# Exotic Decay

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## Outline

- Introduction and physics motivation of exotic decay search.
- Exotic decay searches in PIENU

- Prospects of exotic decay searches in PIONEER
- Summary.

$$
- \pi^+ \to e^+ \nu_H
$$

$$
- \mu^+ \to e^+ X_H
$$

## Introduction and Physics Motivation

- In the SM,,,
	- neutrinos are massless particles.
	- lepton flavor conservation.

## ➡Due to the observation of neutrino oscillation (PRL 81, 1562 (1998)), neutrinos are massive and lepton flavor is not strictly conserved.

- Other problems in the SM,,,
	- Origin of neutrino masses.
	- Candidates of dark matter.
	- Explanation of baryogenesis.
	- Solution of strong CP problem. and so on...







# Introduction and Physics Motivation

- To solve such problems and extend the SM, the existence of massive or massless weakly interacting neutral particles  $(\nu_{H}, X, \cdots)$  has been suggested. - Heavy neutrinos: JHEP 04 (2011) 011. PRD 97, 075016 (2018). - Axions (or axion like particles): PRD 101, 075002 (2020). *νH*, *X*, ⋯
	-
	-
	- Familons: PRL 49, 1549 (1982).
	- Majorons: PLB 99, 411 (1981). PLB 98, 265 (1982). ■Pion and muon (and other mesons such as kaon) decays are potential

sources of new particles.

## Summary of Exotic Decay Searches in PIENU



 $\mathcal{L}$ *\** Sensitive to new interactions such as  $\nu$  −  $\nu$  interaction (PLB 32, 121 (1970)) and 6-fermion interaction(Phys. Rev. 133, B130 (1964))





$$
(\ell = e, \mu)
$$





# Quick Theory of Massive Neutrinos

- Extension of the SM by adding massive neutrinos.
- . The relation between the weak eigenstates  $\nu_e$  and mass eigenstates  $\nu_i$ JHEP 1104 11 (2011)
- with the mixing parameter  $U_{\ell i}$  is  $U_{\ell i}$

$$
\nu_\ell = \sum_{i=1}^{3+k} U_{\ell i} \nu_i \qquad (\ell = e, \mu, \tau, \chi)
$$

• The contribution from massive neutrino to  $\pi^+ \to e^+ \nu_e$  decay is

$$
R_{ei} = \frac{\Gamma(\pi^+ \to e^+ \nu_i)}{\Gamma(\pi^+ \to e^+ \nu_e)} = |U_{ei}|^2 \rho_e
$$

Neutrino Minimal Standard Model  $(k = 3)$ 

 $\left( 1, \chi_2, \cdots, \chi_k \right)$ 

$$
\begin{array}{l} \rho_e=\frac{(\delta_e+\delta_i-(\delta_e-\delta_i)^2)\sqrt{1+\delta_e^2+\delta_i^2-2(\delta_e+\delta_i+\delta_i)}}{\delta_e(1-\delta_i)^2}\\ \delta_e=\frac{m_e^2}{m_\pi^2},\ \delta_i=\frac{m_{\nu_i}^2}{m_\pi^2} \end{array}
$$

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- *π* Dominant  $\pi^+ \to \mu^+ \to e^+$  ( $\pi$ DAR,  $\pi$ DIF,  $\pi$ DAR  $\mu$ DIF) can be removed using
	- early time requirement: × ∼ 1/100
	- energy deposit in the target ( $T_\mu = 4.1 \text{ MeV}$ ):  $\times \sim 1/1000$
	- tracking information called "kink angle" cut:  $x \sim 1/2$



 $\pi^+ \to e^+ \nu_H$  (and  $\pi^+ \to e^+ \nu_X$ ,  $\pi^+ \to e^+ \nu_e \nu_V$  $\bar{\nu})$ 

- *π* Remaining backgrounds are  $\pi$ DIF,  $\mu$ DIF, and  $\pi^+ \rightarrow e^+ \nu_e$  tail.
	- $\blacksquare$  Peak search with the range of  $E_e = 4 \sim 56$  MeV ( $m_{\nu_i} = 60 \sim 135$  MeV/ $c^2$ ). − Set the 90% CL limits on  $U_{ei}^2$  with  $10^{-7} \sim 10^{-8}$ .
- I'm revisiting higher mass region where photo nuclear bump appears.

Hypothetical signal at  $E_e = 40$  MeV with  $U_{ei}$  $^{2} = 1 \times 10^{-8}$ 



Results of  $\pi^+ \rightarrow e^+ \nu_H$ 

 $= 60 \sim 135 \text{ MeV}/c^2$ 



### **Axions and Family Symmetry Breaking**

Institute for Theoretical Physics, University of California at Santa Barbara, Santa Barbara, California 93106 (Received 20 September 1982)

Possible advantages of replacing the Peccei-Quinn  $U(1)$  quasisymmetry by a group of genuine flavor symmetries are pointed out. Characteristic neutral Nambu-Goldstone bosons will arise, which might be observed in rare K or  $\mu$  decays. The formulation of Lagrangians embodying these ideas is discussed schematically.

• This was introduced to solve strong CP problem, baryon asymmetry, and

$$
- \mu^+ \to e^+ X
$$

$$
- K^+ \to \pi^+ X
$$

Phys. Rev. Lett. 49 549 (1982)PACS numbers: 11.30.Qc, 12.35.Cn, 14.80.Gt • Wilczek proposed neutral NG boson called "familon" will arise in the decays

this might be the hint for dark matter, and so on.

Quick Theory of  $\mu^+ \rightarrow e^+ X_H$ 

### Frank Wilczek



- Similar analysis method, but  $\pi^+ \to \mu^+ \to e^+$  events were selected by
	- a requirement of late decay time  $t > 200$  ns.
- Michel spectrum was fitted to the smooth polynomial function and signal shape produced by MC.

Hypothetical signal at  $m_{X_H} = 90$  MeV/ $c^2$  with  $R = 5 \times 10^{-5}$ 

\n- New limits were set on 
$$
R = \Gamma(\mu^+ \to e^+ X_H)/\Gamma(\mu^+ \to e^+ \nu \bar{\nu})
$$
 with an order of 10<sup>-5</sup>. Hypothetical signal at *m*<sub>*X<sub>H</sub>* = 90 MeV/*c*<sup>2</sup> with *R* = 5 × 10<sup>-5</sup></sub>
\n- Another constraint in very light mass region and the nontrivial signal at *m*<sub>*X<sub>H</sub>* = 90 MeV/*c*<sup>2</sup> with *R* = 5 × 10<sup>-5</sup></sub>
\n- Comological study (PRD 103, 52007 202)
\n

*t* > 200 ns



 $\mu^+ \rightarrow e^+ X_H$ 



## Prospects of Exotic Decay Searches in PIONEER

- To improve the sensitivity for exotic decay searches,
	- larger statistics: ×100
	-
	- better energy resolution: × 2?
- Toy MC study in  $\pi^+ \rightarrow e^+ \nu_H$ 
	- $\times$  100 larger statistics than PIENU
	- $\overline{h}$  No  $\pi$ DIF event and 1%  $\pi$ <sup>+</sup>  $\rightarrow$   $e^{+}\nu_{e}$  tail.
	- the same energy resolution with PIENU

Also for other exotic decays!!! 13 An order of magnitude of improvement is expected!!!



# Additional Decay Modes

- In the PIENU, the detector acceptance was limited (~20%).
- . All neutral particles *X* analyzed in PIENU were assumed to be long lifetime.
	- If X is very short lifetime and decays  $X \to \gamma \gamma$ , one  $\gamma$  is missed. *X* is very short lifetime and decays  $X \to \gamma \gamma$ , one  $\gamma$
- For PIONEER, 2γ could be detected due to larger acceptance.







# Summary

- Exotic decay searches may reveal the problem of the SM.
- PIENU and PIONEER have capability to search for the decays involving new particles.
- PIENU has achieved many searched with high sensitivity.
- PIONEER is expected to improve the sensitivity by an order of magnitude. • New decay modes can also be search for.
- 



# Back Up





## • In the SM, the BRs are  $- BR(π<sup>+</sup> → μ<sup>+</sup>ν<sub>μ</sub>νν<sub>ν</sub>ν) ≈ 10<sup>-20</sup>$

 $-BR(\pi^+ \to e^+ \nu_e \nu \bar{\nu}) \approx 10^{-18}$ 

 $\pi^+ \rightarrow \ell^+ \nu_\ell \nu \bar{\nu}$  $\overline{\nu}$ 



### • Similar analysis with  $\pi^+ \rightarrow e^+\nu_H$ , but cuts were optimized.



FIG. 5. Top: The  $E_e$  spectra of  $\pi^+ \to e^+ \nu_e$  decay after  $\pi^+ \to \mu^+ \to e^+$  suppression cuts from data taken before (a) and after (b) 1 November 2010. The black crosses with the statistical uncertainties show the data. Background components illustrated by the dashed and dotted green line, dashed gray line, dotted blue line, and solid red line represent  $\pi^+ \to \mu^+ \to e^+$  decays, low-energy  $\pi^+ \to e^+ \nu_e$  tail,  $\mu$ DIF events, and the sum of those three components, respectively (see text). Bottom: The residual plots (rebinned for visual clarity) shown by the black circles with the statistical error bars and hypothetical neutrino-neutrino interaction  $(I_{\nu\bar{\nu}})$  signals (solid red curves) with the branching ratio  $B_{l_{\nu\bar{\nu}}}^{\pi e_{3\nu}} = 1.0 \times 10^{-6}$  for data taken before (a) and after (b) 2 November 2010. The dashed horizontal red lines represent the residual of 0.

 $\pi^+ \rightarrow e^+ \nu_e \nu \bar{\nu}$  $\overline{\nu}$ 

## • The same analysis method with  $\pi^+ \to \mu^+ \nu_H$ .

FIG. 4. (a) The  $T_{\mu}$  spectra of  $\pi^{+} \rightarrow \mu^{+} \nu_{\mu}$  decay. The black crosses with the statistical uncertainties show the data. The dotted green line, dashed blue line, and solid red line represent the Gaussian distribution at 4.1 MeV,  $\pi^+ \rightarrow \mu^+ \nu_\mu \gamma$  decay, and the sum of those two functions, respectively. (b) Residual plot shown by the black circles with the statistical error bars for the signal region  $T_{\mu} = 1.3-3.4$  MeV. The solid red curve represents the hypothetical neutrino-neutrino interaction  $(I_{\nu\bar{\nu}})$  signal with the branching ratio  $B_{L_x}^{\pi\mu 3\nu} = 6.0 \times 10^{-5}$ . The dashed horizontal red line indicates the residual of 0.

 $\pi^+$   $\rightarrow$   $\mu^+ \nu_{\mu} \nu \bar{\nu}$ 



• Similar analysis with heavy neutrino searches, but the signal spectra are continuous.



 $\pi^+ \rightarrow \ell^+ \nu X$ 



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## • The same analysis method with  $\pi^+ \to \ell^+ \nu_\ell \nu \bar{\nu}$ .



 $\pi^+ \rightarrow \ell^+ \nu X$ 



## Constraint on the Mojoron Model:  $\pi^+ \rightarrow e^+ \nu X$

- **Majoron model can also be constrained using**  $R^{\pi}$ **.**
- . The predicted  $R^{\pi}$  including the massless Majoron  $X_0$  and a light neutral  $Higgs H'(\leq 1 \text{ MeV}/c^2)$  can be written as (PRD 25, 907 (1982)) .  $\Gamma(\pi \to e L^0)/\Gamma(\pi \to \mu L^0)$  $= 1 + 157.5g<sup>2</sup>$ 
	- Final state  $\nu, \nu X_0, \nu H'$  g : Majoron-neutrino coupling constant  $\Gamma(\pi \to e \nu_e)/\Gamma(\pi \to \mu \nu_\mu)$  $L^0$ : Final state  $\nu, \nu X_0, \nu H'$  *g*:
- . Averaged experimental branching ratio:  $R_{\rm exp}^{\pi} = (1.2327 \pm 0.0023) \times 10^{-4}$  $R^{\pi}_{\rm exp}$  $<$  1.0014  $\Rightarrow$   $g^2$  < 9 × 10<sup>-6</sup>
- If 0.01% precision of  $R^{\pi}$  is assumed, then  $g^2 \approx 1 \times 10^{-6}$ . *Rπ* SM  $R^\pi$  is assumed, then  $g^2 \approx 1 \times 10^{-6}$

