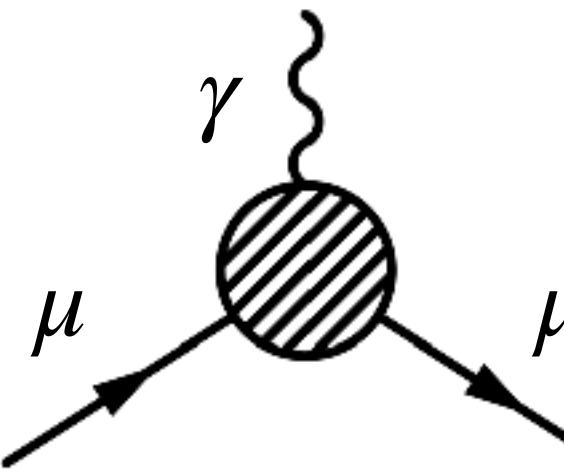
<sup>1</sup>Paul Scherrer Institut, CH-5232 Villigen PSI, Switzerland

<sup>2</sup>Institute for Nuclear Theory, University of Washington, Seattle, Washington 98195-1550, USA

(Received 3 August 2018; published 7 December 2018)

The EFT includes terms for the magnetic ( $\sigma_{\alpha\beta}$ ) and electric ( $\sigma_{\alpha\beta}\gamma^5$ ) dipole moments:

$$H_{\text{eff}} = c_R^{\ell_f \ell_i} \bar{\ell}_f \sigma_{\alpha\beta} P_R \ell_i F^{\alpha\beta} + \text{h.c.}$$

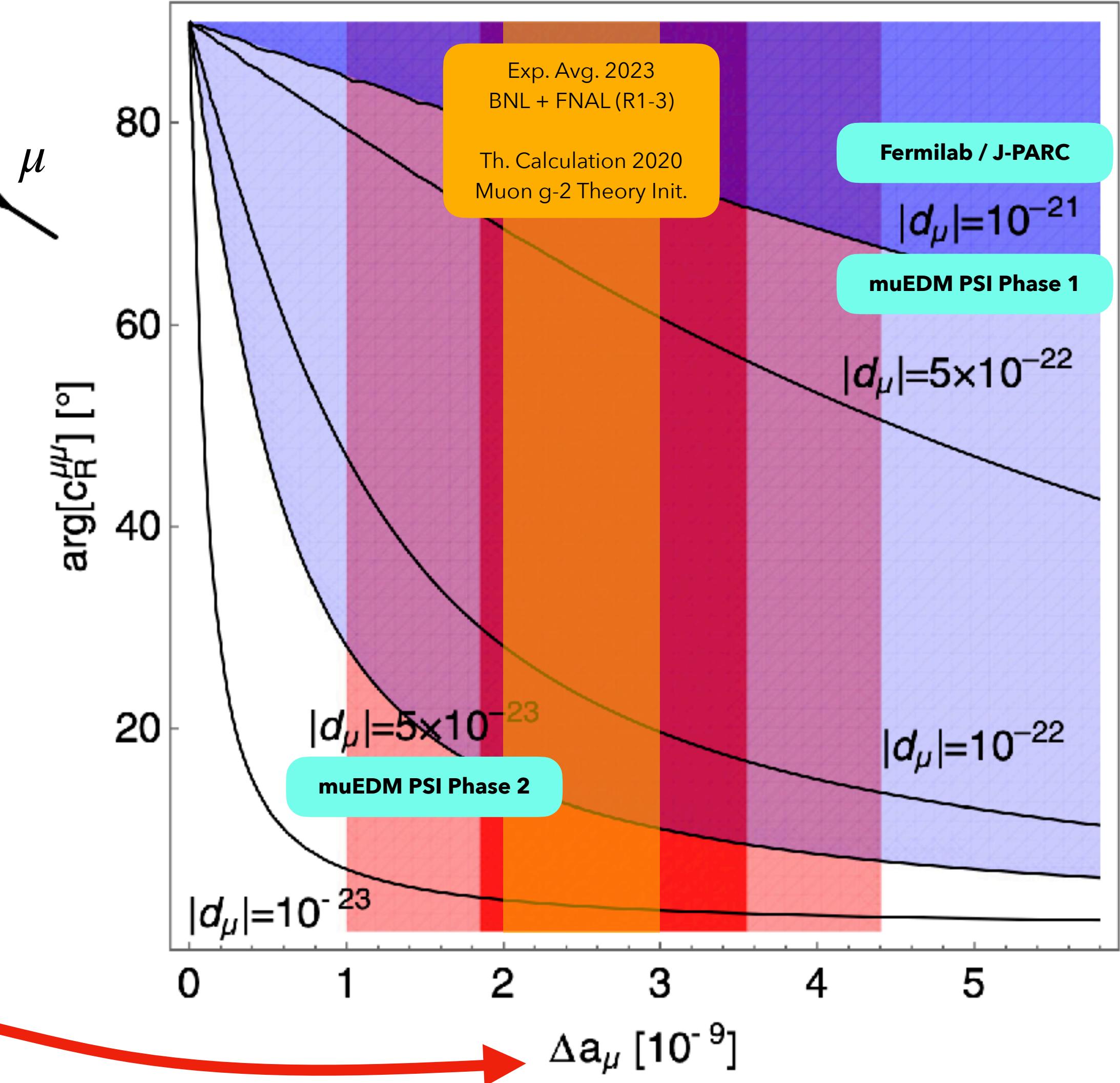


Expanding the terms ( $P_R = (1 + \gamma^5)/2$ ) and reducing to low energy limit (dipole form factors as  $q^2 \rightarrow 0$ ), gives:

$$a_\mu^{(\text{eff})} = \frac{4m_\mu}{e} \text{Re}(c_R^{\mu\mu}) \quad d_\mu^{(\text{eff})} = -2\text{Im}(c_R^{\mu\mu})$$

Precision eEDM searches have constrained  $\text{Im}(c_R^{ee})$ , but that of the muon  $\text{Im}(c_R^{\mu\mu})$  remains largely unconstrained:

$$\arg(c_R^{\mu\mu}) = \arctan \left( \frac{2m_\mu d_\mu^{(\text{eff})}}{e a_\mu^{(\text{eff})}} \right)$$



# Indirect Limits

Atomic EDMs can also constrain CP-violating observables.  
The atomic EDM (due to relativistic corrections that  
counter the Schiff Theorem) is scaling as  $d_{\text{atom}} \propto \alpha^2 Z^3$

PHYSICAL REVIEW LETTERS 128, 131803 (2022)

## Improved Indirect Limits on Muon Electric Dipole Moment

Yohei Ema<sup>1,\*</sup>, Ting Gao<sup>2,†</sup>, and Maxim Pospelov<sup>2,3,‡</sup>

<sup>1</sup>Deutsches Elektronen-Synchrotron DESY, Notkestraße 85, 22607 Hamburg, Germany

<sup>2</sup>School of Physics and Astronomy, University of Minnesota, Minneapolis, Minnesota 55455, USA

<sup>3</sup>William I. Fine Theoretical Physics Institute, School of Physics and Astronomy, University of Minnesota, Minneapolis, Minnesota 55455, USA

(Received 5 October 2021; revised 27 January 2022; accepted 7 March 2022; published 1 April 2022)

# Indirect Limits

PHYSICAL REVIEW LETTERS 128, 131803 (2022)

## Improved Indirect Limits on Muon Electric Dipole Moment

Yohei Ema<sup>1,\*</sup>, Ting Gao<sup>2,†</sup>, and Maxim Pospelov<sup>2,3,‡</sup>

<sup>1</sup>Deutsches Elektronen-Synchrotron DESY, Notkestraße 85, 22607 Hamburg, Germany

<sup>2</sup>School of Physics and Astronomy, University of Minnesota, Minneapolis, Minnesota 55455, USA

<sup>3</sup>William I. Fine Theoretical Physics Institute, School of Physics and Astronomy, University of Minnesota, Minneapolis, Minnesota 55455, USA

(Received 5 October 2021; revised 27 January 2022; accepted 7 March 2022; published 1 April 2022)

Atomic EDMs can also constrain CP-violating observables.

The atomic EDM (due to relativistic corrections that

counter the Schiff Theorem) is scaling as  $d_{\text{atom}} \propto \alpha^2 Z^3$

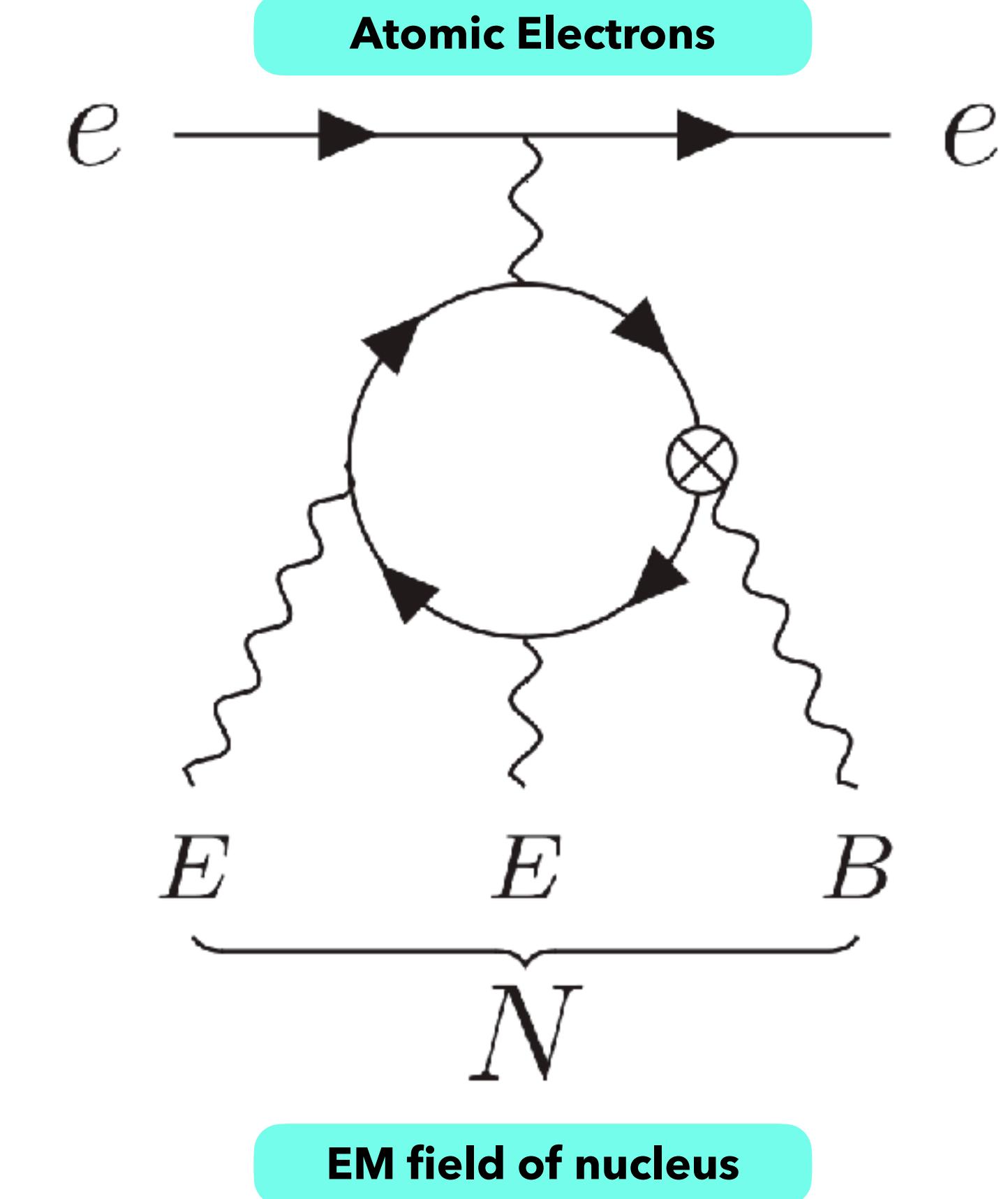
Low photon momenta inside  
nucleus:  $q_\gamma \sim 30 \text{ MeV} < m_\mu$

Treat only the dominant  $\mathbf{E}^3 \mathbf{B}$   
interaction inside the nucleus

$$\mathcal{L} = -\frac{d_\mu e^3}{12\pi^2 m_\mu^3} (\mathbf{E} \cdot \mathbf{B})(\mathbf{E} \cdot \mathbf{E} - \mathbf{B} \cdot \mathbf{B})$$

Nuclear  $\mathbf{E}(r)$  field based on  
collective charge properties

Nuclear  $\mathbf{B}(r, I)$  field estimated  
using shell model



## Improved Indirect Limits on Muon Electric Dipole Moment

Yohei Ema<sup>1,\*</sup>, Ting Gao<sup>2,†</sup>, and Maxim Pospelov<sup>2,3,‡</sup><sup>1</sup>Deutsches Elektronen-Synchrotron DESY, Notkestraße 85, 22607 Hamburg, Germany<sup>2</sup>School of Physics and Astronomy, University of Minnesota, Minneapolis, Minnesota 55455, USA<sup>3</sup>William I. Fine Theoretical Physics Institute, School of Physics and Astronomy, University of Minnesota, Minneapolis, Minnesota 55455, USA

(Received 5 October 2021; revised 27 January 2022; accepted 7 March 2022; published 1 April 2022)

# Indirect Limits

Atomic EDMs can also constrain CP-violating observables.

The atomic EDM (due to relativistic corrections that

counter the Schiff Theorem) is scaling as  $d_{\text{atom}} \propto \alpha^2 Z^3$ Low photon momenta inside nucleus:  $q_\gamma \sim 30 \text{ MeV} < m_\mu$ Treat only the dominant  $\mathbf{E}^3 \mathbf{B}$  interaction inside the nucleus

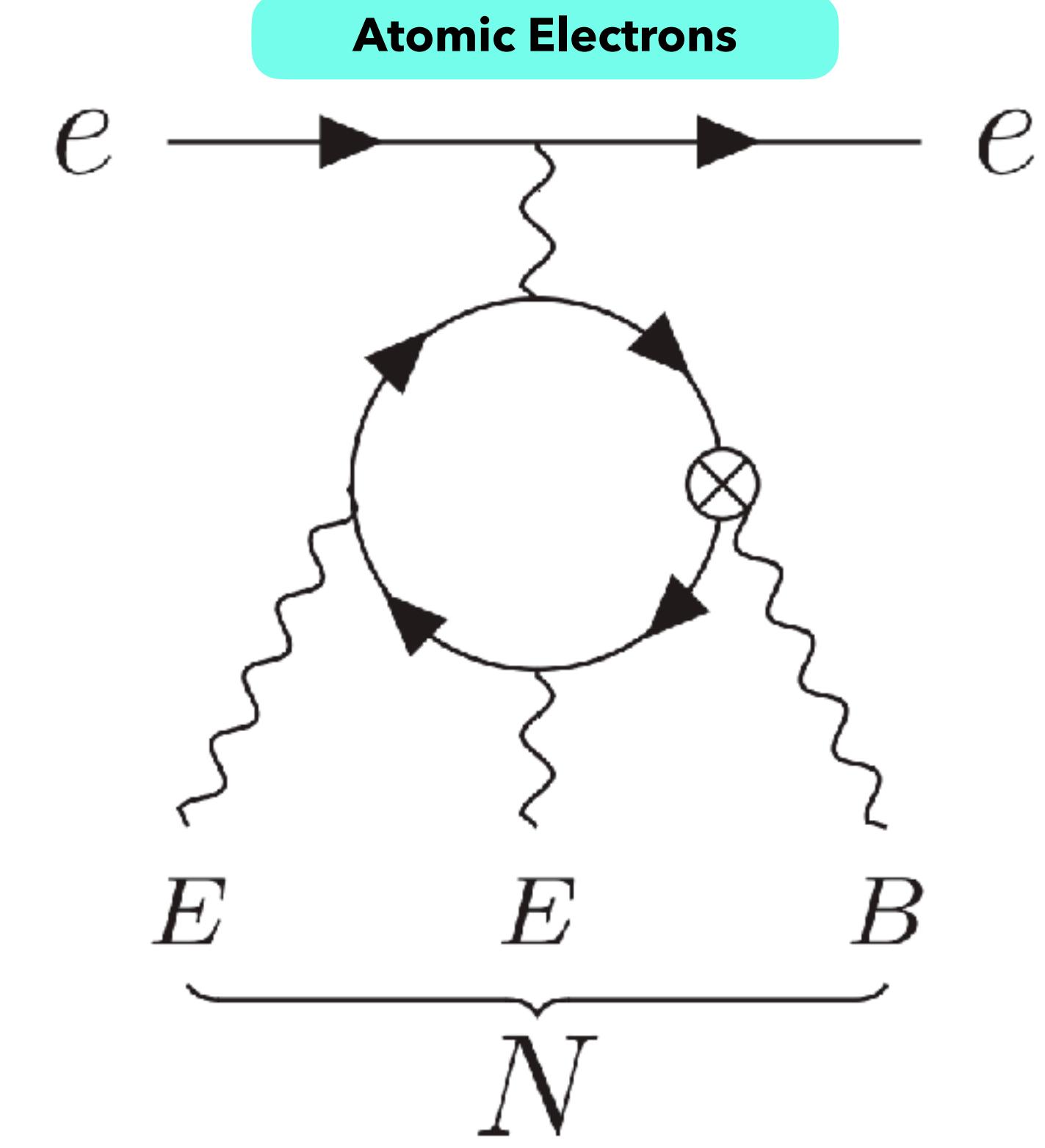
$$\mathcal{L} = -\frac{d_\mu e^3}{12\pi^2 m_\mu^3} (\mathbf{E} \cdot \mathbf{B})(\mathbf{E} \cdot \mathbf{E}) + \dots$$

Nuclear  $\mathbf{E}(r)$  field based on collective charge propertiesNuclear  $\mathbf{B}(r, I)$  field estimated using shell modelCalculate the Schiff moment  $S_N$  of the nucleus in terms of  $d_\mu$  and compare to measurement.

$$S_N = (\text{Nuc. Struct.}) \times -\frac{d_\mu}{m_\mu^3} \frac{Z^2}{\alpha^3} m_p R_N^2$$

$$|S_{^{199}\text{Hg}}^{(\text{exp})}| < 3.1 \times 10^{-13} \text{ efm}^3 \implies d_\mu < 6.4 \times 10^{-20} \text{ ecm}$$

Better than BNL direct limit



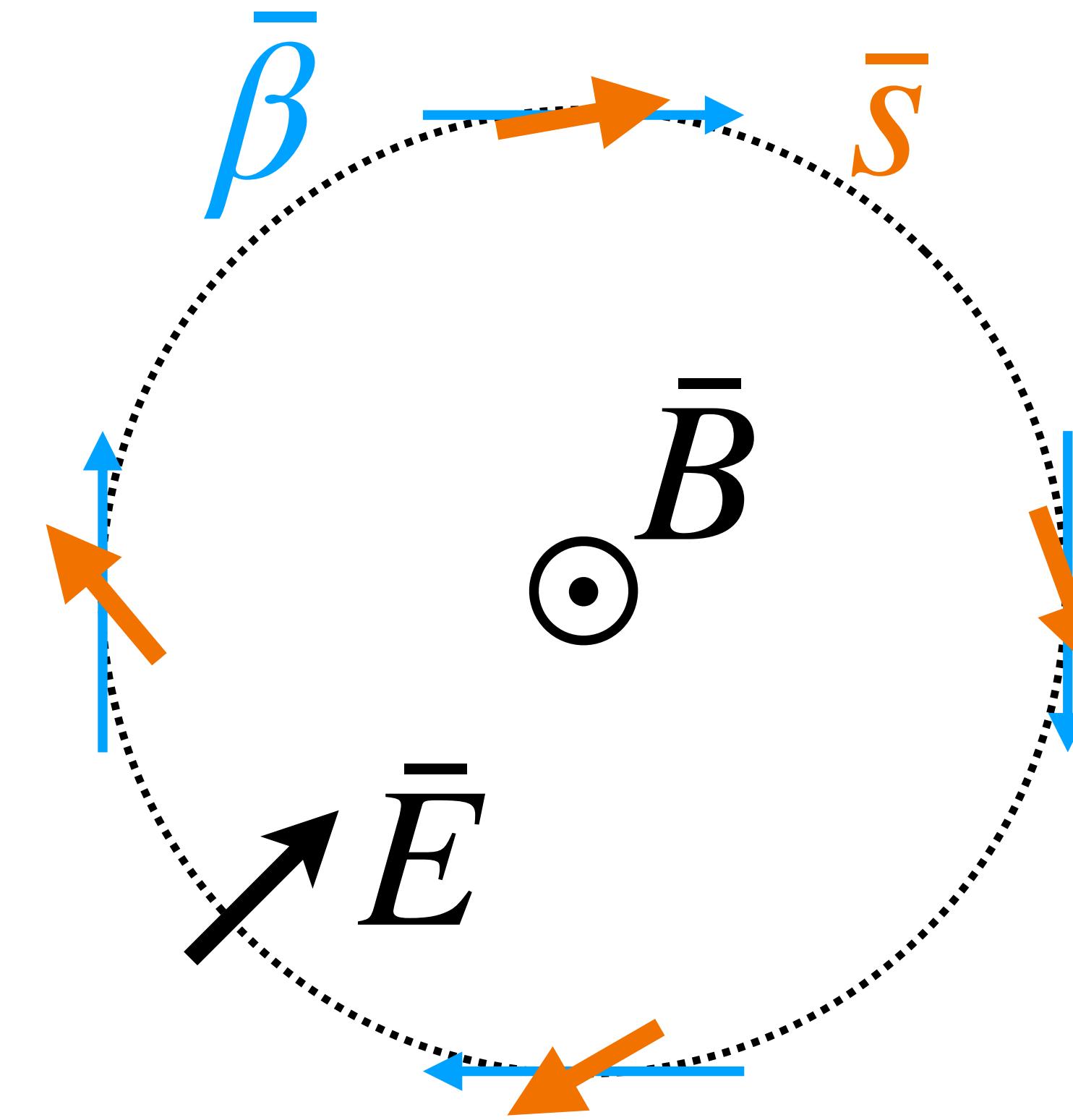
EM field of nucleus

# Spin Precession in a storage ring

$$\left( \vec{d}_\mu = \frac{\eta e}{2mc} \vec{s} \right)$$

Velocity :  $\frac{d\bar{\beta}}{dt} = \bar{\Omega}_c \times \bar{\beta}$        $\bar{\Omega} = \bar{\Omega}_0 - \bar{\Omega}_c$

Spin :  $\frac{d\bar{s}}{dt} = \bar{\Omega}_0 \times \bar{s}$



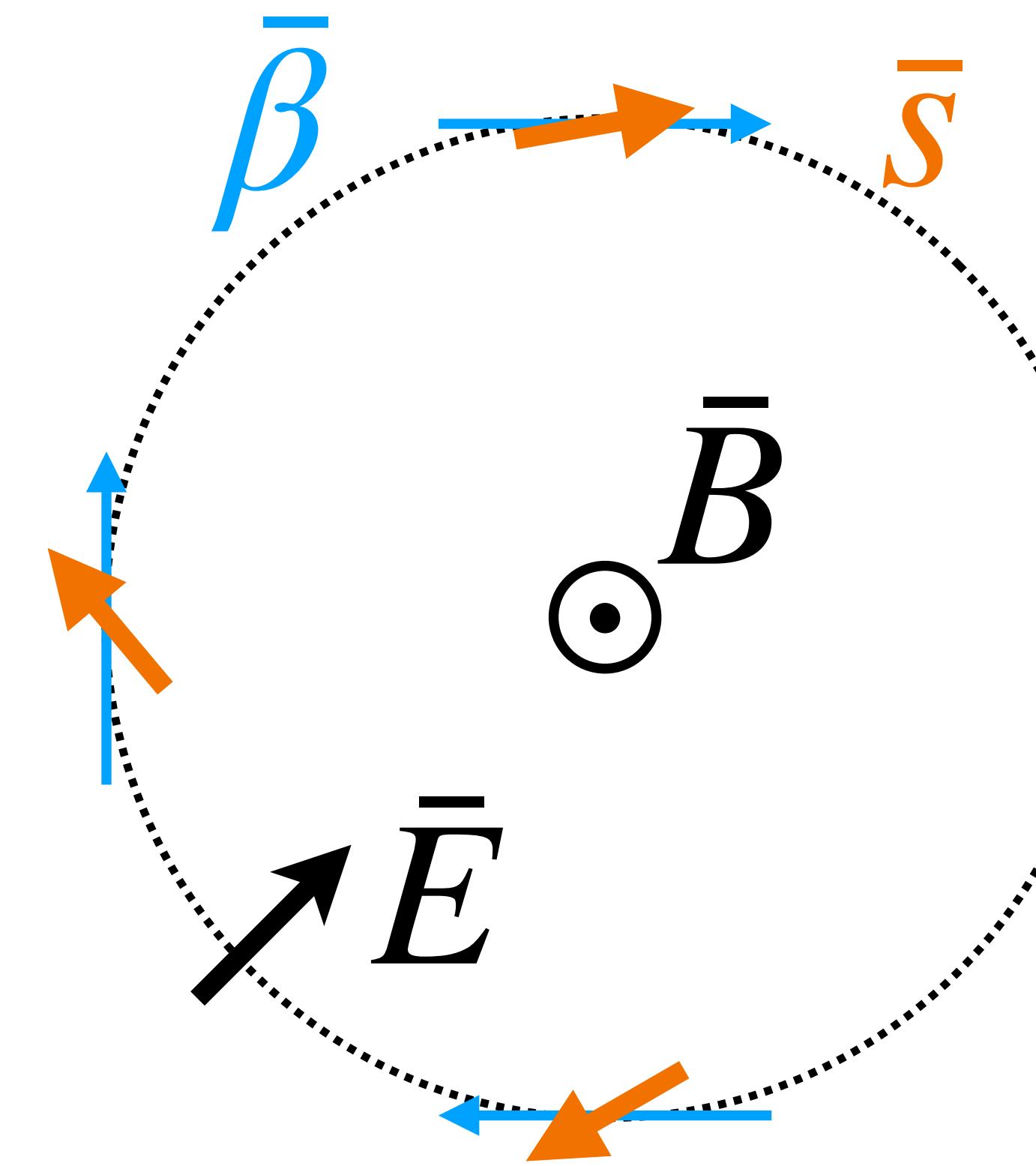
# Spin Precession in a storage ring

$$\left( \vec{d}_\mu = \frac{\eta e}{2mc} \vec{s} \right)$$

Velocity :  $\frac{d\bar{\beta}}{dt} = \bar{\Omega}_c \times \bar{\beta}$        $\bar{\Omega} = \bar{\Omega}_0 - \bar{\Omega}_c$

Spin :  $\frac{d\bar{s}}{dt} = \bar{\Omega}_0 \times \bar{s}$

$$= \frac{aq}{m} \left( \bar{B} - \frac{\gamma}{\gamma + 1} (\bar{\beta} \cdot \bar{B}) \bar{\beta} - \left( 1 + \frac{1}{a(1 - \gamma^2)} \right) \frac{\bar{\beta} \times \bar{E}}{c} \right)$$



$$+ \frac{\eta q}{2m} \left( \bar{\beta} \times \bar{B} + \frac{\bar{E}}{c} - \frac{\gamma/c}{\gamma + 1} (\bar{\beta} \cdot \bar{E}) \bar{\beta} \right)$$

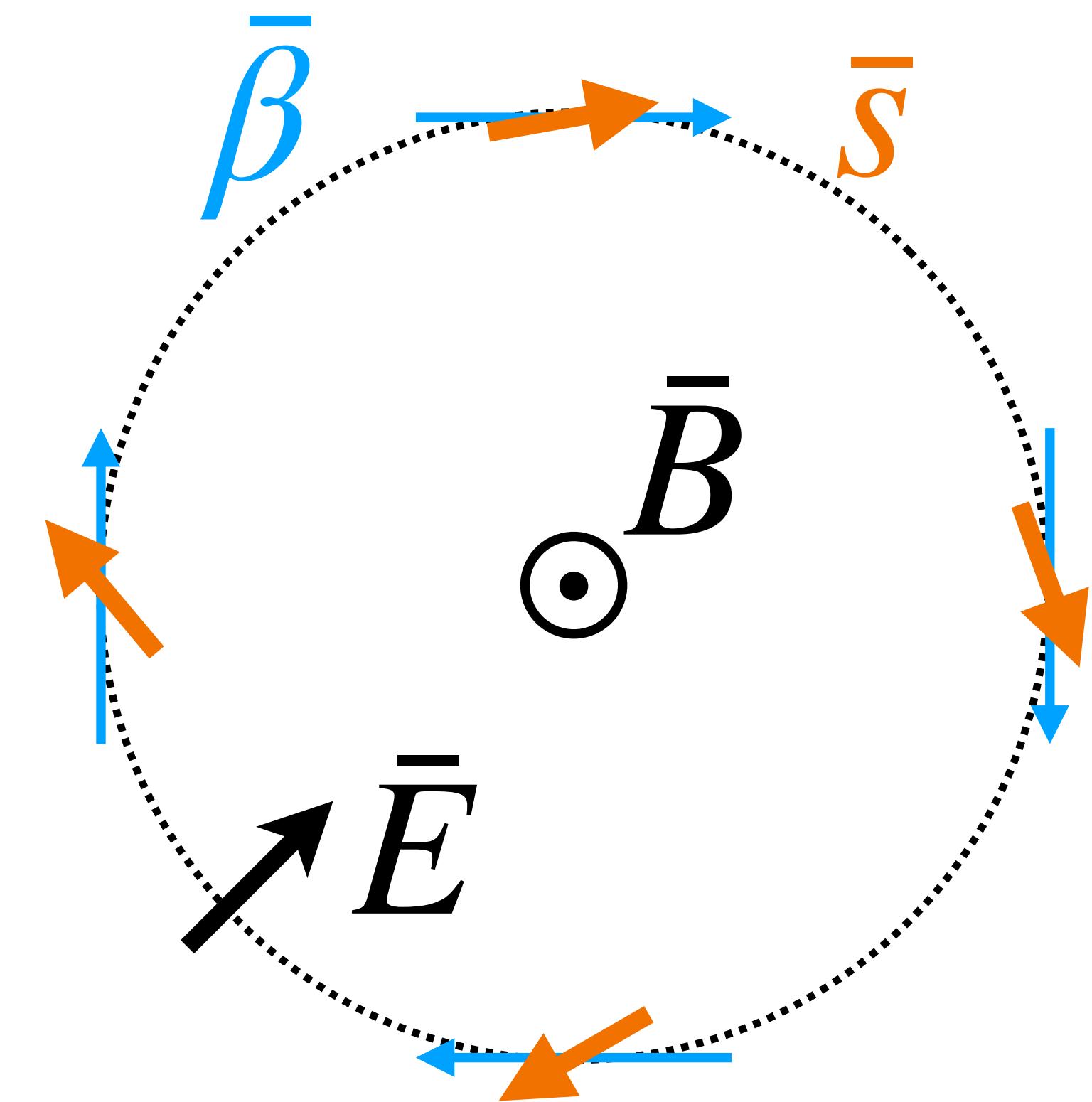
# Spin Precession in a storage ring

$$\left( \vec{d}_\mu = \frac{\eta e}{2mc} \vec{s} \right)$$

Velocity :  $\frac{d\bar{\beta}}{dt} = \bar{\Omega}_c \times \bar{\beta}$        $\bar{\Omega} = \bar{\Omega}_0 - \bar{\Omega}_c$

Spin :  $\frac{d\bar{s}}{dt} = \bar{\Omega}_0 \times \bar{s}$

$$= \frac{aq}{m} \left( \bar{B} - \frac{\gamma}{\gamma + 1} (\bar{\beta} \cdot \bar{B}) \bar{\beta} - \left( 1 + \frac{1}{a(1 - \gamma^2)} \right) \frac{\bar{\beta} \times \bar{E}}{c} \right)$$



$$+ \frac{\eta q}{2m} \left( \bar{\beta} \times \bar{B} + \frac{\bar{E}}{c} - \frac{\gamma/c}{\gamma + 1} (\bar{\beta} \cdot \bar{E}) \bar{\beta} \right)$$

$$\bar{\Omega} = \frac{aq}{m} \left( \bar{B} - \left( 1 + \frac{1}{a(1 - \gamma^2)} \right) \frac{\bar{\beta} \times \bar{E}}{c} \right) + \frac{\eta q}{2m} \left( \bar{\beta} \times \bar{B} + \frac{\bar{E}}{c} \right)$$

# Experimental Approaches

**Fermilab**    
$$\bar{\Omega} = \frac{aq}{m} \left( \bar{B} - \left( 1 + \frac{1}{a(1-\gamma^2)} \right) \frac{\bar{\beta} \times \bar{E}}{c} \right) + \frac{\eta q}{2m} \left( \bar{\beta} \times \bar{B} + \frac{\bar{E}}{c} \right)$$

“magic momentum”

**J-PARC**    
$$\bar{\Omega} = \frac{aq}{m} \left( \bar{B} - \left( 1 + \frac{1}{a(1-\gamma^2)} \right) \frac{\bar{\beta} \times \bar{E}}{c} \right) + \frac{\eta q}{2m} \left( \bar{\beta} \times \bar{B} + \frac{\bar{E}}{c} \right)$$

“E field shielding”

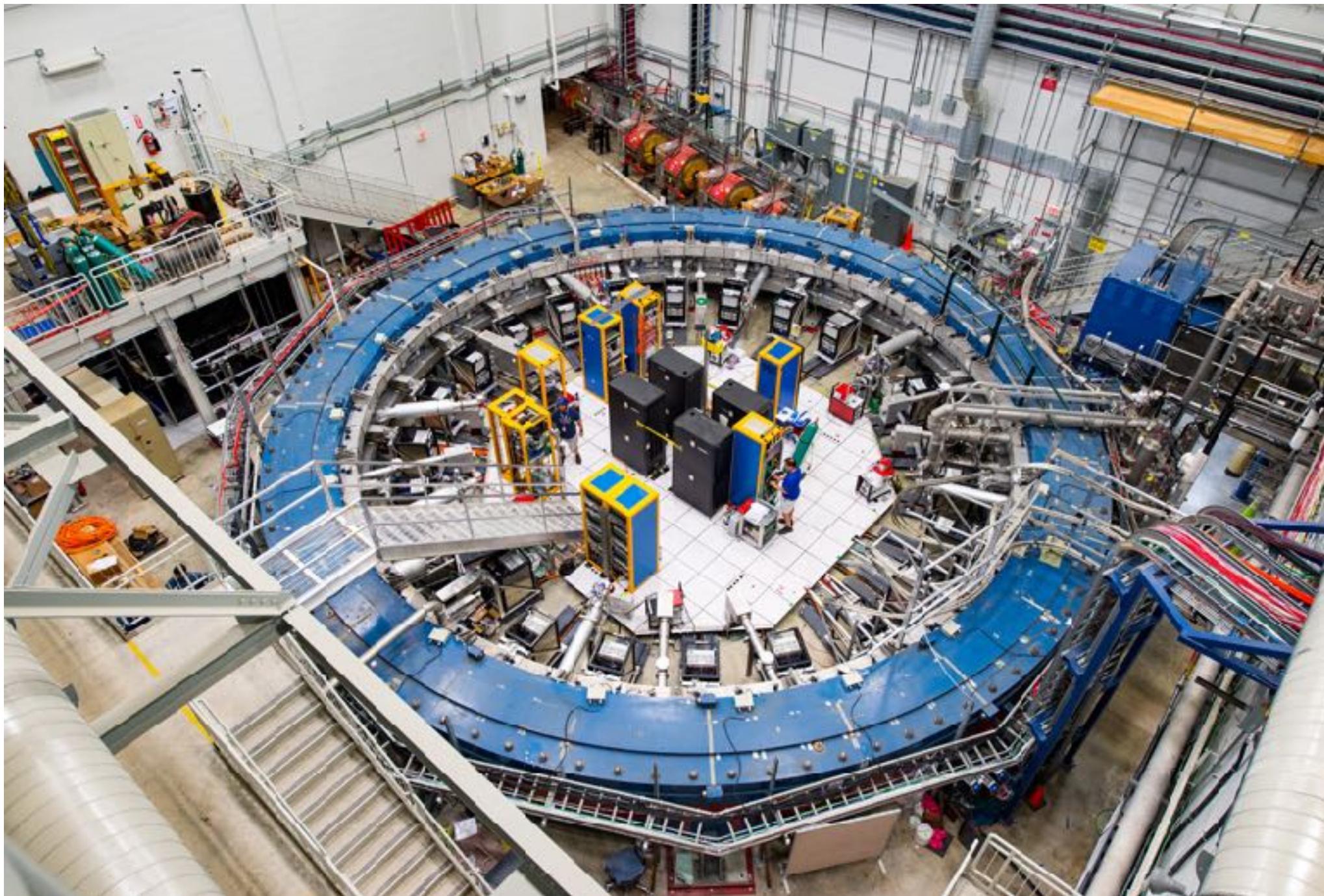
**PSI**    
$$\bar{\Omega} = \frac{aq}{m} \left( \bar{B} - \left( 1 + \frac{1}{a(1-\gamma^2)} \right) \frac{\bar{\beta} \times \bar{E}}{c} \right) + \frac{\eta q}{2m} \left( \bar{\beta} \times \bar{B} + \frac{\bar{E}}{c} \right)$$

“frozen spin”

# Muon g-2 at BNL & Fermilab

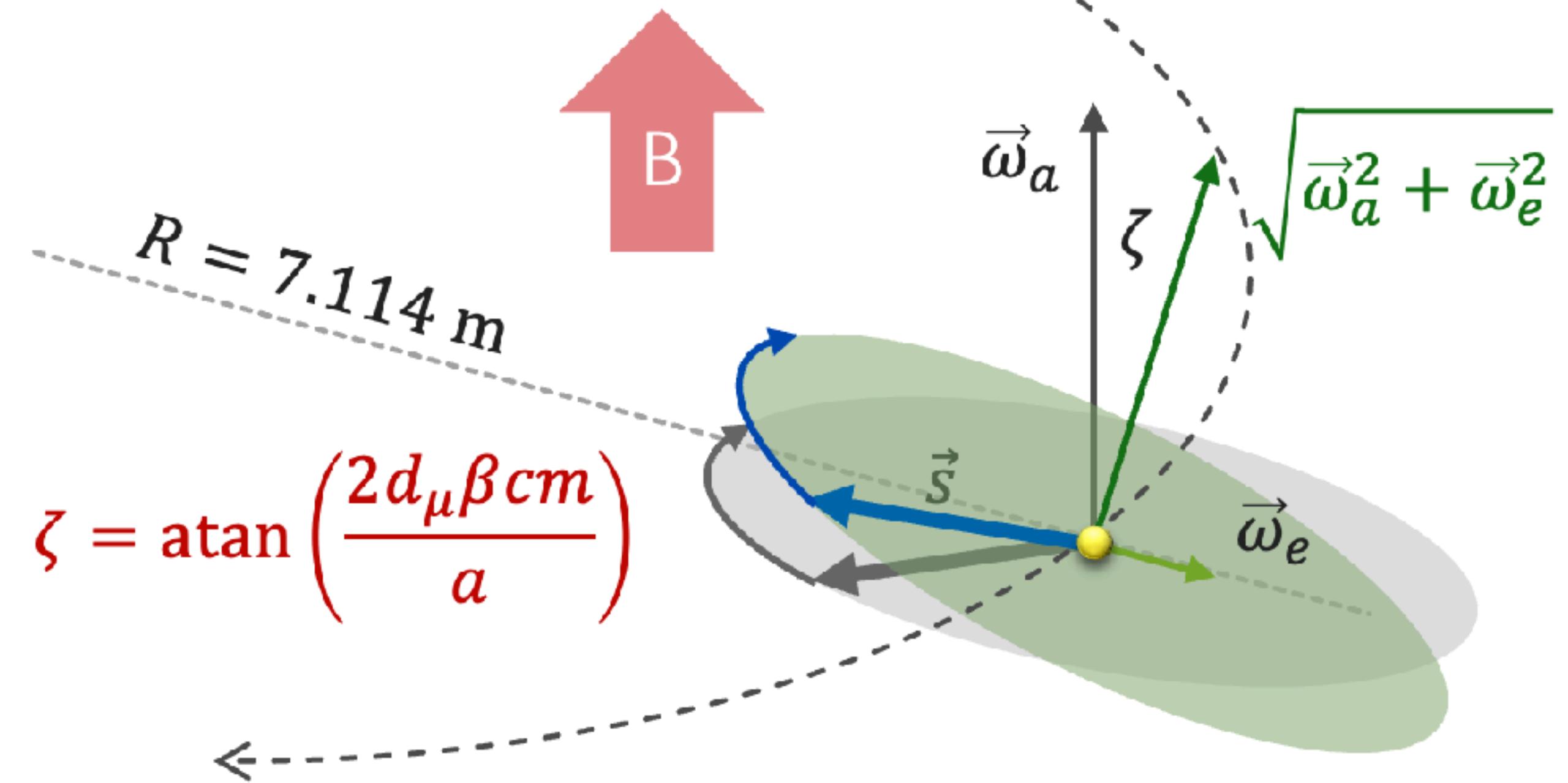
Pulsed (avg. bunch freq.  $11.4\text{ s}^{-1}$ ) muon beam  
with magic momentum  $3.1\text{ GeV}/c$

$$1 + \frac{1}{a_\mu(1 - \gamma^2)} \stackrel{!}{=} 0 \implies \gamma = 29.3$$



$$\bar{\Omega} \approx \frac{aq}{m}\bar{B} + \frac{\eta q}{2m}\vec{\beta} \times \bar{B}$$

**2 simultaneous precession axes!**



Credit: P. Schmidt-Wellenburg

# Frozen Spin Technique

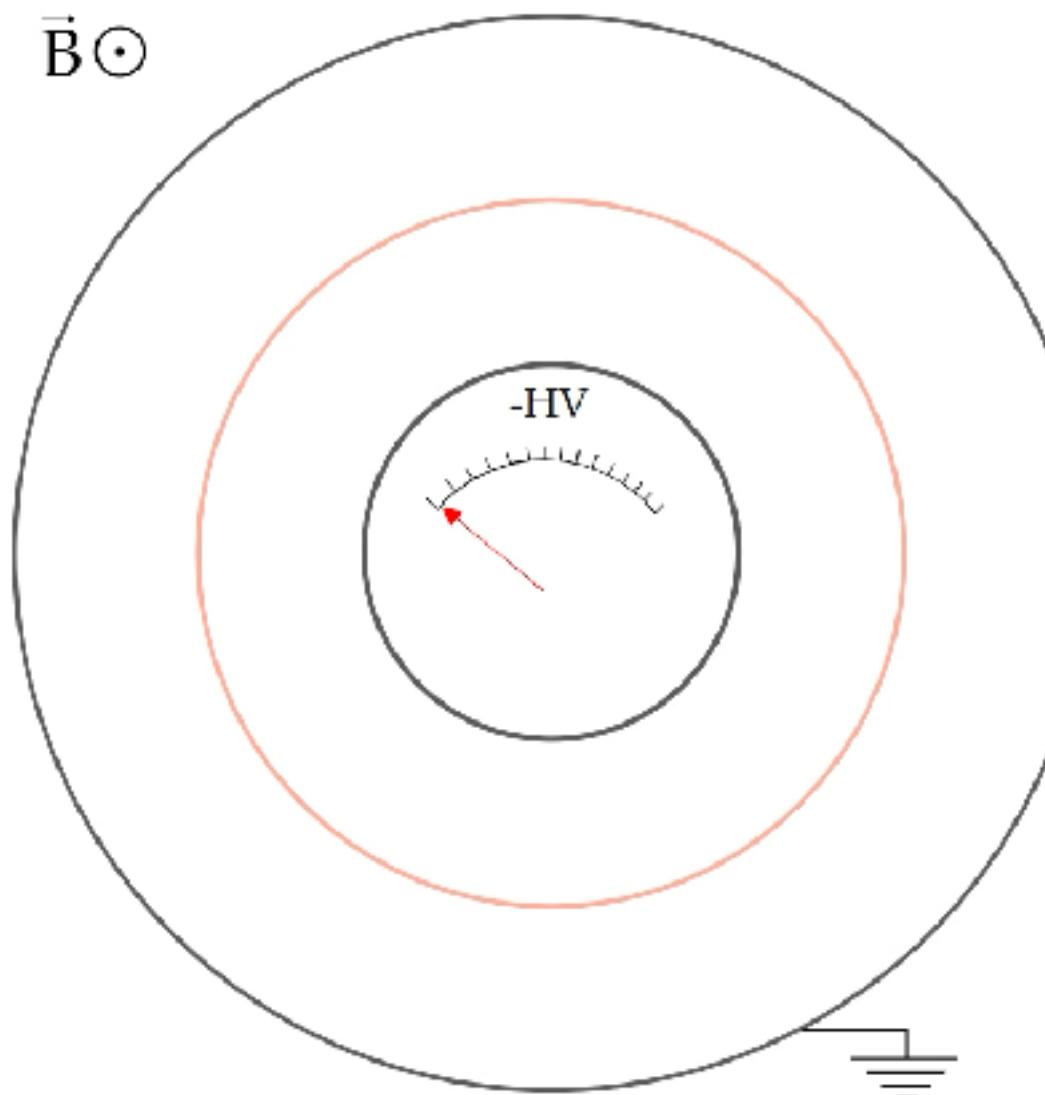
**Goal:** Configure E, B fields such that spin follows velocity vector and EDM is the only inherent source of spin precession.

$$\bar{\Omega} = \bar{\Omega}_0 - \bar{\Omega}_c = \frac{aq}{m} \left( \bar{B} - \frac{\gamma}{\gamma + 1} (\bar{\beta} \cdot \bar{B}) \bar{\beta} - \left( 1 + \frac{1}{a(1 - \gamma^2)} \right) \frac{\bar{\beta} \times \bar{E}}{c} \right) + \frac{\eta q}{2m} \left( \bar{\beta} \times \bar{B} + \frac{\bar{E}}{c} - \frac{\gamma/c}{\gamma + 1} (\bar{\beta} \cdot \bar{E}) \bar{\beta} \right)$$

# Frozen Spin Technique

**Goal:** Configure E, B fields such that spin follows velocity vector and EDM is the only inherent source of spin precession.

$$\bar{\Omega} = \bar{\Omega}_0 - \bar{\Omega}_c = \frac{aq}{m} \left( \bar{B} - \frac{\gamma}{\gamma + 1} (\bar{\beta} \cdot \bar{B}) \bar{\beta} - \left( 1 + \frac{1}{a(1 - \gamma^2)} \right) \frac{\bar{\beta} \times \bar{E}}{c} \right) + \frac{\eta q}{2m} \left( \bar{\beta} \times \bar{B} + \frac{\bar{E}}{c} - \frac{\gamma/c}{\gamma + 1} (\bar{\beta} \cdot \bar{E}) \bar{\beta} \right)$$

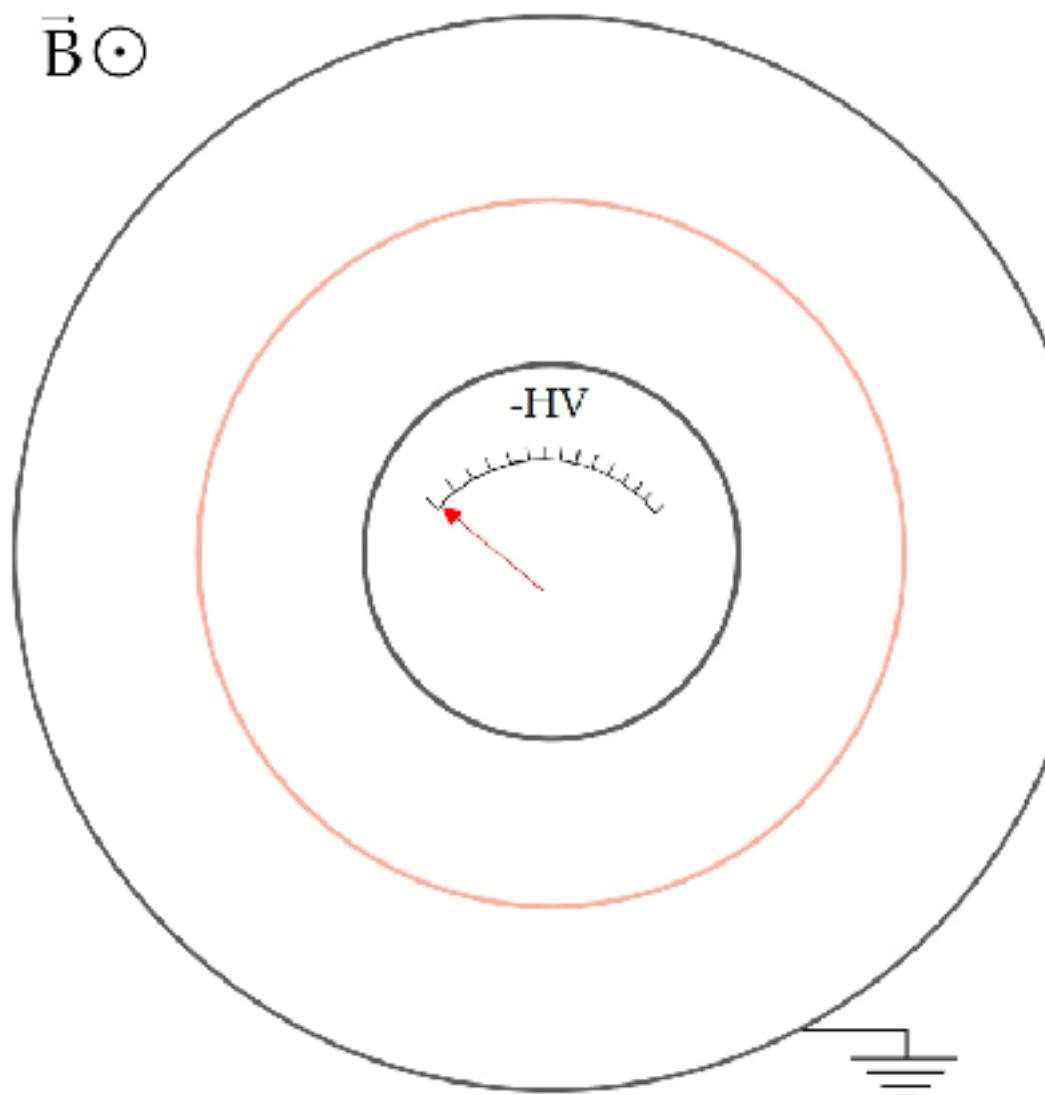


Frozen Spin Condition :  $E_f \stackrel{a \ll 1}{\approx} aB\beta c\gamma^2$

# Frozen Spin Technique

**Goal:** Configure E, B fields such that spin follows velocity vector and EDM is the only inherent source of spin precession.

$$\bar{\Omega} = \bar{\Omega}_0 - \bar{\Omega}_c = \frac{aq}{m} \left( \bar{B} - \frac{\gamma}{\gamma + 1} (\bar{\beta} \cdot \bar{B}) \bar{\beta} - \left( 1 + \frac{1}{a(1 - \gamma^2)} \right) \frac{\bar{\beta} \times \bar{E}}{c} \right) + \frac{\eta q}{2m} \left( \bar{\beta} \times \bar{B} + \frac{\bar{E}}{c} - \frac{\gamma/c}{\gamma + 1} (\bar{\beta} \cdot \bar{E}) \bar{\beta} \right)$$



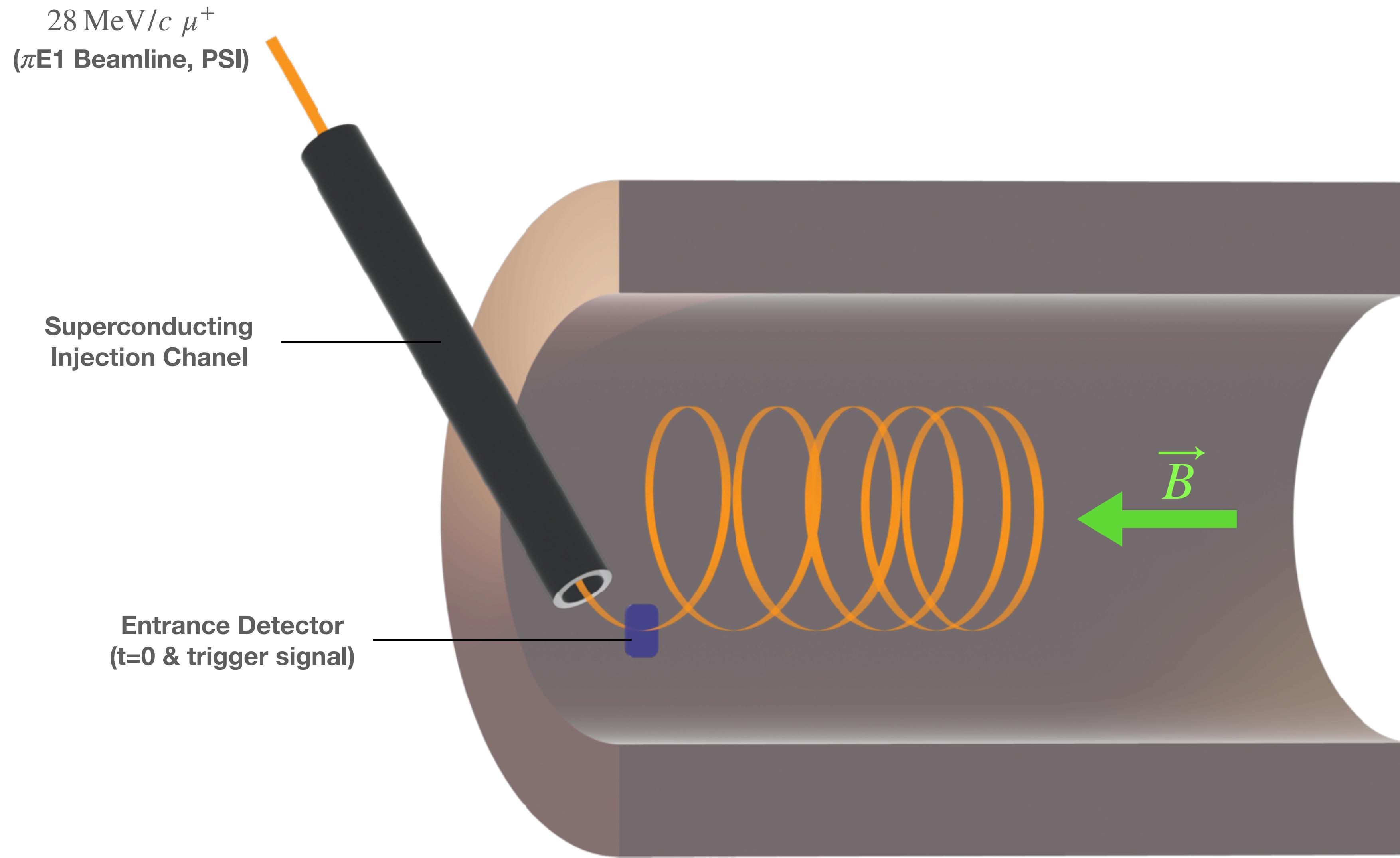
Frozen Spin Condition :  $E_f \stackrel{a < < 1}{\approx} aB\beta c\gamma^2$

## Experimental Requirements:

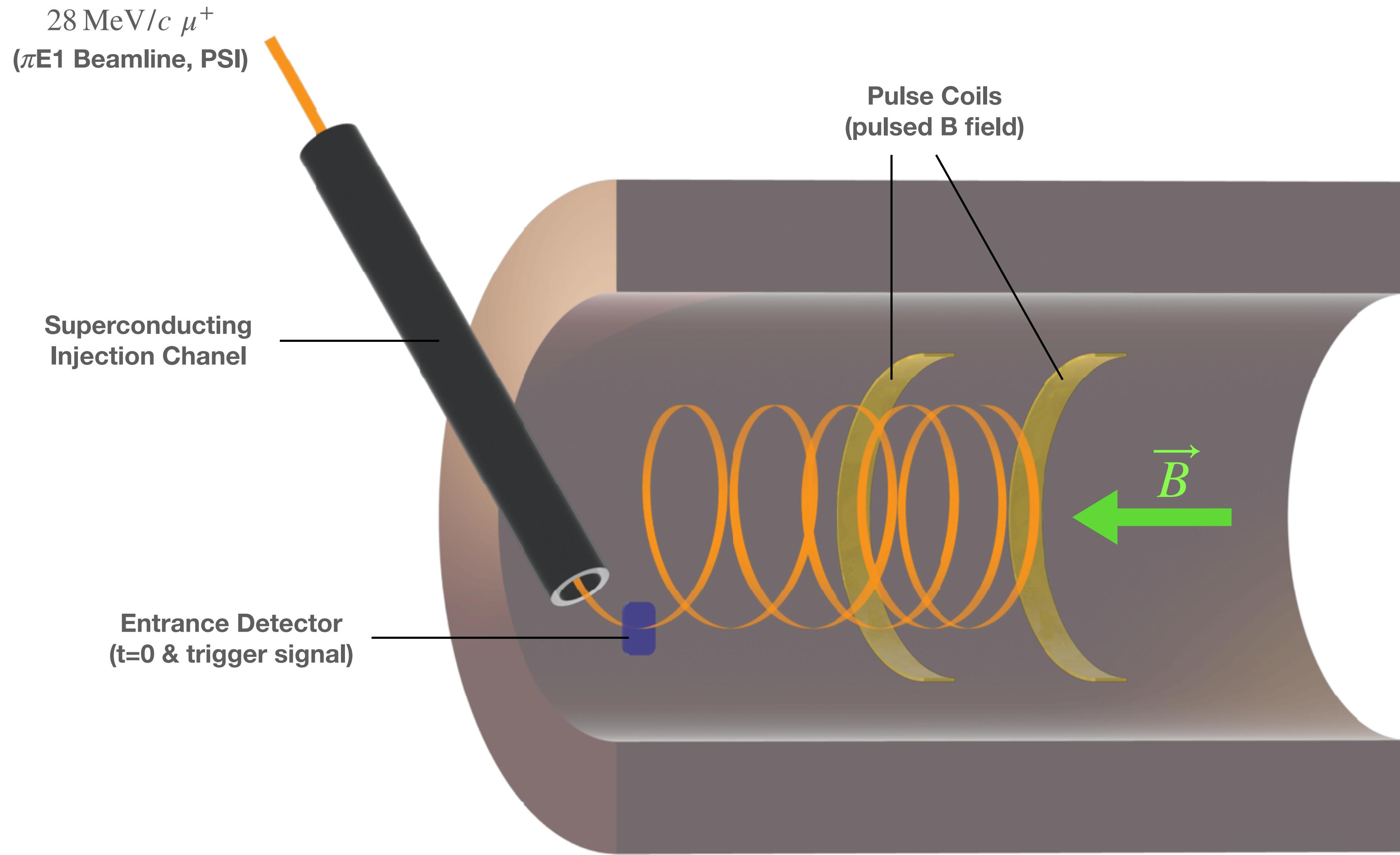
1. Fields  $\perp$  Velocity
2. Precisely tuned  $E = E_f$
3. Constrained  $B_r$  (radial),  $E_z$  (axial)

Any periodic deviations must be stable over the timescale of  $\tau_\mu$ .

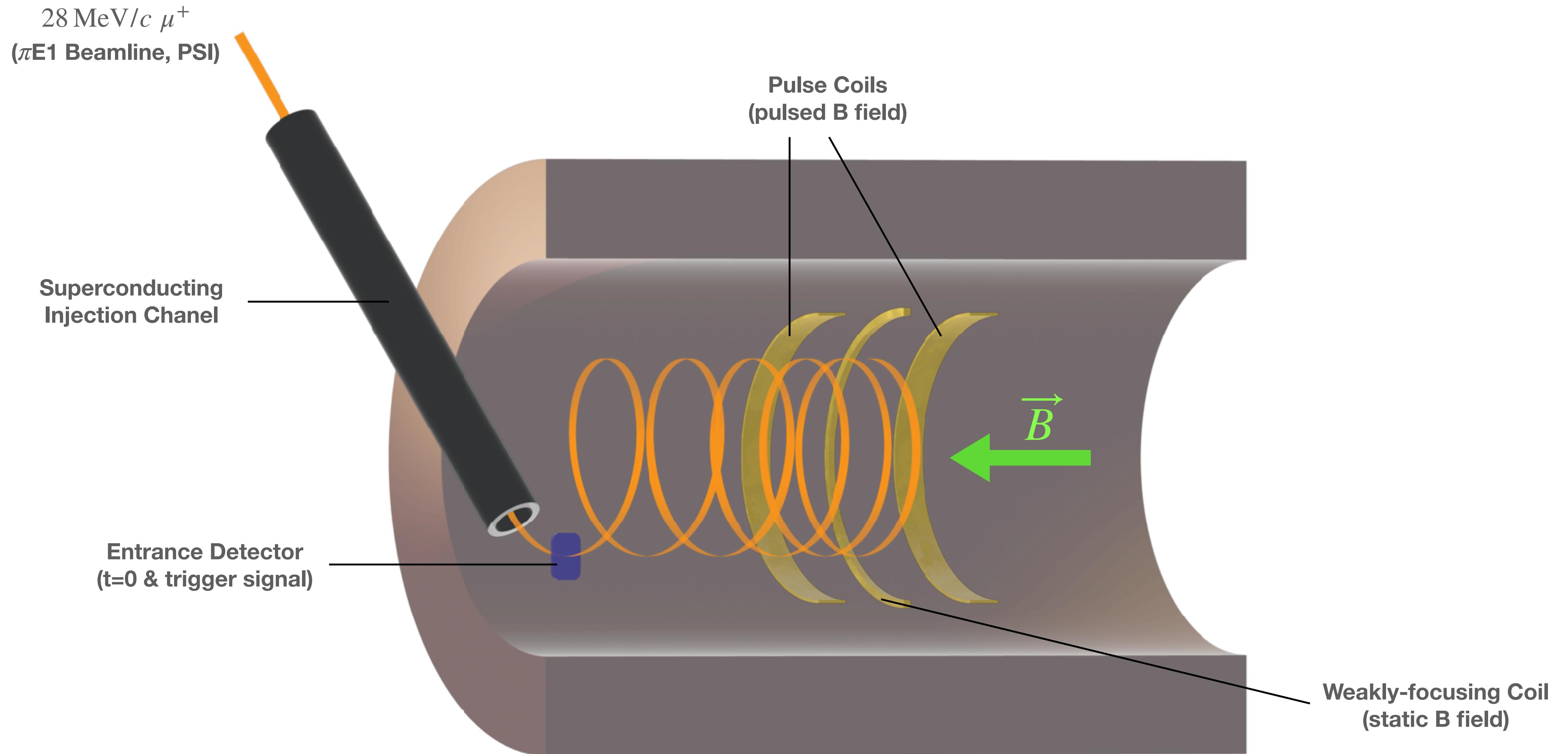
# muEDM Experiment at PSI



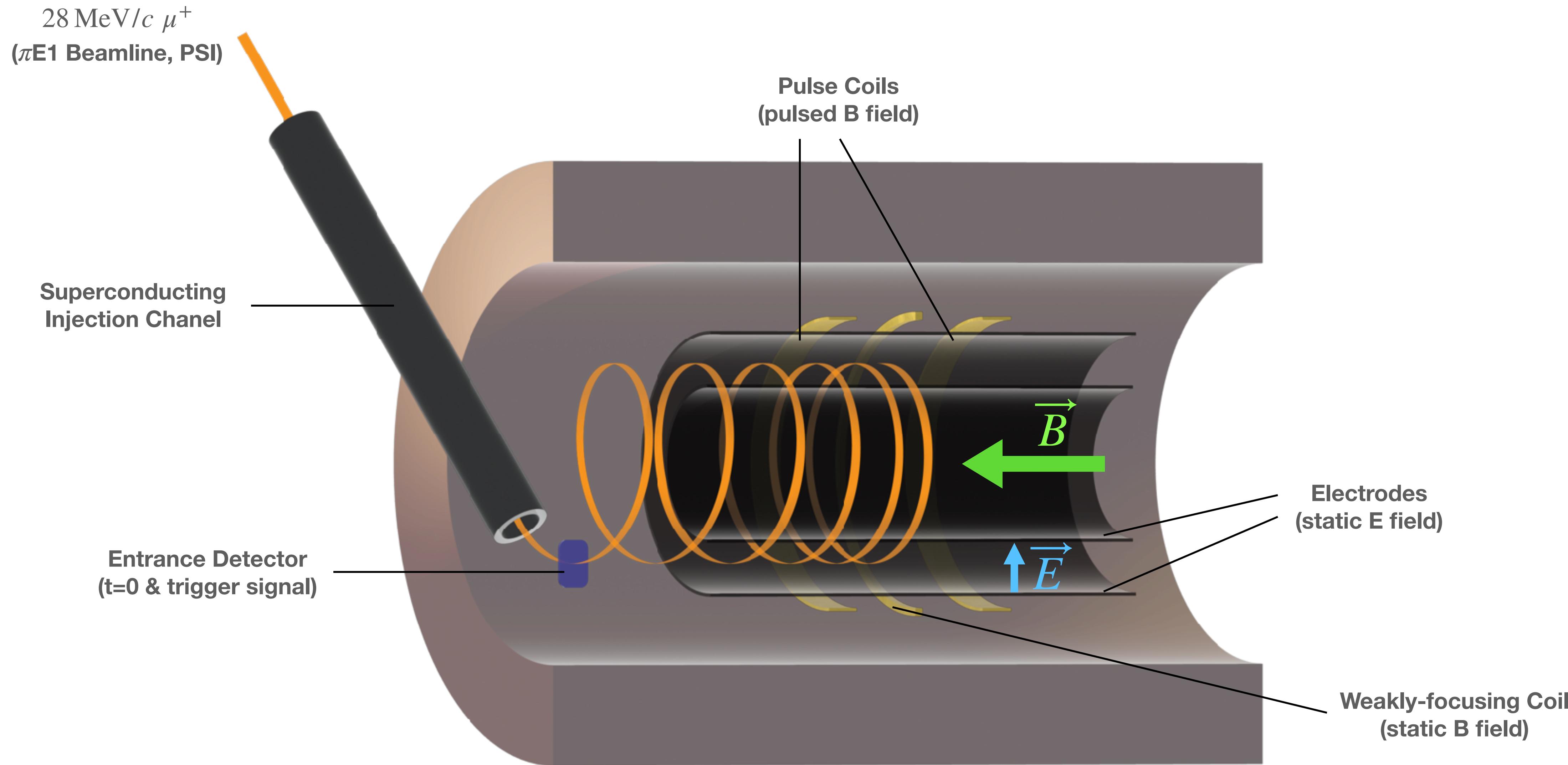
# muEDM Experiment at PSI



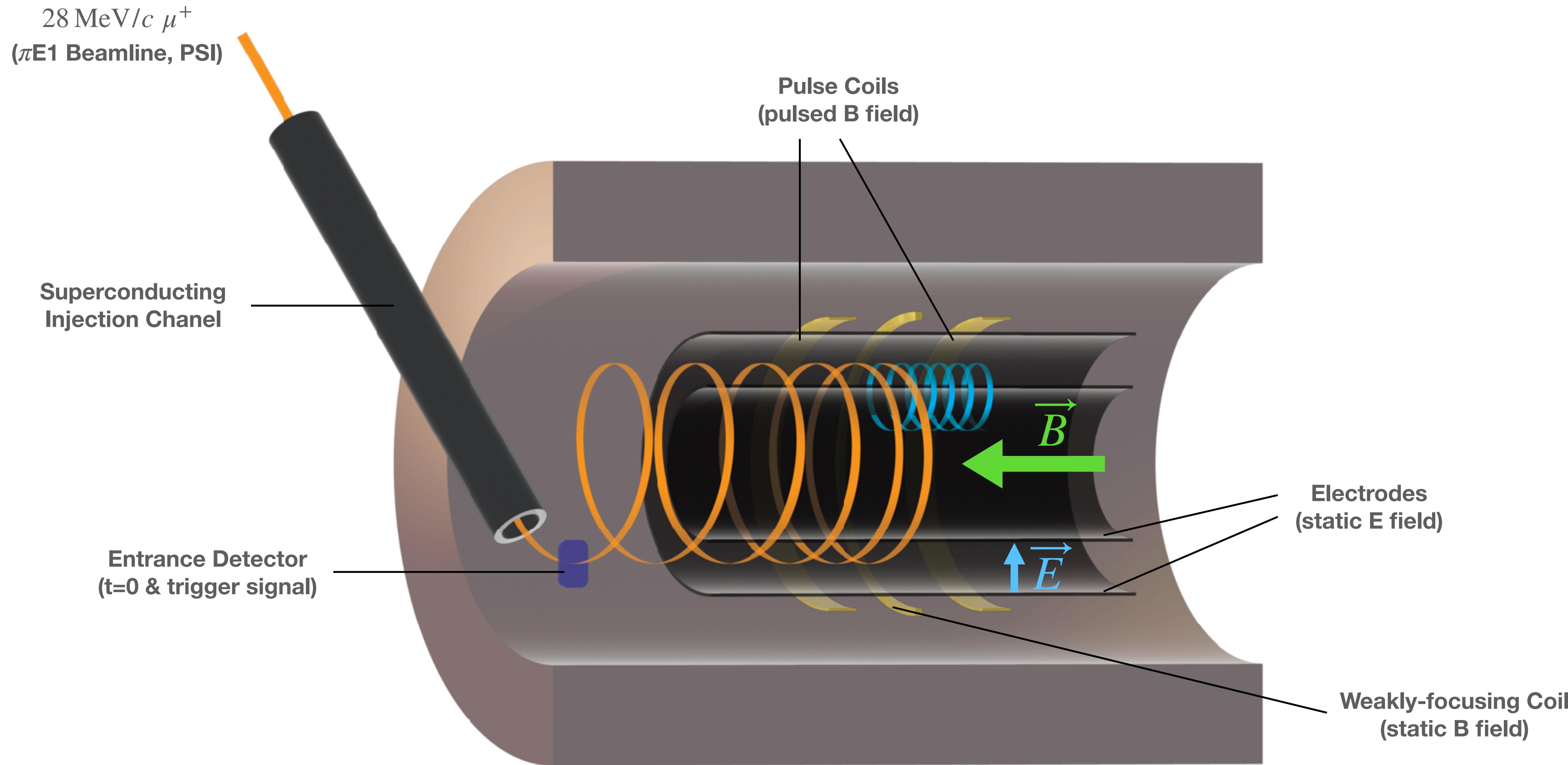
# muEDM Experiment at PSI



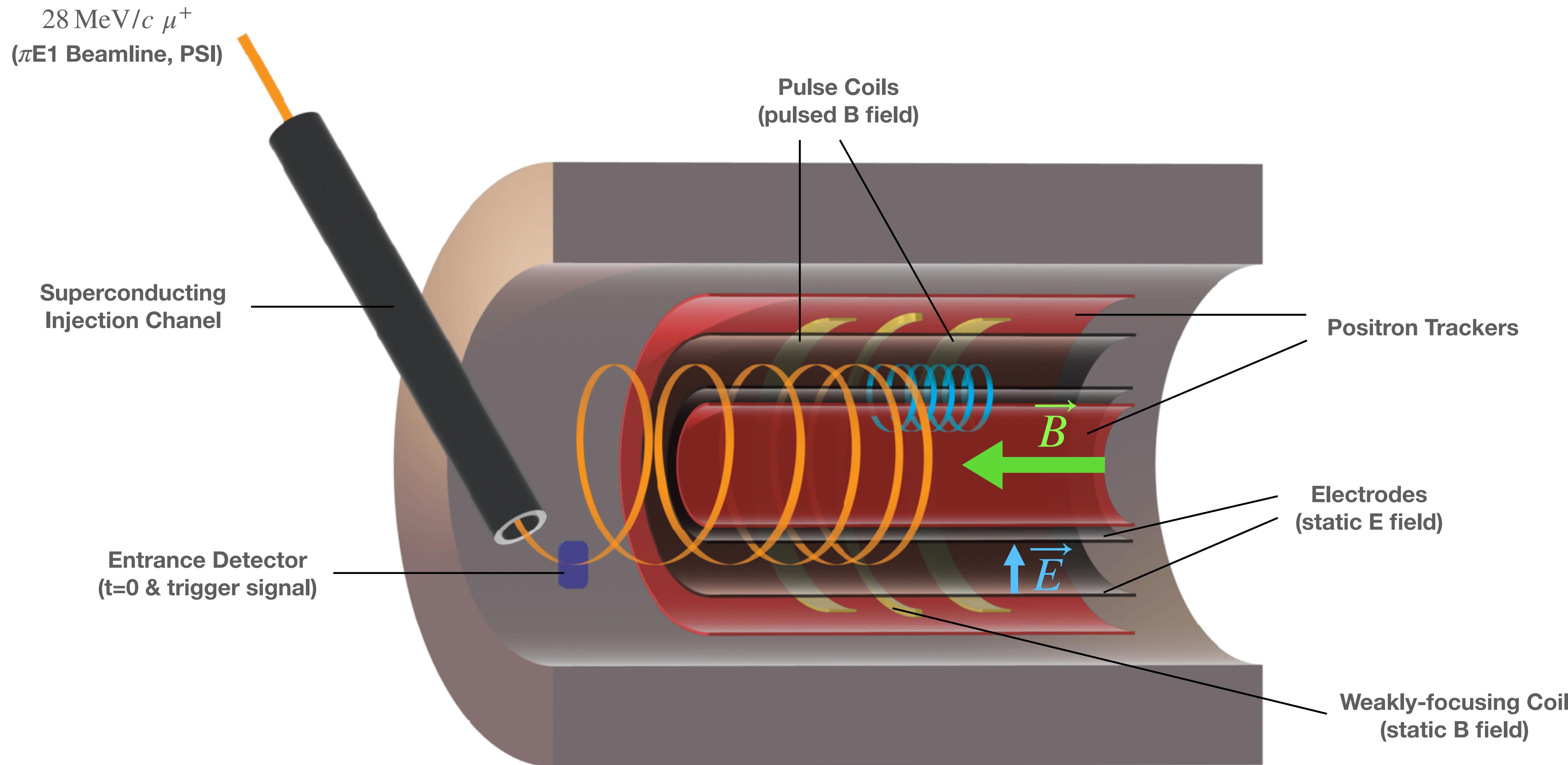
# muEDM Experiment at PSI



# muEDM Experiment at PSI

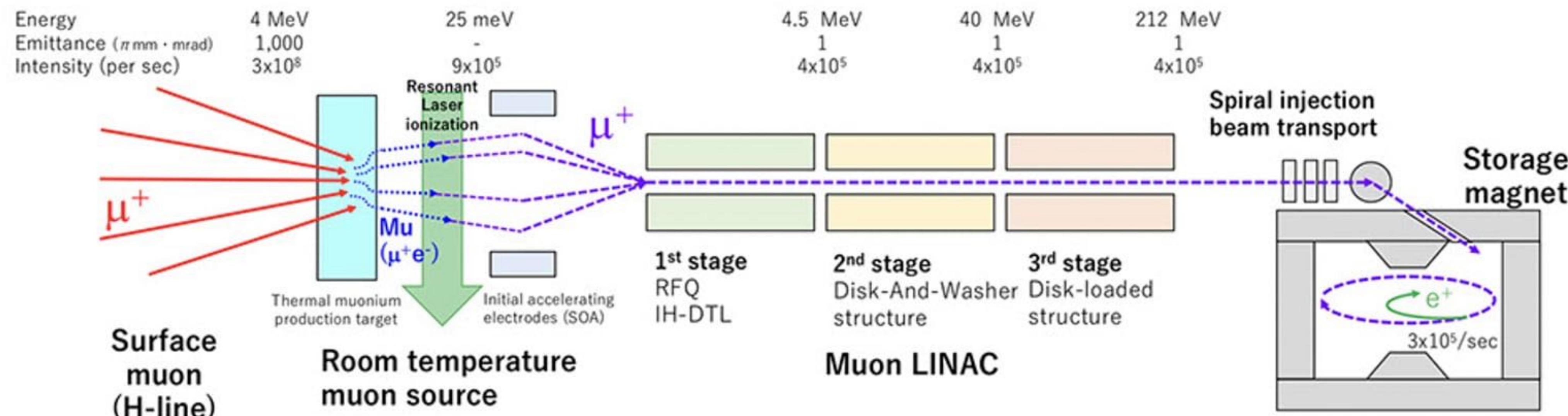


# muEDM Experiment at PSI



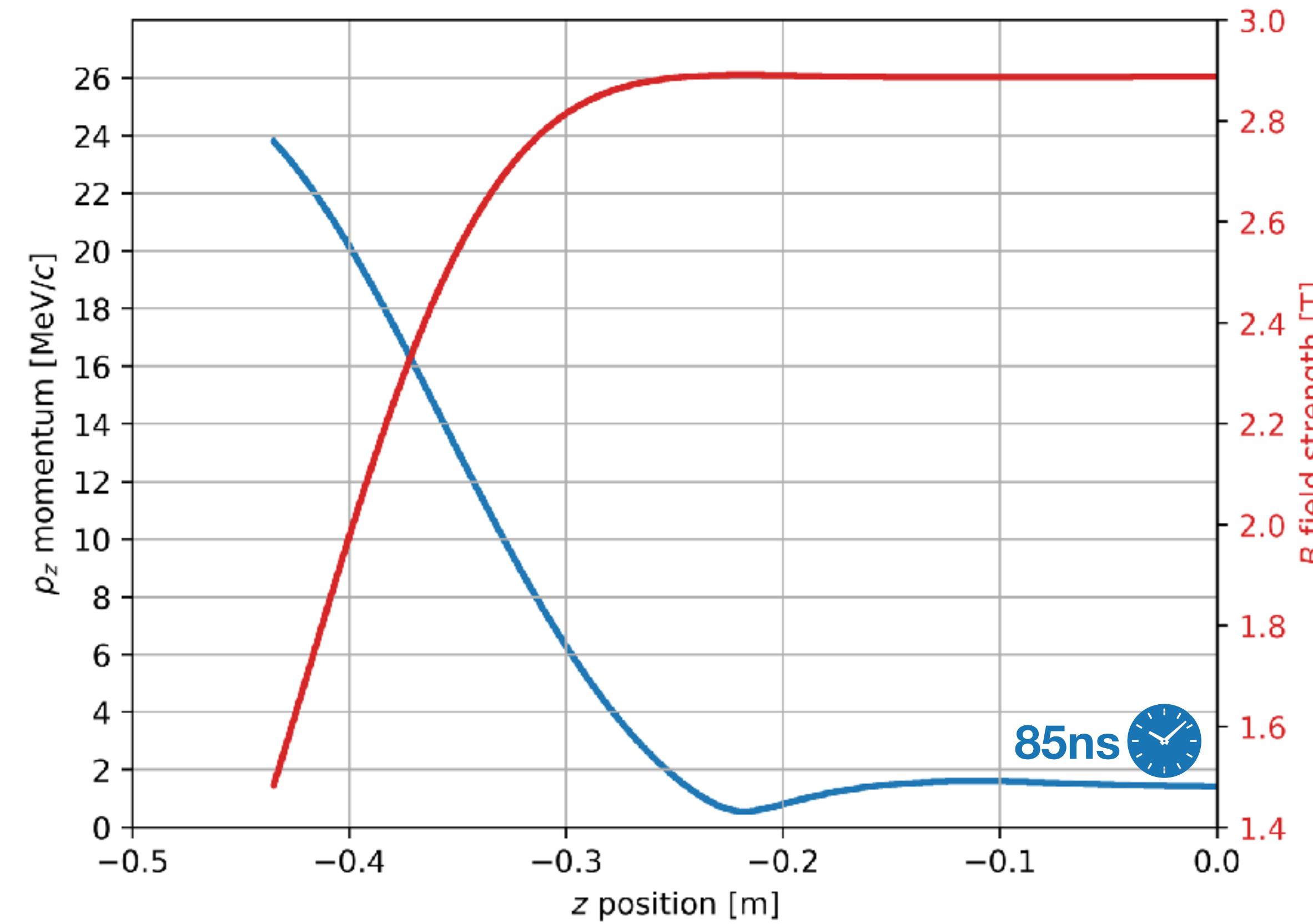
# Muon g-2/EDM at JPARC

- Pulsed (bunch rate 25 Hz) muon beam with momentum  $30 \text{ MeV}/c$ .
- Longitudinal injection (similar to PSI)
- Electric field cancellation means narrow acceptance phase space.
- Pulsed beam demands bunched injection and measurement
- Ultra-cold muon beam developed for phase space compression.

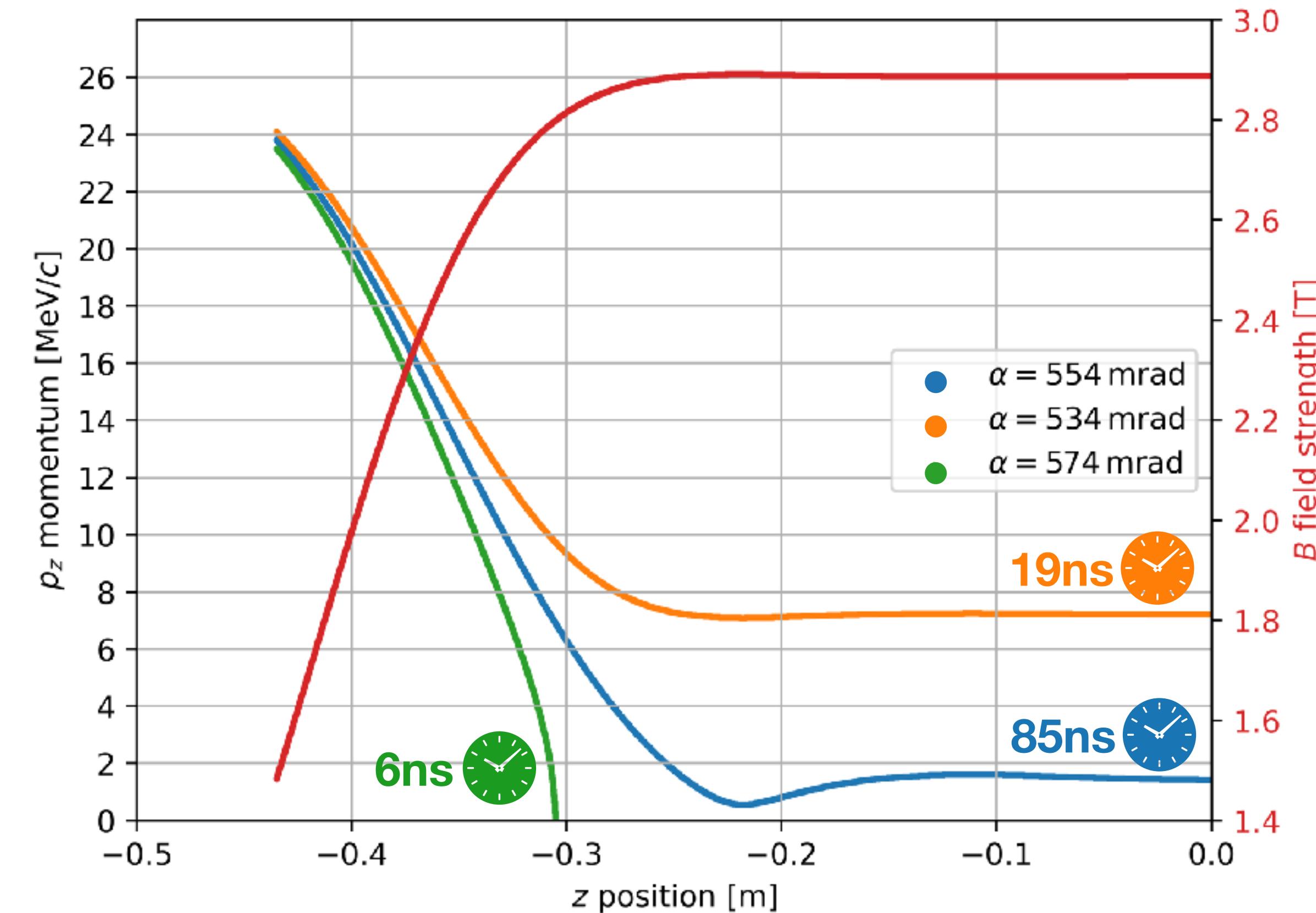


Credit: J-PARC g-2.kek.jp

# Challenges: Injection & Storage



# Challenges: Injection & Storage



# Outlook for muEDM Phase I

