

REBECA GONZALEZ SUAREZ - UPPSALA UNIVERSITY

PSI PARTICLE PHYSICS SUMMER SCHOOL – FROM LOW TO HIGH: PARTICLE PHYSICS AT THE FRONTIER

THE STANDARD MODEL AND BEYOND

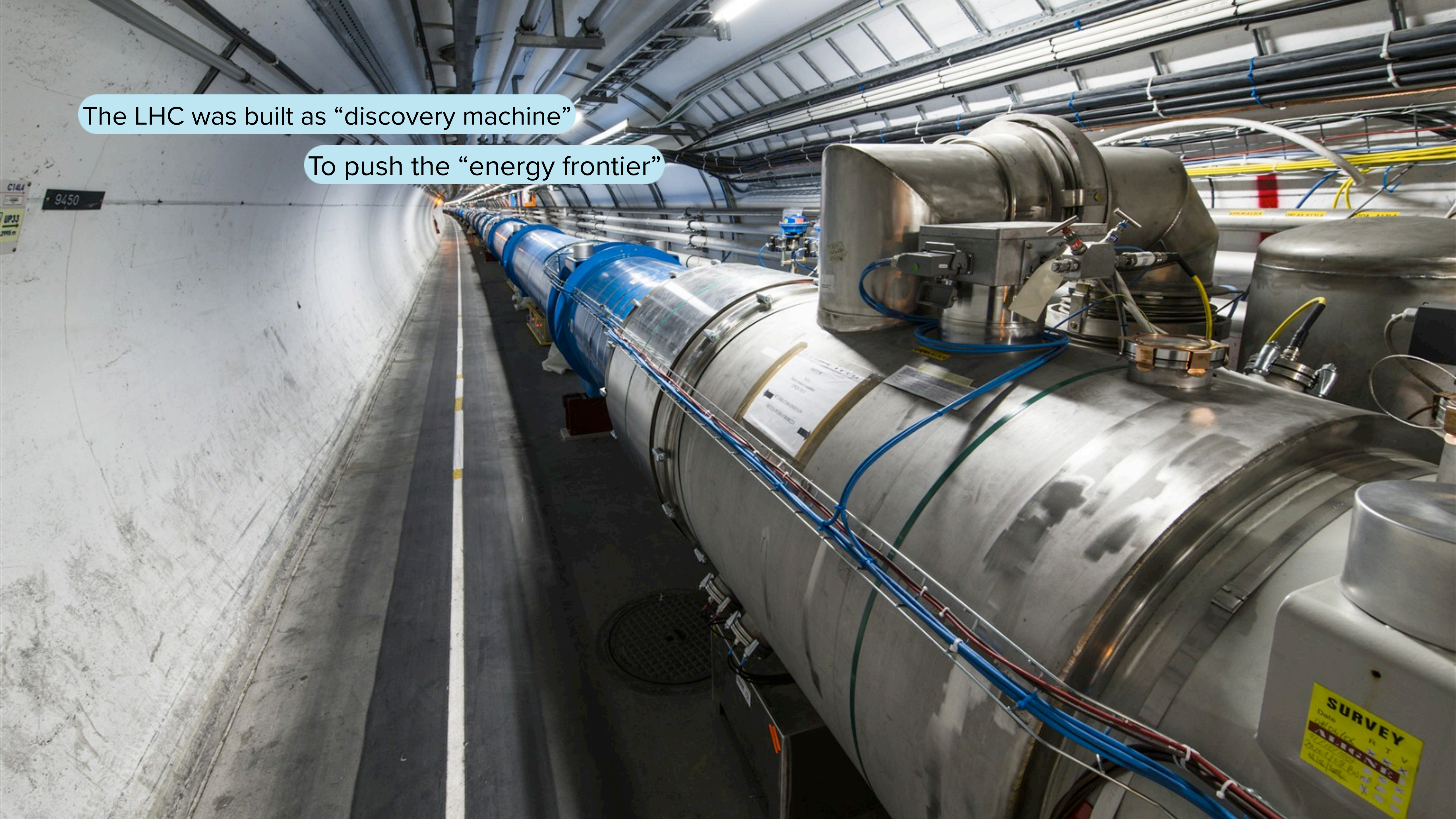
Part II

THE STANDARD MODEL OF PARTICLE PHYSICS

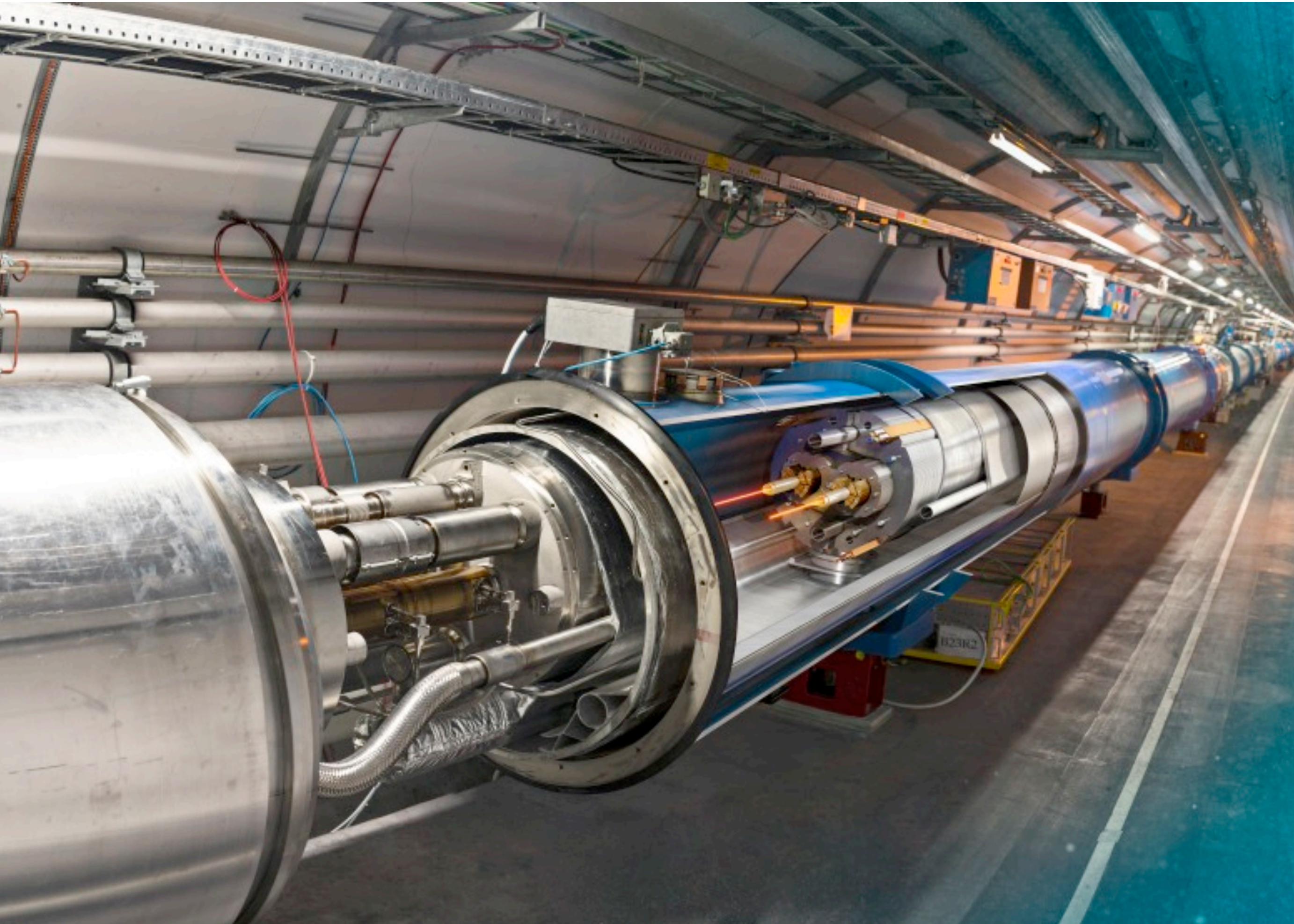
- It is both true that:
 - The SM works extremely well → it is highly predictive and robust, and deviations are few and far in between
 - Last lecture
 - The SM is broken → we know it is hiding something, and it is hiding it very well
 - This one

The LHC was built as “discovery machine”

To push the “energy frontier”



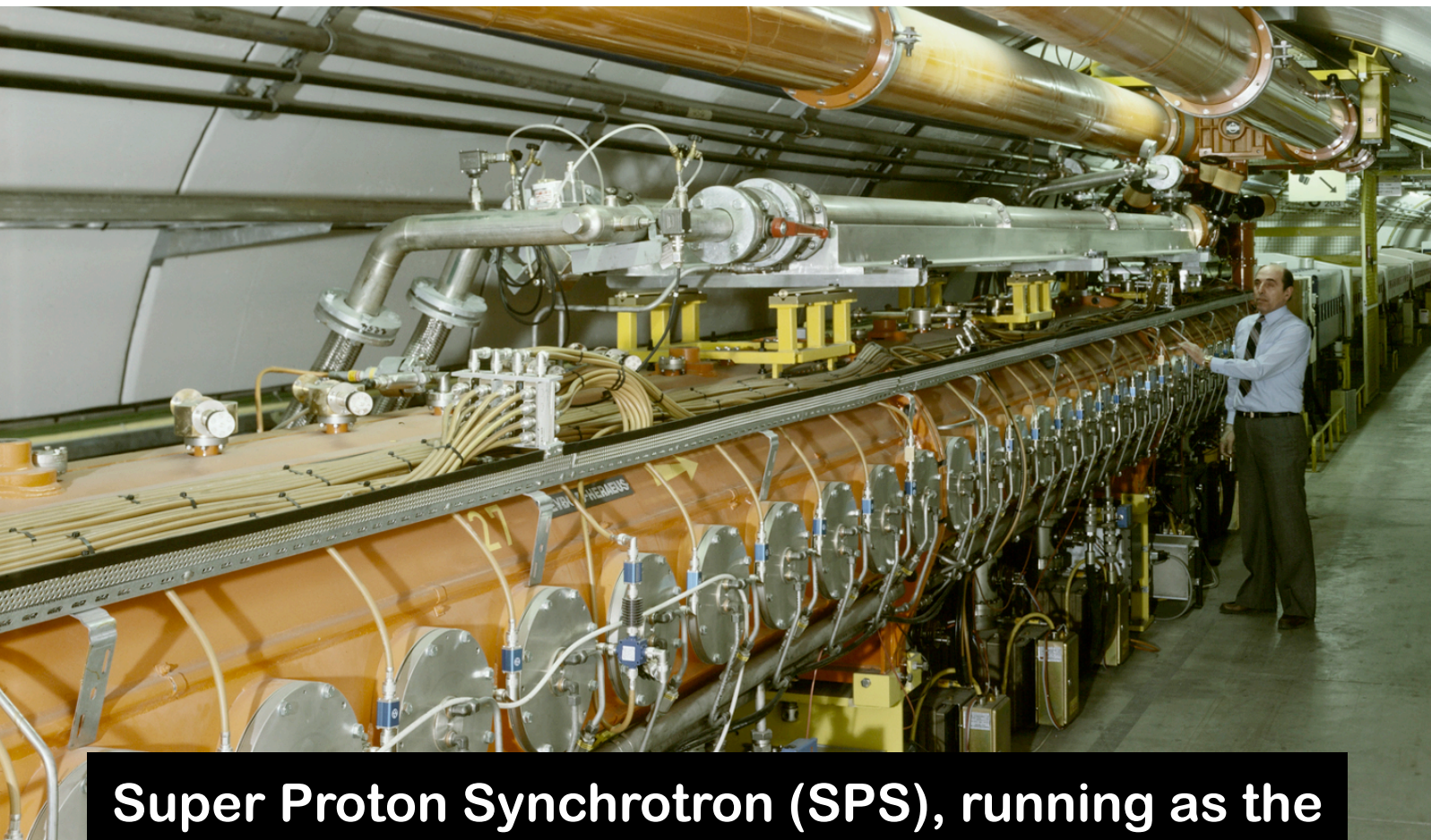
WHEN INCREASING THE ENERGY



- When particles accelerated are close to the speed of light and collide, the collision energy becomes available mass to produce new particles ($E=mc^2$)
- Heavier particles than those colliding can be produced
- When colliding protons (composite particles) we increase the chances to see something new

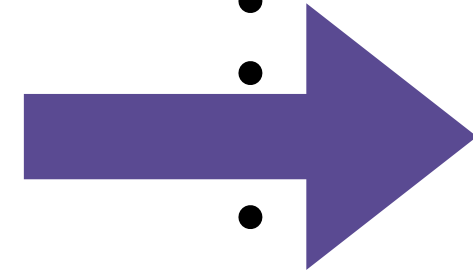
1983

W and Z bosons
80 and 90 GeV



Super Proton Synchrotron (SPS), running as the "Proton–Antiproton Collider" at CERN

540 GeV



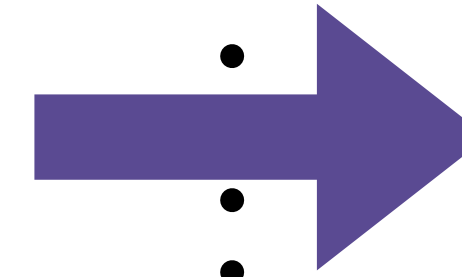
1995

Top quark
173 GeV



Tevatron at Fermilab

1 TeV



We did find a heavy particle!

2012

Higgs boson
125 GeV



Large Hadron collider at CERN

13.6 TeV

THE LHC'S ACHIEVEMENTS

Not every collider finds
new particles
Every collider gets us
one step ahead

- After 14 years we have:
 - The Higgs boson
 - Found after decades of searches in multiple colliders
 - First observations of tons of never-seen-before SM processes
 - mostly multiple production of bosons and/or top quarks
 - About 3000 scientific papers
 - Approximately half of those being precision measurements
 - Un unparalleled battery of searches for BSM phenomena
 - Because the SM is great but...

WE STILL NEED TO GO

BEYOND THE STANDARD MODEL

LEON LEDERMAN'S LAUNDRY LIST

Nature Insight: The Large Hadron Collider
Vol. 448, No. 7151 pp 269-312 (2007)



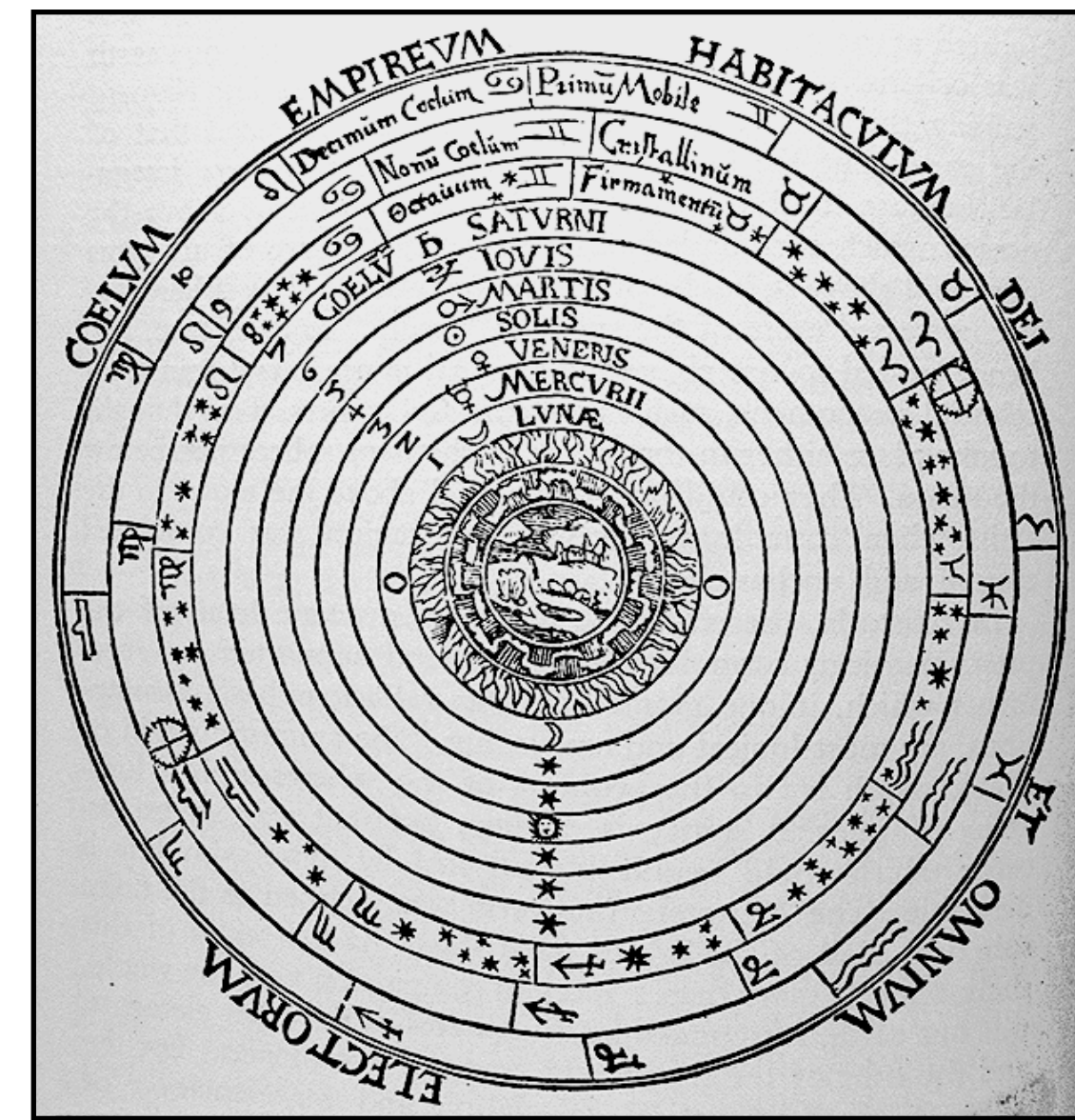
- *“To me, these three factors —the Higgs particles, supersymmetric particles and new dimensions— are the discoveries most likely to emerge from the first five or so years of LHC operations. But there is a long, more speculative laundry list of objects that might be illuminated by the powerful beams of the LHC.”*

Expectations were high for the LHC
That is the “guaranteed discovery” people talk about

- ☑ Higgs boson
- ☐ SUSY
- ☐ Extra dimensions
- ☐ The origin of dark matter
- ☐ The origin of dark energy
- ☐ Compositeness
- ☐ Technicolor (new strong force)
- ☐ WW scattering
- ☐ Additional Higgs bosons
- ☐ Right-handed neutrinos
- ☐ Mini black holes

THE MANY PROBLEMS OF THE SM

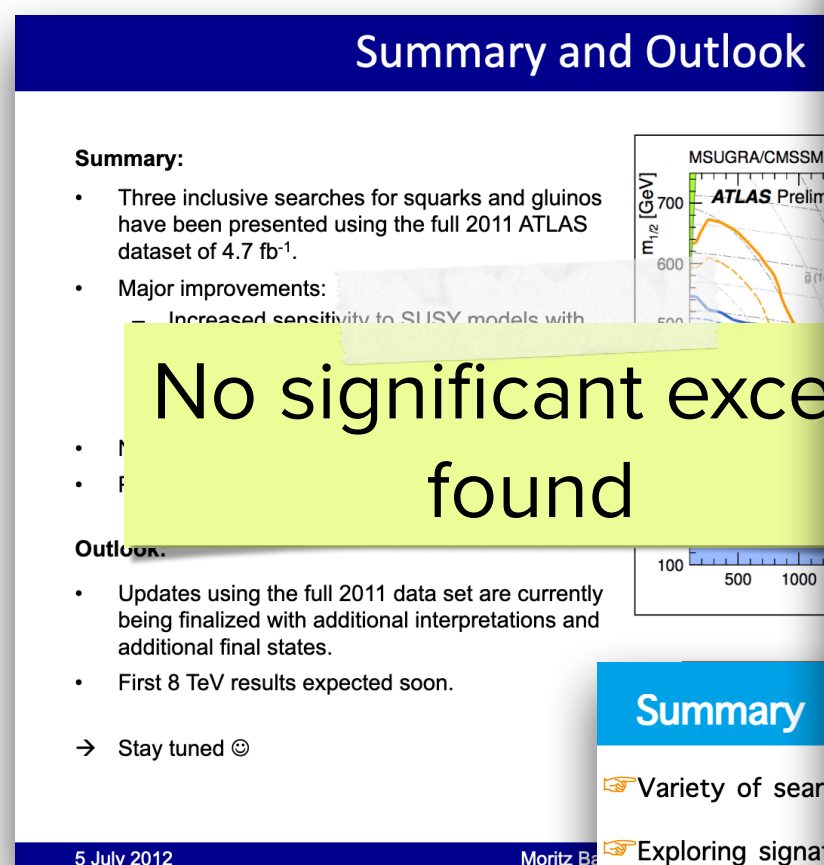
The SM: the Aristotelian Ptolemaic system of our times



- The inexplicable **neutrino masses**
- The fact that **dark matter** exists but we have no idea what it is
- The fact that **dark energy** exists but we have even less idea
- The fact that we are made of **matter** and there is very little **antimatter** around
- We don't know **why** we have **3 families** of particles
- We don't understand their **masses**
- There are a few parameters that **cannot be predicted**, we have to measure them and put them by hand
- The fact that we cannot connect **gravity** from the macroscopic to the fundamental
- ...

FROM THE VERY BEGINNING

■ SUSY was the



Summary and Outlook

Summary:

- Three inclusive searches for squarks and gluinos have been presented using the full 2011 ATLAS dataset of 4.7 fb^{-1} .
- Major improvements:
 - Increased sensitivity to SUSY models with

No significant excess found

Outlook:

- Updates using the full 2011 data set are currently being finalized with additional interpretations and additional final states.
- First 8 TeV results expected soon.

→ Stay tuned ☺

5 July 2012

Just two examples of many

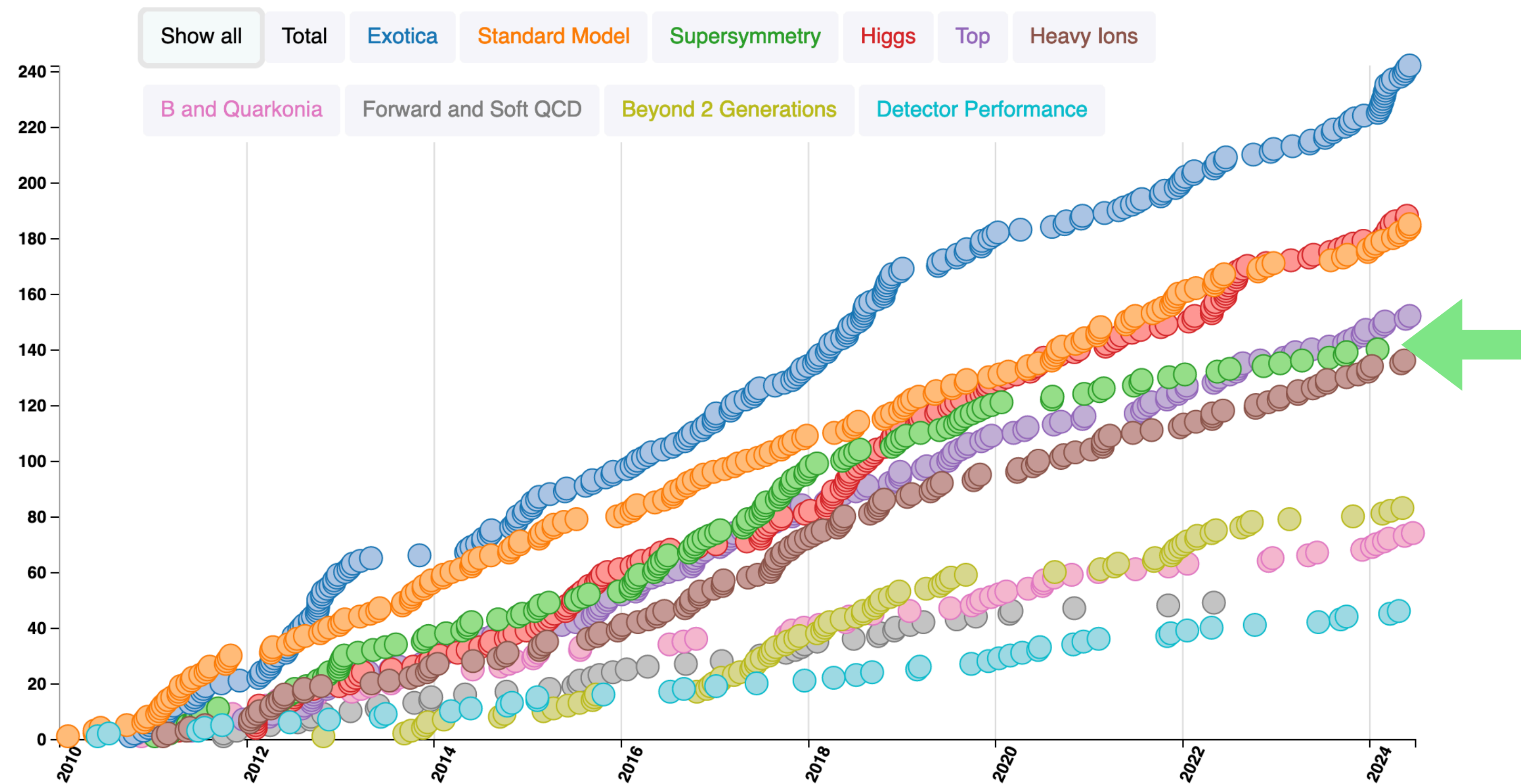
IS SUSY DEAD?
NO, BUT SOMEWHAT HUMBLLED...

ad?



SUSY IS STILL NOT FULLY DEAD

- It will always be a fun ground to test interesting things
- CMS and ATLAS have about 300 SUSY papers



Model	Signature	$\int \mathcal{L} dt$ [fb ⁻¹]	Mass limit	Reference						
Inclusive Searches	$\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$	0 e, μ mono-jet	2-6 jets 1-3 jets	E_T^{miss} E_T^{miss}	140 140	\tilde{q} [1x, 8x Degen.] \tilde{q} [8x Degen.]	1.0 0.9	1.85	$m(\tilde{\chi}_1^0) < 400$ GeV $m(\tilde{q}) - m(\tilde{\chi}_1^0) = 5$ GeV	210.14293 2102.10874
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$	0 e, μ	2-6 jets	E_T^{miss}	140	\tilde{g} \tilde{g}	Forbidden 1.15-1.95	2.3	$m(\tilde{\chi}_1^0) = 0$ GeV $m(\tilde{\chi}_1^0) = 1000$ GeV	210.14293 210.14293
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}W\tilde{\chi}_1^0$	1 e, μ	2-6 jets	E_T^{miss}	140	\tilde{g}		2.2	$m(\tilde{\chi}_1^0) < 600$ GeV	2101.01629
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}(\ell\ell)\tilde{\chi}_1^0$	$ee, \mu\mu$	2 jets	E_T^{miss}	140	\tilde{g}		2.2	$m(\tilde{\chi}_1^0) < 700$ GeV	2204.13072
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow qqWZ\tilde{\chi}_1^0$	0 e, μ	7-11 jets	E_T^{miss}	140	\tilde{g}		1.97	$m(\tilde{\chi}_1^0) < 600$ GeV	2008.06032
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow qqWZ\tilde{\chi}_1^0$	SS e, μ	6 jets	E_T^{miss}	140	\tilde{g}		1.15	$m(\tilde{g}) - m(\tilde{\chi}_1^0) = 200$ GeV	2307.01094
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\tilde{\chi}_1^0$	0-1 e, μ	3 b	E_T^{miss}	140	\tilde{g}		2.45	$m(\tilde{\chi}_1^0) < 500$ GeV	2211.08028
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow t\tilde{\chi}_1^0$	SS e, μ	6 jets	E_T^{miss}	140	\tilde{g}		1.25	$m(\tilde{g}) - m(\tilde{\chi}_1^0) = 300$ GeV	1909.08457
3 rd gen. squarks direct production	$\tilde{b}_1\tilde{b}_1$	0 e, μ	2 b	E_T^{miss}	140	\tilde{b}_1 \tilde{b}_1		1.255 0.68	$m(\tilde{\chi}_1^0) < 400$ GeV 10 GeV $< \Delta m(\tilde{b}_1, \tilde{\chi}_1^0) < 20$ GeV	2101.12527 2101.12527
	$\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_2^0 \rightarrow bh\tilde{\chi}_1^0$	0 e, μ 2 τ	6 b 2 b	E_T^{miss} E_T^{miss}	140 140	\tilde{b}_1 \tilde{b}_1	Forbidden 0.13-0.85	0.23-1.35	$\Delta m(\tilde{\chi}_2^0, \tilde{\chi}_1^0) = 130$ GeV, $m(\tilde{\chi}_1^0) = 100$ GeV $\Delta m(\tilde{\chi}_2^0, \tilde{\chi}_1^0) = 130$ GeV, $m(\tilde{\chi}_1^0) = 0$ GeV	1908.03122 2103.08189
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$	0-1 e, μ	≥ 1 jet	E_T^{miss}	140	\tilde{t}_1		1.25	$m(\tilde{\chi}_1^0) = 1$ GeV	2004.14060, 2012.03799
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$	1 e, μ	3 jets/1 b	E_T^{miss}	140	\tilde{t}_1	Forbidden	1.05	$m(\tilde{\chi}_1^0) = 500$ GeV	2012.03799, ATLAS-CONF-2023-043
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow \tilde{\tau}_1 b\nu, \tilde{\tau}_1 \rightarrow \tau\tilde{G}$	1-2 τ	2 jets/1 b	E_T^{miss}	140	\tilde{t}_1	Forbidden	1.4	$m(\tilde{\tau}_1) = 800$ GeV	2108.07665
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0 / \tilde{c}\tilde{c}, \tilde{c} \rightarrow c\tilde{\chi}_1^0$	0 e, μ 0 e, μ	2 c mono-jet	E_T^{miss} E_T^{miss}	36.1 140	\tilde{c} \tilde{t}_1		0.85 0.55	$m(\tilde{\chi}_1^0) = 0$ GeV $m(\tilde{t}_1, \tilde{c}) - m(\tilde{\chi}_1^0) = 5$ GeV	1805.01649 2102.10874
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow t\tilde{\chi}_2^0, \tilde{\chi}_2^0 \rightarrow Z/h\tilde{\chi}_1^0$	1-2 e, μ	1-4 b	E_T^{miss}	140	\tilde{t}_1		0.067-1.18	$m(\tilde{\chi}_2^0) = 500$ GeV	2006.05880
	$\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$	3 e, μ	1 b	E_T^{miss}	140	\tilde{t}_2	Forbidden	0.86	$m(\tilde{\chi}_1^0) = 360$ GeV, $m(\tilde{t}_1) - m(\tilde{\chi}_1^0) = 40$ GeV	2006.05880
EW direct	$\tilde{\chi}_1^+\tilde{\chi}_2^0$ via WZ	Multiple ℓ /jets $ee, \mu\mu$	≥ 1 jet	E_T^{miss} E_T^{miss}	140 140	$\tilde{\chi}_1^+/\tilde{\chi}_2^0$ $\tilde{\chi}_1^+/\tilde{\chi}_2^0$		0.96 0.205	$m(\tilde{\chi}_1^0) = 0$, wino-bino $m(\tilde{\chi}_1^+) - m(\tilde{\chi}_1^0) = 5$ GeV, wino-bino	2106.01676, 2108.07586 1911.12606
	$\tilde{\chi}_1^+\tilde{\chi}_1^+$ via WW	2 e, μ		E_T^{miss}	140	$\tilde{\chi}_1^+$		0.42	$m(\tilde{\chi}_1^0) = 0$, wino-bino	1908.08215
	$\tilde{\chi}_1^+\tilde{\chi}_2^0$ via Wh	Multiple ℓ /jets		E_T^{miss}	140	$\tilde{\chi}_1^+/\tilde{\chi}_2^0$	Forbidden	1.06	$m(\tilde{\chi}_1^0) = 70$ GeV, wino-bino	2004.10894, 2108.07586
	$\tilde{\chi}_1^+\tilde{\chi}_1^+$ via $\tilde{\ell}_L/\tilde{\nu}$	2 e, μ		E_T^{miss}	140	$\tilde{\chi}_1^+$		1.0	$m(\tilde{\ell}, \tilde{\nu}) = 0.5(m(\tilde{\chi}_1^+) + m(\tilde{\chi}_1^0))$	1908.08215
	$\tilde{\tau}\tilde{\tau}, \tilde{\tau} \rightarrow \tau\tilde{\chi}_1^0$	2 τ		E_T^{miss}	140	$\tilde{\tau}$ [$\tilde{\tau}_R, \tilde{\tau}_{R,L}$]		0.34 0.48	$m(\tilde{\chi}_1^0) = 0$	ATLAS-CONF-2023-029
	$\tilde{\ell}_{L,R}\tilde{\ell}_{L,R}, \tilde{\ell} \rightarrow \ell\tilde{\chi}_1^0$	2 e, μ $ee, \mu\mu$	0 jets ≥ 1 jet	E_T^{miss} E_T^{miss}	140 140	$\tilde{\ell}$ $\tilde{\ell}$		0.7 0.26	$m(\tilde{\chi}_1^0) = 0$ $m(\tilde{\ell}) - m(\tilde{\chi}_1^0) = 10$ GeV	1908.08215 1911.12606
	$\tilde{H}\tilde{H}, \tilde{H} \rightarrow h\tilde{G}/Z\tilde{G}$	0 e, μ 4 e, μ 0 e, μ 2 e, μ	≥ 3 b 0 jets ≥ 2 large jets ≥ 2 jets	E_T^{miss} E_T^{miss} E_T^{miss} E_T^{miss}	140 140 140 140	\tilde{H} \tilde{H} \tilde{H} \tilde{H}		0.94 0.55 0.45-0.93 0.77	$\text{BR}(\tilde{\chi}_1^0 \rightarrow h\tilde{G}) = 1$ $\text{BR}(\tilde{\chi}_1^0 \rightarrow Z\tilde{G}) = 1$ $\text{BR}(\tilde{\chi}_1^0 \rightarrow Z\tilde{G}) = 1$ $\text{BR}(\tilde{\chi}_1^0 \rightarrow Z\tilde{G}) = \text{BR}(\tilde{\chi}_1^0 \rightarrow h\tilde{G}) = 0.5$	To appear 2103.11684 2108.07586 2204.13072
	Long-lived particles	Direct $\tilde{\chi}_1^+\tilde{\chi}_1^-$ prod., long-lived $\tilde{\chi}_1^\pm$	Disapp. trk	1 jet	E_T^{miss}	140	$\tilde{\chi}_1^\pm$ $\tilde{\chi}_1^\pm$		0.66 0.21	Pure Wino Pure higgsino
Stable \tilde{g} R-hadron		pixel dE/dx		E_T^{miss}	140	\tilde{g}		2.05		2205.06013
Metastable \tilde{g} R-hadron, $\tilde{g} \rightarrow qq\tilde{\chi}_1^0$		pixel dE/dx		E_T^{miss}	140	\tilde{g} [$\tau(\tilde{g}) = 10$ ns]		2.2	$m(\tilde{\chi}_1^0) = 100$ GeV	2205.06013
$\tilde{\ell}\tilde{\ell}, \tilde{\ell} \rightarrow \ell\tilde{G}$		Displ. lep		E_T^{miss}	140	$\tilde{\ell}, \tilde{\mu}$ $\tilde{\tau}$		0.7 0.34 0.36	$\tau(\tilde{\ell}) = 0.1$ ns $\tau(\tilde{\ell}) = 0.1$ ns $\tau(\tilde{\ell}) = 10$ ns	2011.07812 2011.07812 2205.06013
RPV	$\tilde{\chi}_1^+\tilde{\chi}_1^+/\tilde{\chi}_1^0, \tilde{\chi}_1^+ \rightarrow Z\ell \rightarrow \ell\ell\ell$	3 e, μ		E_T^{miss}	140	$\tilde{\chi}_1^+/\tilde{\chi}_1^0$ [$\text{BR}(Z\tau)=1, \text{BR}(Ze)=1$]		0.625 1.05	Pure Wino	2011.10543
	$\tilde{\chi}_1^+\tilde{\chi}_1^+/\tilde{\chi}_2^0 \rightarrow WW/Z\ell\ell\nu\nu$	4 e, μ	0 jets	E_T^{miss}	140	$\tilde{\chi}_1^+/\tilde{\chi}_2^0$ [$\lambda_{33} \neq 0, \lambda_{12k} \neq 0$]		0.95 1.55	$m(\tilde{\chi}_1^0) = 200$ GeV	2103.11684
	$\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow qqq$		≥ 8 jets	E_T^{miss}	140	\tilde{g} [$m(\tilde{\chi}_1^0) = 50$ GeV, 1250 GeV]		1.6 2.25	Large λ'_{112}	To appear
	$\tilde{u}, \tilde{t} \rightarrow t\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow tbs$		Multiple	E_T^{miss}	36.1	\tilde{t} [$\lambda'_{323} = 2e-4, 1e-2$]		0.55 1.05	$m(\tilde{\chi}_1^0) = 200$ GeV, bino-like	ATLAS-CONF-2018-003
	$\tilde{u}, \tilde{t} \rightarrow b\tilde{\chi}_1^+, \tilde{\chi}_1^+ \rightarrow bbs$		$\geq 4b$	E_T^{miss}	140	\tilde{t}	Forbidden	0.95	$m(\tilde{\chi}_1^+) = 500$ GeV	2010.01015
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow bs$		2 jets + 2 b	E_T^{miss}	36.7	\tilde{t}_1 [qq, bs]		0.42 0.61		1710.07171
	$\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow q\ell$	2 e, μ 1 μ	2 b DV	E_T^{miss}	36.1 136	\tilde{t}_1		0.4-1.45 1.0 1.6	$\text{BR}(\tilde{t}_1 \rightarrow b\ell/b\mu) > 20\%$ $\text{BR}(\tilde{t}_1 \rightarrow q\mu) = 100\%, \cos\theta_{12} = 1$	1710.05544 2003.11956
	$\tilde{\chi}_1^+/\tilde{\chi}_2^0/\tilde{\chi}_1^0, \tilde{\chi}_{1,2}^0 \rightarrow tbs, \tilde{\chi}_1^+ \rightarrow bbs$	1-2 e, μ	≥ 6 jets	E_T^{miss}	140	$\tilde{\chi}_1^0$		0.2-0.32	Pure higgsino	2106.09609

*Only a selection of the available mass limits on new states or phenomena is shown. Many of the limits are based on simplified models, c.f. refs. for the assumptions made.

10⁻¹ 1 Mass scale [TeV]

THERE IS SO MUCH MORE THAN SUSY

Resonance searches
“bump hunt”

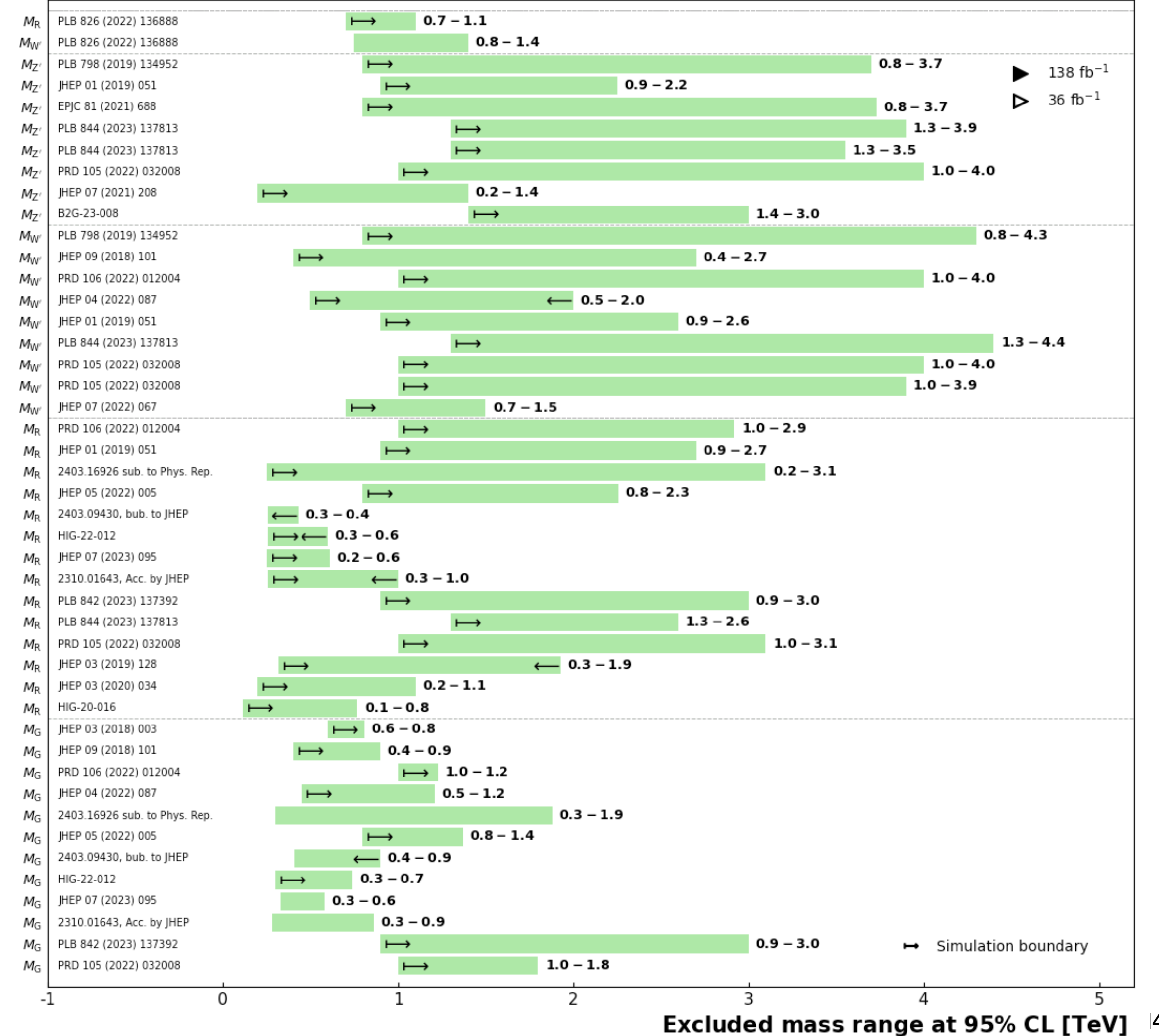
- VV/VH/HH/HV resonances**
- HST**
 - ▶ $R \rightarrow q\bar{q}\gamma \rightarrow W\gamma$ ($g_m = 0.1, \Lambda = 4M_X$)
 - ▶ $W' \rightarrow q\bar{q}\gamma \rightarrow W\gamma$ ($g_m = 0.1, \Lambda = 4M_X$)
 - ▷ Z' (2016 combination)
 - ▷ $Z' \rightarrow ZH \rightarrow q\bar{q}\tau\bar{\tau}$
 - Z', HVT B**
 - ▶ $Z' \rightarrow ZH \rightarrow (\ell\ell, \nu\nu)b\bar{b}$
 - ▶ $Z' \rightarrow ZH \rightarrow q\bar{q}q\bar{q}$
 - ▶ $Z' \rightarrow WW \rightarrow q\bar{q}q\bar{q}$
 - ▶ $Z' \rightarrow WW \rightarrow \ell\nu q\bar{q}$
 - ▶ $Z' \rightarrow \ell\ell$
 - ▶ $Z' \rightarrow ZH \rightarrow \ell\nu\nu, cc/4q$
 - W', HVT B**
 - ▷ W' (2016 combination)
 - ▷ $W' \rightarrow WZ \rightarrow \ell\ell q\bar{q}$
 - ▶ $W' \rightarrow WZ \rightarrow \nu\nu q\bar{q}$
 - ▶ $W' \rightarrow WZ \rightarrow \ell\ell q\bar{q}$
 - ▷ $W' \rightarrow WH \rightarrow q\bar{q}\tau\bar{\tau}$
 - ▶ $W' \rightarrow WZ \rightarrow q\bar{q}q\bar{q}$
 - ▶ $W' \rightarrow WH \rightarrow \ell\nu q\bar{q}$
 - ▶ $W' \rightarrow WZ \rightarrow \ell\nu q\bar{q}$
 - ▶ $W' \rightarrow \ell\nu$
 - Radion, $\Lambda_R = 3\text{TeV}$**
 - ▶ $R \rightarrow ZZ \rightarrow \nu\nu q\bar{q}$
 - ▷ $R \rightarrow HH \rightarrow q\bar{q}\tau\bar{\tau}$
 - ▶ $R \rightarrow HH$ (combination)
 - ▶ $R \rightarrow HH \rightarrow b\bar{b}WW$ (lep.) merged-jet
 - ▶ $R \rightarrow HH \rightarrow b\bar{b}WW$ (lep.)
 - ▶ $R \rightarrow HH \rightarrow \tau\tau\gamma\gamma$ (not in HH Comb.)
 - ▶ $R \rightarrow HH \rightarrow \text{multi-leptons}$
 - ▶ $R \rightarrow HH \rightarrow \gamma\gamma b\bar{b}$
 - ▶ $R \rightarrow HH \rightarrow b\bar{b}b\bar{b}$ merged-jet
 - ▶ $R \rightarrow VV \rightarrow q\bar{q}q\bar{q}$
 - ▶ $R \rightarrow WW \rightarrow \ell\nu q\bar{q}$
 - ▷ $R \rightarrow ZZ$
 - ▷ $R \rightarrow WW$
 - ▶ $R \rightarrow WW$
 - Bulk $G, \kappa/M_{Pl} = 0.5$**
 - ▷ $G \rightarrow ZZ \rightarrow \ell\ell\nu\nu$
 - ▷ $G \rightarrow ZZ \rightarrow \ell\ell q\bar{q}$
 - ▶ $G \rightarrow ZZ \rightarrow \nu\nu q\bar{q}$
 - ▶ $G \rightarrow ZZ \rightarrow \ell\ell q\bar{q}$
 - ▶ $G \rightarrow HH$ (combination)
 - ▶ $G \rightarrow HH \rightarrow b\bar{b}WW$ (lep.) merged-jet
 - ▶ $G \rightarrow HH \rightarrow b\bar{b}WW$ (lep.)
 - ▶ $G \rightarrow HH \rightarrow \tau\tau\gamma\gamma$ (not in HH Comb.)
 - ▶ $G \rightarrow HH \rightarrow \text{multi-leptons}$
 - ▶ $G \rightarrow HH \rightarrow \gamma\gamma b\bar{b}$
 - ▶ $G \rightarrow HH \rightarrow b\bar{b}b\bar{b}$ merged-jet
 - ▶ $G \rightarrow WW \rightarrow \ell\nu q\bar{q}$

Overview of CMS B2G Results

June 2024

CMS Preliminary

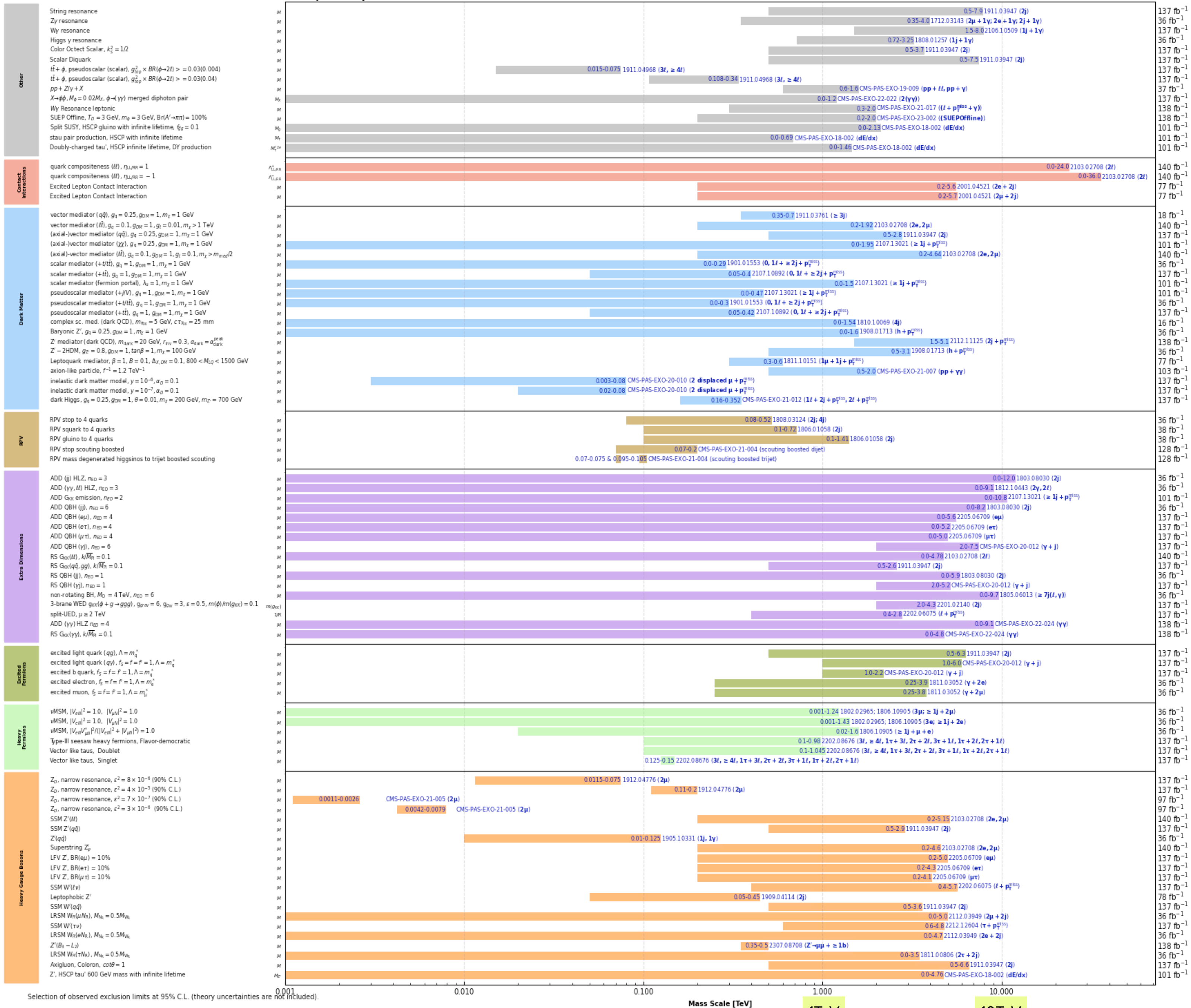
36 – 138 fb⁻¹ (13 TeV)



Overview of CMS EXO results

CMS preliminary

March 2024



We are well above the TeV scale in heavy particle searches

Selection of observed exclusion limits at 95% C.L. (theory uncertainties are not included).

ATLAS Heavy Particle Searches* - 95% CL Upper Exclusion Limits

Status: March 2023

ATLAS Preliminary

$\int \mathcal{L} dt = (3.6 - 139) \text{ fb}^{-1}$

$\sqrt{s} = 13 \text{ TeV}$

Model	ℓ, γ	Jets [†]	E_T^{miss}	$\int \mathcal{L} dt [\text{fb}^{-1}]$	Limit	Reference	
Extra dimen.	ADD $G_{KK} + g/q$	$0 e, \mu, \tau, \gamma$	$1 - 4 j$	Yes	139	M_D 11.2 TeV $n = 2$	2102.10874
	ADD non-resonant $\gamma\gamma$	2γ	-	-	36.7	M_S 8.6 TeV $n = 3$ HLZ NLO	1707.04147
	ADD QBH	-	$2 j$	-	139	M_{th} 9.4 TeV $n = 6$	1910.08447
	ADD BH multijet	-	$\geq 3 j$	-	3.6	M_{th} 9.55 TeV $n = 6, M_D = 3 \text{ TeV}$, rot BH	1512.02586
	RS1 $G_{KK} \rightarrow \gamma\gamma$	2γ	-	-	139	G_{KK} mass 4.5 TeV $k/\overline{M}_{Pl} = 0.1$	2102.13405
	Bulk RS $G_{KK} \rightarrow WW/ZZ$	multi-channel	-	-	36.1	G_{KK} mass 2.3 TeV $k/\overline{M}_{Pl} = 1.0$	1808.02380
	Bulk RS $g_{KK} \rightarrow tt$	$1 e, \mu$	$\geq 1 b, \geq 1 J/2j$	Yes	36.1	g_{KK} mass 3.8 TeV $\Gamma/m = 15\%$	1804.10823
	2UED / RPP	$1 e, \mu$	$\geq 2 b, \geq 3 j$	Yes	36.1	KK mass 1.8 TeV Tier (1,1), $\mathcal{B}(A^{(1,1)} \rightarrow tt) = 1$	1803.09678
Gauge bosons	SSM $Z' \rightarrow \ell\ell$	$2 e, \mu$	-	-	139	Z' mass 5.1 TeV	1903.06248
	SSM $Z' \rightarrow \tau\tau$	2τ	-	-	36.1	Z' mass 2.42 TeV	1709.07242
	Leptophobic $Z' \rightarrow bb$	-	$2 b$	-	36.1	Z' mass 2.1 TeV	1805.09299
	Leptophobic $Z' \rightarrow tt$	$0 e, \mu$	$\geq 1 b, \geq 2 J$	Yes	139	Z' mass 4.1 TeV $\Gamma/m = 1.2\%$	2005.05138
	SSM $W' \rightarrow \ell\nu$	$1 e, \mu$	-	Yes	139	W' mass 6.0 TeV	1906.05609
	SSM $W' \rightarrow \tau\nu$	1τ	-	Yes	139	W' mass 5.0 TeV	ATLAS-CONF-2021-025
	SSM $W' \rightarrow tb$	-	$\geq 1 b, \geq 1 J$	-	139	W' mass 4.4 TeV	ATLAS-CONF-2021-043
	HVT $W' \rightarrow WZ$ model B	$0-2 e, \mu$	$2 j / 1 J$	Yes	139	W' mass 4.3 TeV	2004.14636
	HVT $W' \rightarrow WZ \rightarrow \ell\nu \ell'\ell'$ model C	$3 e, \mu$	$2 j$ (VBF)	Yes	139	W' mass 340 GeV $g_V = 3$ $g_V c_H = 1, g_f = 0$	2207.03925
	HVT $Z' \rightarrow WW$ model B	$1 e, \mu$	$2 j / 1 J$	Yes	139	Z' mass 3.9 TeV $g_V = 3$	2004.14636
LRSM $W_R \rightarrow \mu N_R$	2μ	$1 J$	-	80	W_R mass 5.0 TeV $m(N_R) = 0.5 \text{ TeV}, g_L = g_R$	1904.12679	
CI	CI $qqqq$	-	$2 j$	-	37.0	Λ 21.8 TeV η_{LL}^-	1703.09127
	CI $\ell\ell qq$	$2 e, \mu$	-	-	139	Λ 35.8 TeV η_{LL}^-	2006.12946
	CI $eebs$	$2 e$	$1 b$	-	139	Λ 1.8 TeV $g_s = 1$	2105.13847
	CI $\mu\mu bs$	2μ	$1 b$	-	139	Λ 2.0 TeV $g_s = 1$	2105.13847
	CI $tttt$	$\geq 1 e, \mu$	$\geq 1 b, \geq 1 j$	Yes	36.1	Λ 2.57 TeV $ C_{4t} = 4\pi$	1811.02305
DM	Axial-vector med. (Dirac DM)	-	$2 j$	-	139	m_{med} 3.8 TeV	ATL-PHYS-PUB-2022-036
	Pseudo-scalar med. (Dirac DM)	$0 e, \mu, \tau, \gamma$	$1 - 4 j$	Yes	139	m_{med} 376 GeV	2102.10874
	Vector med. Z' -2HDM (Dirac DM)	$0 e, \mu$	$2 b$	Yes	139	$m_{Z'}$ 3.0 TeV	2108.13391
	Pseudo-scalar med. 2HDM+a	multi-channel	-	-	139	m_a 800 GeV	ATLAS-CONF-2021-036
LQ	Scalar LQ 1 st gen	$2 e$	$\geq 2 j$	Yes	139	LQ mass 1.8 TeV $\beta = 1$	2006.05872
	Scalar LQ 2 nd gen	2μ	$\geq 2 j$	Yes	139	LQ mass 1.7 TeV $\beta = 1$	2006.05872
	Scalar LQ 3 rd gen	1τ	$2 b$	Yes	139	LQ_3^u mass 1.49 TeV $\mathcal{B}(LQ_3^u \rightarrow b\tau) = 1$	2303.01294
	Scalar LQ 3 rd gen	$0 e, \mu$	$\geq 2 j, \geq 2 b$	Yes	139	LQ_3^d mass 1.24 TeV $\mathcal{B}(LQ_3^d \rightarrow t\nu) = 1$	2004.14060
	Scalar LQ 3 rd gen	$\geq 2 e, \mu, \geq 1 \tau$	$\geq 1 j, \geq 1 b$	-	139	LQ_3^d mass 1.43 TeV $\mathcal{B}(LQ_3^d \rightarrow t\tau) = 1$	2101.11582
	Scalar LQ 3 rd gen	$0 e, \mu, \geq 1 \tau$	$0 - 2 j, 2 b$	Yes	139	LQ_3^d mass 1.26 TeV $\mathcal{B}(LQ_3^d \rightarrow b\nu) = 1$	2101.12527
	Vector LQ mix gen	multi-channel	$\geq 1 j, \geq 1 b$	Yes	139	LQ_3^V mass 2.0 TeV $\mathcal{B}(\tilde{U}_1 \rightarrow t\mu) = 1$, Y-M coupl.	ATLAS-CONF-2022-052
	Vector LQ 3 rd gen	$2 e, \mu, \tau$	$\geq 1 b$	Yes	139	LQ_3^V mass 1.96 TeV $\mathcal{B}(LQ_3^V \rightarrow b\tau) = 1$, Y-M coupl.	2303.01294
Vector-like fermions	VLQ $TT \rightarrow Zt + X$	$2e/2\mu/\geq 3e, \mu$	$\geq 1 b, \geq 1 j$	-	139	T mass 1.46 TeV	SU(2) doublet
	VLQ $BB \rightarrow Wt/Zb + X$	multi-channel	-	-	36.1	B mass 1.34 TeV	SU(2) doublet
	VLQ $T_{5/3} T_{5/3} T_{5/3} \rightarrow Wt + X$	$2(SS)/\geq 3 e, \mu$	$\geq 1 b, \geq 1 j$	Yes	36.1	$T_{5/3}$ mass 1.64 TeV $\mathcal{B}(T_{5/3} \rightarrow Wt) = 1, c(T_{5/3} Wt) = 1$	1807.11883
	VLQ $T \rightarrow Ht/Zt$	$1 e, \mu$	$\geq 1 b, \geq 3 j$	Yes	139	T mass 1.8 TeV	SU(2) singlet, $\kappa_T = 0.5$
	VLQ $Y \rightarrow Wb$	$1 e, \mu$	$\geq 1 b, \geq 1 j$	Yes	36.1	Y mass 1.85 TeV	$\mathcal{B}(Y \rightarrow Wb) = 1, c_R(Wb) = 1$
	VLQ $B \rightarrow Hb$	$0 e, \mu$	$\geq 2b, \geq 1j, \geq 1J$	-	139	B mass 2.0 TeV	SU(2) doublet, $\kappa_B = 0.3$
	VLL $\tau' \rightarrow Z\tau/H\tau$	multi-channel	$\geq 1 j$	Yes	139	τ' mass 898 GeV	SU(2) doublet
Exctd ferm.	Excited quark $q^* \rightarrow qg$	-	$2 j$	-	139	q^* mass 6.7 TeV	only u^* and d^* , $\Lambda = m(q^*)$
	Excited quark $q^* \rightarrow q\gamma$	1γ	$1 j$	-	36.7	q^* mass 5.3 TeV	only u^* and d^* , $\Lambda = m(q^*)$
	Excited quark $b^* \rightarrow bg$	-	$1 b, 1 j$	-	139	b^* mass 3.2 TeV	$\Lambda = 4.6 \text{ TeV}$
	Excited lepton τ^*	2τ	$\geq 2 j$	-	139	τ^* mass 4.6 TeV	
	Other	Type III Seesaw	$2,3,4 e, \mu$	$\geq 2 j$	Yes	139	N^0 mass 910 GeV
LRSM Majorana ν		2μ	$2 j$	-	36.1	N_R mass 3.2 TeV	DY production
Higgs triplet $H^{\pm\pm} \rightarrow W^\pm W^\pm$		$2,3,4 e, \mu$ (SS)	various	Yes	139	$H^{\pm\pm}$ mass 350 GeV	DY production
Higgs triplet $H^{\pm\pm} \rightarrow \ell\ell$		$2,3,4 e, \mu$ (SS)	-	-	139	$H^{\pm\pm}$ mass 1.08 TeV	DY production, $ q = 5e$
Multi-charged particles		-	-	-	139	multi-charged particle mass 1.59 TeV	DY production, $ g = 1g_D$
Magnetic monopoles		-	-	-	34.4	monopole mass 2.37 TeV	

$\sqrt{s} = 13 \text{ TeV}$
partial data $\sqrt{s} = 13 \text{ TeV}$
full data

We look for many other different models, extra dimensions, extra bosons...

*Only a selection of the available mass limits on new states or phenomena is shown.

† Small-radius (large-radius) jets are denoted by the letter j (J).

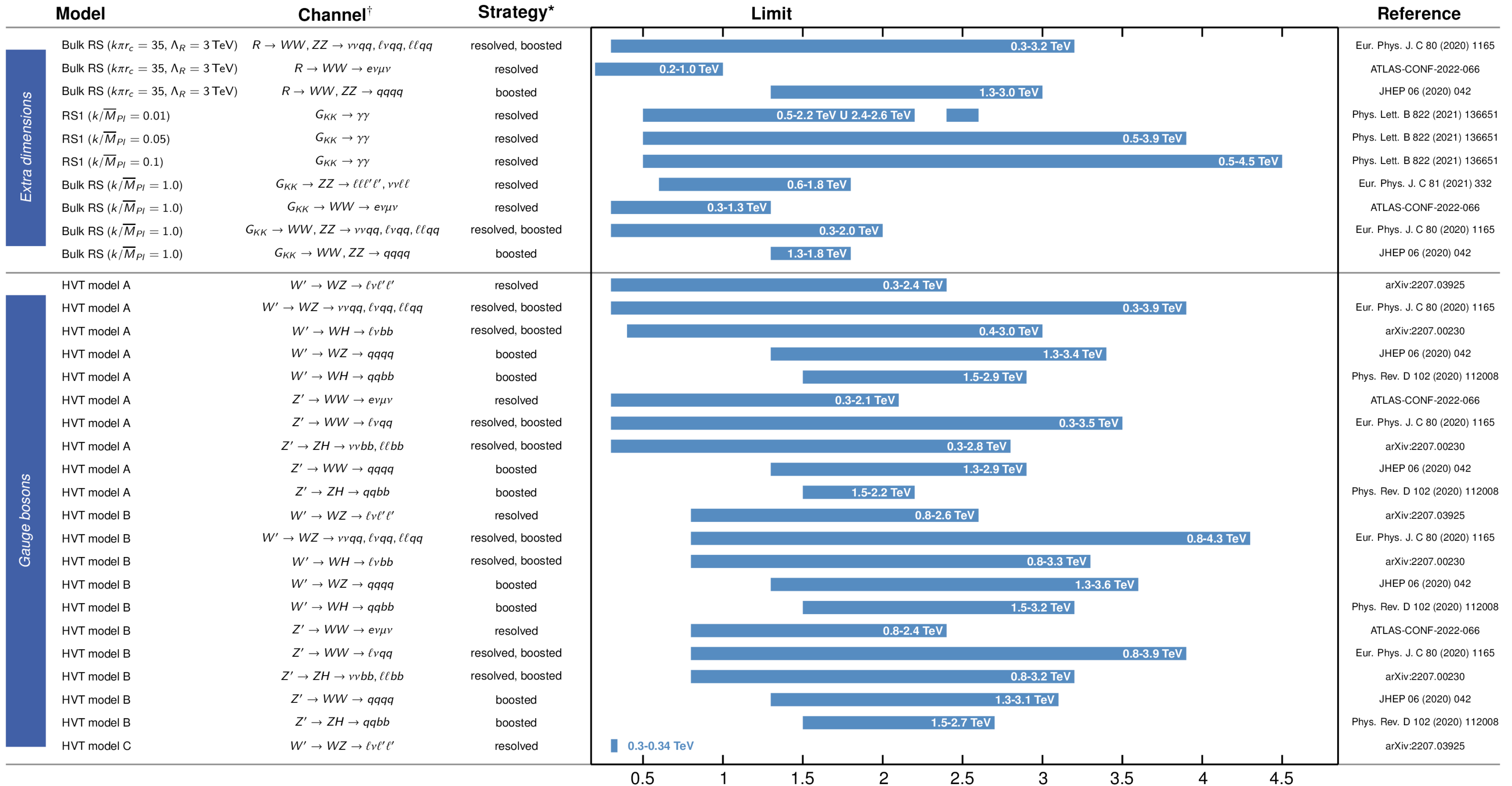
ATLAS Diboson Searches - 95% CL Exclusion Limits

Status: March 2023

$\mathcal{L} = 139 \text{ fb}^{-1}$

ATLAS Preliminary

$\sqrt{s} = 13 \text{ TeV}$



HVT model A: $g_F = -0.55, g_H = -0.56$

HVT model B: $g_F = 0.14, g_H = -2.9$

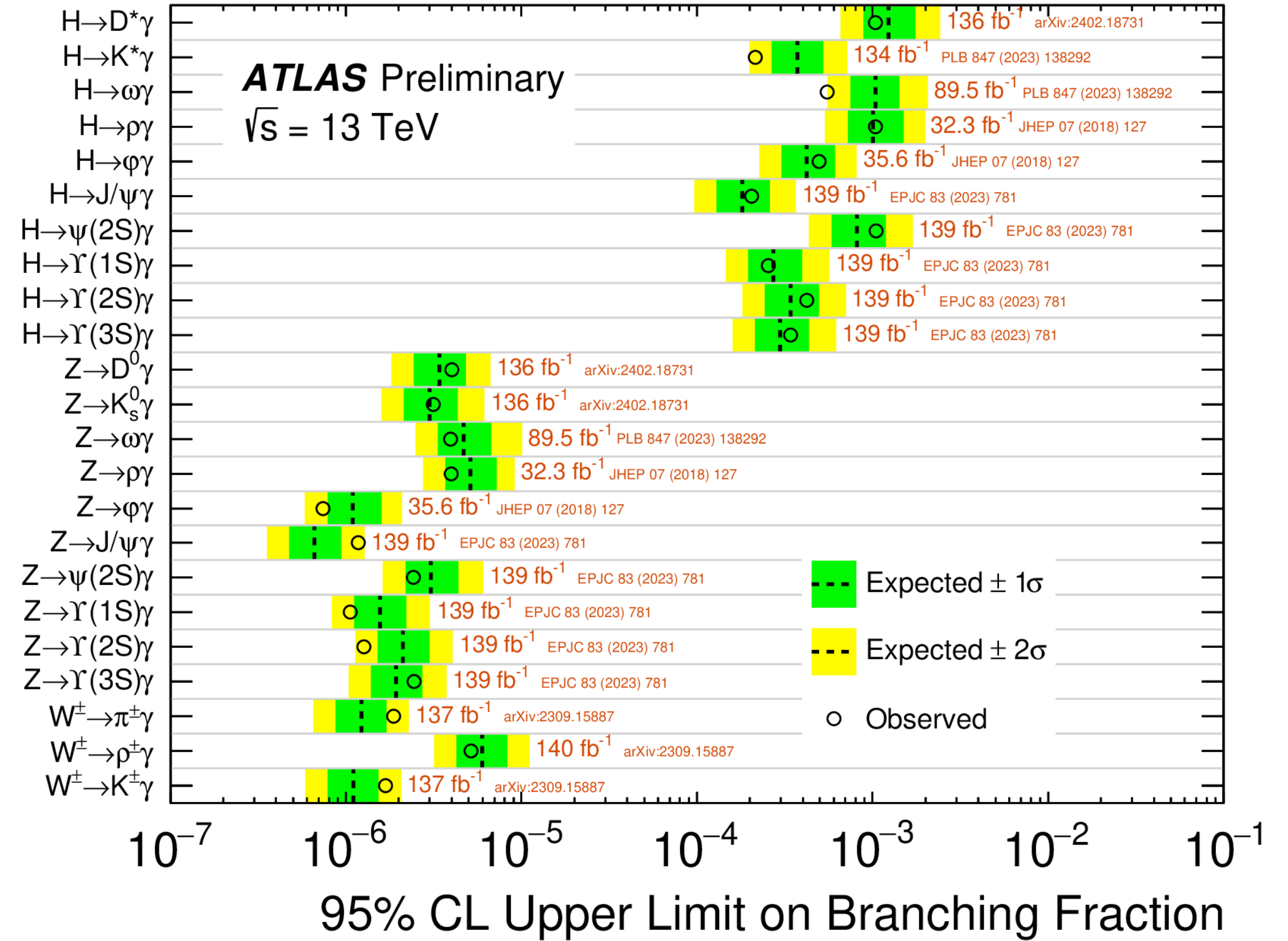
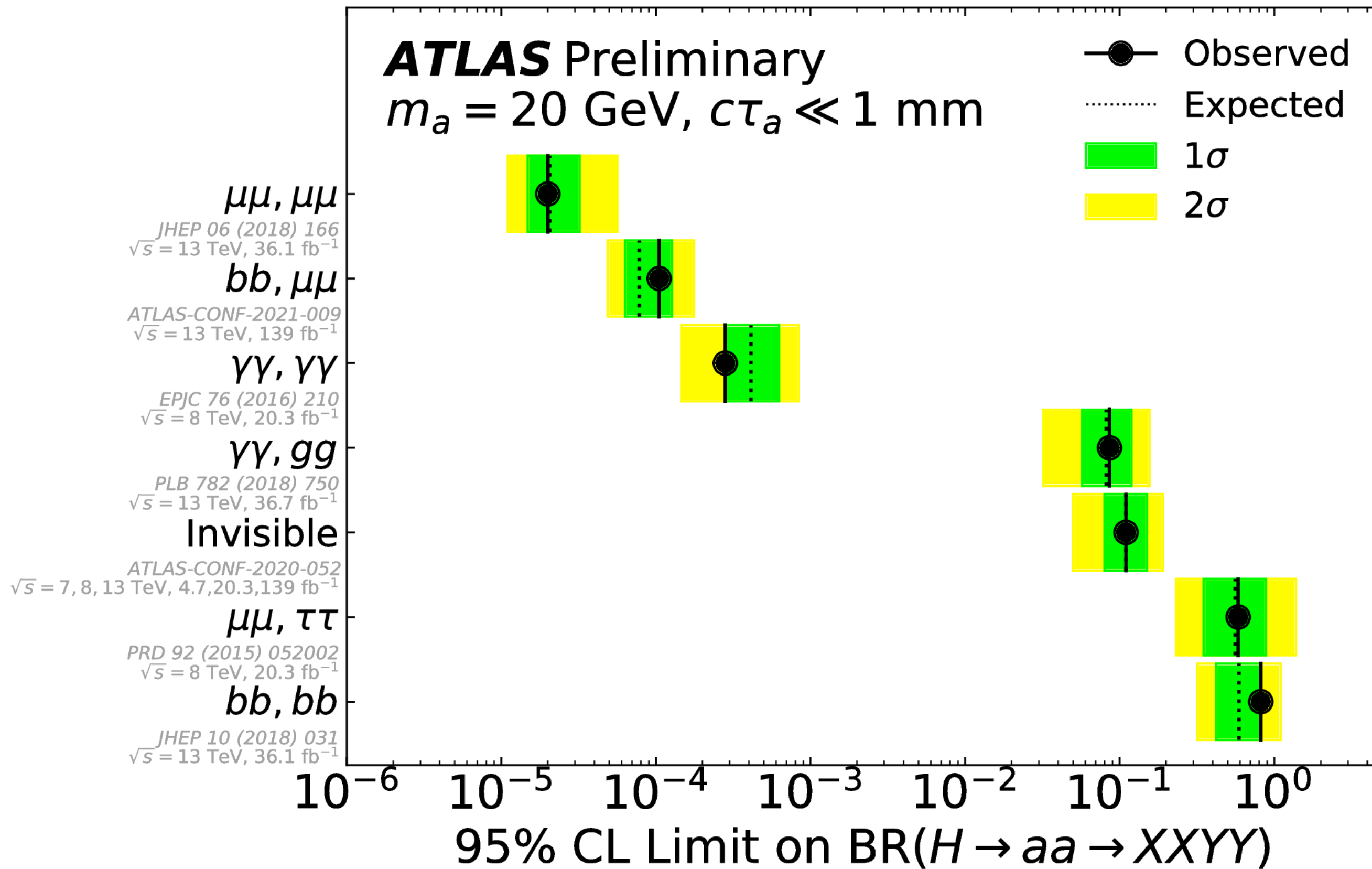
HVT model C: $g_F = 0, g_H = 1$

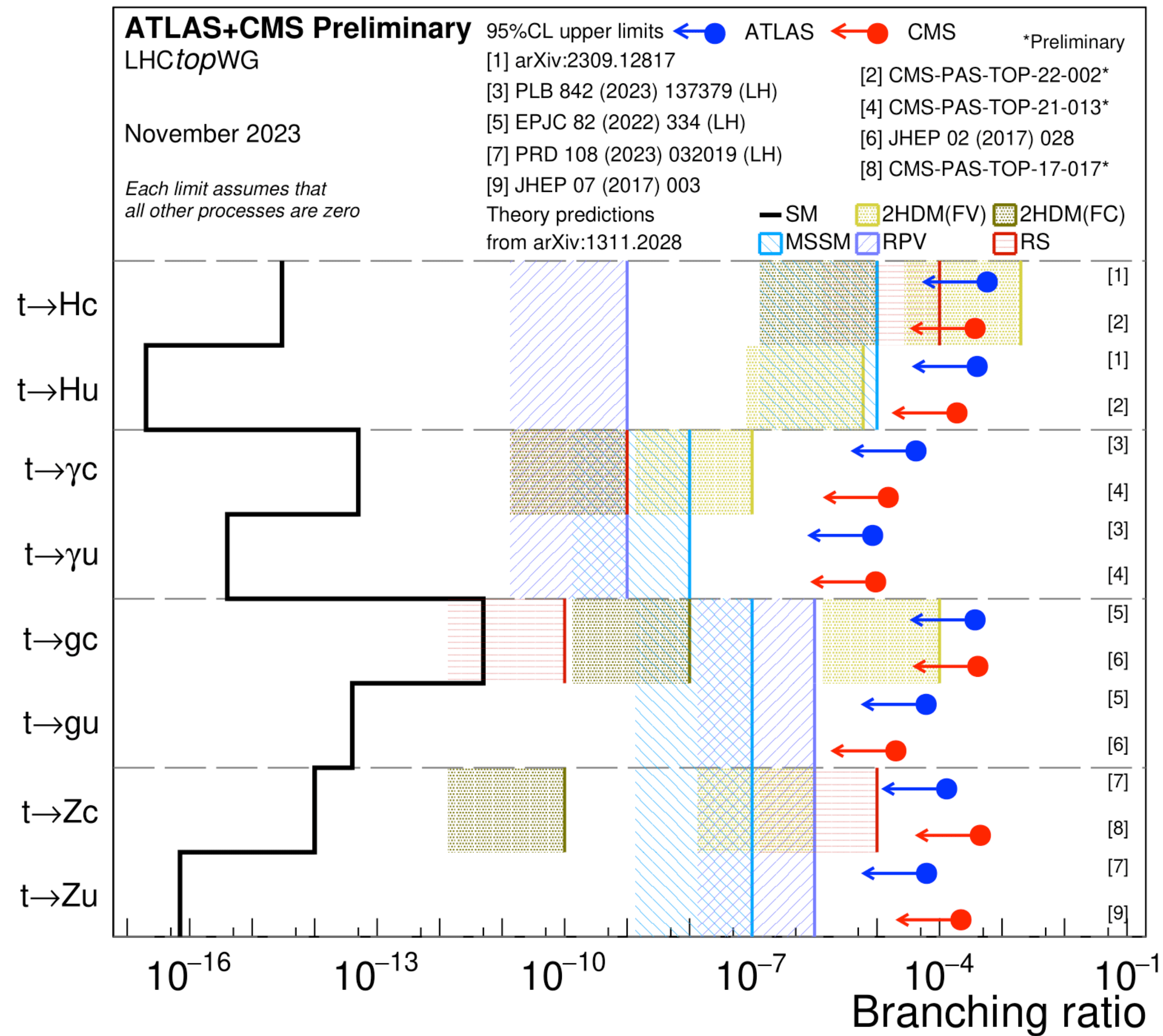
*small-radius (large-radius) jets are used in resolved (boosted) events

[†]with $\ell = \mu, e$

We look for exotic decays of the Higgs boson

March 2021

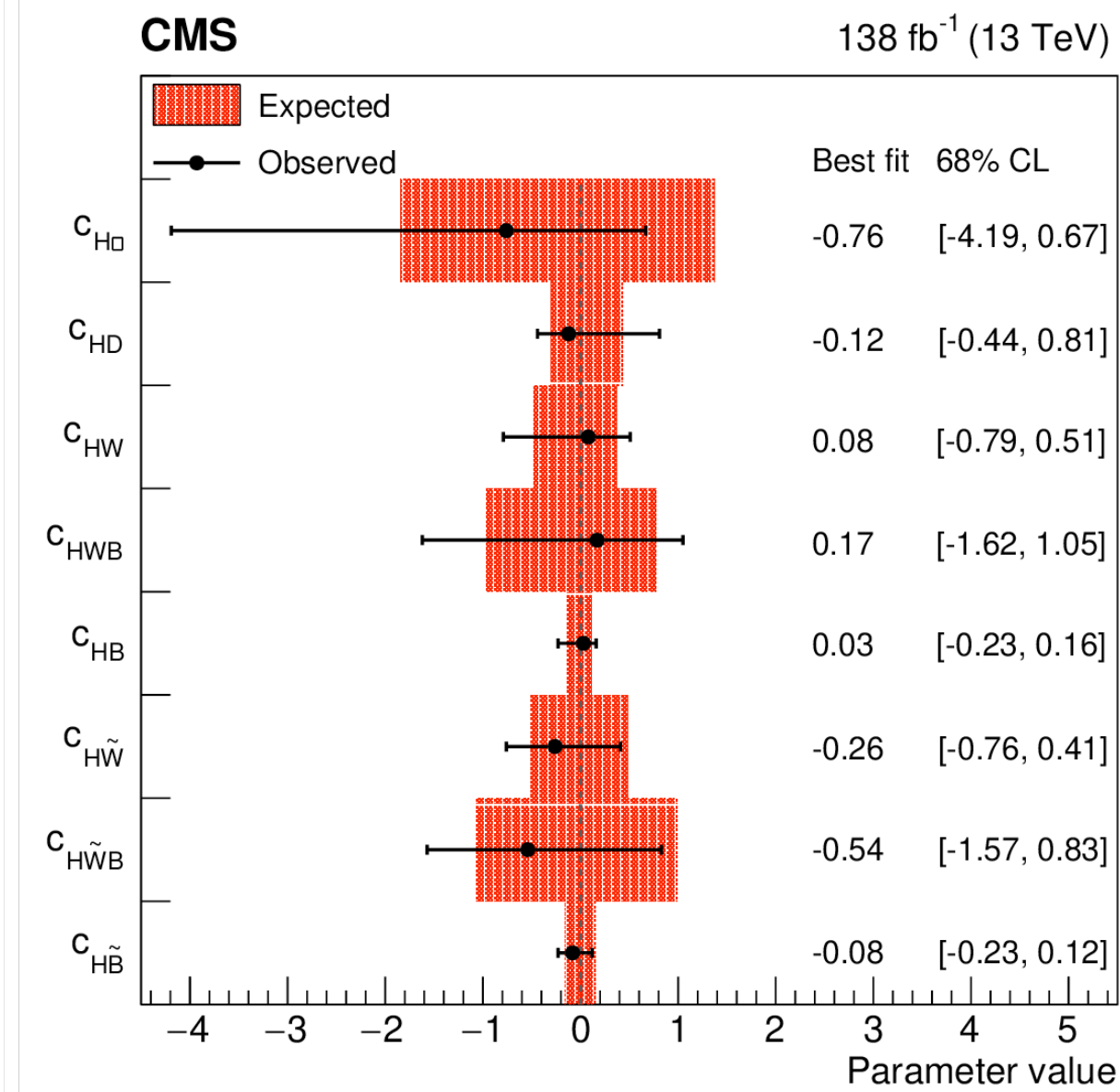
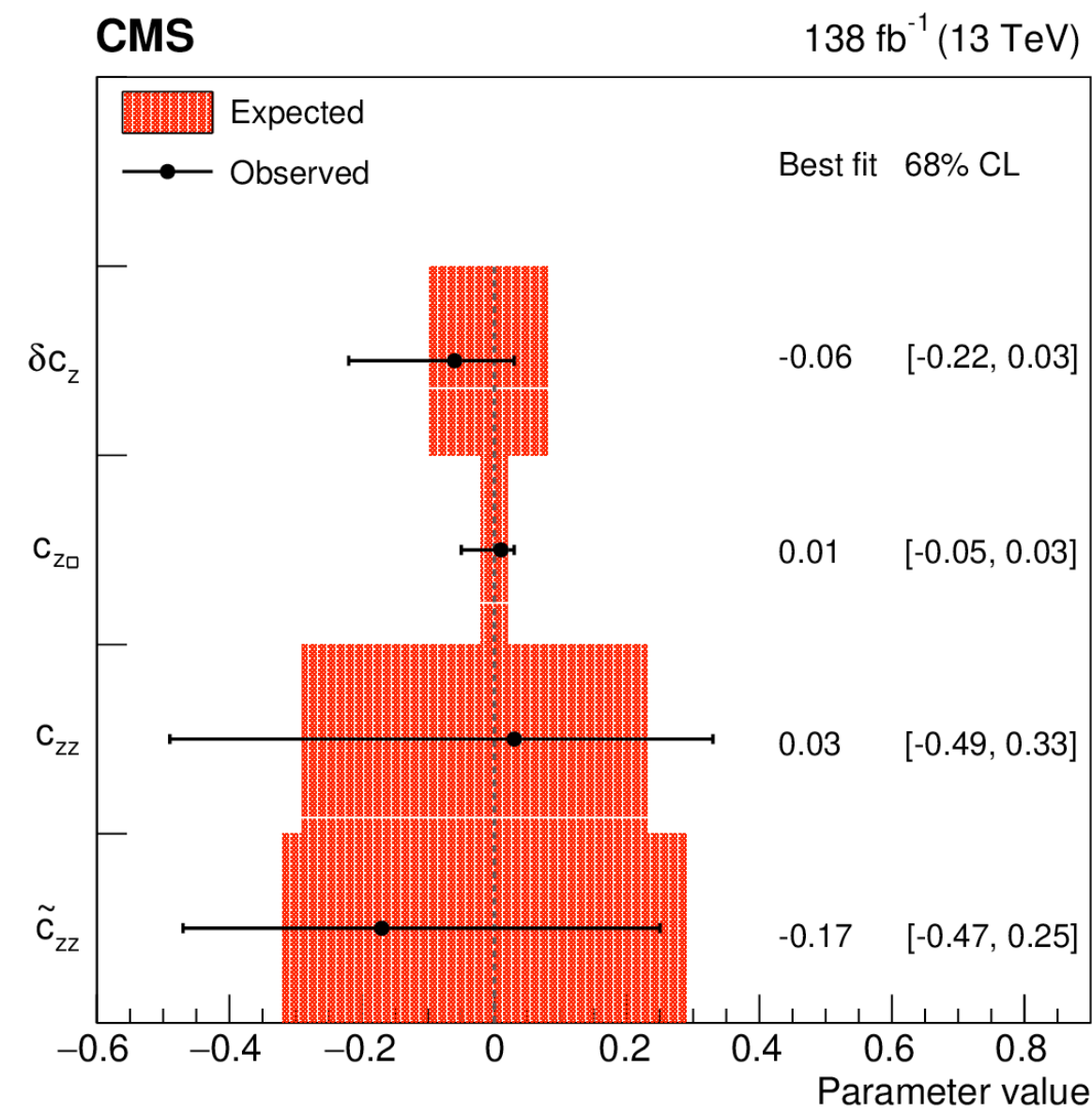




And of the top quark

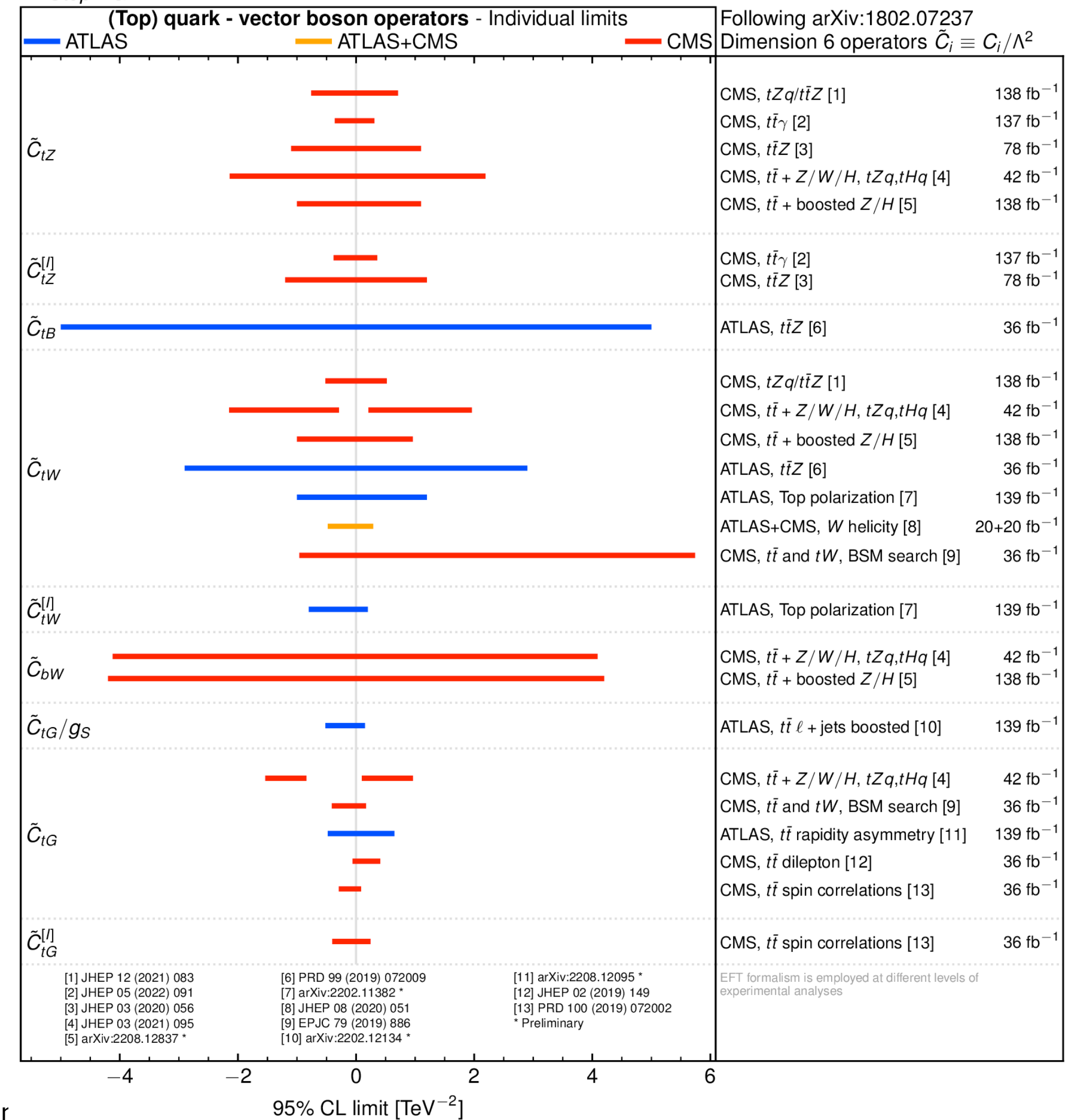
WE LOOK AT PRECISION IN A DIFFERENT WAY

- Interpreting the LHC data in the context of effective field theories (EFTs)

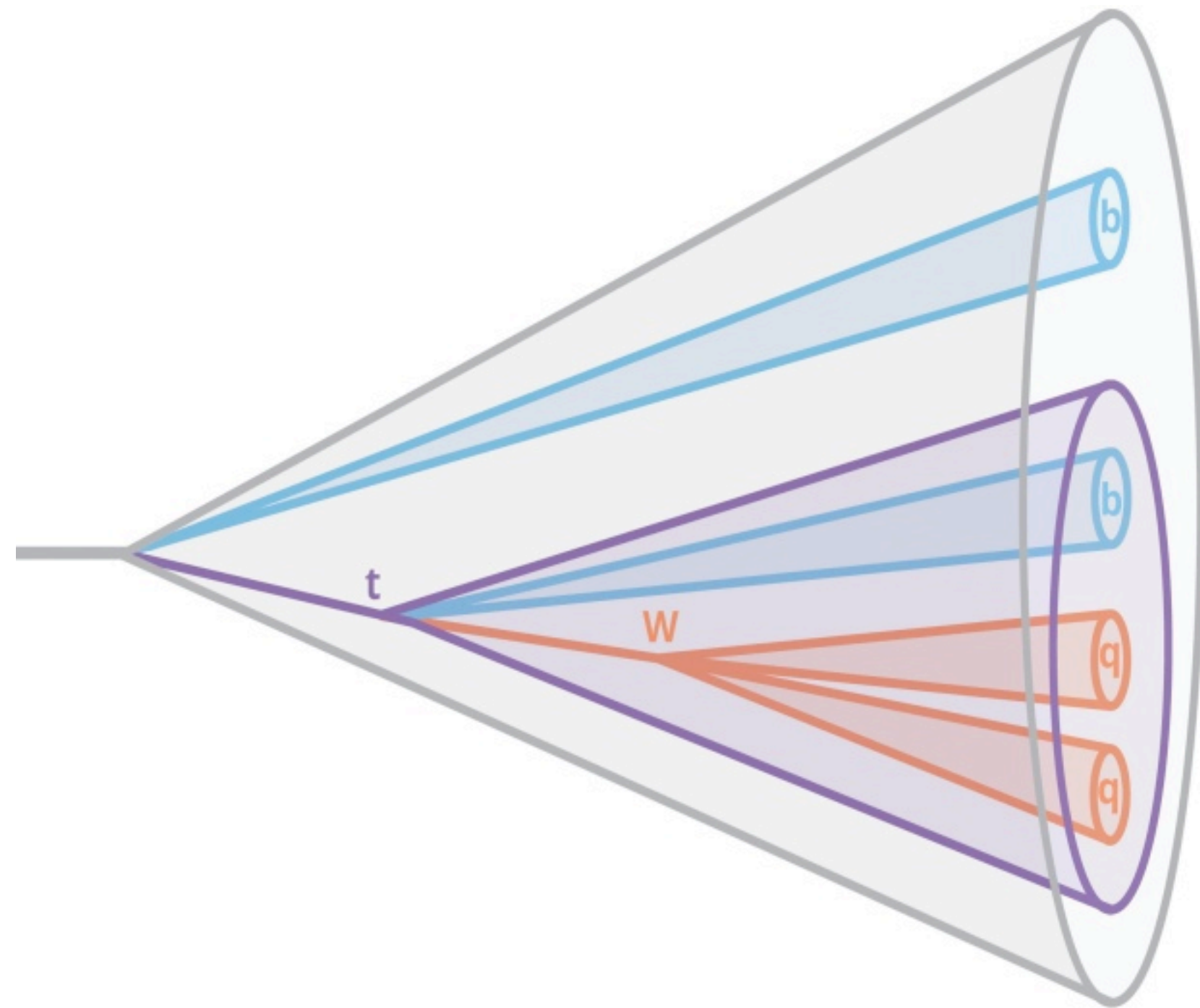


ATLAS+CMS Preliminary
LHCtopWG

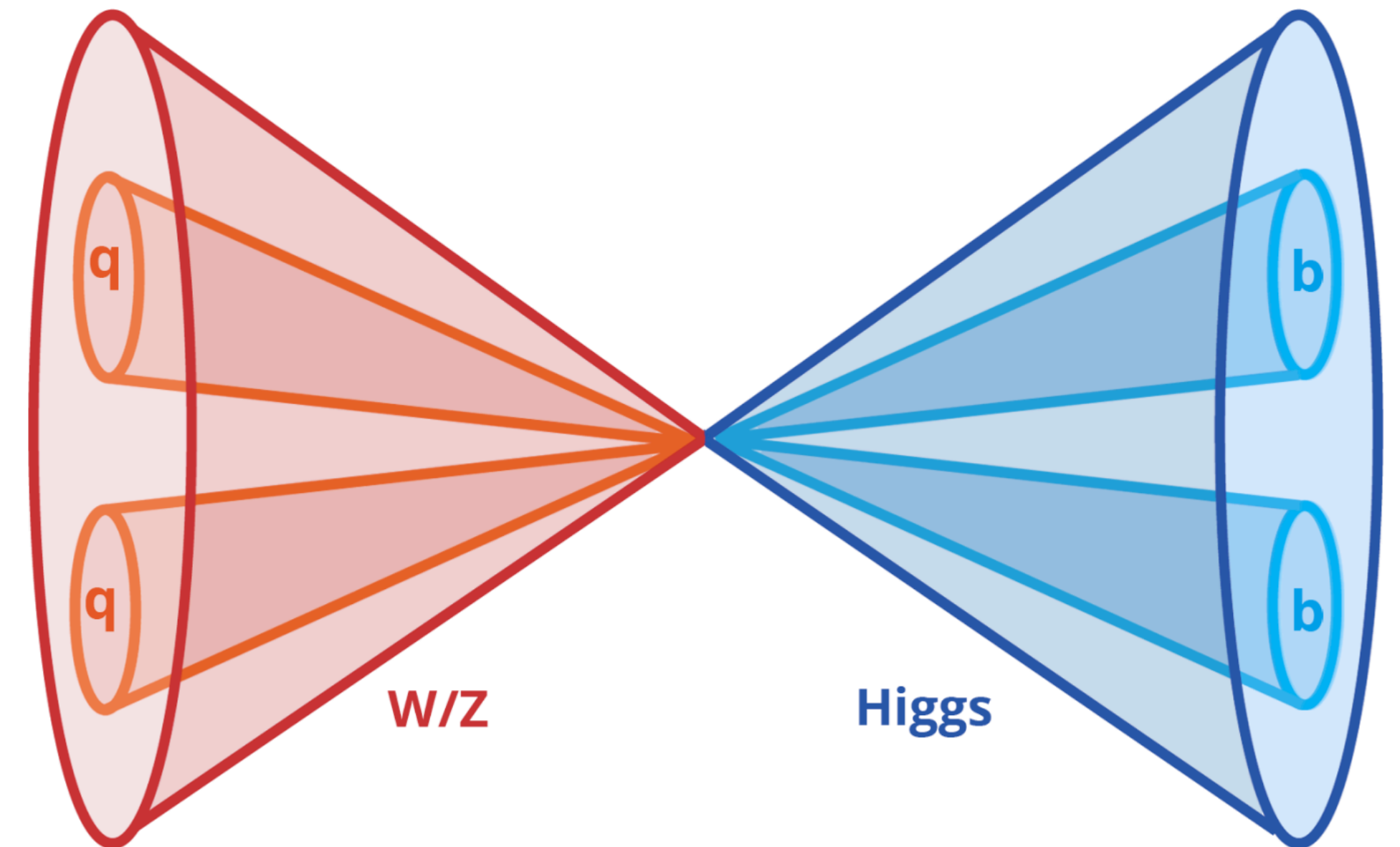
November 2022



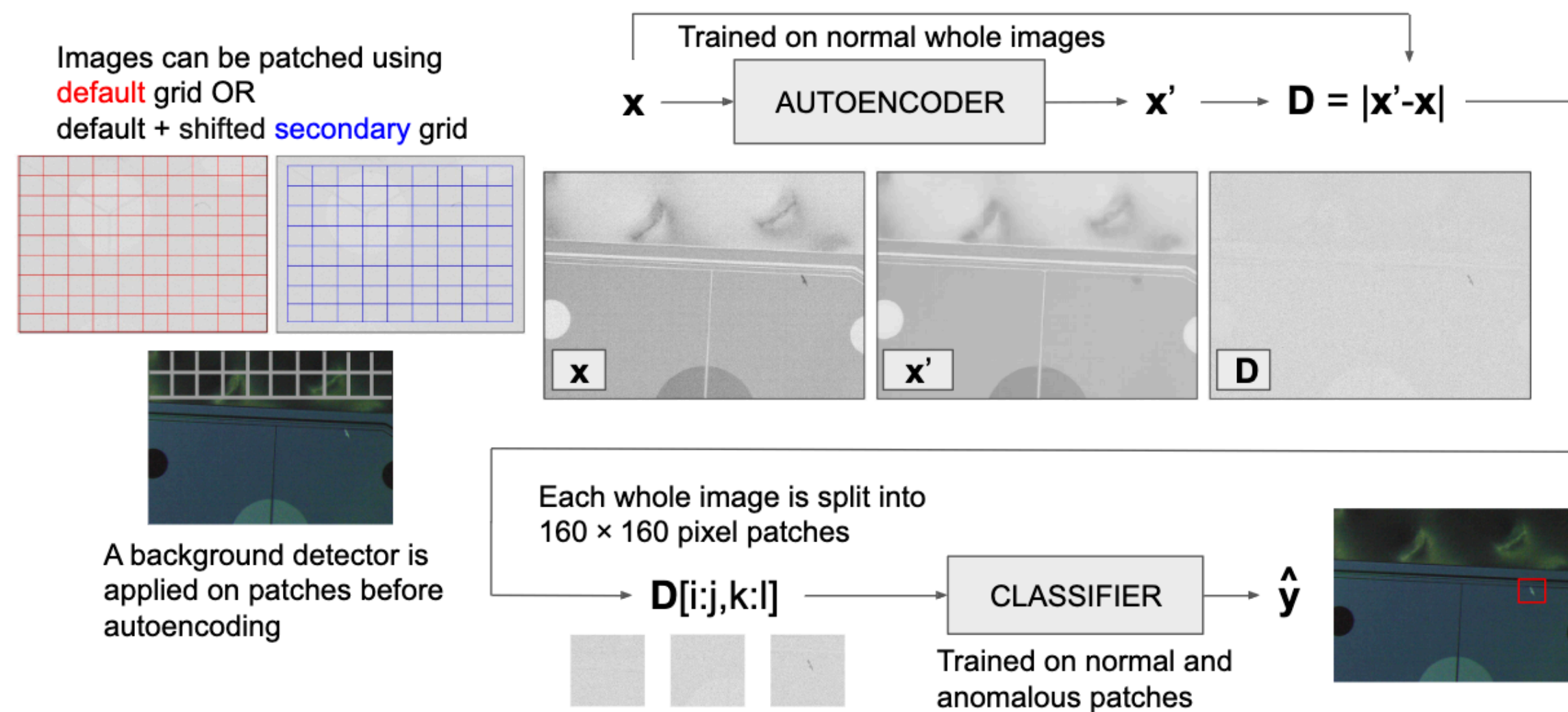
WE INCLUDE INCREASINGLY COMPLICATED INGREDIENTS



Like boosted objects, collimated decays, that can have substructure



AND INCREASINGLY COMPLICATED METHODS



<https://arxiv.org/pdf/2303.15319>

We've come a long way from the cut-based analysis days

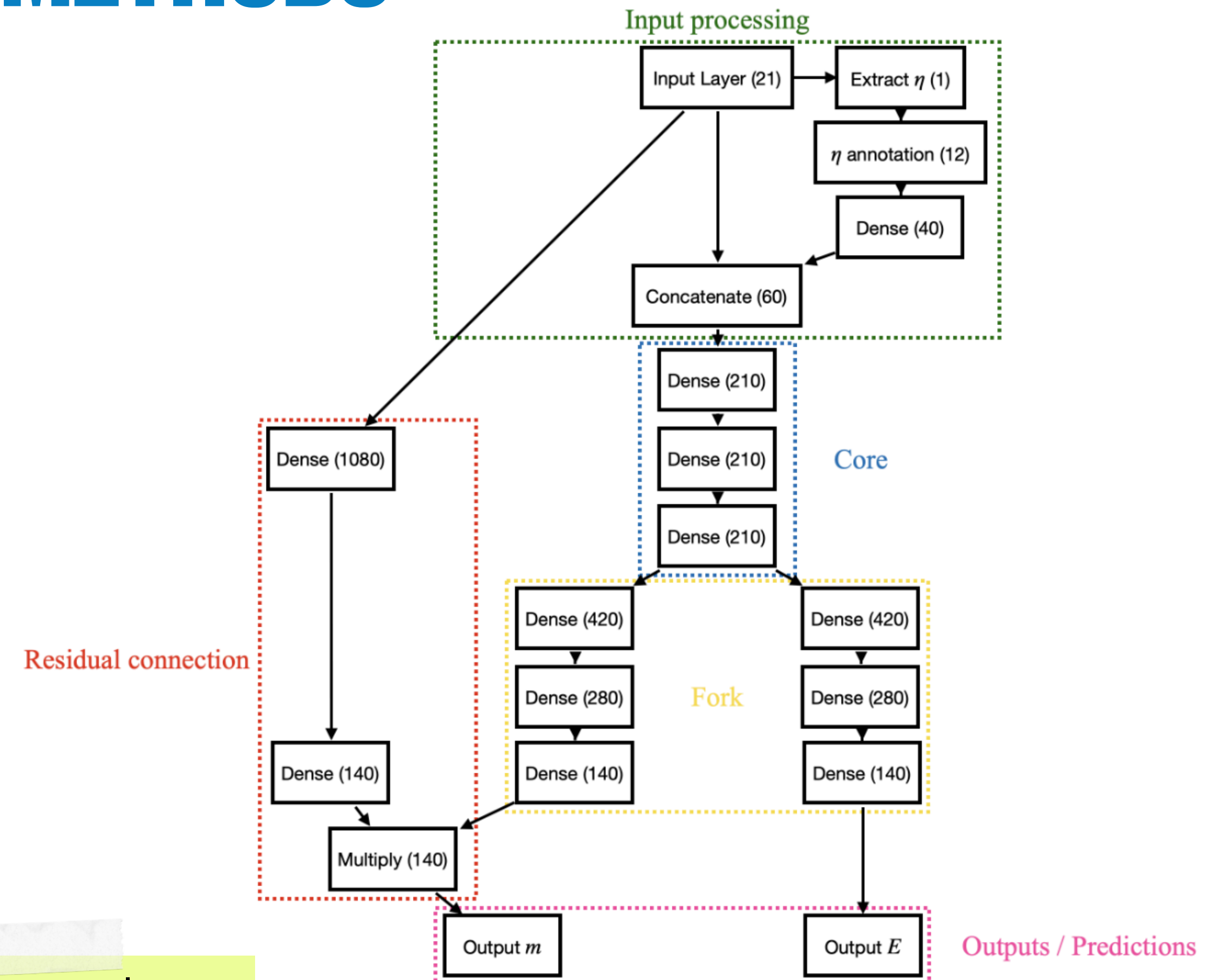
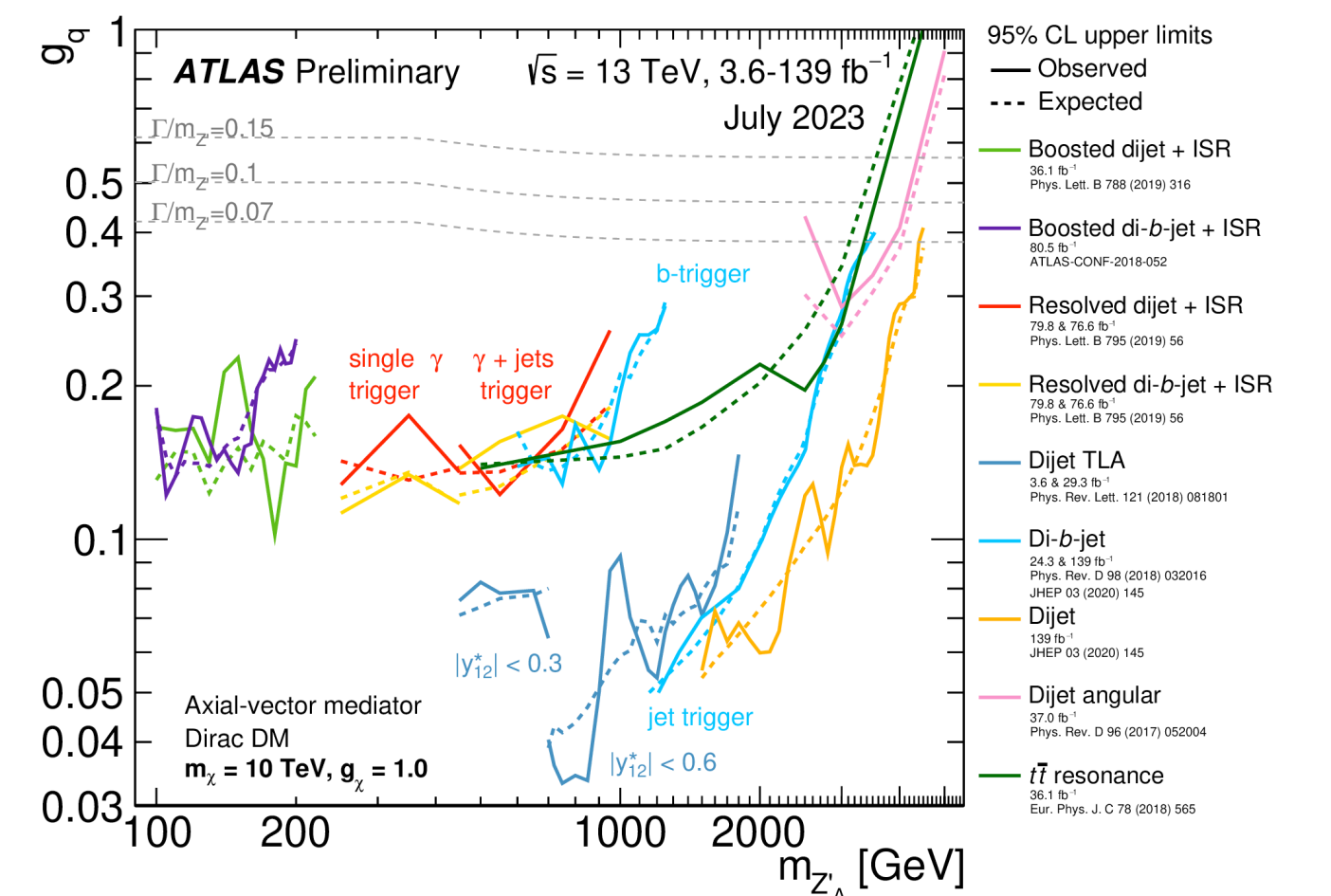
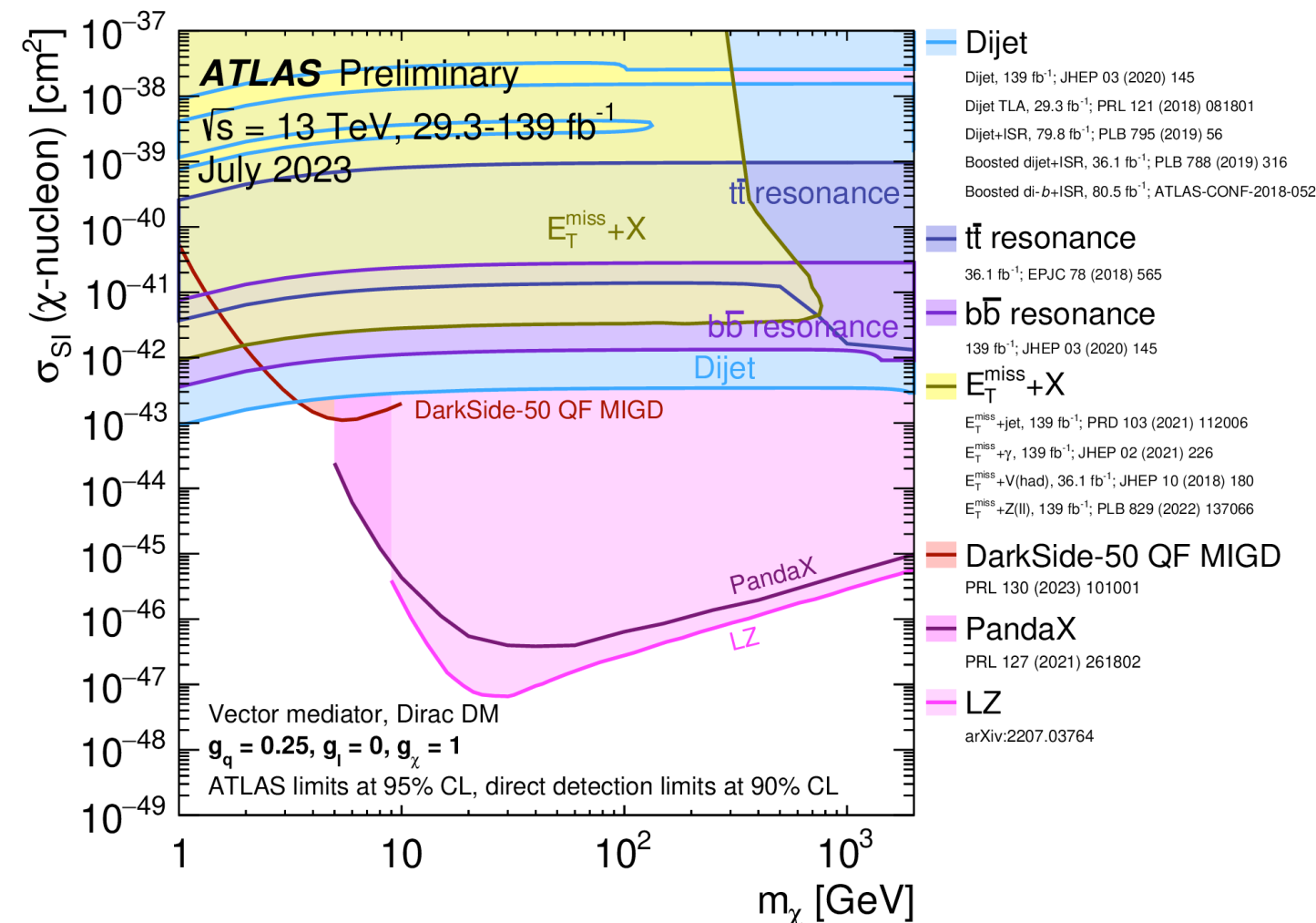
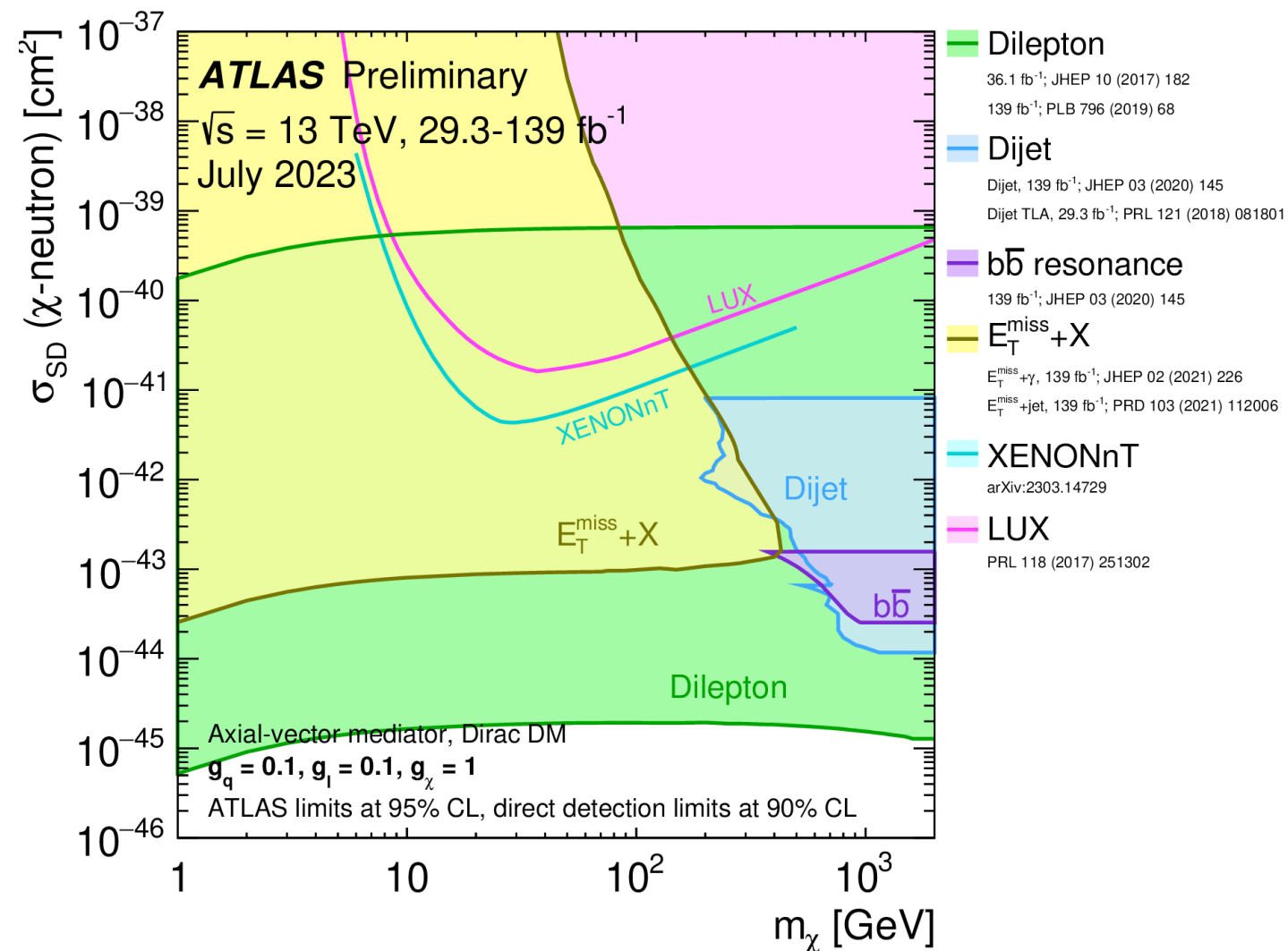
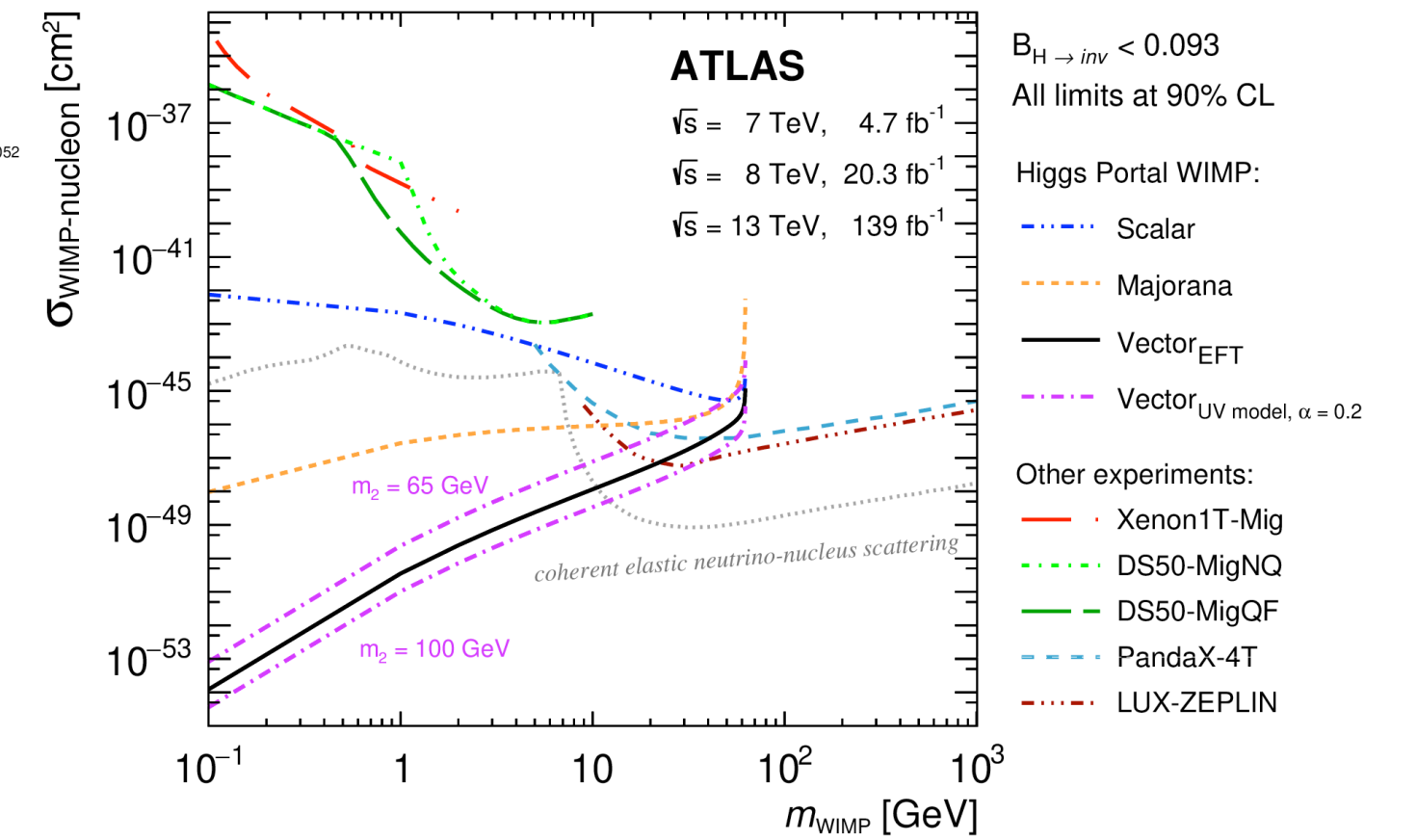
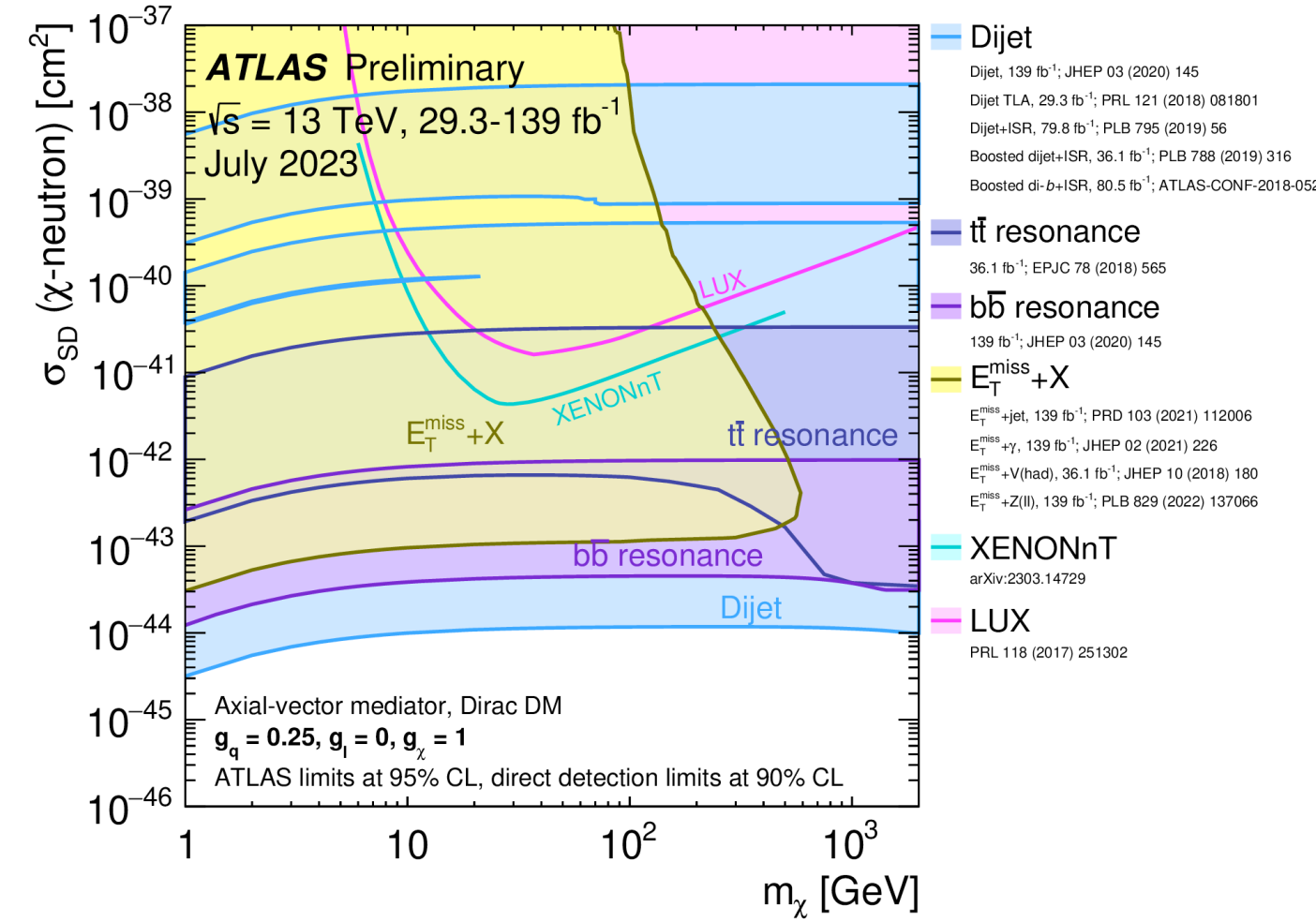
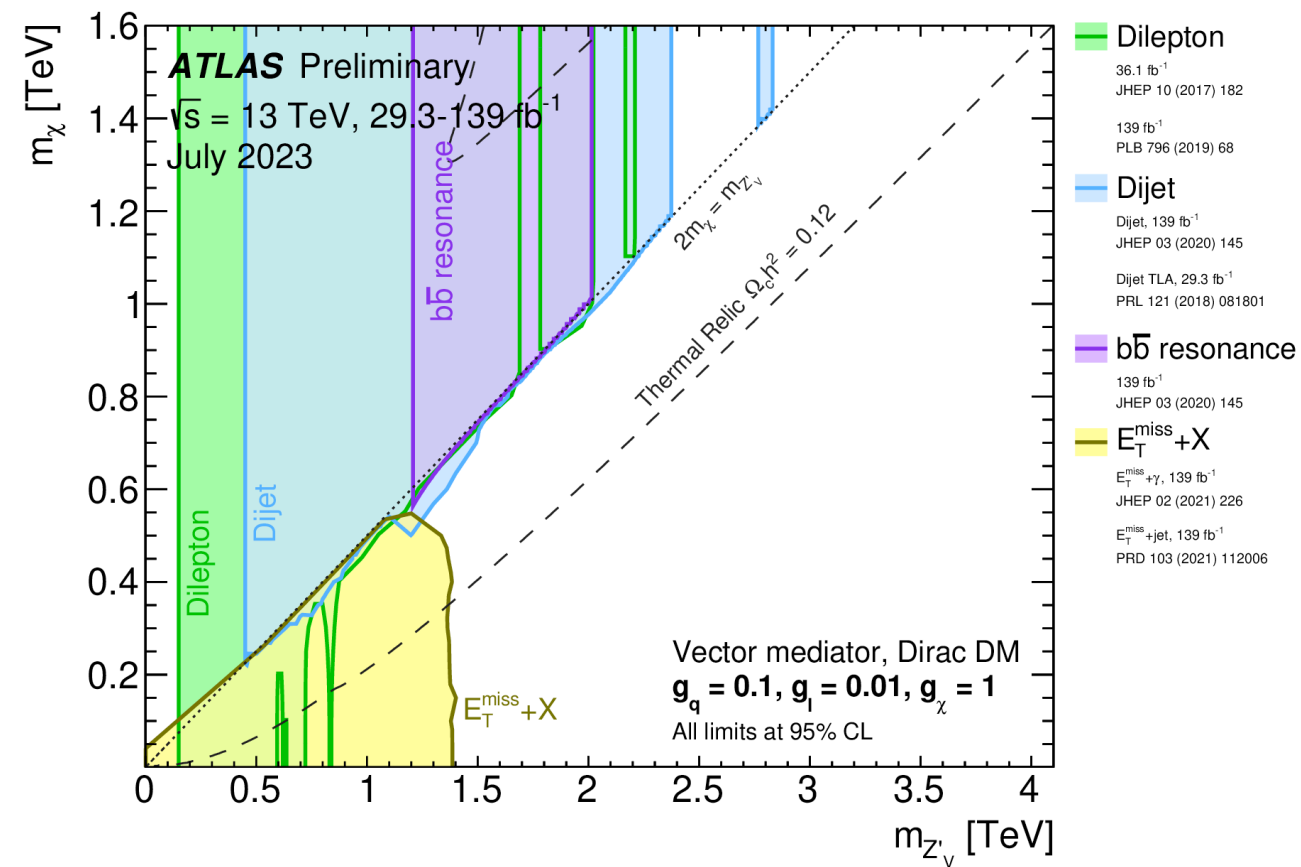


Figure 2: The architecture of the DNN.

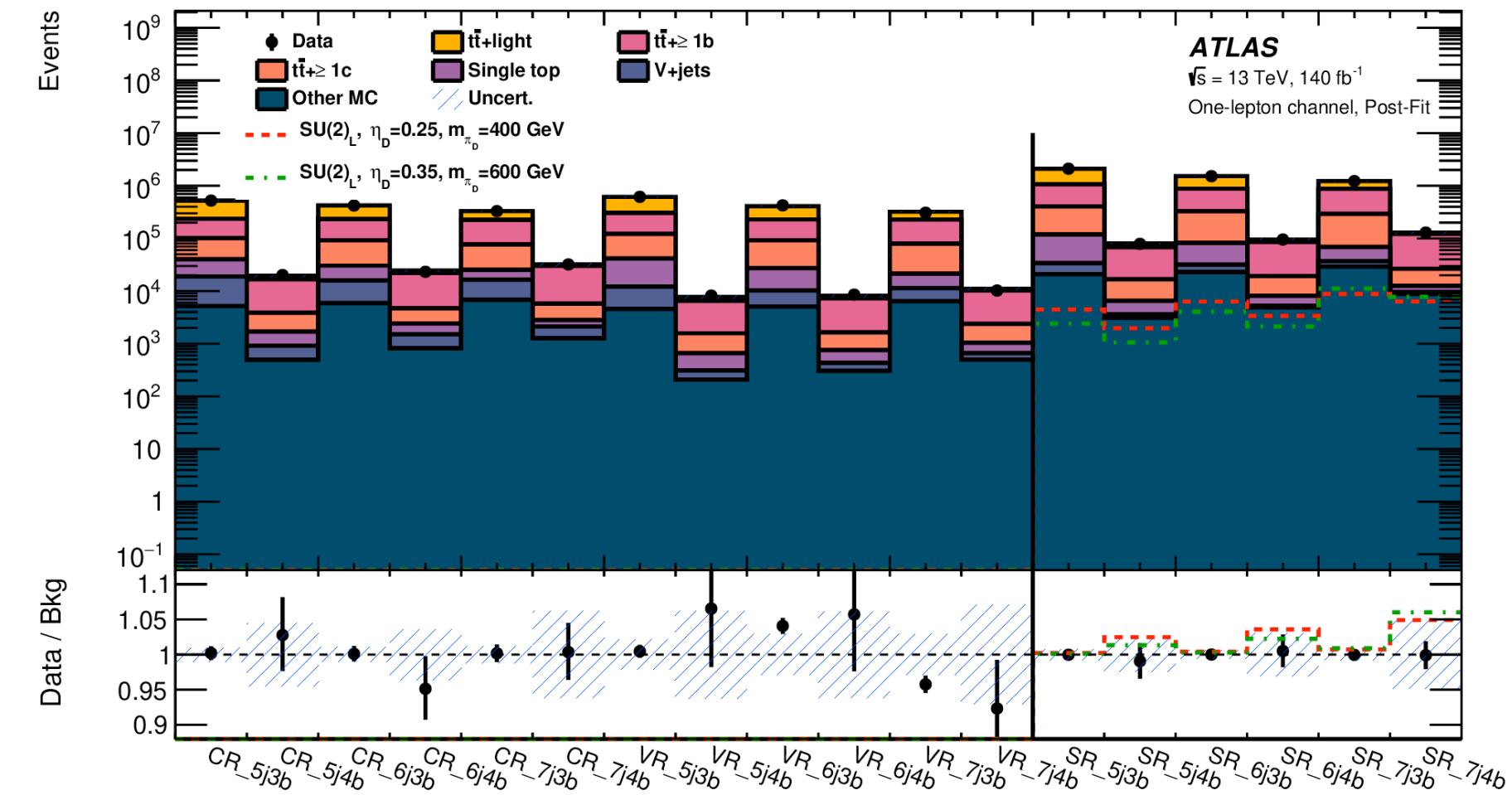
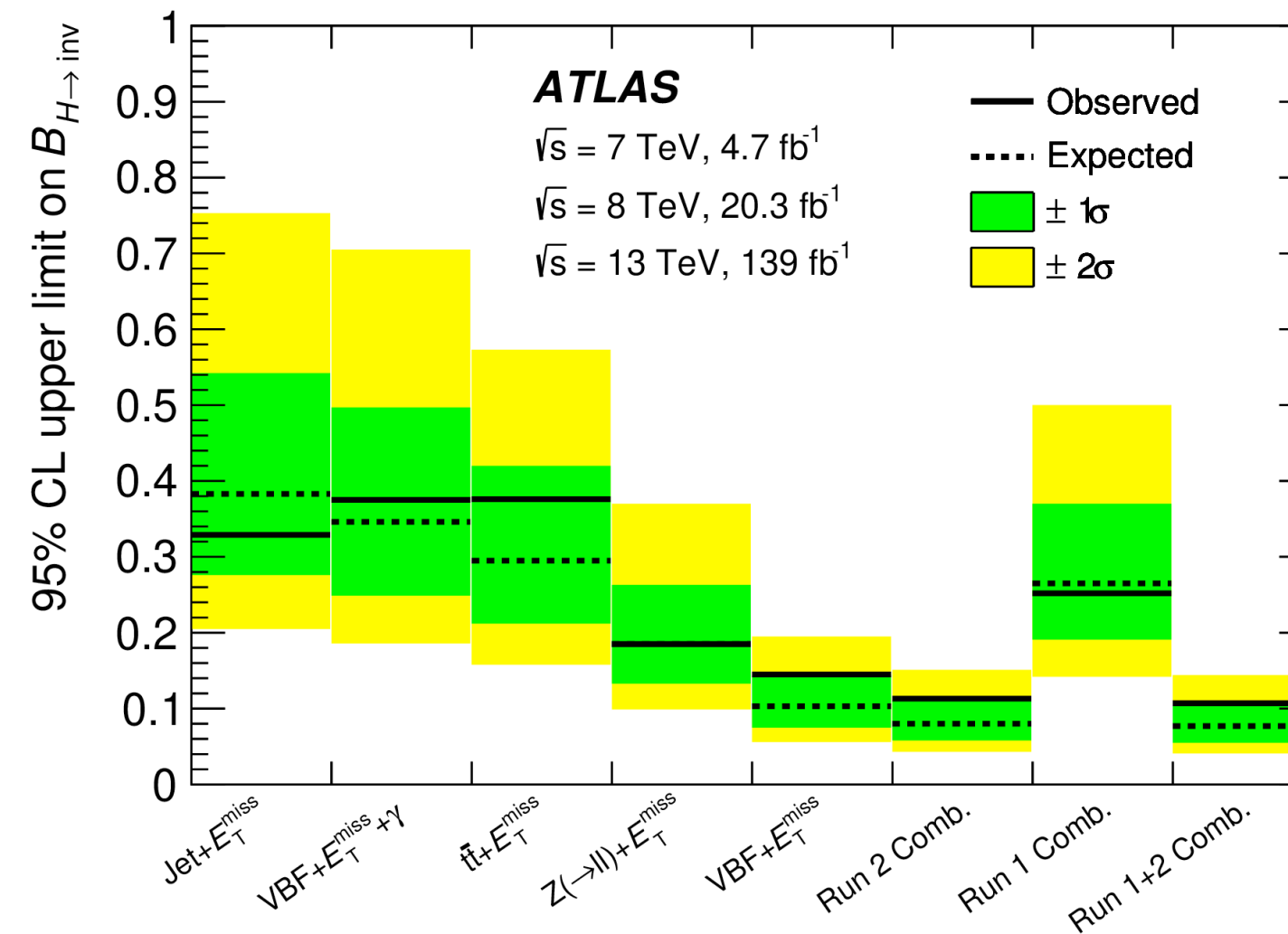
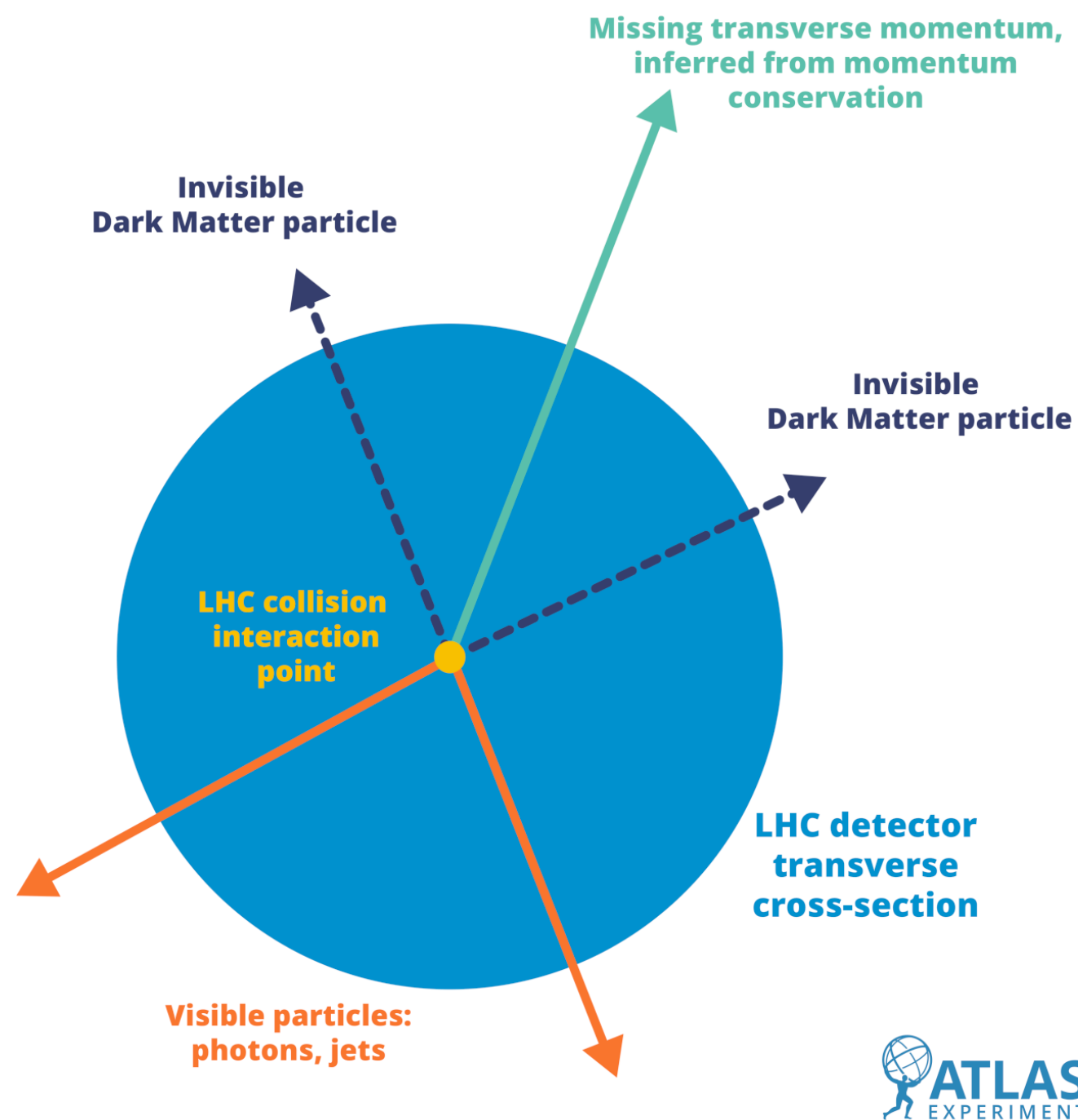
<https://arxiv.org/pdf/2311.08885>

WE GO ACROSS PHYSICS FIELDS

We look for dark matter in many different ways



ALSO WITH INCREASED COMPLEXITY



At the very beginning, our dark matter searches were mostly “mono” things (a SM object + MET) or invisible decays of the Higgs boson but after many searches we have expanded the criteria to more complex models



🍏 Even more complicated signatures, complicated background, needs for new techniques, new triggers, machine learning, new ideas

🍏🍏 More complicated signatures, larger background, smaller signals

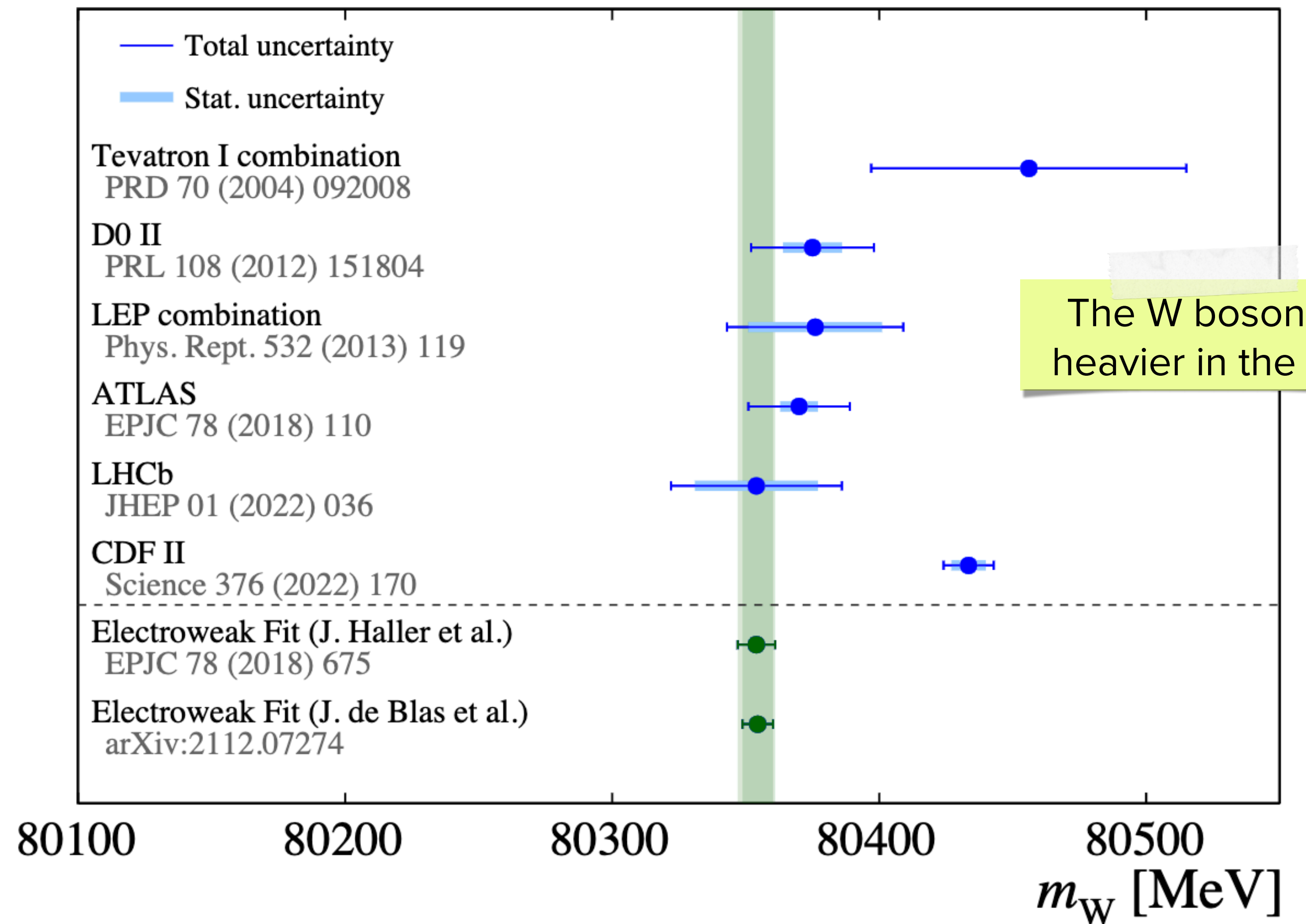
🍏🍏🍏 Simple signatures, large cross sections, simplified models



Time

MORE TO INVESTIGATE

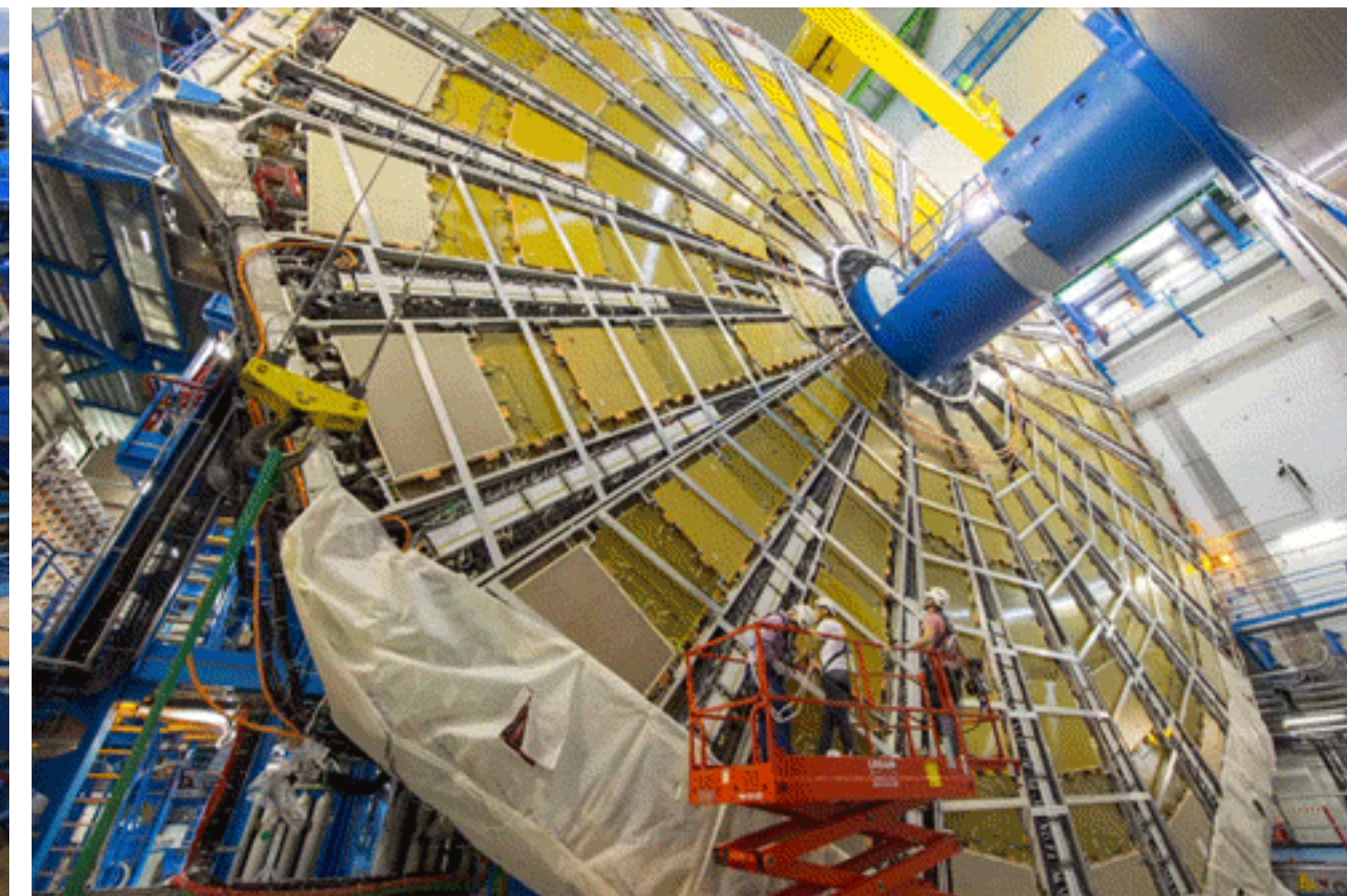
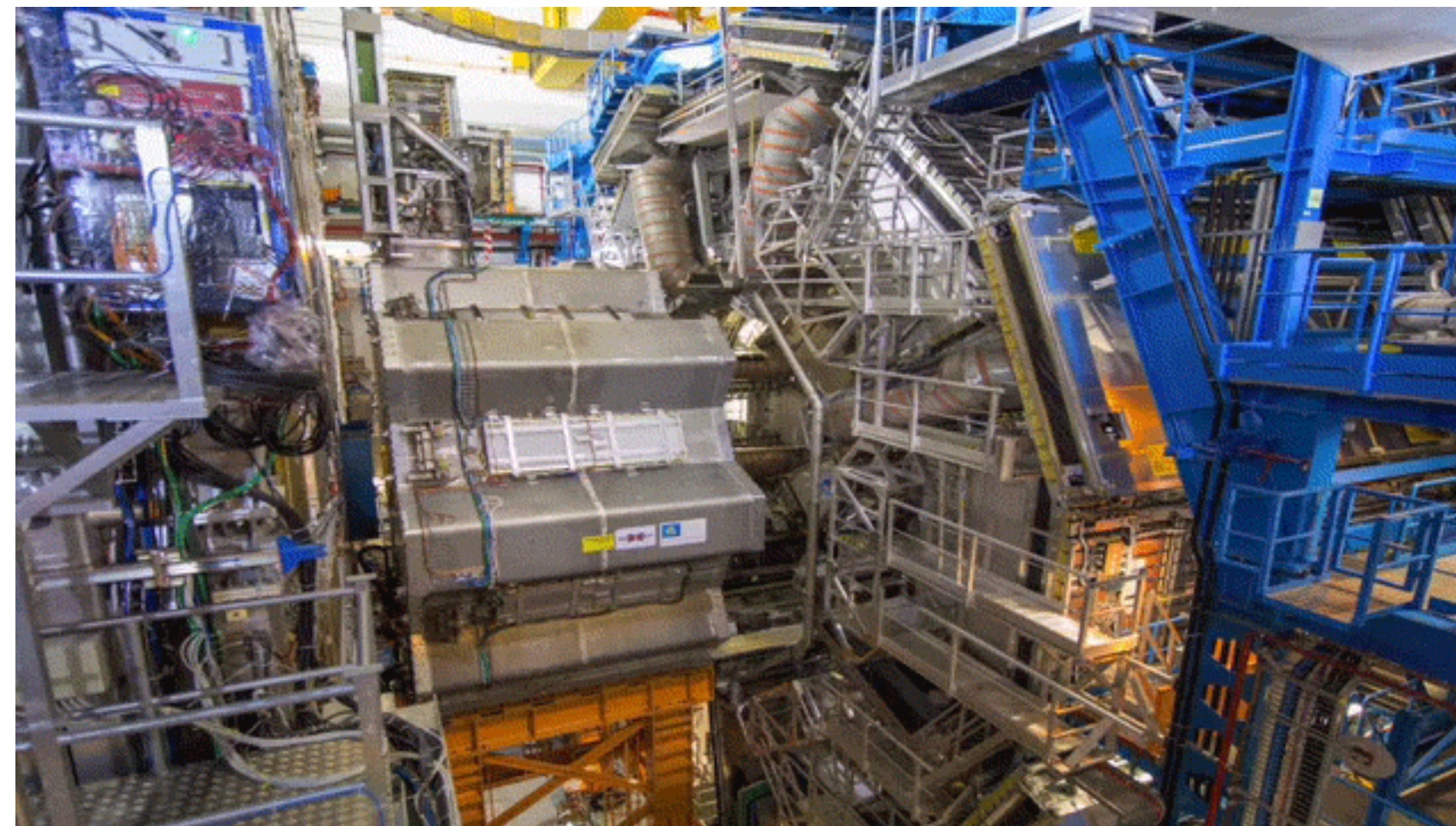
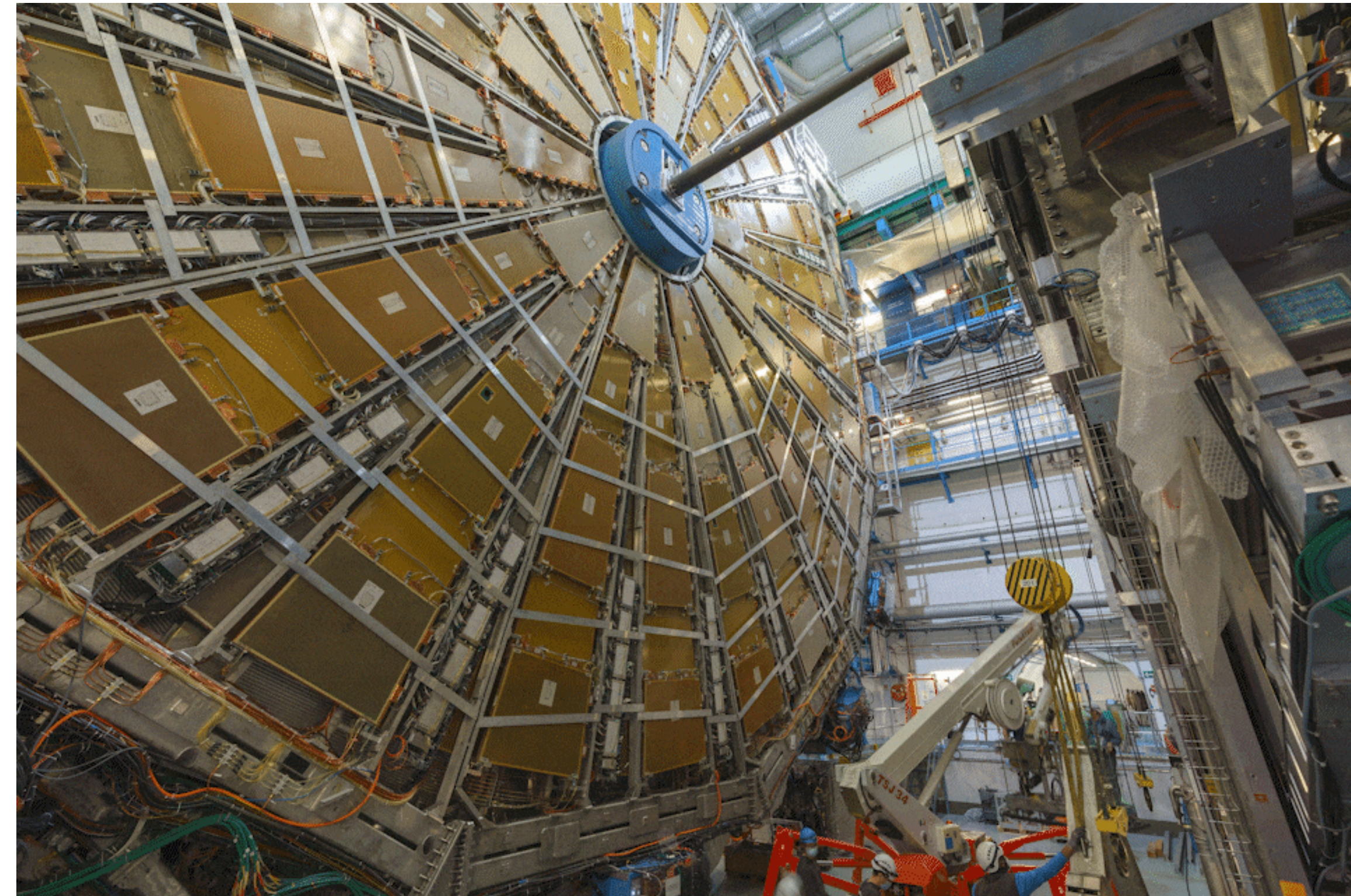
- Not only we have the general SM shortcomings we already discussed but we have additional interesting effects to investigate



RUN-3

FEELS LIKE RUN 3 JUST STARTED

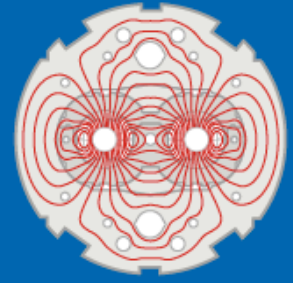
- Run-3 started after the second long shutdown (LS 2) in 2022
- Most ATLAS and CMS analysis use Run-2 data



BUT RUN-3 IS THE LAST RUN OF THE LHC AS WE KNOW IT

- The goal for Run-3 is 250 fb^{-1} at 13.6 TeV
- Run-3 will finish next year (or maybe the one after?)
- Will be followed by a Long shut-down (LS3)
- To fully upgrade the LHC and the experiments
 - To the High-Luminosity LHC

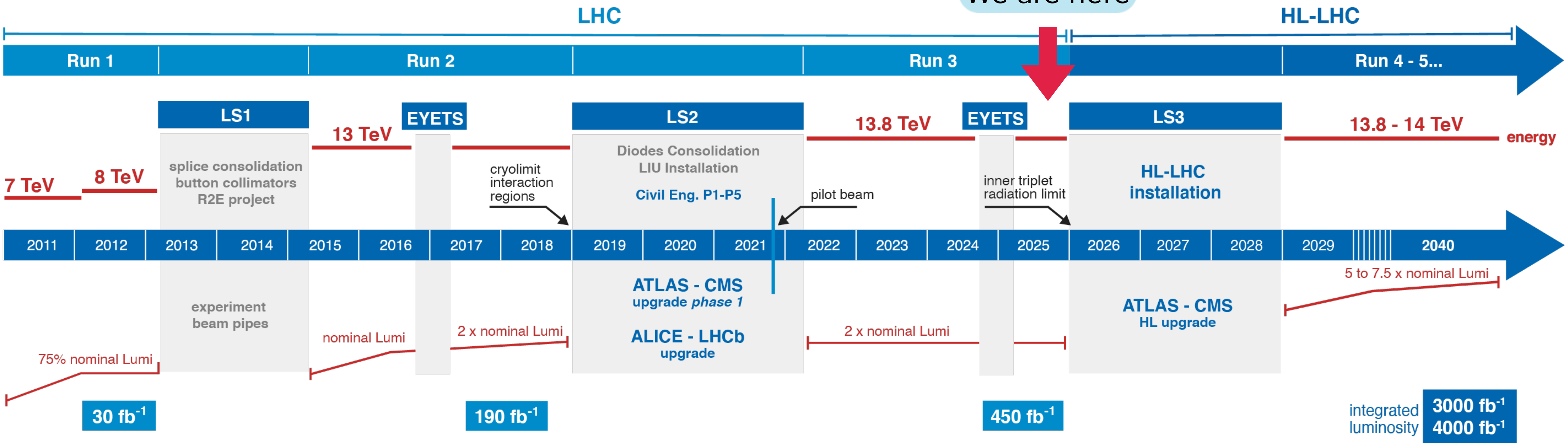




LHC / HL-LHC Plan



We are here

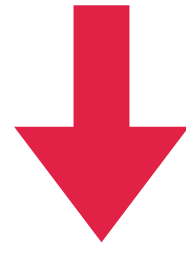


HL-LHC TECHNICAL EQUIPMENT:



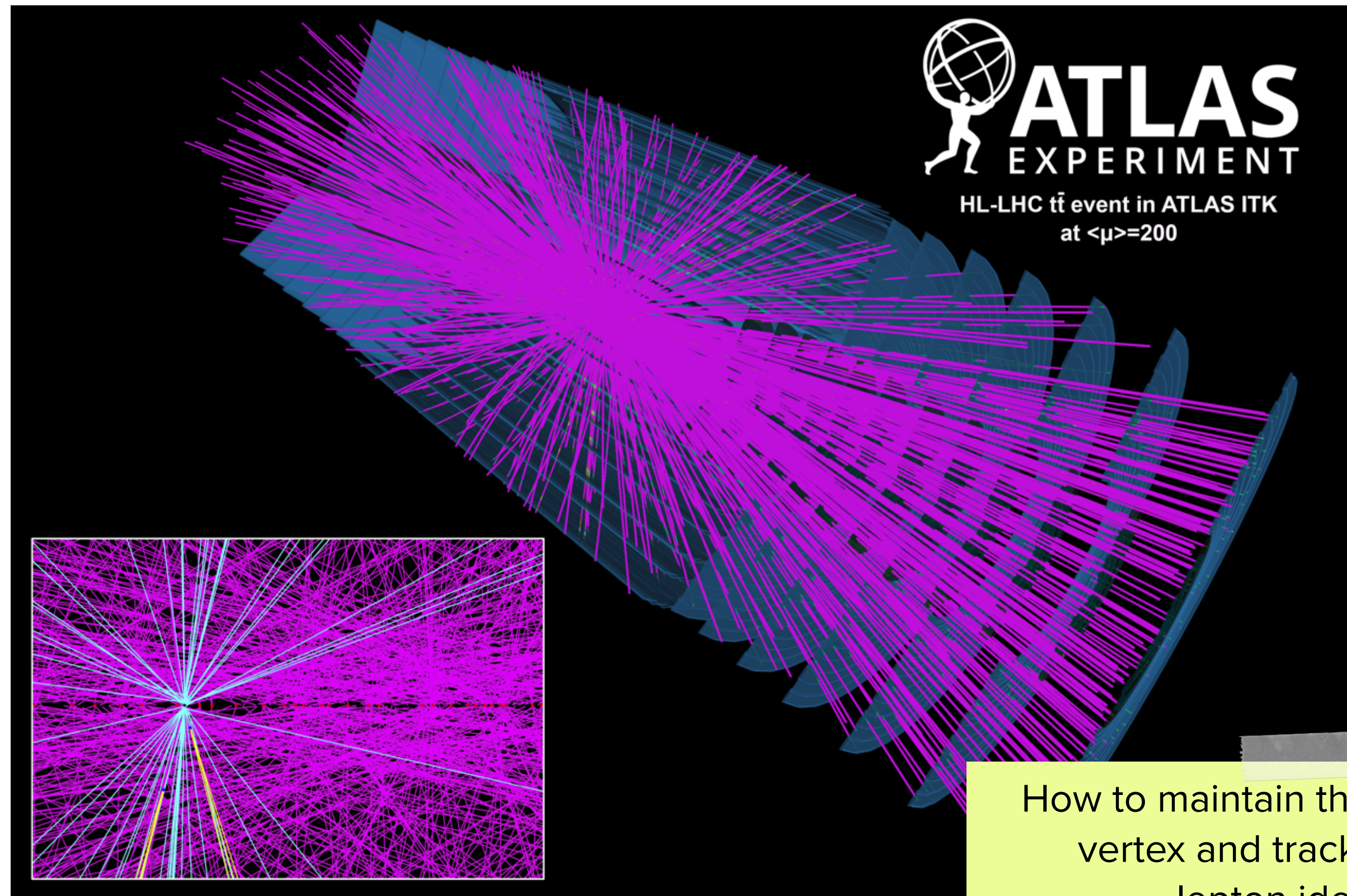
HL-LHC CIVIL ENGINEERING:





Last update: April 2023

THE HL-LHC WILL BE TOUGH FOR THE DETECTORS



How to maintain the performance in
vertex and track reconstruction?
lepton identification?
heavy flavor tagging?

EVERY EXPERIMENT HAS A BATTERY OF UPGRADES

- E.g.: Three major detector upgrades for ATLAS
 - complete replacement of the inner tracking system
 - new radiation-tolerant read-out electronics for the tracking, calorimeter and muon systems
 - Trigger and Data Acquisition (TDAQ) architecture



THE DATA OF THE HL-LHC

- More than 10 extra years of running:
 - Targeting 3ab^{-1}
 - 3000fb^{-1} → to be compared with about 300fb^{-1} from Runs 1 to 3
 - 10x more events
 - Same center of mass energy, 13.6 TeV

0.5M Higgs bosons per experiment in Run 1
8M in Run 2
10M expected in Run 3
200M in HL-LHC

The High-Luminosity LHC

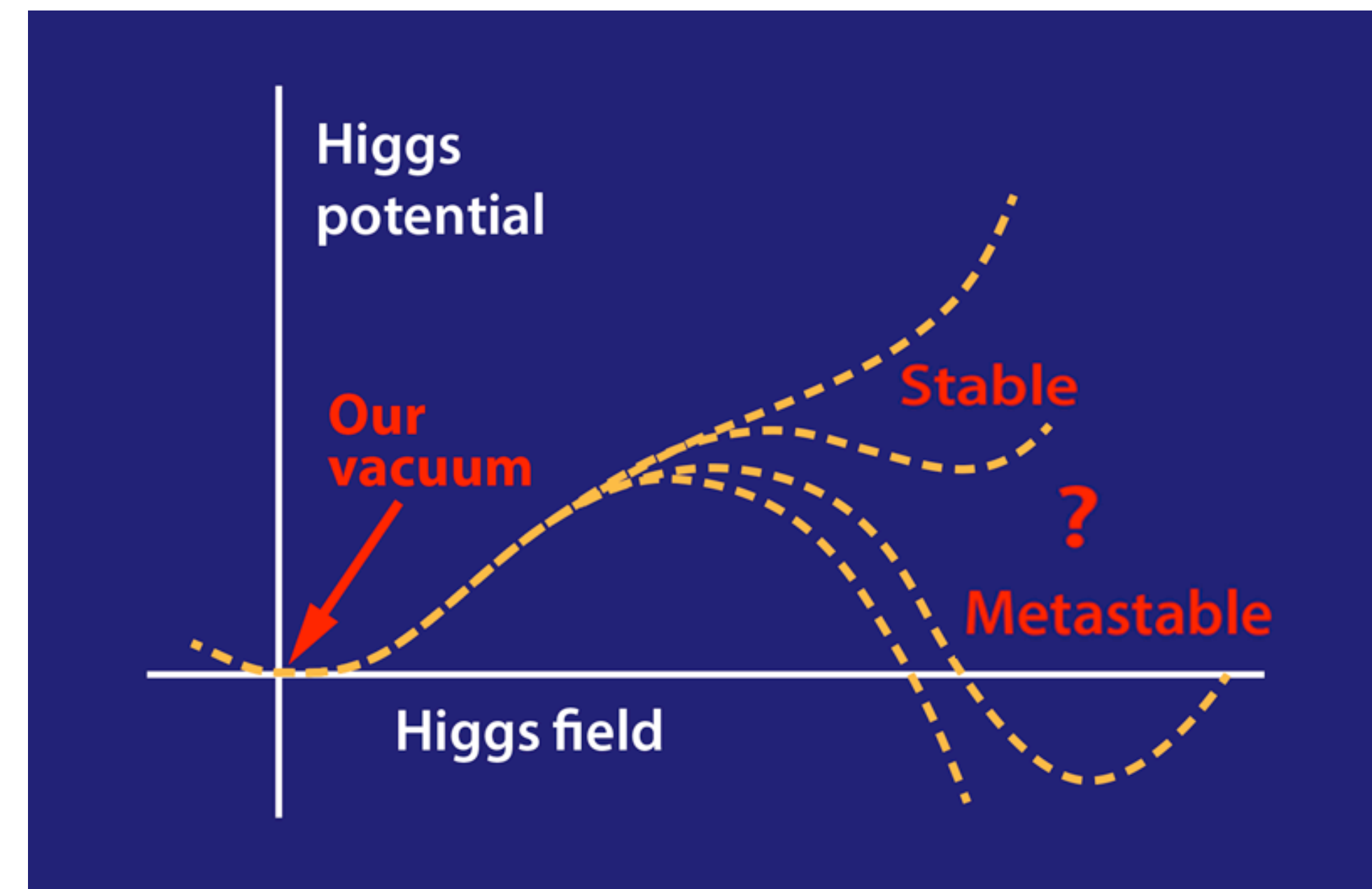
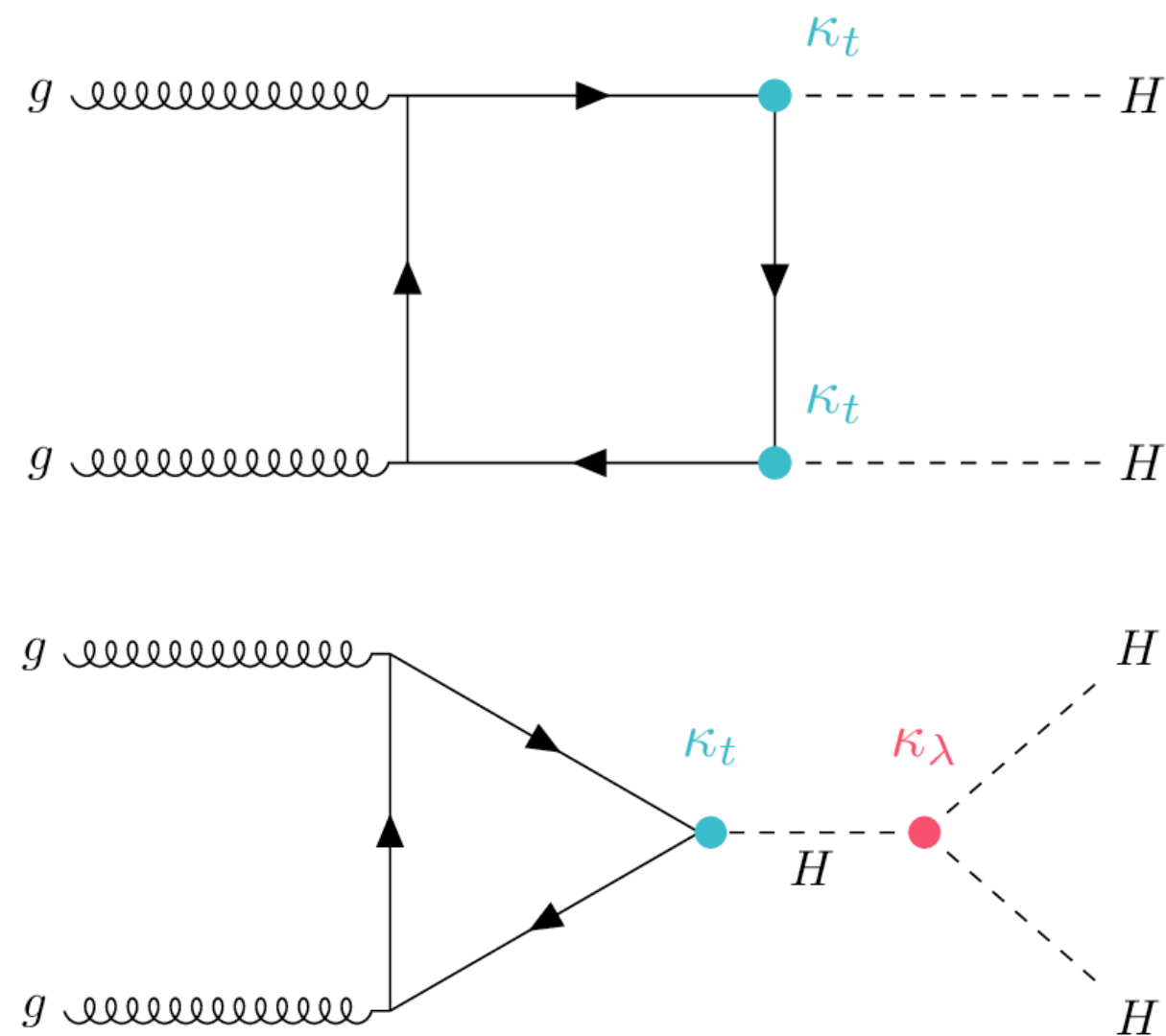




THE PHYSICS CASE

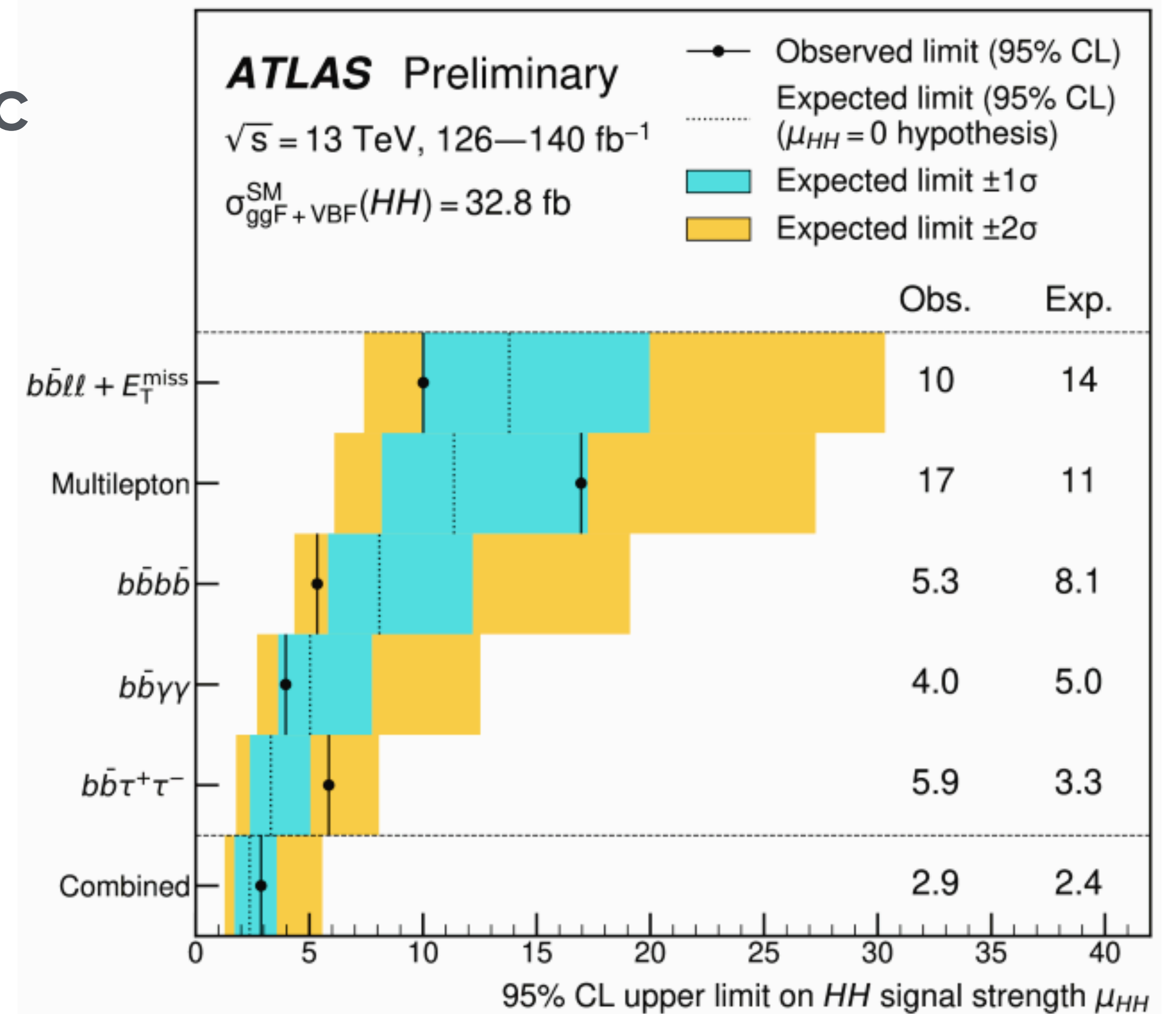
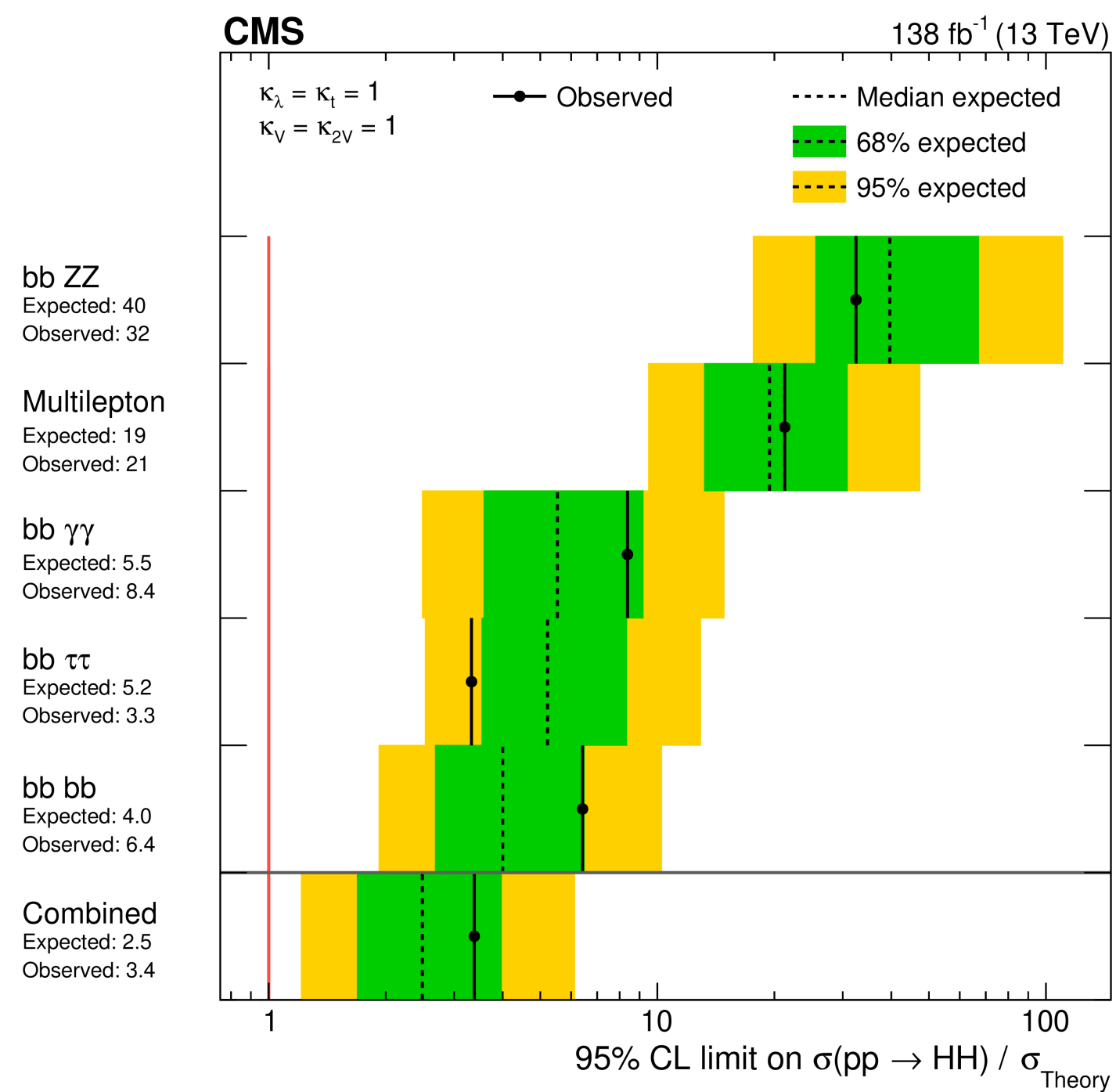
OUR FIRST SHOT TO UNDERSTAND DOUBLE HIGGS PRODUCTION

- All Higgs properties are interesting
- But di-Higgs production is particularly interesting
 - Shape of the Higgs potential (cosmological implications)



DI-HIGGS IS ONE OF THOSE THINGS

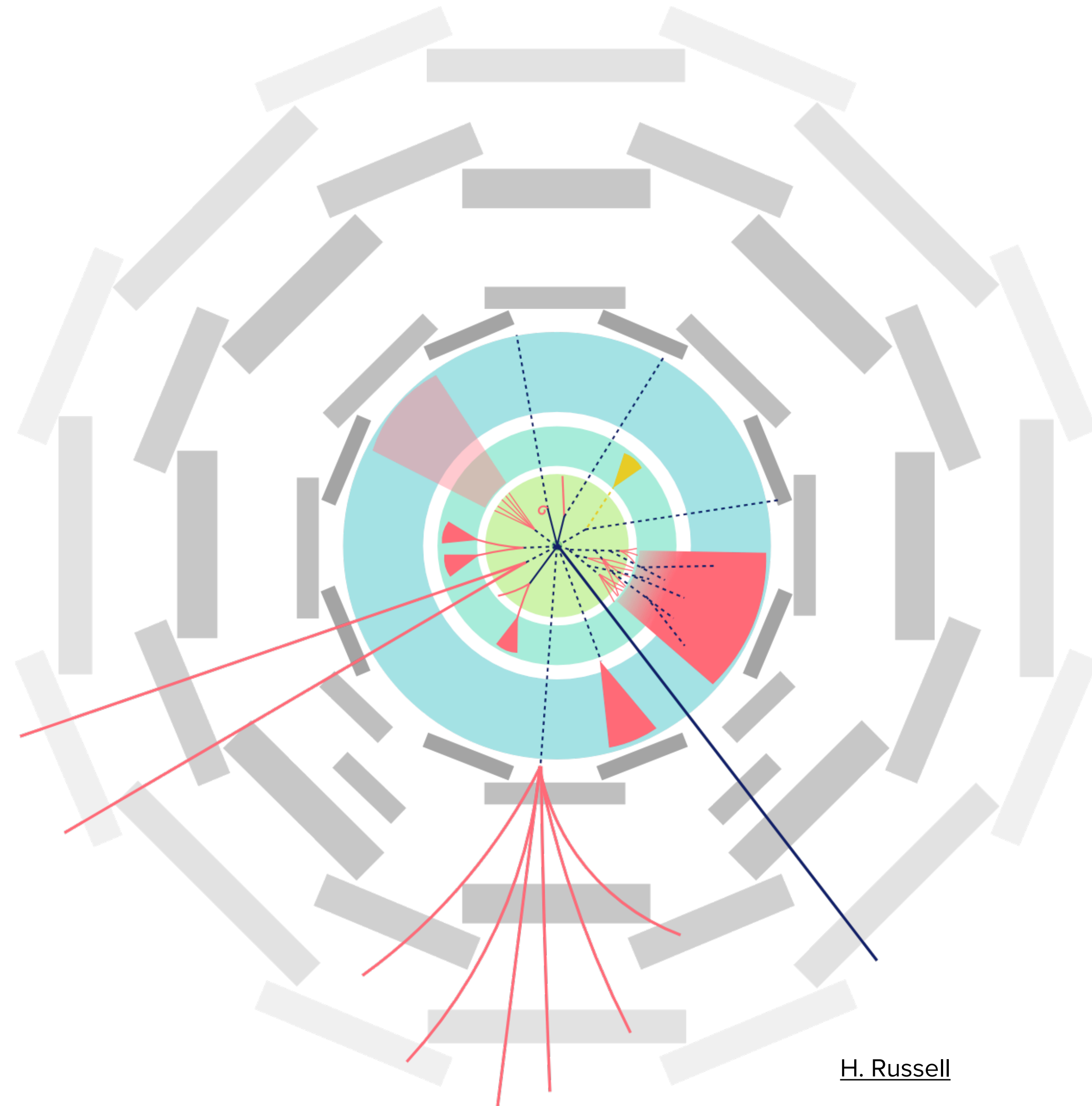
- That we thought was not possible at the LHC
- But we are getting closer every day



UNEXPLORED AREAS

- At the HL-LHC we will have more data
 - The opportunity to directly discover new particles and phenomena will increase
 - Especially background 0 searches

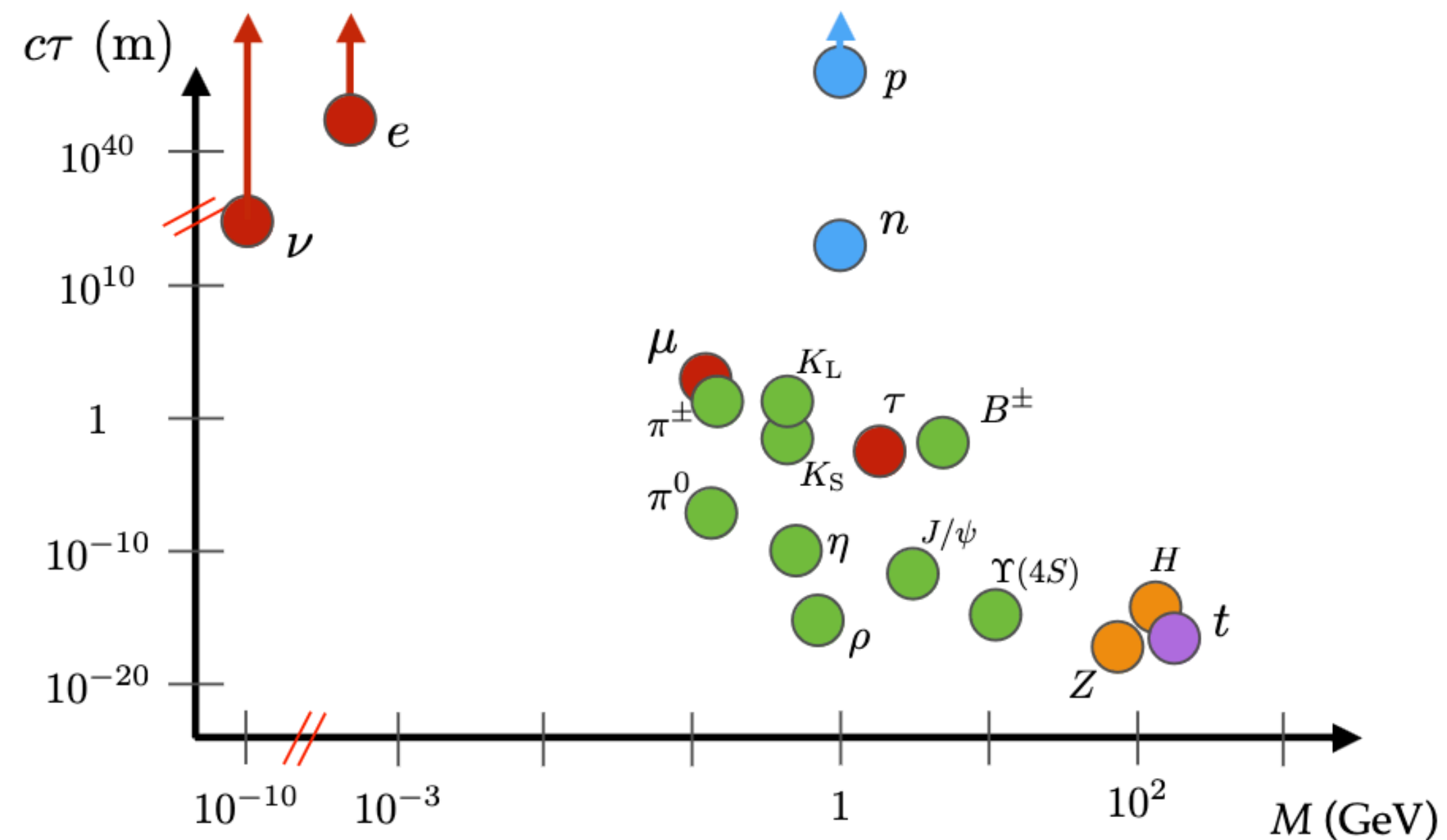
My favorite kind are searches for long-lived particles



H. Russell

ANOTHER DETOUR: LONG-LIVED PARTICLES

LONG-LIVED PARTICLES

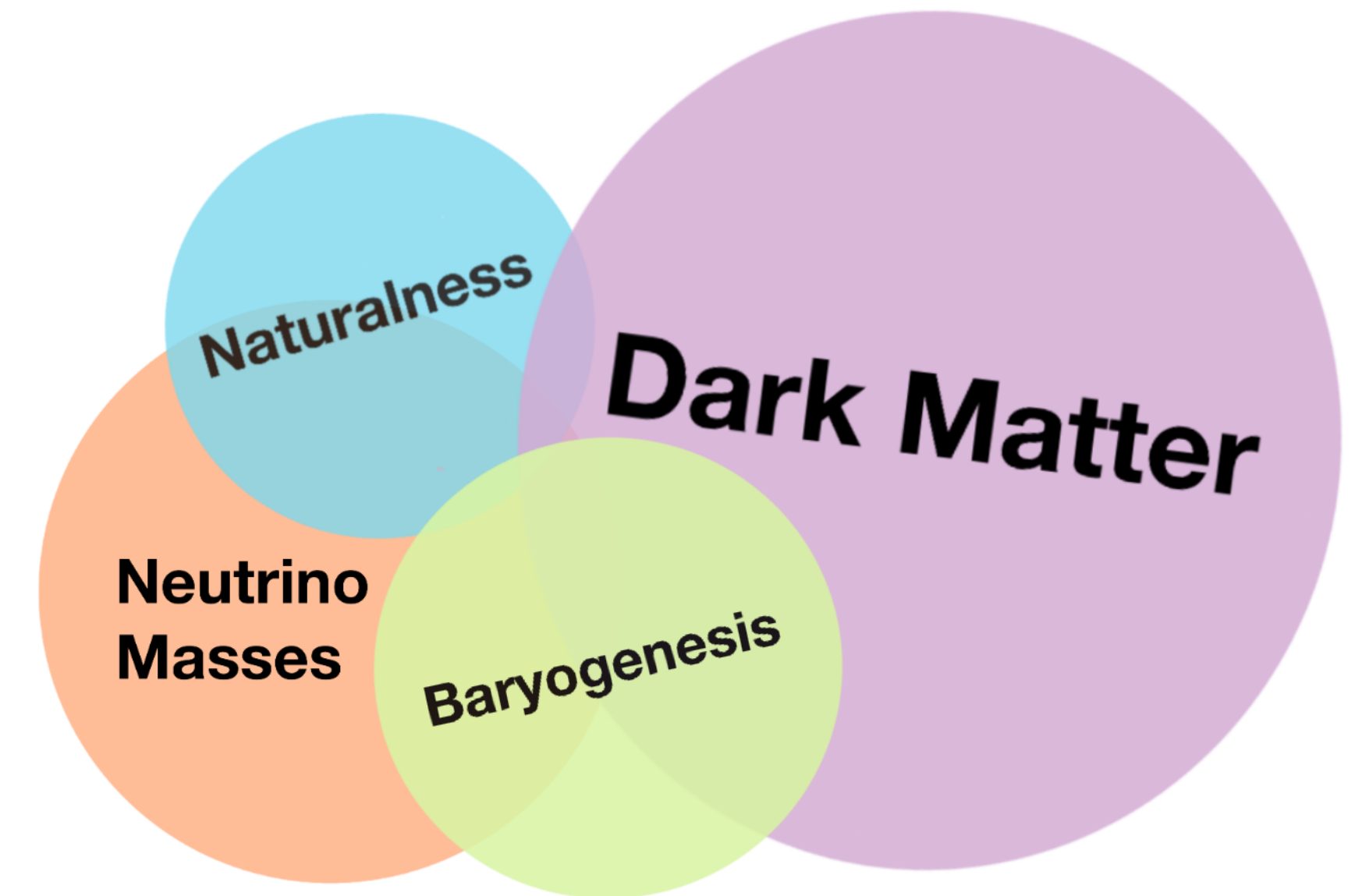



[arXiv:1903.04497](https://arxiv.org/abs/1903.04497)

- SM particles all have different lifetimes, even with similar masses
- Many of them are long-lived
 - Due to e.g. small couplings or a suppressed decay phase space
- But we use Long-lived particles (LLPs) as an umbrella term
 - New particles, that we have not discovered yet, with lifetimes long enough to travel measurable distances inside the detectors before decaying

WHERE DO WE GET LLPS?

- LLPs are a generic signature of BSM physics, connected to central questions
 - R parity violating (RPV) and conserving (RPC) SUSY
 - Heavy Neutral leptons (right handed/sterile neutrinos)
 - Exotic Higgs decays
 - or new scalars, e.g. dark photon or Axion-Like Particles (ALPs)...
- In general, LLPs feature extensively in hidden sectors
- If light (<1 GeV) new particles exist, they must be very weakly coupled \rightarrow LLPs



A bright sun is centered in the upper half of the frame, surrounded by a soft, glowing aura. The sky is a uniform, warm orange color. Below the sun, a thick layer of white and orange-tinted clouds stretches across the horizon. The overall scene is hazy and atmospheric.

**LLP searches have been going on for years at colliders in different ways
Looking for them is nothing new**

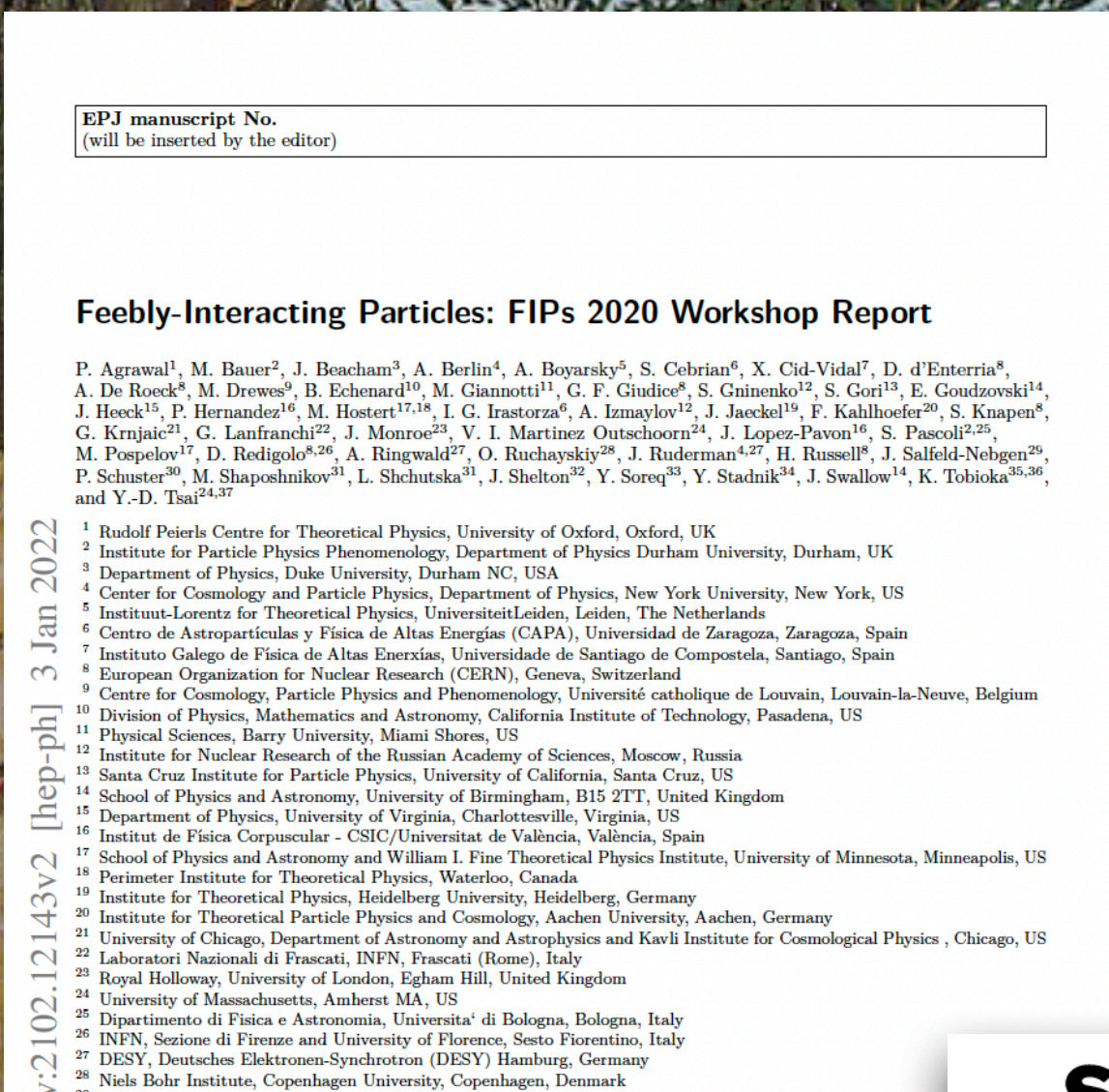
But one could say that Long-Lived Particles are living a Renaissance

arXiv:1806.07396v2 [hep-ph] 5 Mar 2019

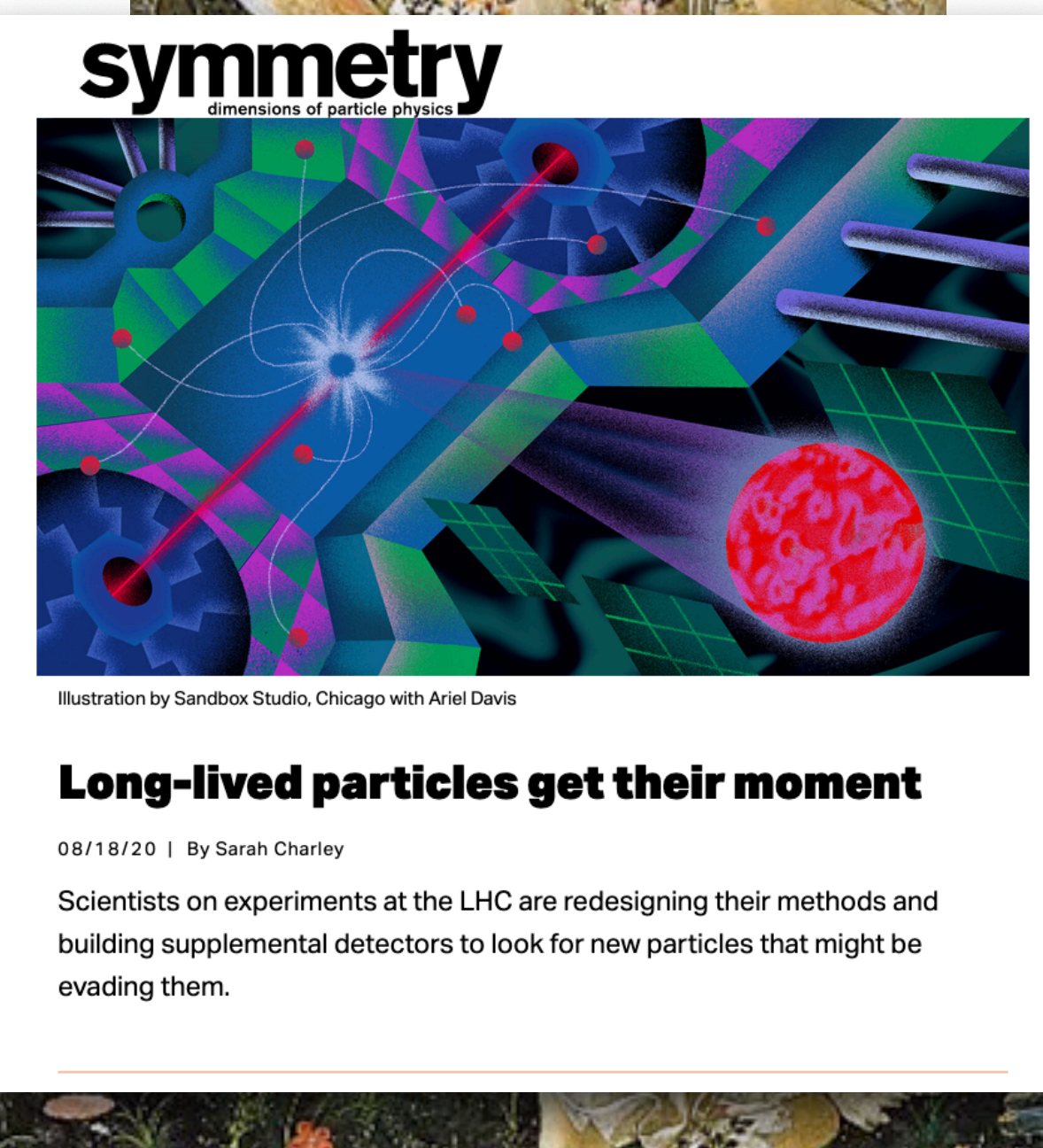


21 Department of Physics, University of Toronto, Toronto, ON M5S 1A7, Canada
 22 Centre for Cosmology, Particle Physics and Phenomenology, Université catholique de Louvain, Louvain-la-Neuve, B-1348, Belgium
 23 CERN, TH Department, CH-1211 Geneva, Switzerland
 24 C.N. Yang Institute for Theoretical Physics, Stony Brook University, Stony Brook, NY 11794, USA
 25 Maryland Center for Fundamental Physics, MD 20742-4111 USA
 26 Department of Physics, University of Illinois at Chicago, Chicago, IL 60607-7199, USA
 27 Harvey Mudd College, 301 Platt Blvd., Claremont, CA 91711, USA
 28 SLAC National Accelerator Laboratory, Menlo Park, CA 94025, USA

matches for “long-lived” in the indico agenda
 ICHEP 2012: 11
 ICHEP 2022: 53



31 Institute of Physics, University of Oxford, Oxford, UK
 32 Department of Physics, Durham University, Durham, UK
 33 Department of Physics, Department of Physics, New York University, New York, NY, USA
 34 Center for Cosmology and Particle Physics, Universität Leiden, Leiden, The Netherlands
 35 Institut-Lorentz for Theoretical Physics, Universiteit Leiden, Leiden, The Netherlands
 36 Centro de Astroparticulas y Física de Altas Energías (CAPA), Universidad de Zaragoza, Zaragoza, Spain
 37 Instituto Galego de Física de Altas Enerxías, Universidade de Compostela, Santiago, Spain
 38 European Organization for Nuclear Research (CERN), Geneva, Switzerland
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 45 Department of Physics, University of Virginia, Charlottesville, VA, USA
 46 Perimeter Institute for Theoretical Physics, Waterloo, ON, Canada
 47 Institut de Física Corpuscular - CSIC/Universitat de València, València, Spain
 48 School of Physics and Astronomy and William I. Fine Theoretical Physics Institute, University of Minnesota, Minneapolis, MN, USA
 49 Perimeter Institute for Theoretical Physics, Heidelberg, Germany
 50 Institute for Theoretical Particle Physics and Cosmology, Aachen University, Aachen, Germany
 51 University of Chicago, Department of Astronomy and Astrophysics and Kavli Institute for Cosmological Physics, Chicago, IL, USA
 52 Laboratori Nazionali di Frascati, INFN, Frascati (Rome), Italy
 53 Royal Holloway, University of London, Egham Hill, United Kingdom
 54 University of Massachusetts, Amherst MA, USA
 55 Dipartimento di Fisica e Astronomia, Università di Bologna, Bologna, Italy
 56 INFN, Sezione di Firenze and University of Florence, Sesto Fiorentino, Italy
 57 DESY, Deutsches Elektronen-Synchrotron (DESY) Hamburg, Germany
 58 Niels Bohr Institute, Copenhagen University, Copenhagen, Denmark
 59 Department of Physics, Princeton University, Princeton, NJ, USA
 60 SLAC National Accelerator Laboratory Menlo Park, CA, USA
 61 Institute of Physics, Ecole Polytechnique Fédérale de Lausanne (EPFL), Lausanne, Switzerland
 62 University of Illinois, Urbana, IL, USA
 63 Physics Department, Technion, Institute of Technology, Haifa 3200003, Israel
 64 Kavli Institute for the Physics and Mathematics of the Universe (KIPMU), University of Tokyo, Japan
 65 Department of Physics, Florida State University, Tallahassee, FL, USA
 66 High Energy Accelerator Research Organization (KEK), Tsukuba, Japan
 67 Kavli Institute for Cosmological Physics, University of Chicago, Chicago, IL, USA



Search for long-lived particles beyond the Standard Model at the Large Hadron Collider
 March 6, 2019
 Particles beyond the Standard Model (SM) can generically have lifetimes that are long compared to SM particles at the weak scale. When produced at experiments such as the Large Hadron Collider (LHC) at CERN, these long-lived particles (LLPs) can decay far from the interaction vertex of the primary proton-proton collision. Such LLP signatures are distinct from those of promptly decaying particles that are targeted by the majority of searches for new physics at the LHC, often requiring customized techniques to identify, for example, significantly displaced decay vertices, tracks with atypical properties, and short track segments. Given their non-standard nature, a comprehensive overview of LLP signatures at the LHC is beneficial to ensure that possible avenues of the discovery of new physics are not overlooked. Here we report on the joint work of a community of theorists and experimentalists with the ATLAS, CMS, and LHCb experiments — as well as those working on dedicated experiments such as MoEDAL, milliQan, MATHUSLA, CODEX-b, and FASER — to survey the current state of LLP searches at the LHC, and to chart a path for the development of LLP searches into the future, both in the upcoming Run 3 and at the High-Luminosity LHC. The work is organized around the current and future potential capabilities of LHC experiments to generally discover new LLPs, and takes a signature-based approach to surveying classes of models that give rise to LLPs rather than emphasizing any particular theory motivation. We develop a set of simplified models; assess the coverage of current searches; document known, often unexpected backgrounds; explore the capabilities of proposed detector upgrades; provide recommendations for the presentation of search results; and look towards the newest frontiers, namely high-multiplicity “dark showers”, highlighting opportunities for expanding the LHC reach for these signals.
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WHY IS THIS HAPPENING?

- There are a few reasons for why LLPs are so interesting nowadays:
 - Searches for LLPs cover intermediate areas, gaps of sensitivity between experiments (eg. dark matter searches between colliders and astro)
 - They address the lack of prompt BSM signals → providing accessible new areas where BSM could be hiding
 - LLP searches offer us the opportunity to think outside the box, to be creative and to propose new ways to solve problems
 - **Innovation**: in methods and experimental setups

AT HIGHER ENERGIES

- We gain access to more massive particles that in turn tend to be shorter-lived

Main offenders



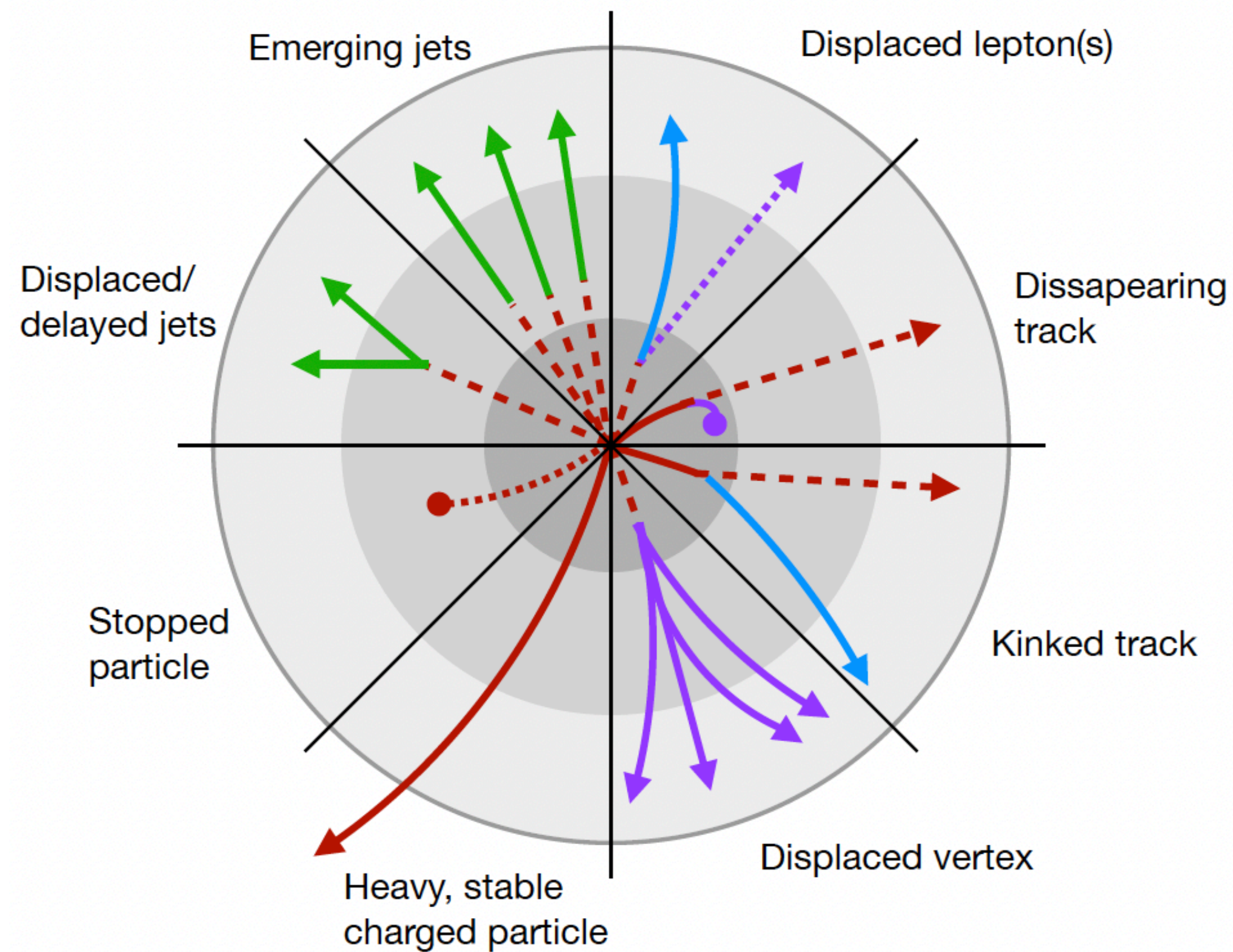
The Higgs boson
2012 - LHC
Sort-Lived
 10^{-22} seconds
You blink and you miss it!



The top quark
1995 - Tevatron
Sort-Lived
 10^{-25} seconds
So short-lived it does not even have
time to form hadrons!

- And we naturally optimize our detectors, trigger, and reconstruction methods to find them
 - LLPs could be regularly produced in collisions and we wouldn't know it
- LLPs produce unconventional signatures in colliders
 - clearly different from other processes (easy to spot!), but potentially invisible to current data-acquisition methods → we could be throwing them away

NON-STANDARD EXPERIMENTAL SIGNATURES



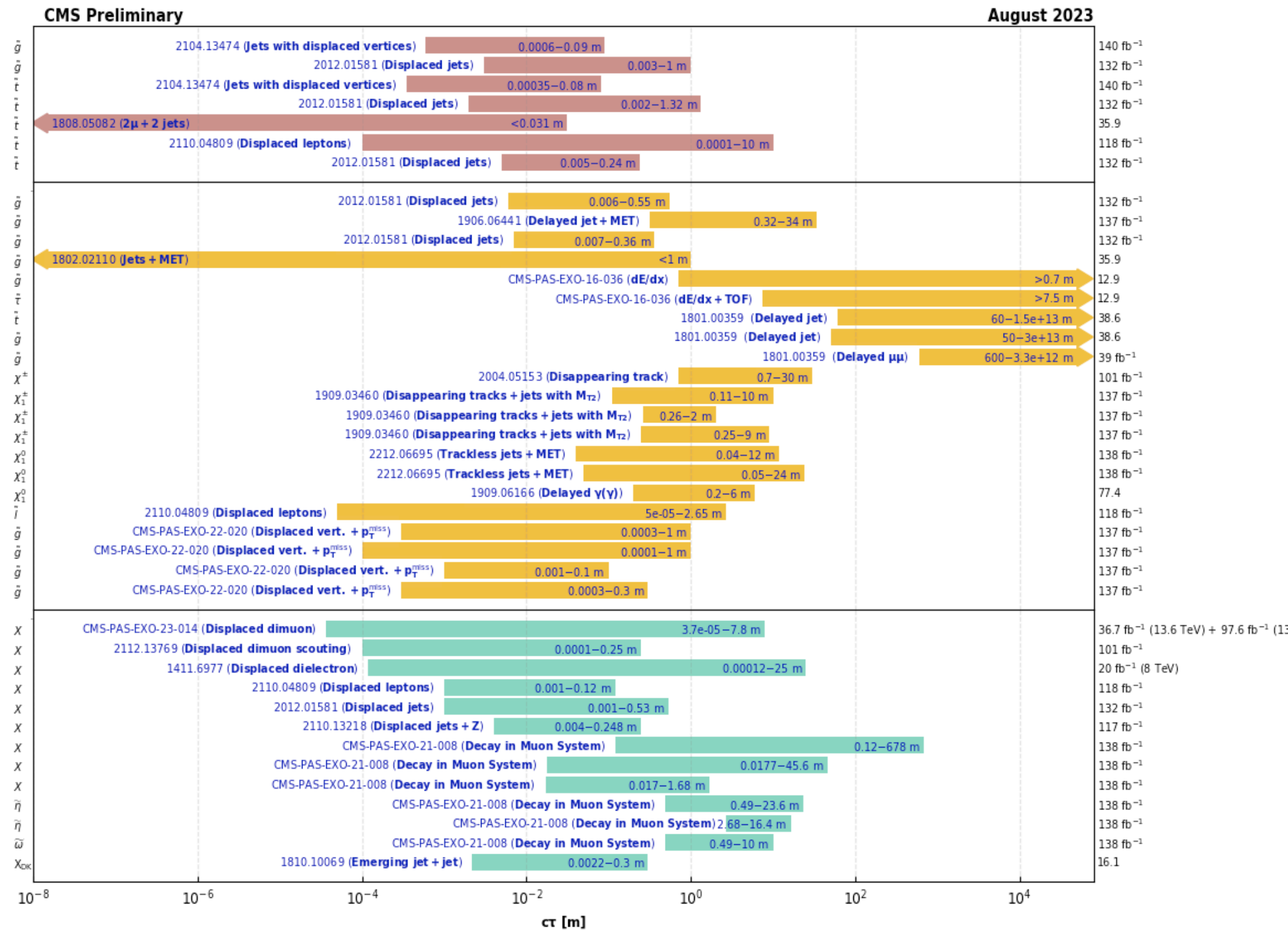
- | | | | |
|---------|---------|---|----------|
| — | Charged | ■ | BSM |
| - - - - | Neutral | ■ | Lepton |
| | Either | ■ | Quark |
| | | ■ | Anything |

- We are talking about
 - displaced and/or delayed objects (leptons, photons, jets); disappearing tracks; nonstandard tracks produced by monopoles, quirks or heavy stable charged particles (HSCPs); nonstandard jets produced in dark showers...
- LLP analyses at the LHC experiments:
 - require customisation: dedicated triggers, object reconstruction, background estimation and in general analysis methods
 - are affected by challenging backgrounds near the collision points → motivate dedicated experiments

ATLAS AND CMS

Vibrant scene of long-lived searches in the exotics and SUSY groups

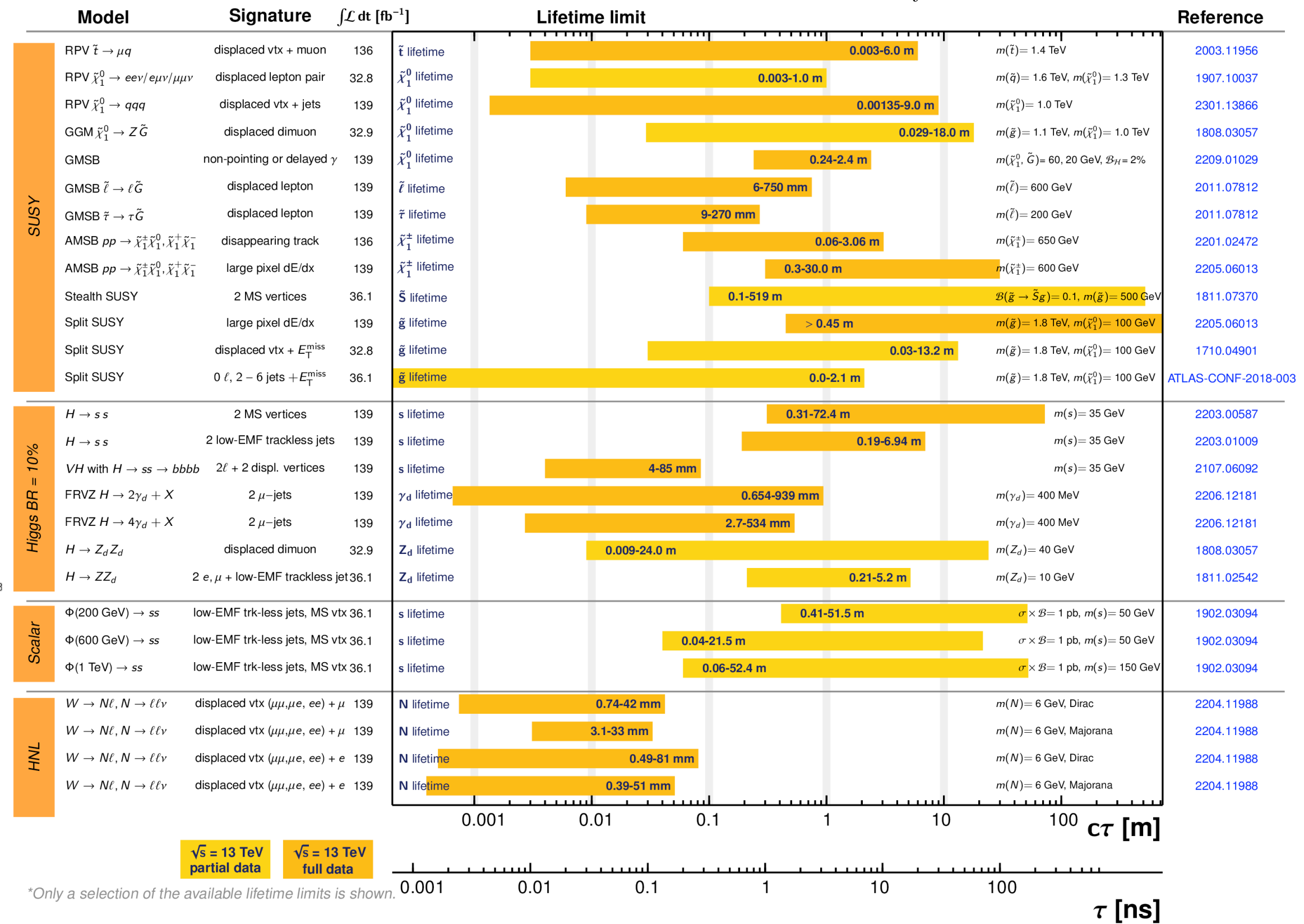
Overview of CMS long-lived particle searches



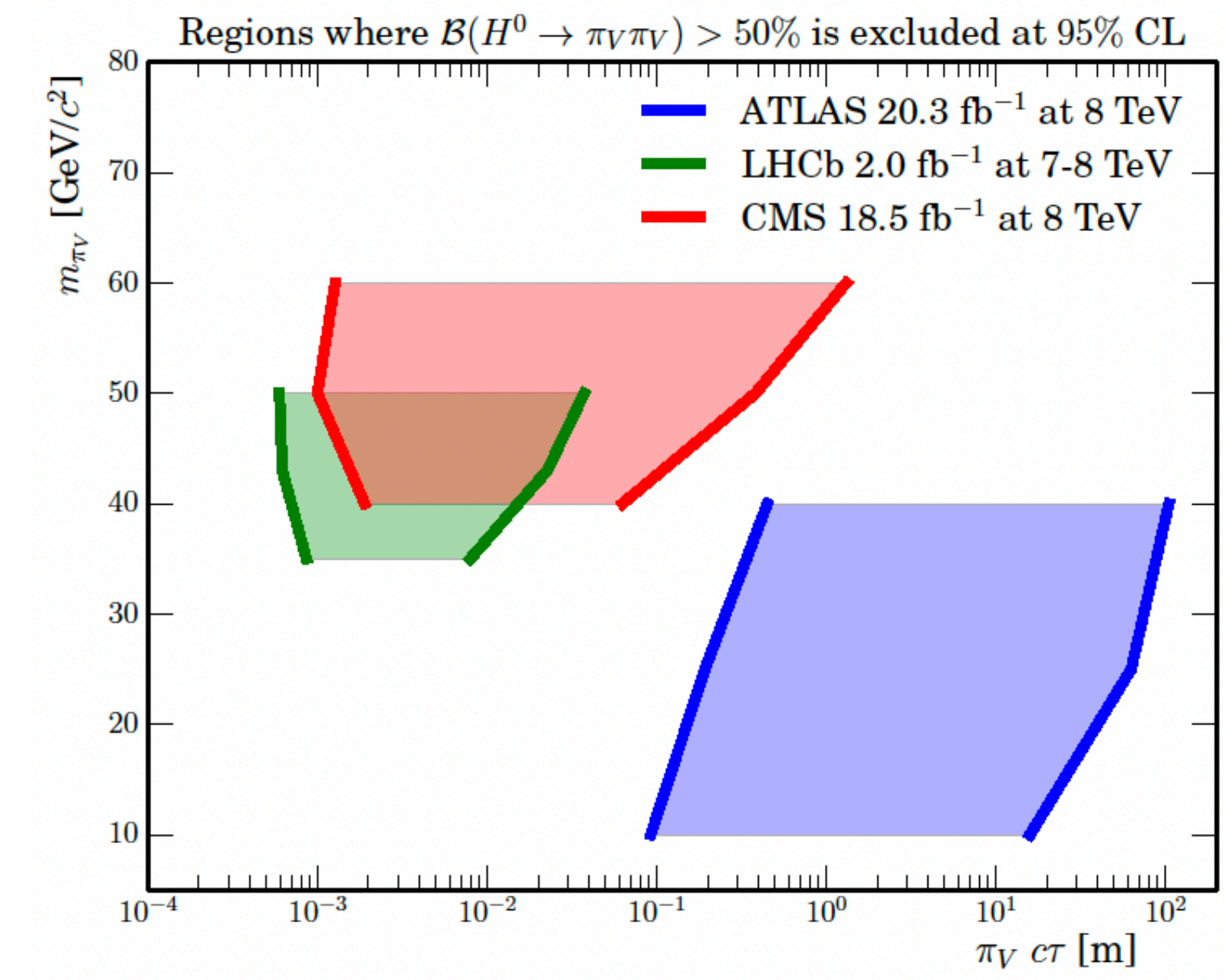
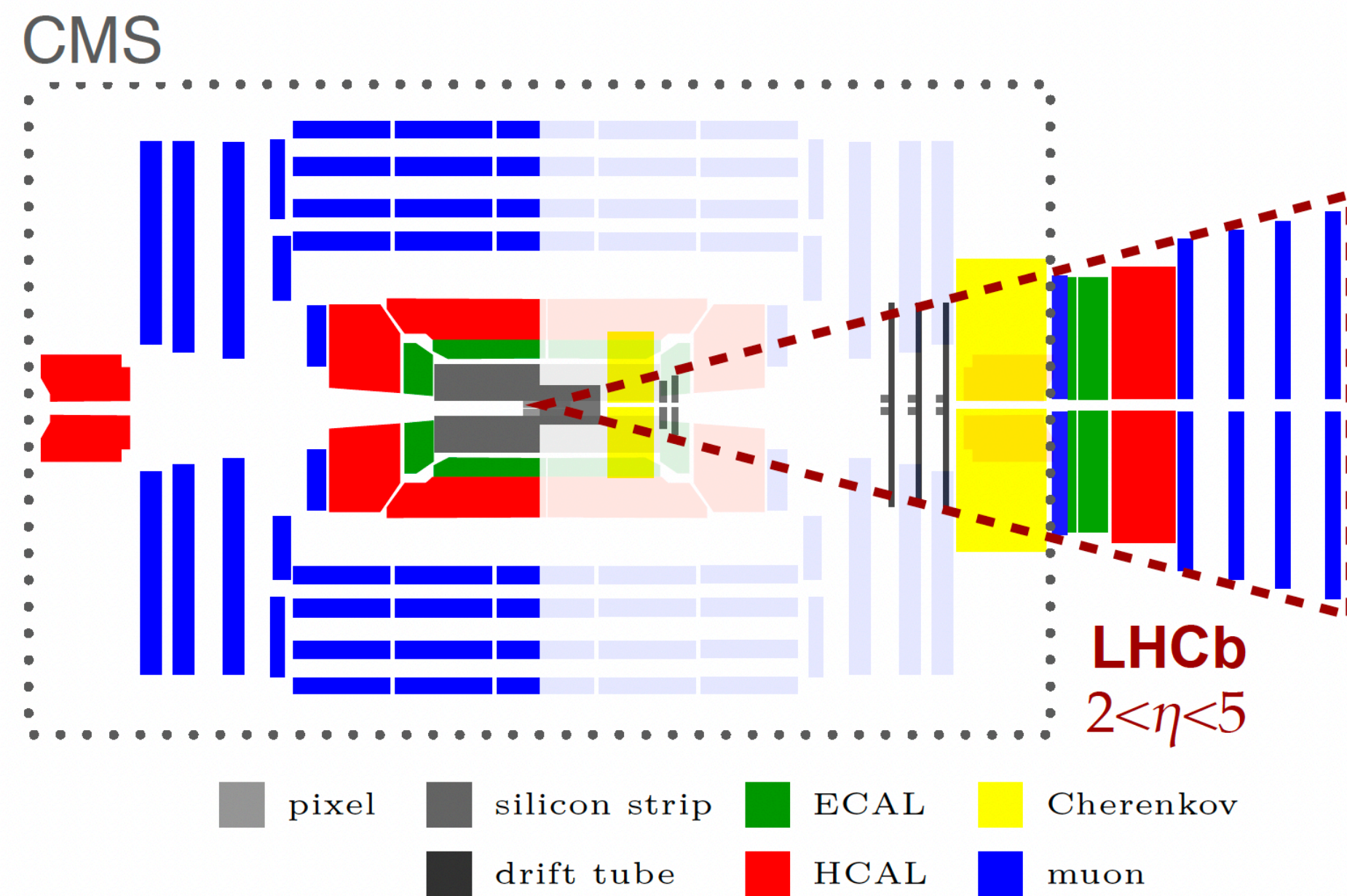
ATLAS Long-lived Particle Searches* - 95% CL Exclusion

Status: March 2023

ATLAS Preliminary
 $\int \mathcal{L} dt = (32.8 - 139) \text{ fb}^{-1}$
 $\sqrt{s} = 13 \text{ TeV}$

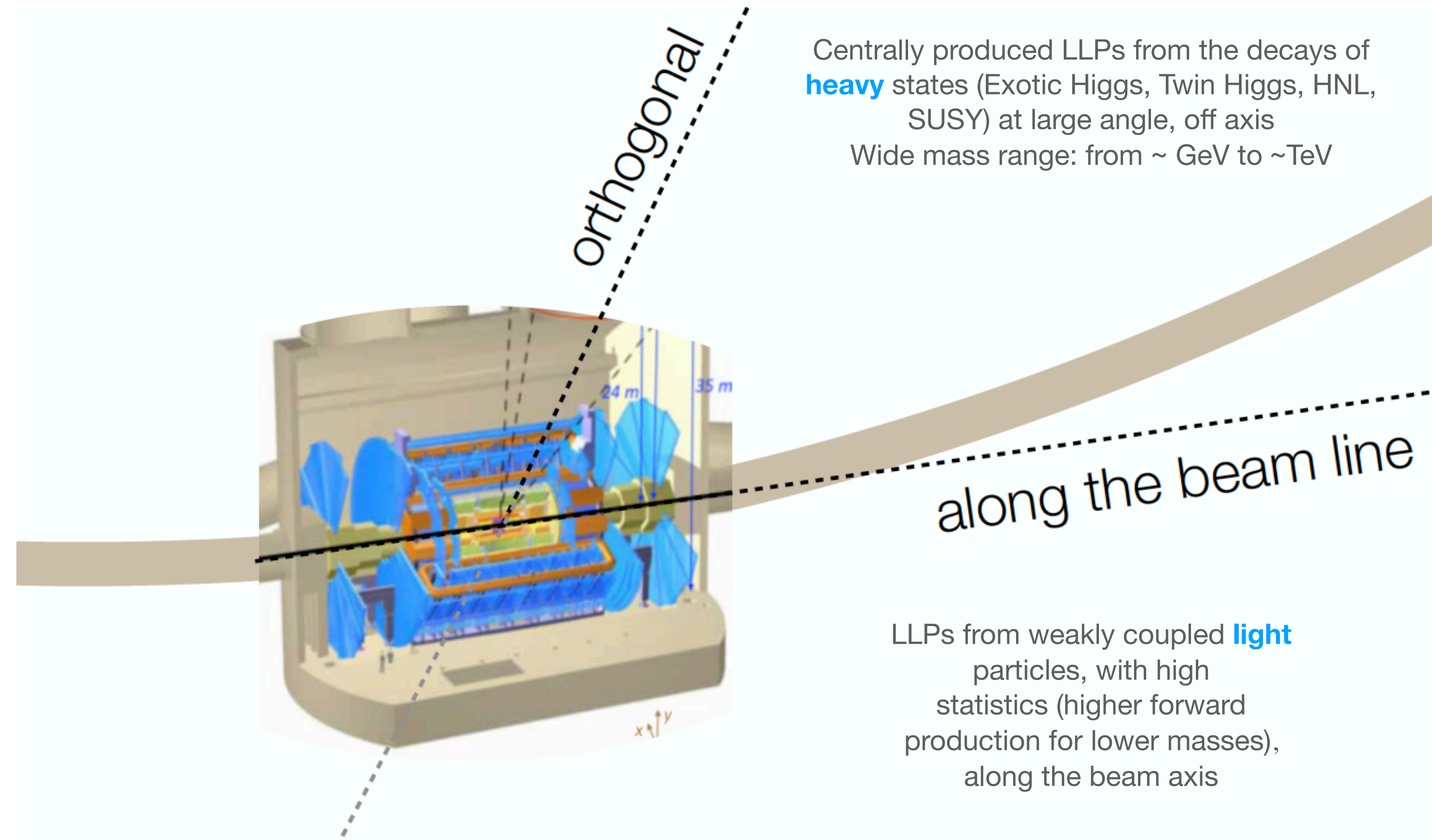
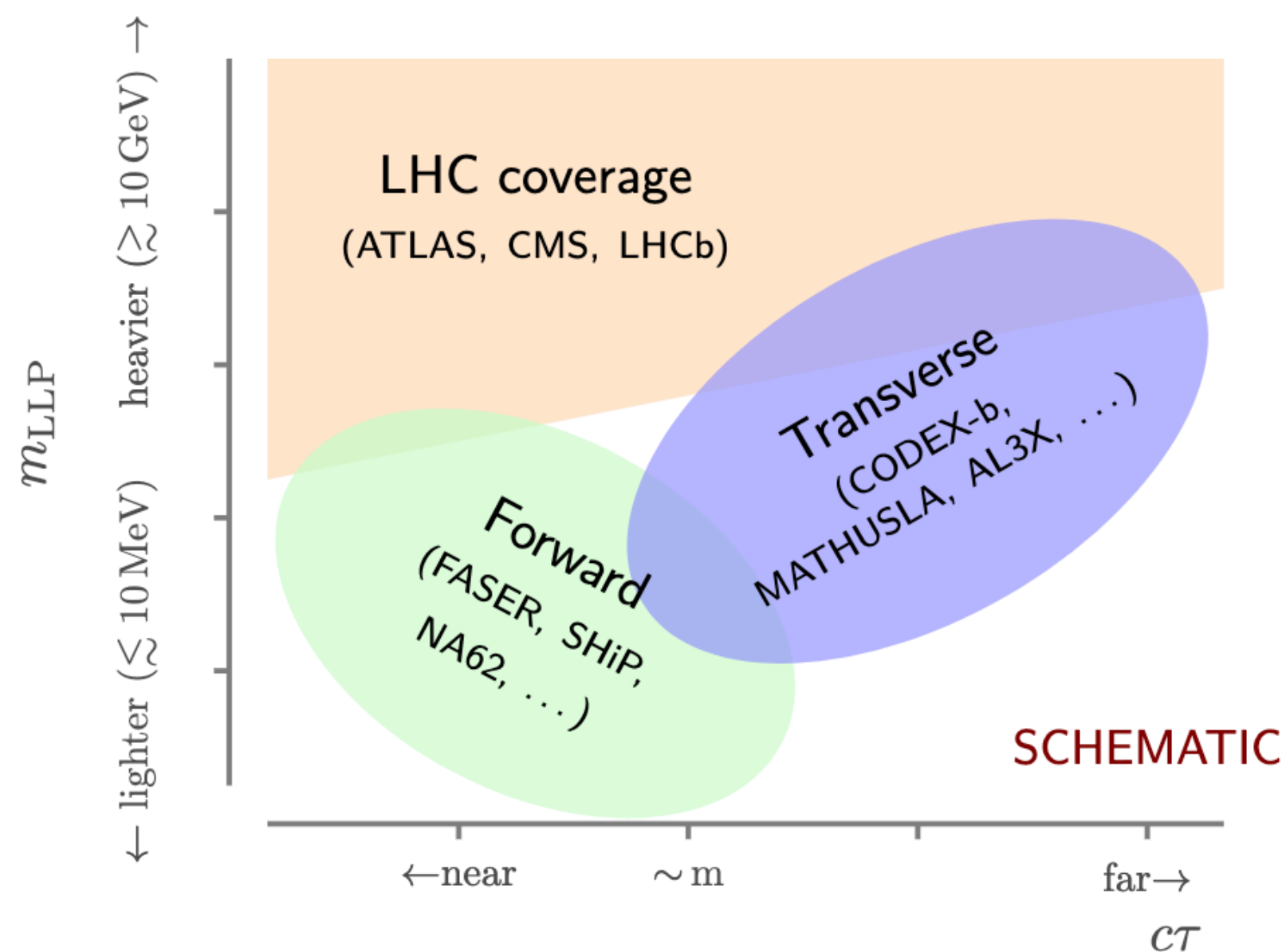


- Existing results on dark photons, exotic Higgs decays, HNLs...
- Complementary coverage to ATLAS and CMS
- First fully GPU trigger in HEP opens new possibility for LLPs



THINKING OUTSIDE THE LHC DETECTORS

- We can supplement them with external detectors optimized for LLPs
 - Access to longer decay lengths
 - Less background (shielding)
 - Easy trigger (or trigger-less)



Centrally produced LLPs from the decays of **heavy** states (Exotic Higgs, Twin Higgs, HNL, SUSY) at large angle, off axis
Wide mass range: from $\sim \text{GeV}$ to $\sim \text{TeV}$

LLPs from weakly coupled **light** particles, with high statistics (higher forward production for lower masses), along the beam axis

THERE ARE MANY EXPERIMENTS

Running

FASER

MoEDAL

NA62

SND@LHC
Scattering and Neutrino Detector at the LHC

Proposed

MATISSE A

SHiP
Search for Hidden Particles

CODEX-b

FPF
forward physics facility

AL3X

ANUBIS

SHADOWS

FACET

FORMOSA

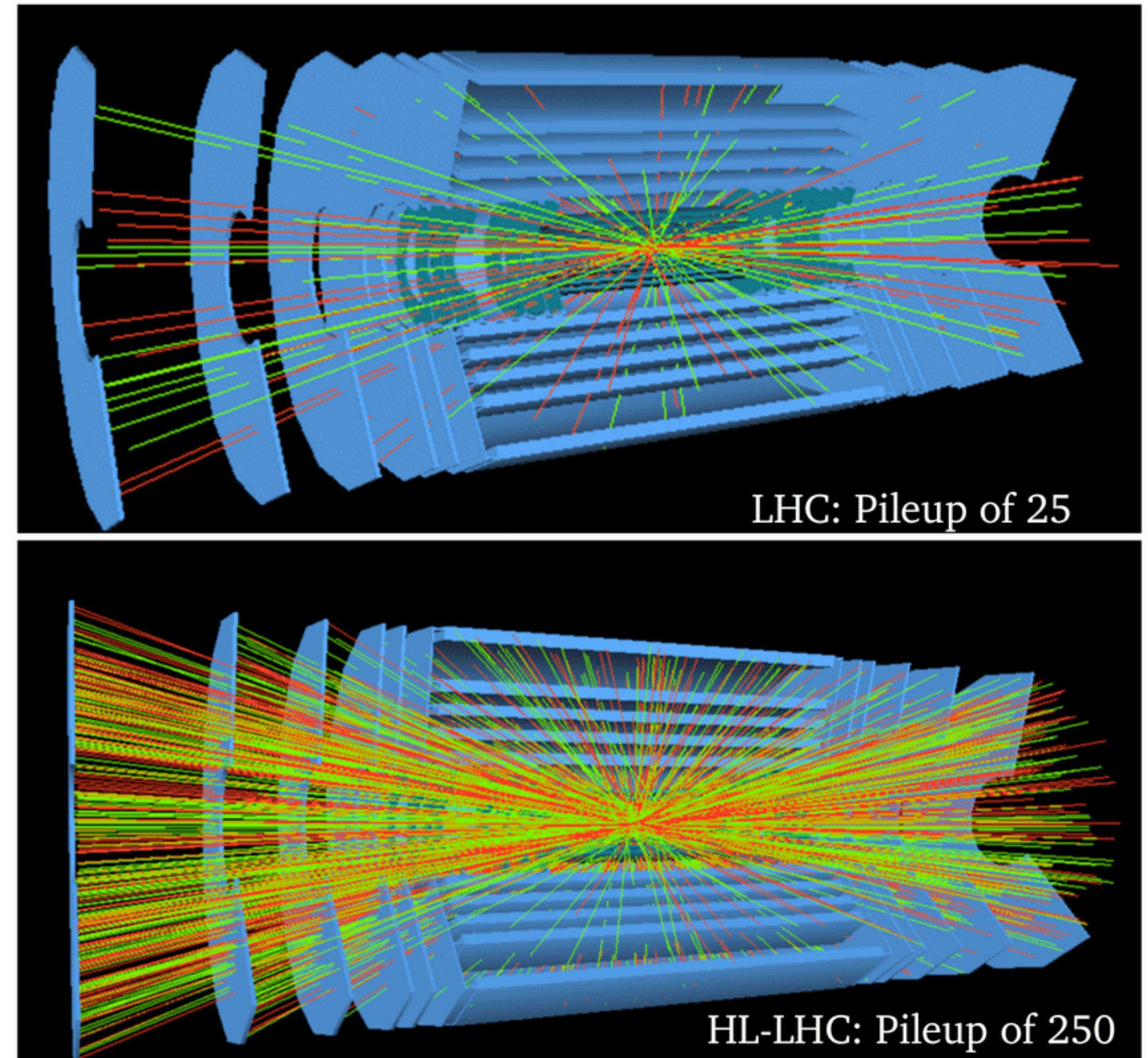
Your experiment here



THE HL-LHC WON'T BE THE END OF THE STORY

- At the HL-LHC we will have more data
 - We will keep probing SM processes, with higher precision
 - Higher chance to find deviations with respect to the predictions
- Most of our current analyses are NOT statistically limited
- No increase of the center of mass energy, HL-LHC will stay at 13.6 TeV
- The pileup will be brutal

There is potential for a surprise but a surprise will be surprising



<https://arxiv.org/abs/1201.5469>

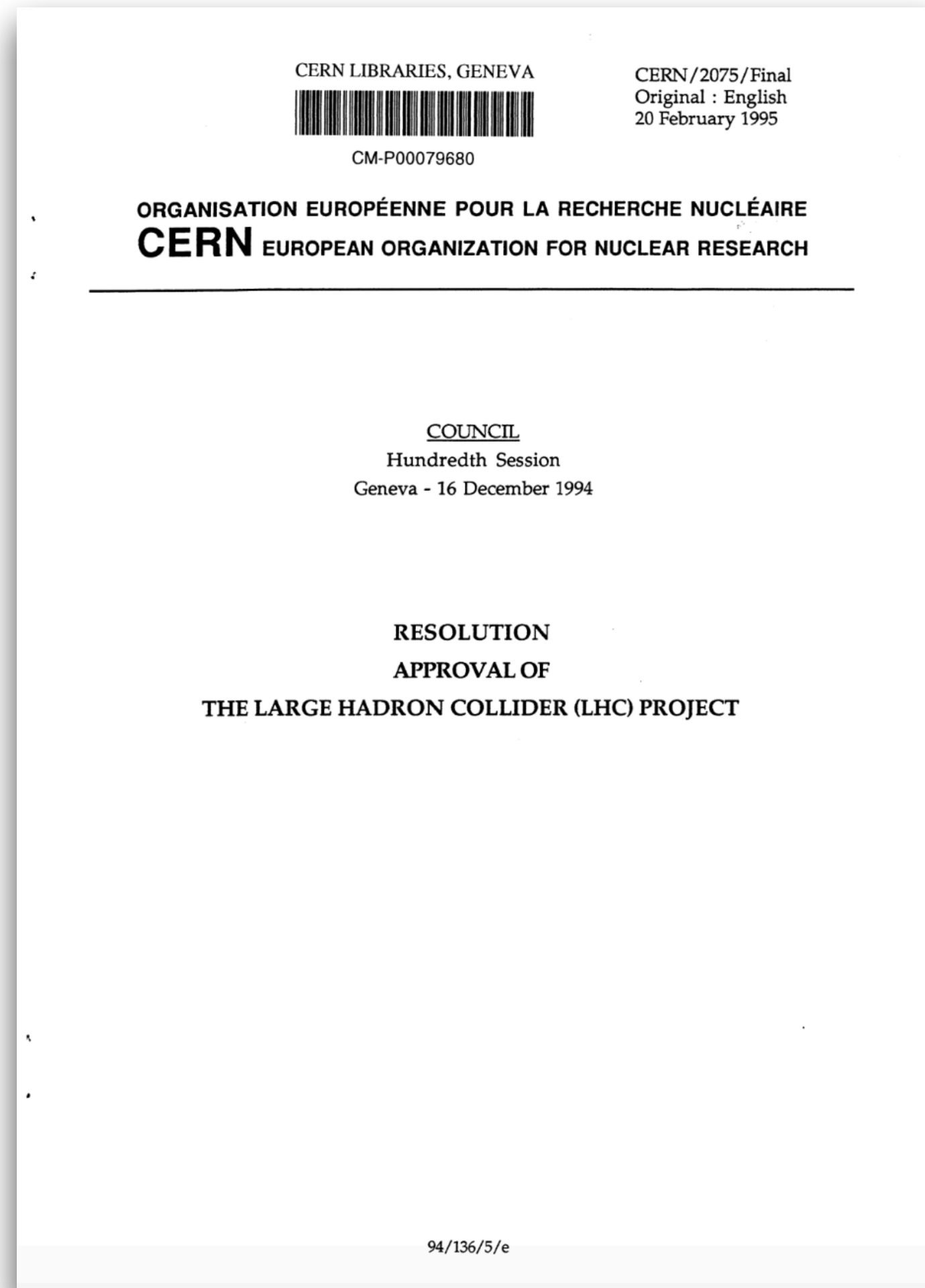
SO, WHAT'S NEXT?

WELL, WE DON'T KNOW

BUT THE CYCLES OF COLLIDER PHYSICS ARE LONG

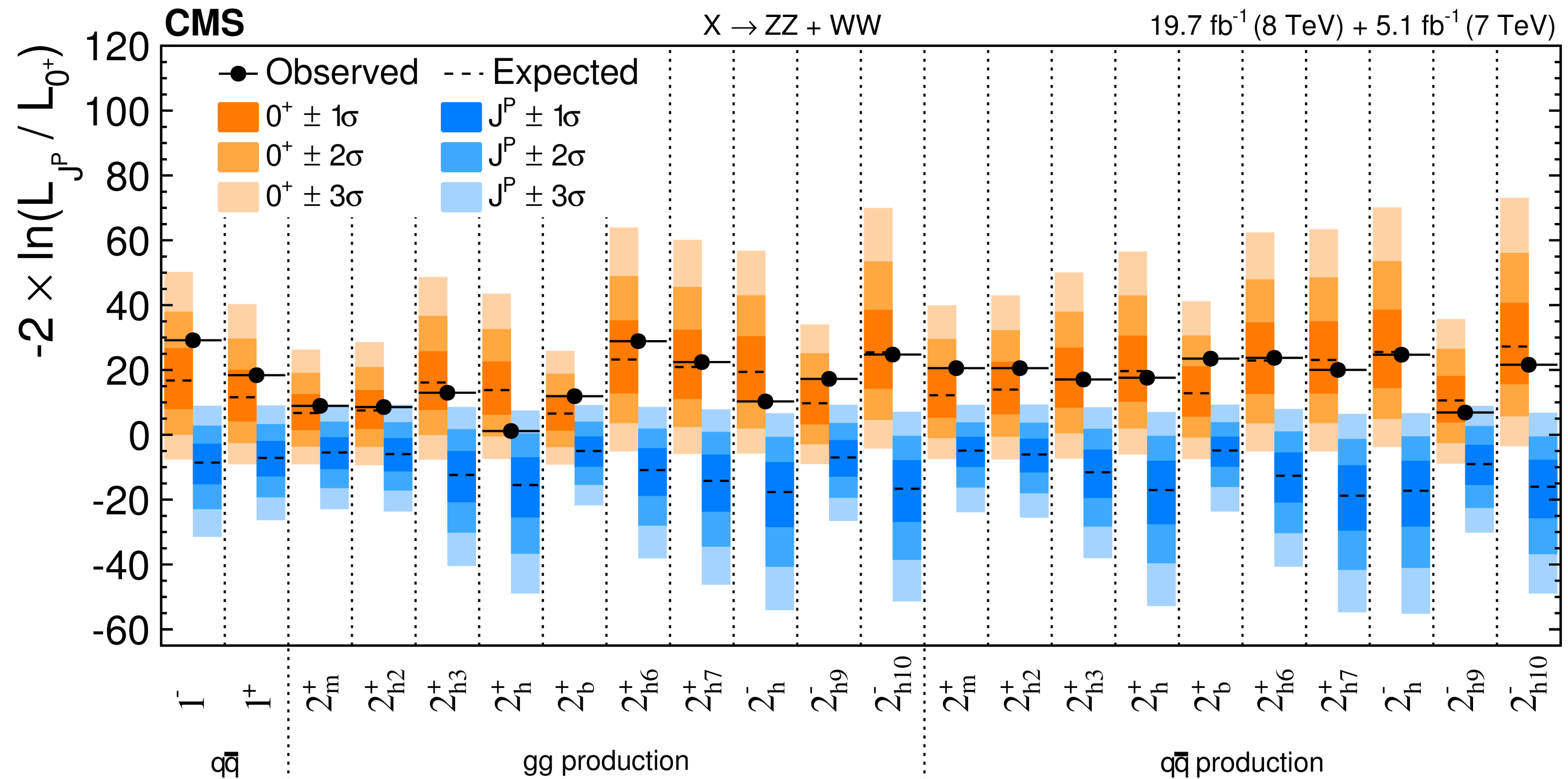
- The LHC was approved in 1994 → started running in 2009 (15 years) → will run until 2041 (32 years)
- The Tevatron was approved in 1978 → started running in 1985 (7 years) → run until 2009 (14 years)
- If we want something to be running shortly after the HL-LHC, now is the time to approve it

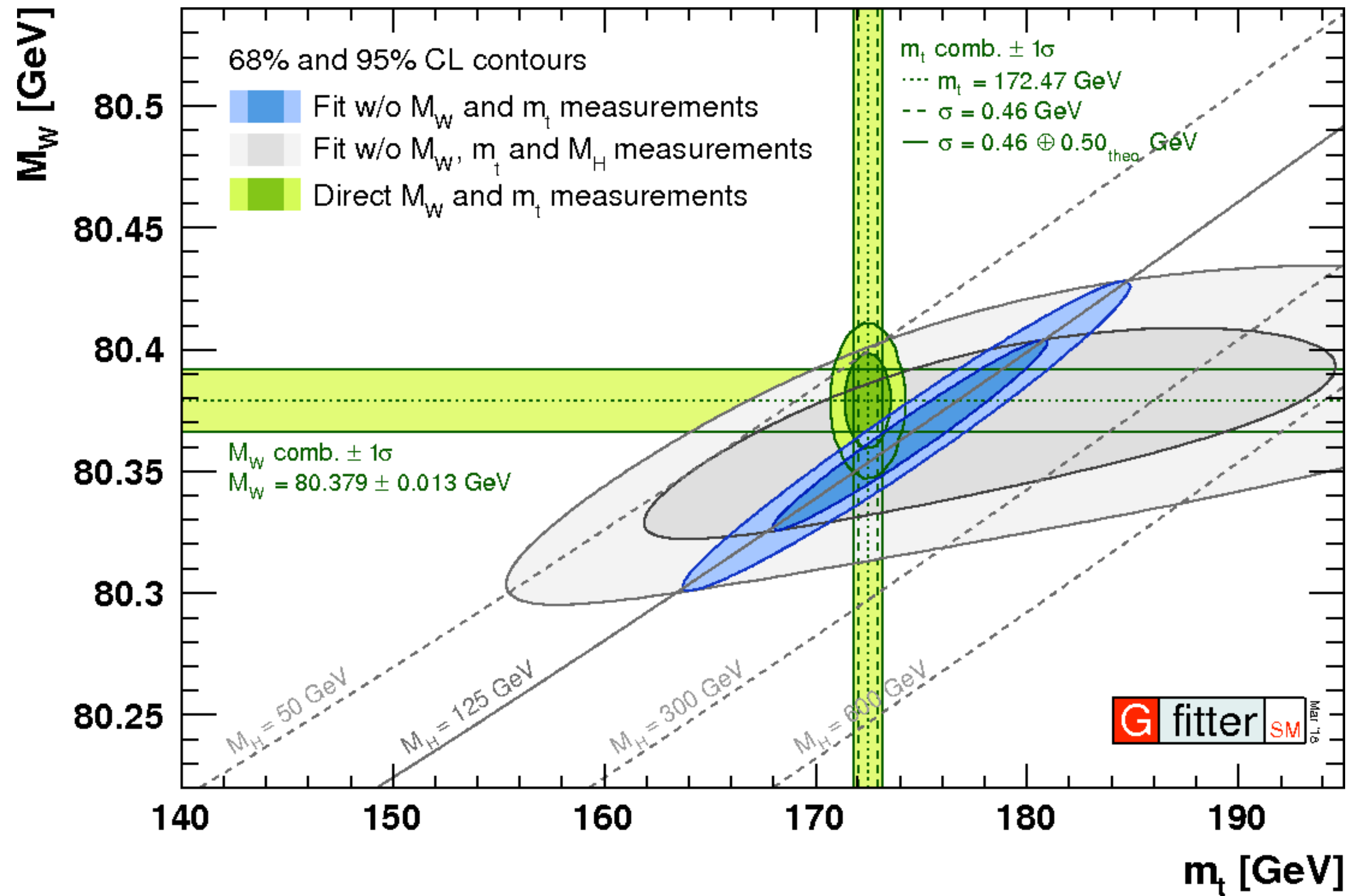
The spokespersons of the future colliders are walking among us, already in the field!
They could be any of you



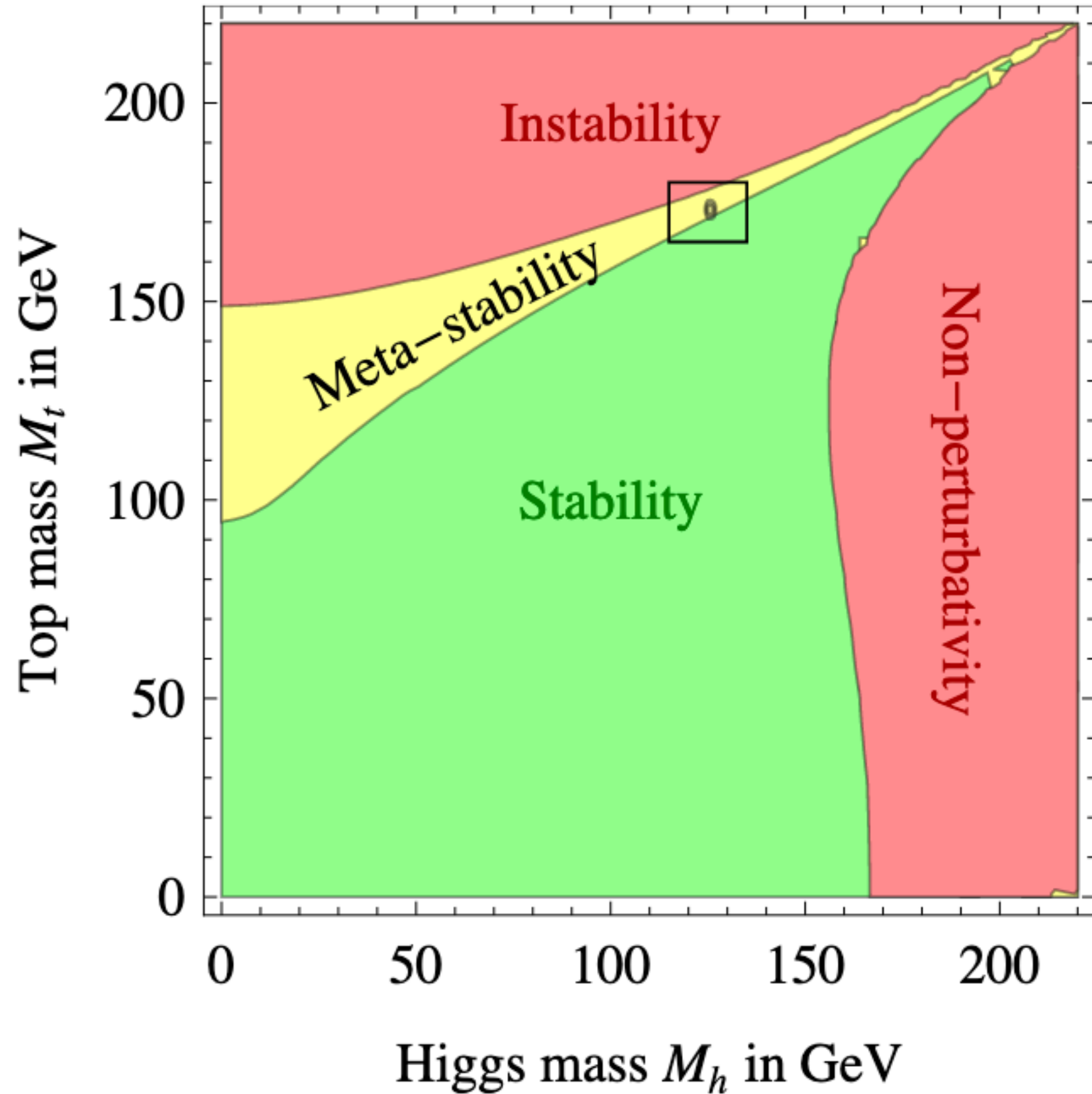
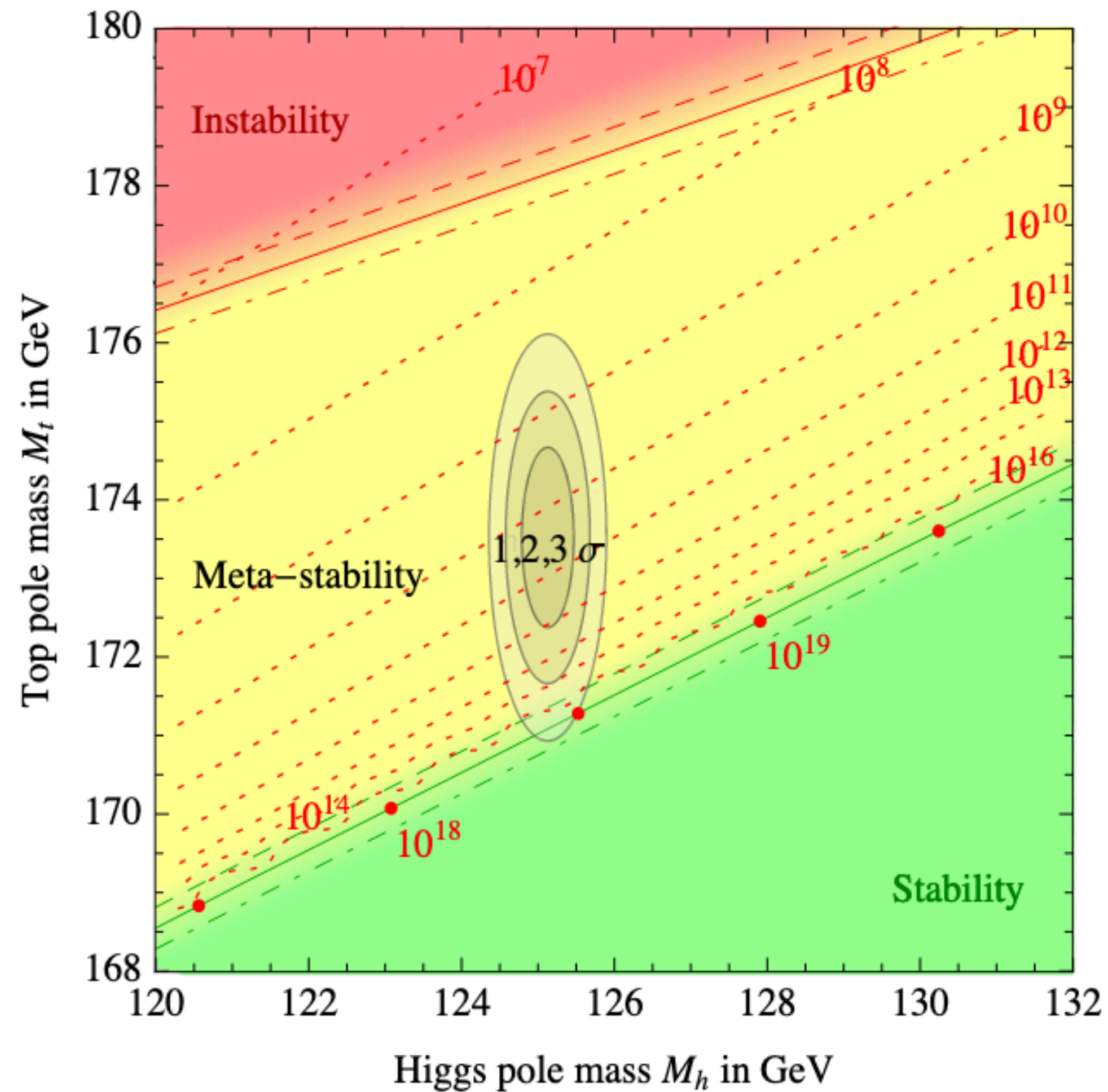
WHAT DO WE HAVE TODAY?

- A relatively new particle that is very special, an exploration tool.





The value of the Higgs mass of 125GeV is very interesting. When combined with the masses of the top quark and the W boson, it hints at something beyond the standard model.



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AND WE ALSO HAVE

- Various options on the table.
- Decades of collider expertise.
- The largest community we ever had.
- Priorities



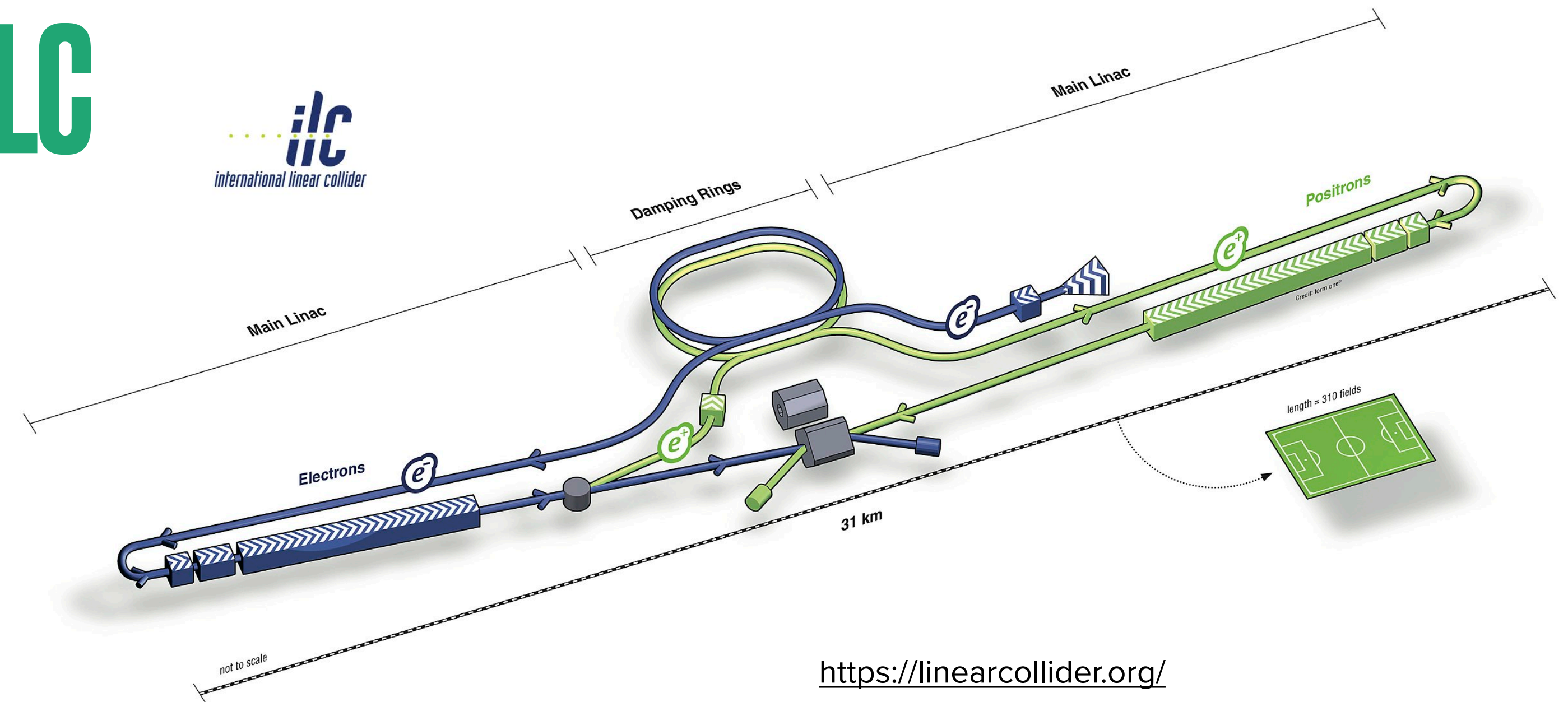
2020 update of the European Strategy for Particle Physics
“An electron-positron Higgs factory is the highest-priority next collider. For the longer term, the European particle physics community has the ambition to operate a proton-proton collider at the highest achievable energy.”

Working towards the 2025 strategy update

2023 US P5 Report
Advocates for substantial US participation and effort to support development of an offshore Higgs factory, with the goal of leading and potentially hosting a muon (!) collider beyond it.

LET'S DIVE A LITTLE DEEPER

LINEAR e^+e^- COLLIDERS: ILC



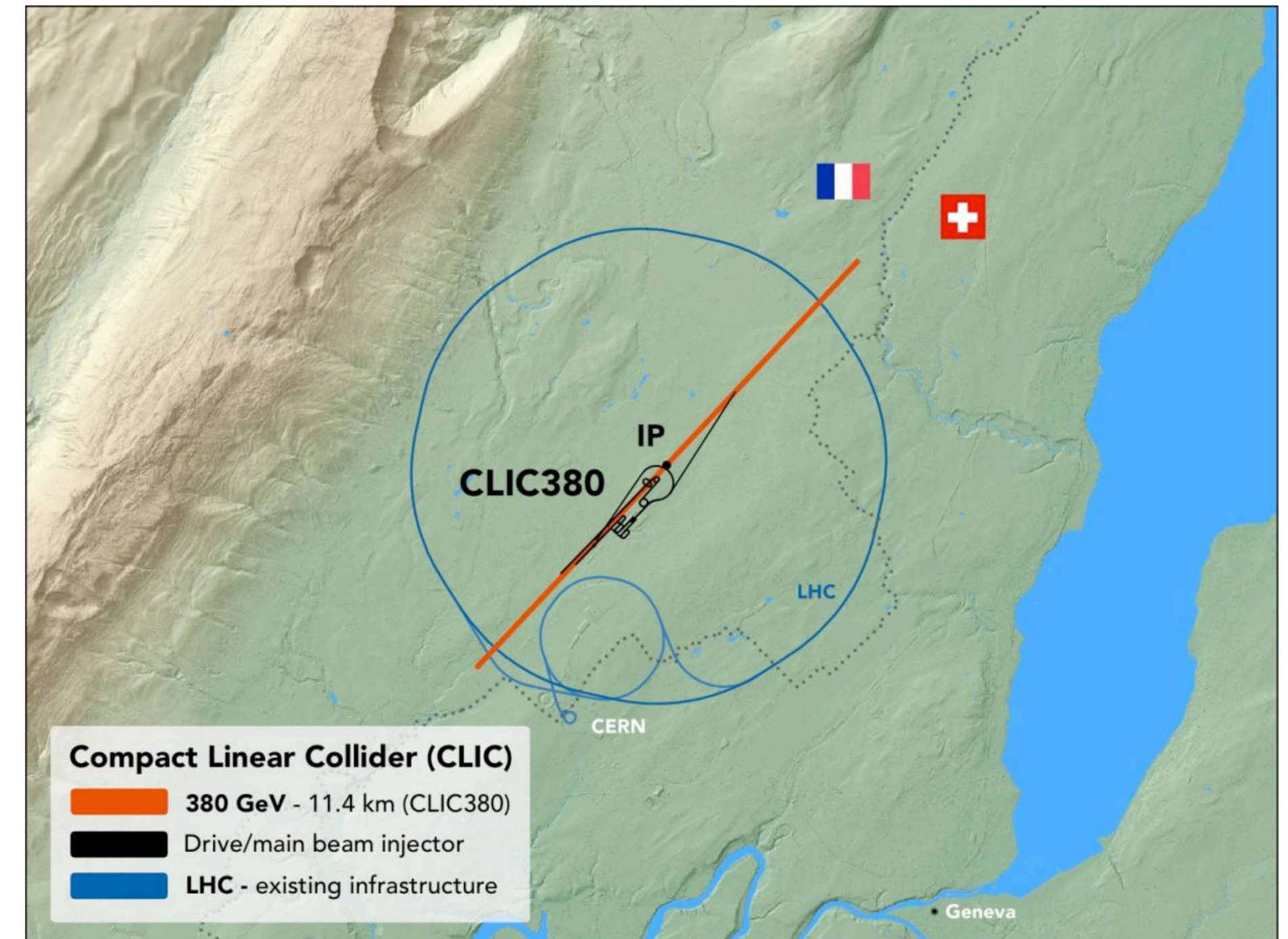
- 250 GeV (could be upgraded to 500)
- 30-km-long
- TDR in 2012
- In Japan originally (before that Germany), currently: location pending
- One collision zone: up to two experiments

LINEAR e^+e^- COLLIDERS: CLIC

<https://clicdp.web.cern.ch/>

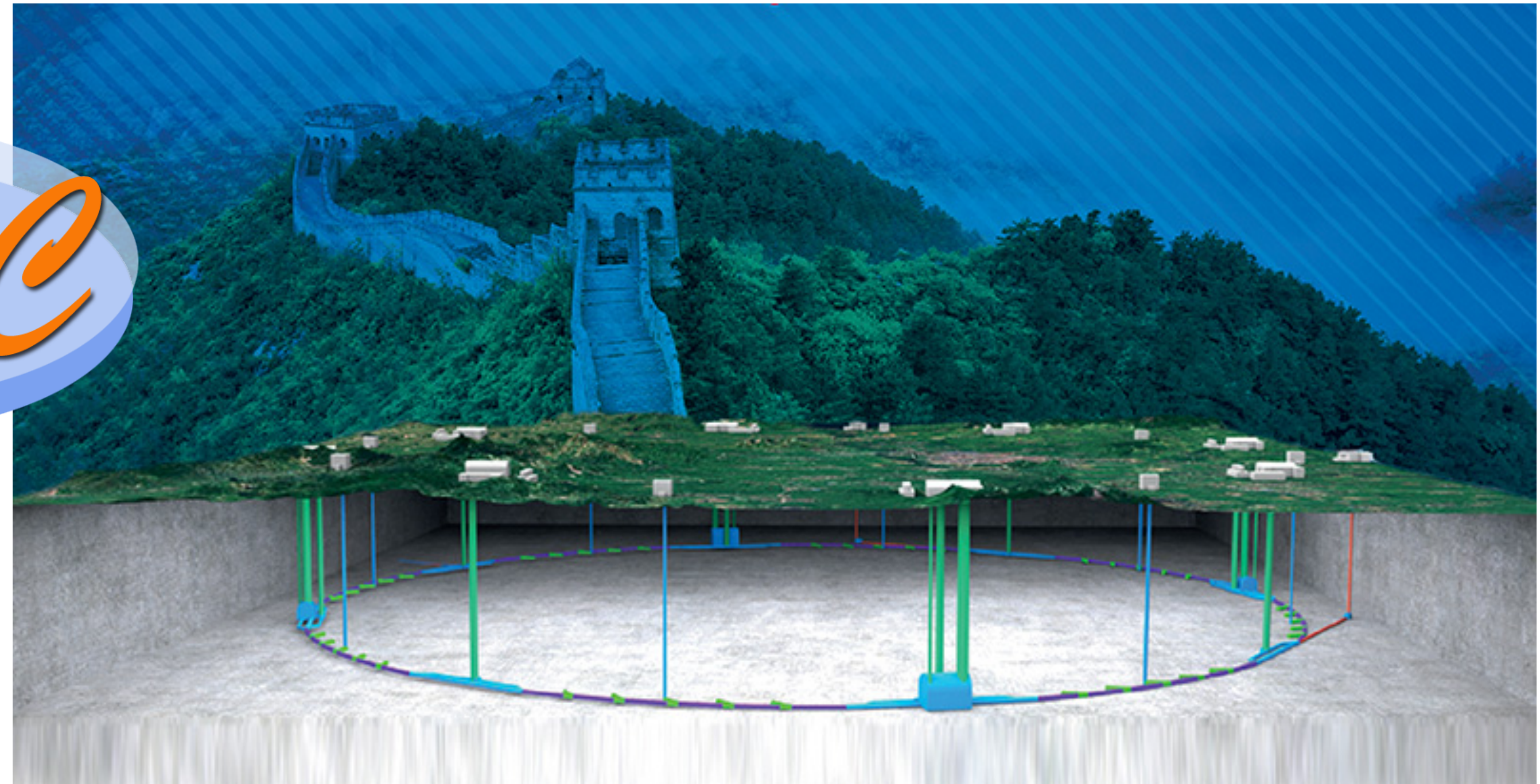


- 380 GeV (up to 3TeV)
- compact: 11 km total length
- ~LHC diameter
- CDR in 2013
- at CERN
- One interaction region



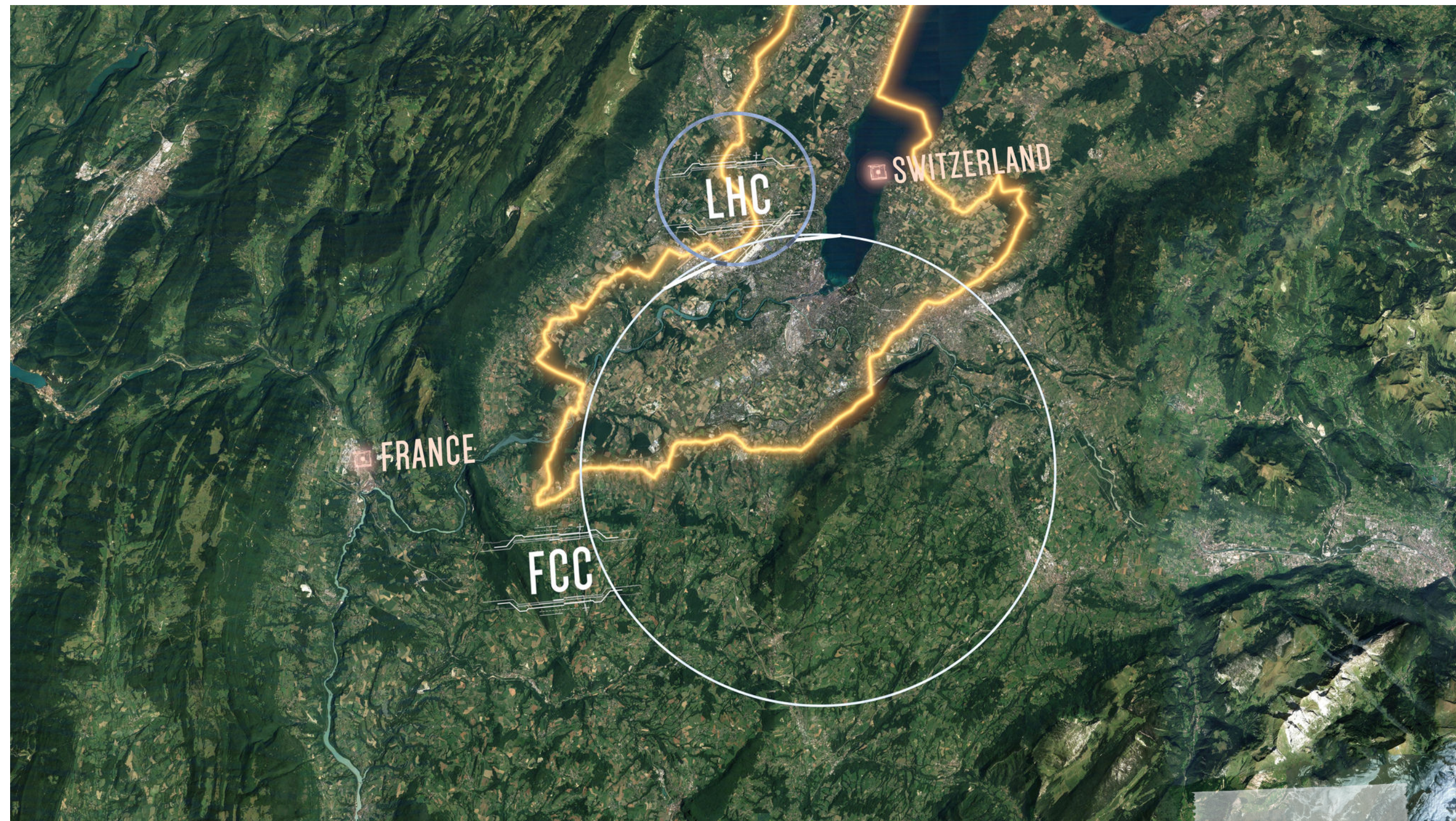
CIRCULAR e^+e^- COLLIDERS: CEPC

<http://cepc.ihep.ac.cn/>



- 90-365 GeV
- underground tunnel of 100 km of circumference, double-ring collider
- TDR in 2023
- To be hosted in China, exact location tbc
- Two (4?) interaction points (IPs)

CIRCULAR e^+e^- COLLIDERS: FCC-ee



More on this
one in a
minute

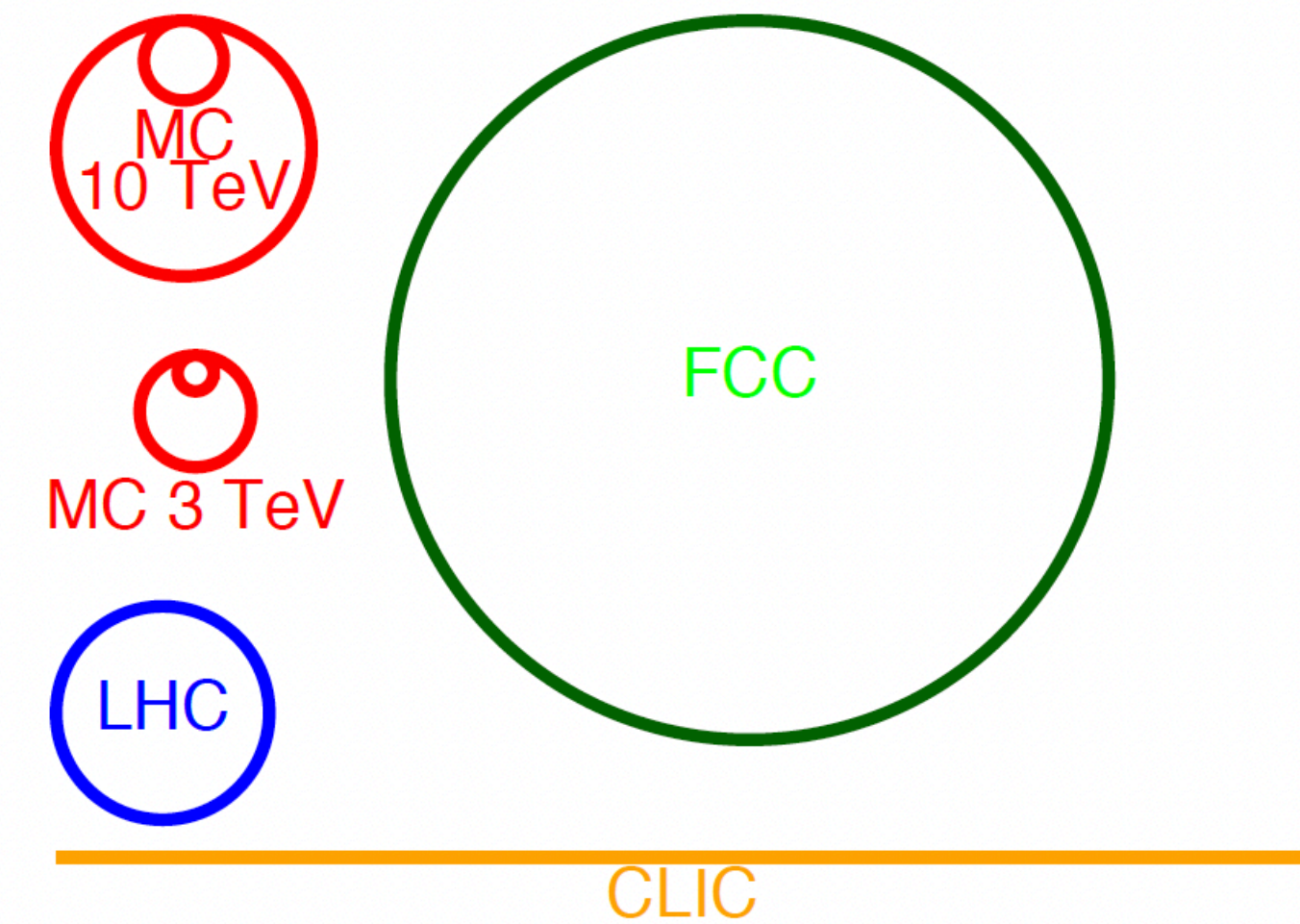


<https://fcc-ee.web.cern.ch/>

THE WILDCARD: A MUON COLLIDER



- The coolest kid on the block
- Requires a lot of R&D in the coming years
- US (or Europe?)

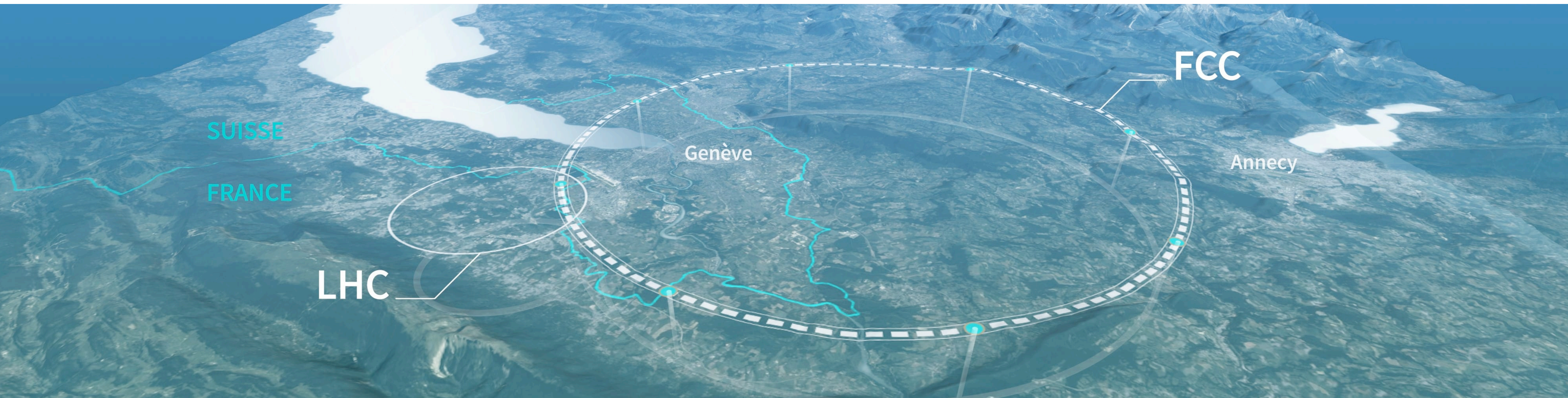


THE OPTION I AM WORKING ON



FUTURE CIRCULAR COLLIDER (FCC) AT CERN

- A versatile, next-generation particle collider housed in a 90km underground ring
- Linked to the LHC accelerator chain
- Implemented in stages, one e^+e^- machine, followed by a high-energy hadron collider



FCC PUSHES TWO FRONTIERS

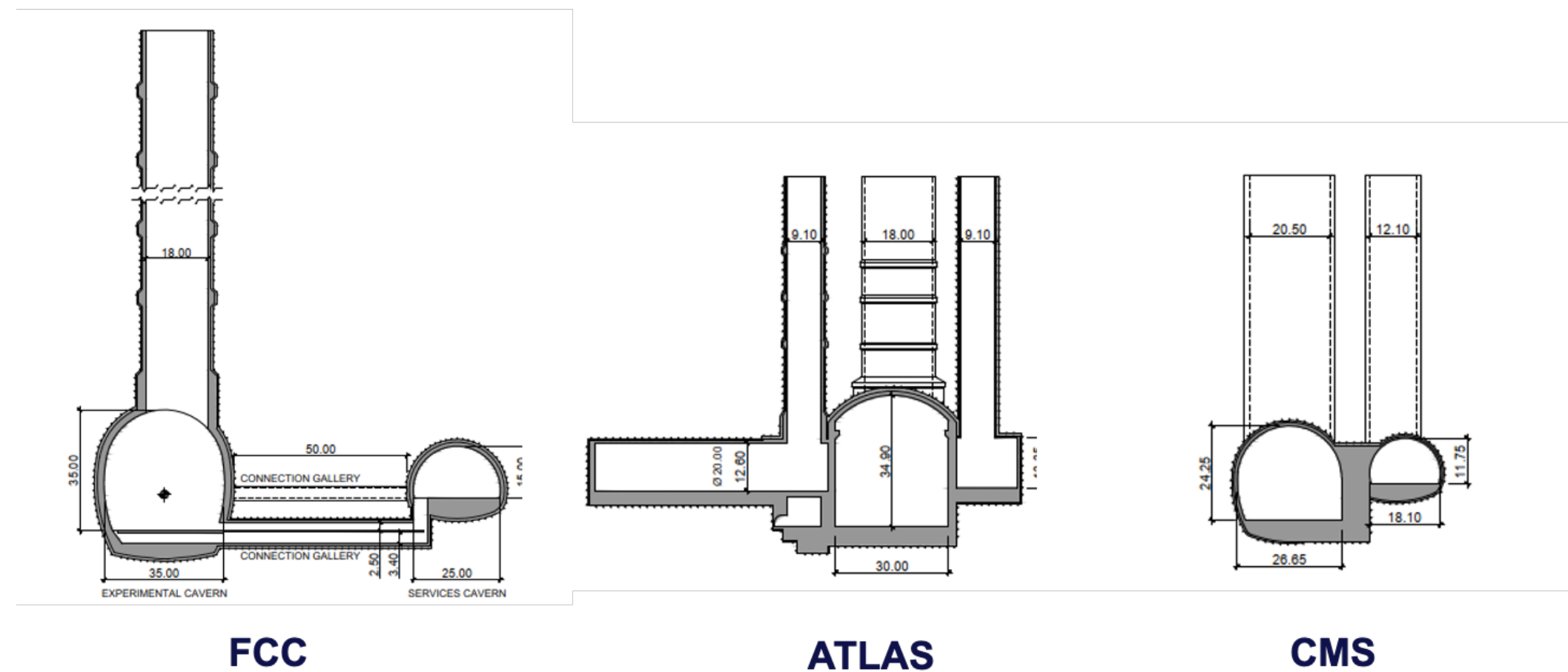
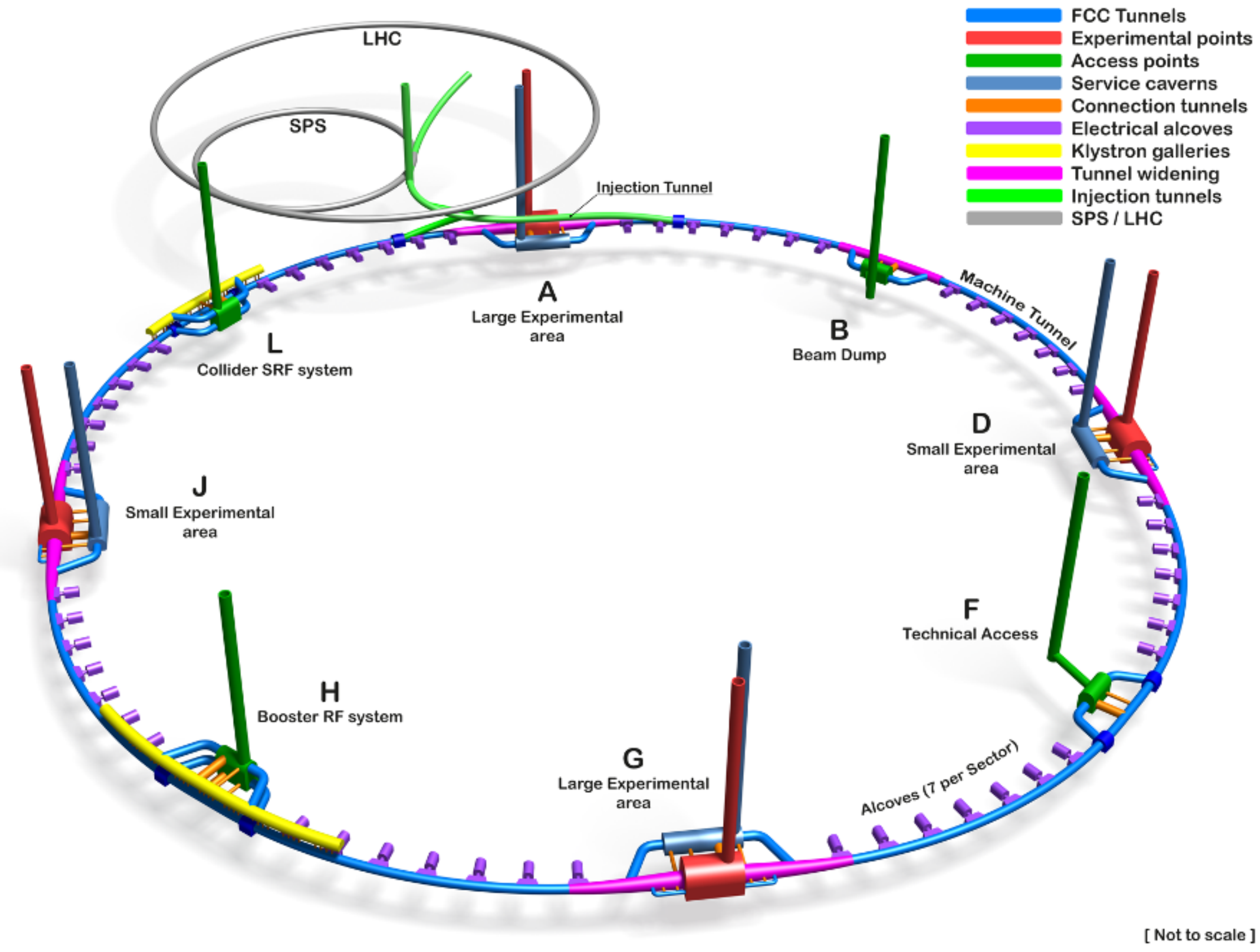
- **1st stage collider, e⁺e⁻**
 - FCC-ee
 - electron-positron collisions
 - 90-365 GeV
 - Higgs, EW, top factory
 - Construction starts: **2033**
 - Physics starts: **2048 (45 if accelerated)**
- **2nd stage collider, pp**
 - FCC-hh
 - proton-proton collisions
 - ≥ 100 TeV
 - Discovery machine
 - Physics operation: **~ 2070**

Additional modes
supported
Heavy ions, eh

Complementary
Synergetic
All-in-one facility

COMMON INFRASTRUCTURE

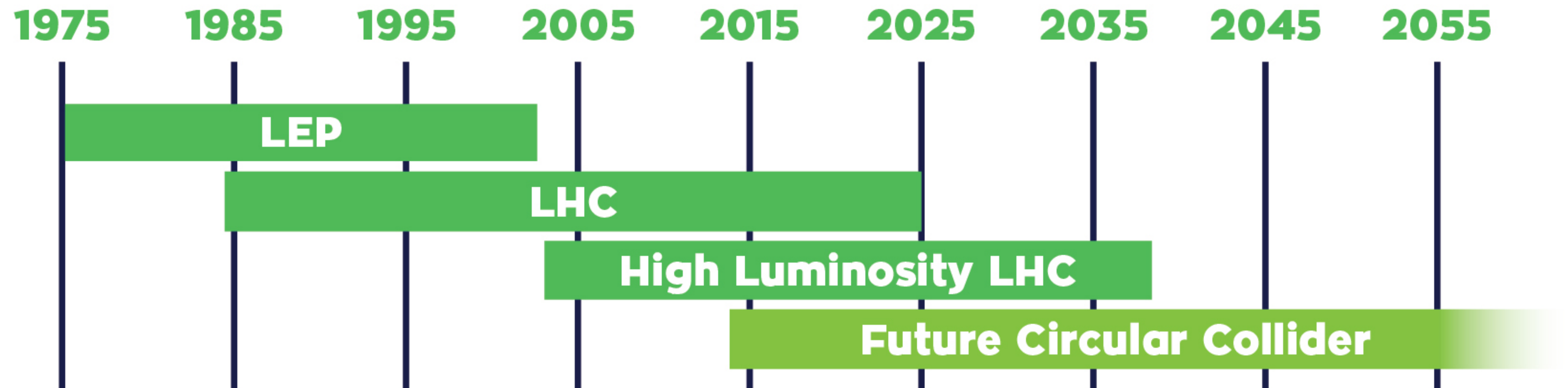
- Making use of the current acceleration chain
- Using one tunnel (and one set of caverns) for both stages
 - 90.7 km ring, 8 surface points



- 4 Experimental areas 2 large (> ATLAS) & 2 small (~CMS)
- Deepest shaft: 400m
- Average shaft depth: 243m

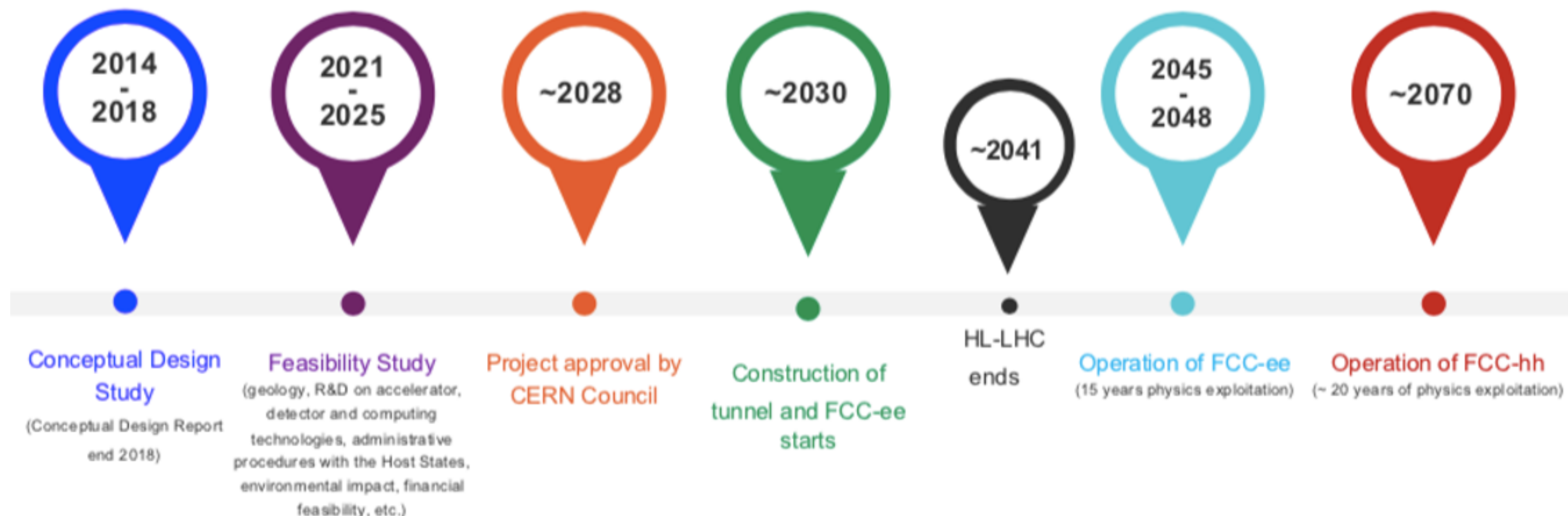
CONTINUITY

- FCC-ee technology is mature and ready → construction in parallel to HL-LHC operation
- Physics a few years after the HL-LHC
 - Guarantees continuity for generations of high energy physicists
 - Can accommodate the size of the CERN community



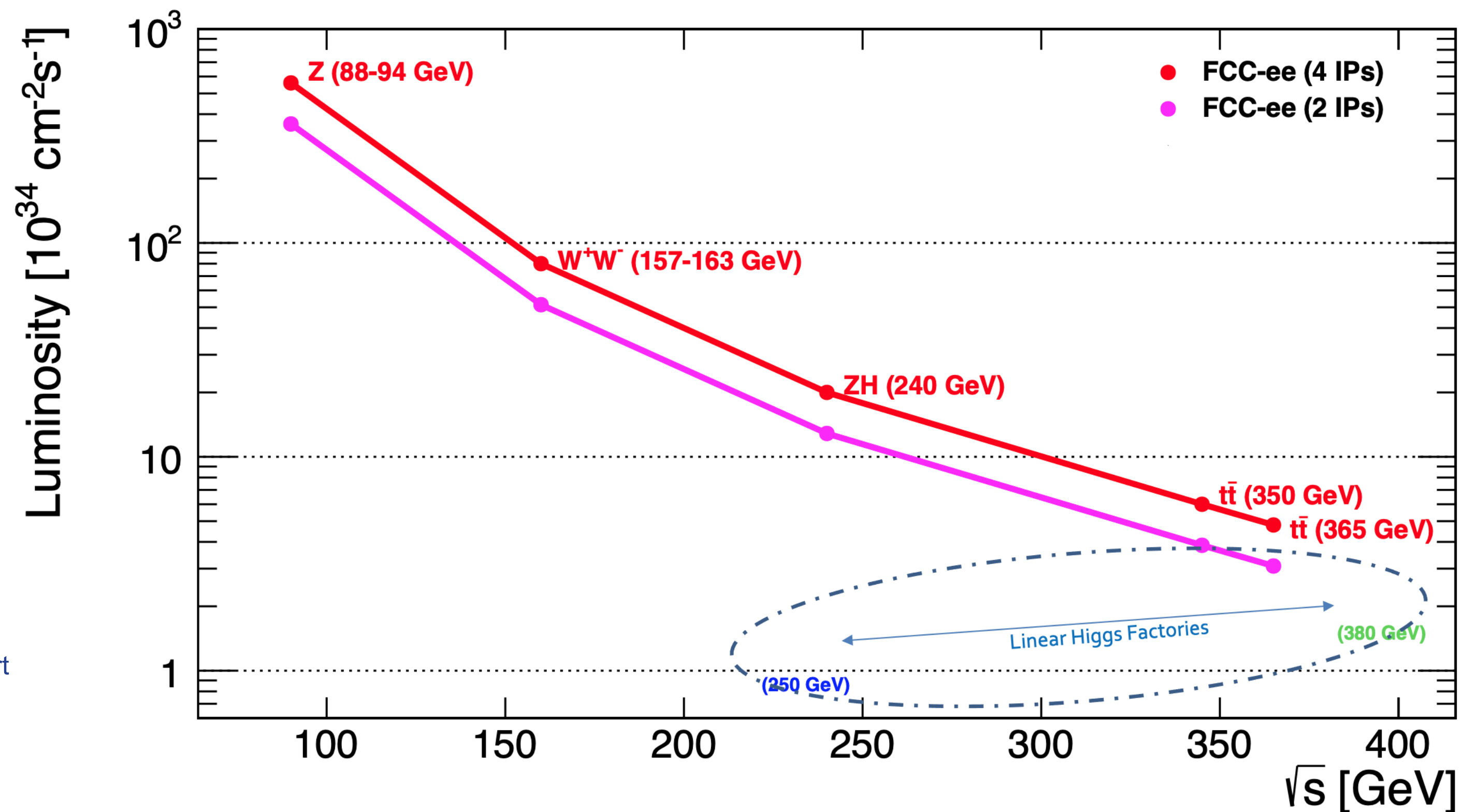
STAGED APPROACH

- Two-stage approach
 - Allows to spread the cost of the more expensive FCC-hh over more years
 - 20 years of R&D work towards optimal and affordable magnets
 - Optimization of overall investment by reusing civil engineering and large part of the technical infrastructure



LUMINOSITY

- FCC-ee: highest luminosities of all proposed Higgs and EW factories, clean experimental conditions, and a range of energies that cover Z, WW, ZH, and tt.

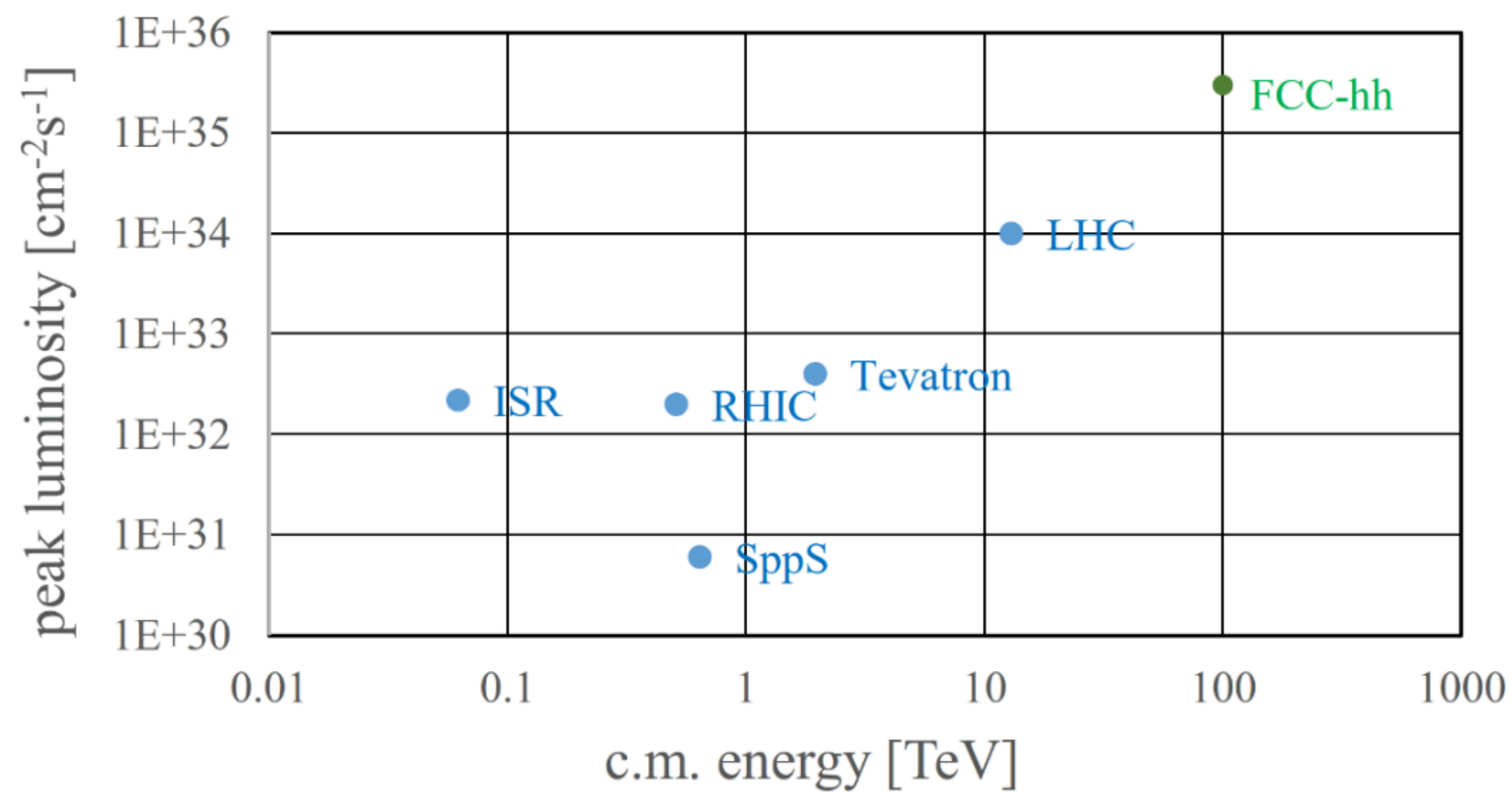


combining successful ingredients of several recent colliders → highest luminosities & energies

FCC feasibility Mid-term report
- Deliverable #8, physics and Experiments

ENERGY

- FCC-hh: Able to directly reach the next energy frontier (\sim x10 LHC)
 - order of magnitude performance increase in both energy & luminosity wrt the LHC



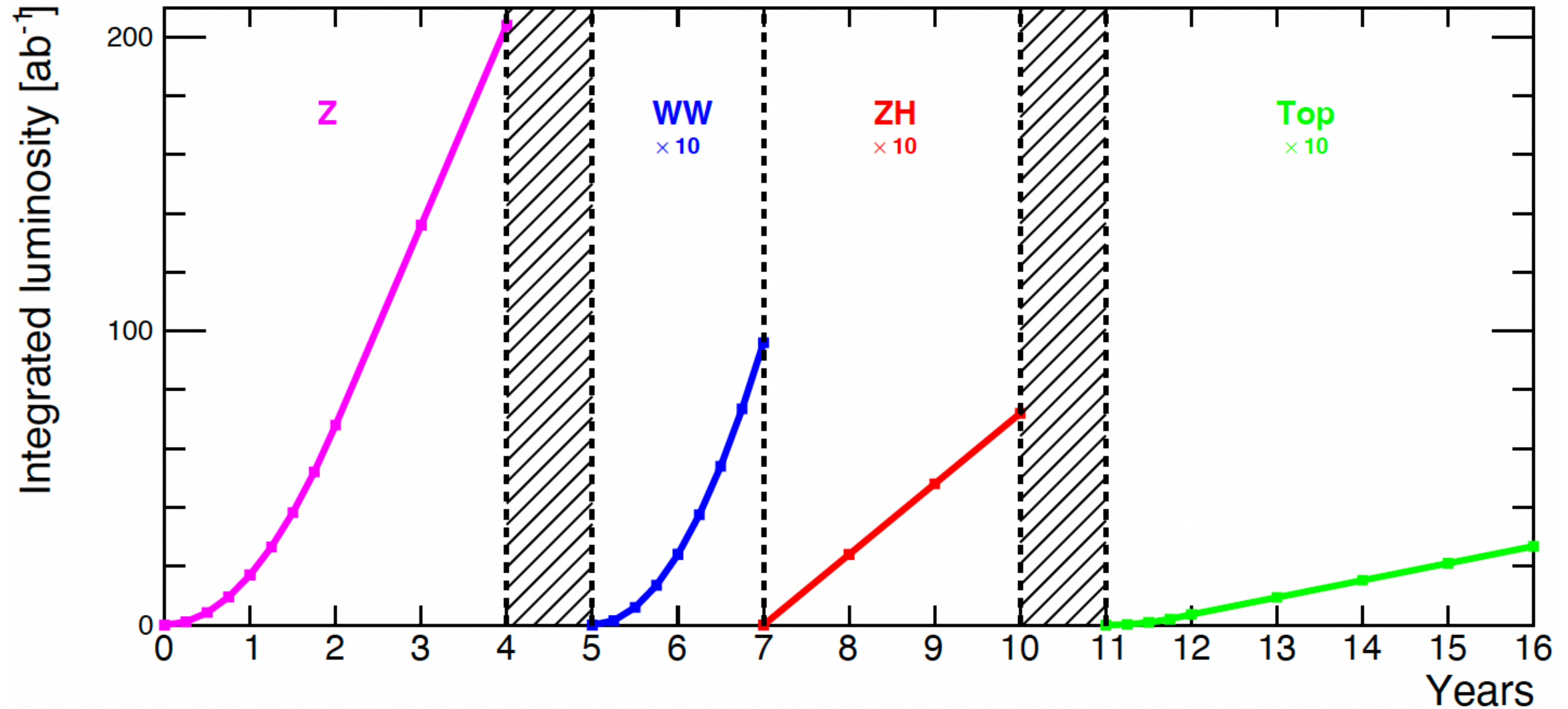
parameter	FCC-hh	HL-LHC	LHC
collision energy cms [TeV]	84 - 119	14	
dipole field [T]	14 - 20	8.33	
circumference [km]	90.7	26.7	
Integrated luminosity/main IP [fb^{-1}]	20000	3000	300

[Frank Zimmermann - ICFA Seminar 2023](#)

That's not all!
 Heavy-ion collisions and, possibly, ep/e-ion collisions, additional experiments and e.g. a FPF could be there from the beginning

FCC-ee

- 16 years, 4 IPs
- Flexibility in the run scenario:
in order and operation periods.
 - Additional runs, e.g. 125GeV possible
- Stringent experimental requirements

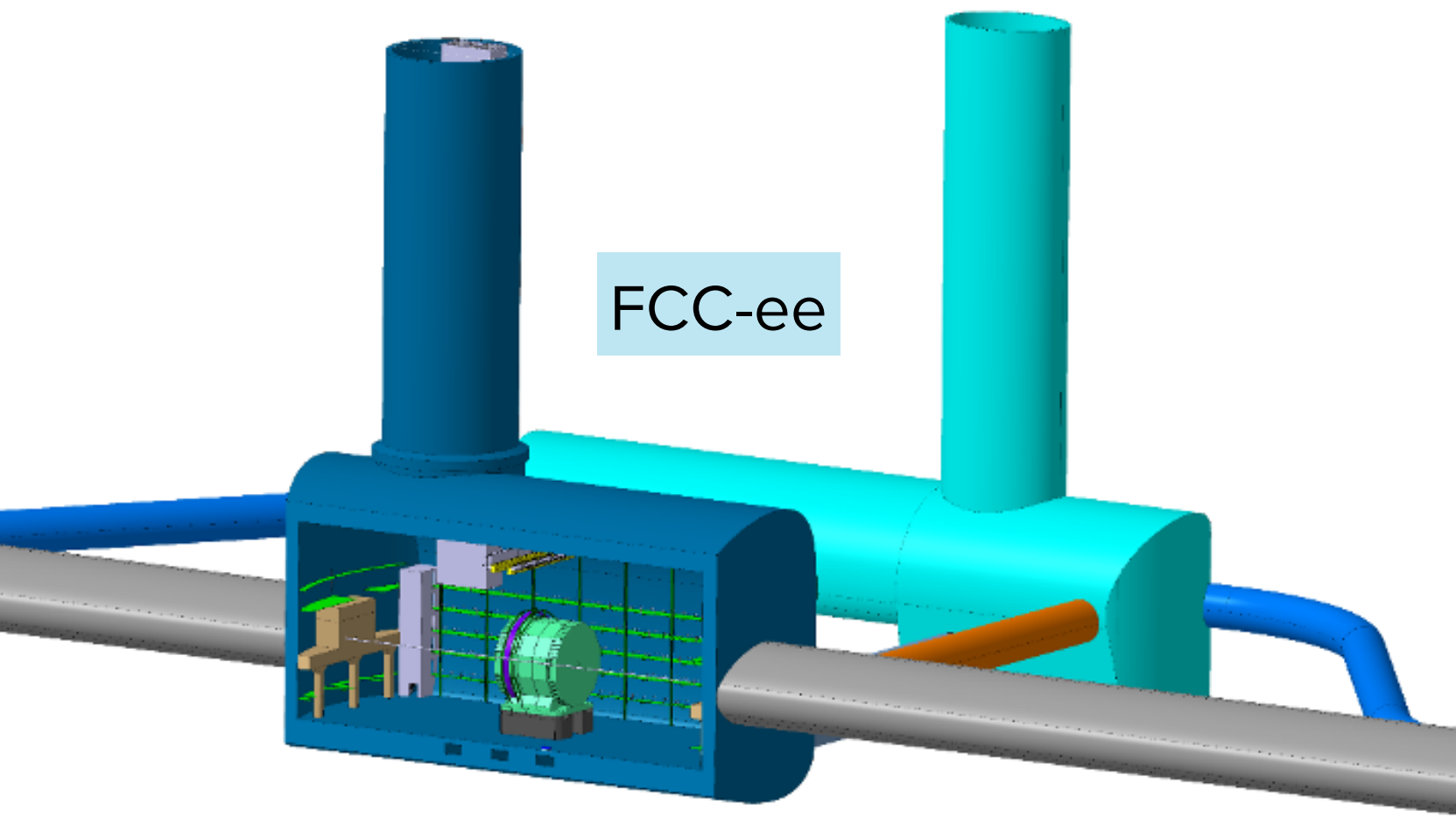


integrated
luminosity per year
summed over 4 IPs
corresponding
to 185 days of
physics per year
and 75% efficiency

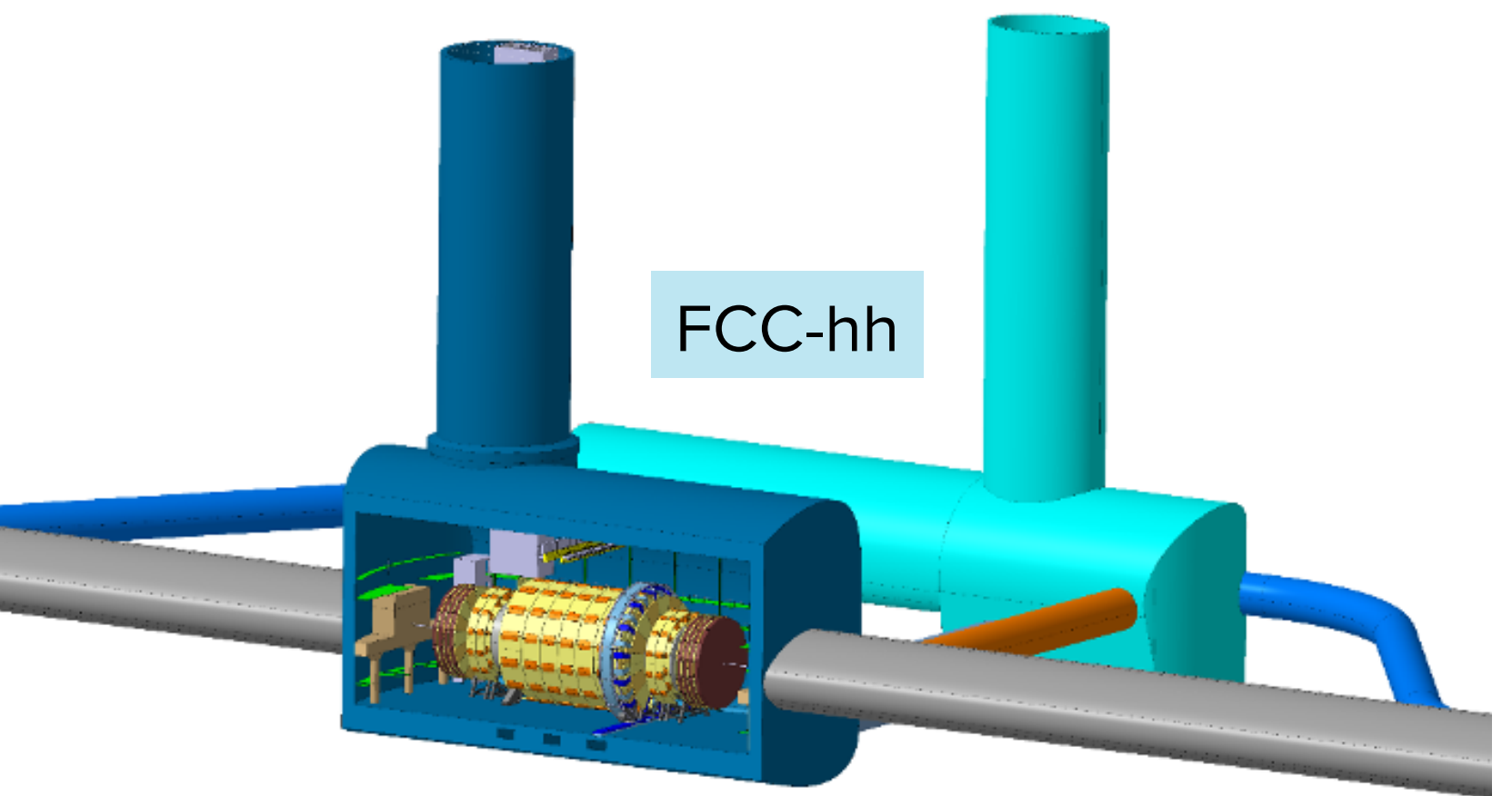
Working point	Z, years 1-2	Z, later	WW, years 1-2	WW, later	ZH	$t\bar{t}$
\sqrt{s} (GeV)	88, 91, 94		157, 163		240	340–350
Lumi/IP ($10^{34} \text{ cm}^{-2} \text{ s}^{-1}$)	70	140	10	20	5.0	0.75
Lumi/year (ab^{-1})	34	68	4.8	9.6	2.4	0.36
Run time (year)	2	2	2	–	3	1
Number of events	6×10^{12} Z		2.4×10^8 WW		1.45×10^6 ZH + 45k WW \rightarrow H	1.9×10^6 $t\bar{t}$ +330k ZH +80k WW \rightarrow H

all the data of
LEP1 in minutes

DETECTOR CONCEPTS



FCC-ee



FCC-hh

CLD	IDEA	ALLEGRO
<p>Based on CLIC detector design, arXiv:1911.12230</p> <p>Full silicon vertex detector and tracker 3D-imaging highly-granular calorimeter system Coil outside calorimeter system</p>	<p>Innovative, possibly cheaper than CLD https://pos.sissa.it/390/819 Baseline in many ongoing studies</p> <p>Silicon vertex detector Short-drift, ultra-light wire chamber Dual-readout calorimeter Thin and light solenoid coil inside calorimeter system</p>	<p>GranuLAr WS, IJCLab 2022 – Martin Aleksa</p> <p>Highly granular noble-liquid calorimeter Thin 2T solenoid in the calorimeter cryostat.</p>

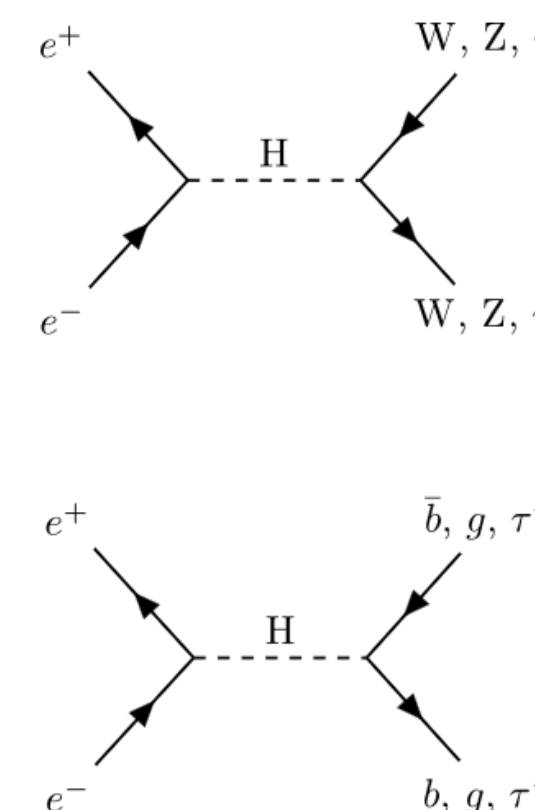
More complementary options possible (4 IP!) → Can we optimize detector designs for the complete physics program? Yes! opportunities to contribute

HIGGS UNDER THE MICROSCOPE

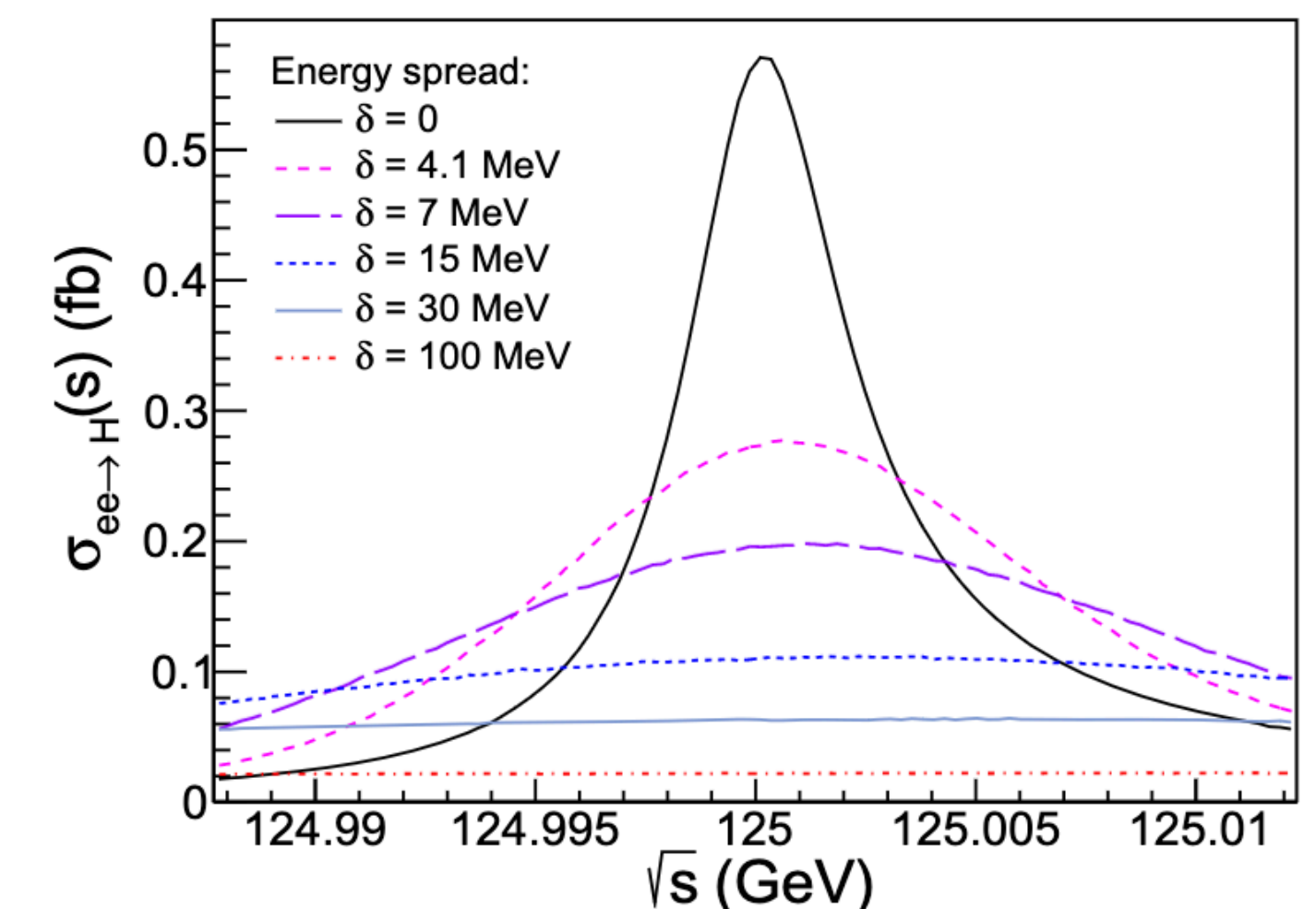
FCC feasibility Mid-term report -
Deliverable #8, physics and Experiments

- We will be able to map all the Higgs properties with accuracy
- Improvement in the precision of several Higgs boson couplings of about one order of magnitude wrt the end of the HL-LHC
 - sub-% measurement of couplings to W, Z, b, τ , % to gluon and charm
- absolute measurement width/couplings
- Recoil method
- Access to direct Higgs production at 125GeV

Coupling	HL-LHC	FCC-ee (240–365 GeV) 2 IPs / 4 IPs
κ_W [%]	1.5*	0.43 / 0.33
κ_Z [%]	1.3*	0.17 / 0.14
κ_g [%]	2*	0.90 / 0.77
κ_γ [%]	1.6*	1.3 / 1.2
$\kappa_{Z\gamma}$ [%]	10*	10 / 10
κ_c [%]	–	1.3 / 1.1
κ_t [%]	3.2*	3.1 / 3.1
κ_b [%]	2.5*	0.64 / 0.56
κ_μ [%]	4.4*	3.9 / 3.7
κ_τ [%]	1.6*	0.66 / 0.55
BR _{inv} (<%, 95% CL)	1.9*	0.20 / 0.15
BR _{unt} (<%, 95% CL)	4*	1.0 / 0.88



[Eur. Phys. J. Plus \(2022\) 137:201](#)



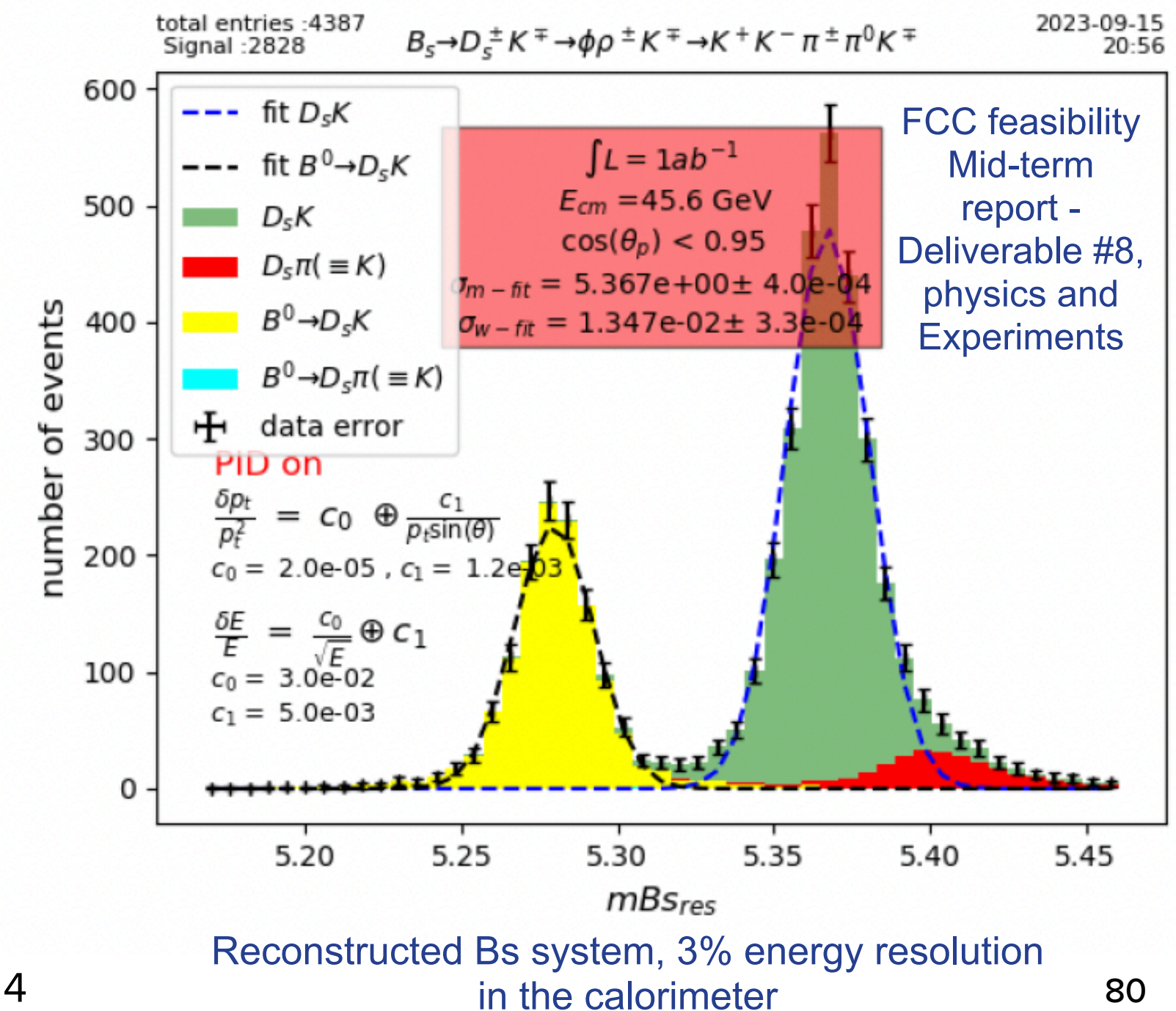
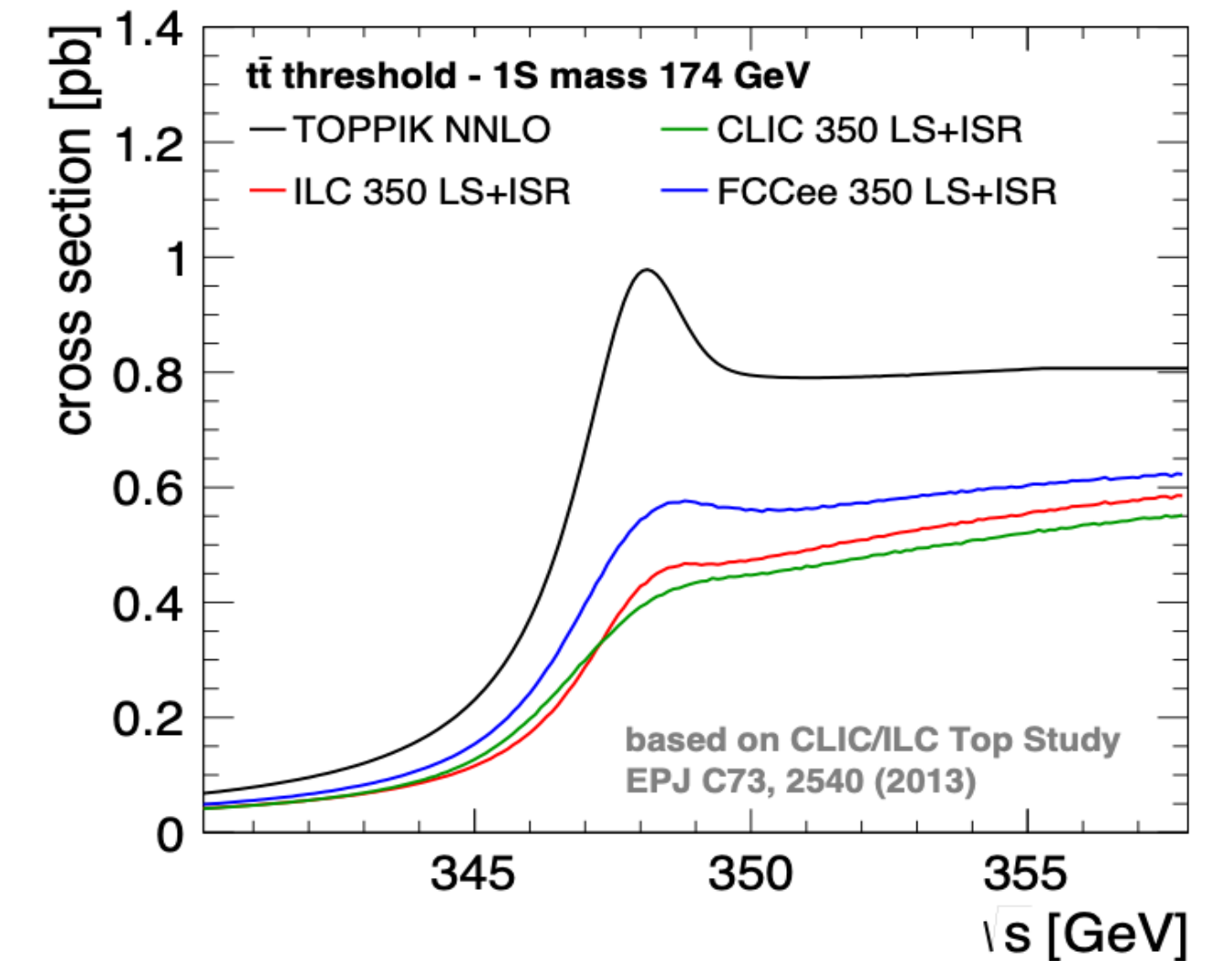
ELECTROWEAK

- Dedicated W and Z runs with unprecedented statistics
 - Z pole run → LEP Statistical uncertainties divided by ~1000
- Comprehensive measurements of the Z lineshape and many Electroweak Precision Observables
 - 50x improved precision → 7x jump in indirect sensitivity to BSM effects
- Direct and uniquely precise determinations of $\alpha_{\text{QED}}(m_Z)$ (for the first time) and $\alpha_s(m_Z)$

Observable	present value	present \pm error	FCC-ee Stat.	FCC-ee Syst.	Comment and leading error
m_Z (keV)	91186700	\pm 2200	4	100	From Z line shape scan Beam energy calibration
Γ_Z (keV)	2495200	\pm 2300	4	25	From Z line shape scan Beam energy calibration
$\sin^2 \theta_W^{\text{eff}} (\times 10^6)$	231480	\pm 160	2	2.4	From $A_{\text{FB}}^{\mu\mu}$ at Z peak Beam energy calibration
$1/\alpha_{\text{QED}}(m_Z^2) (\times 10^3)$	128952	\pm 14	3	small	From $A_{\text{FB}}^{\mu\mu}$ off peak QED&EW errors dominate
$R_\ell^Z (\times 10^3)$	20767	\pm 25	0.06	0.2-1	Ratio of hadrons to leptons Acceptance for leptons
$\alpha_s(m_Z^2) (\times 10^4)$	1196	\pm 30	0.1	0.4-1.6	From R_ℓ^Z
$\sigma_{\text{had}}^0 (\times 10^3)$ (nb)	41541	\pm 37	0.1	4	Peak hadronic cross section Luminosity measurement
$N_\nu (\times 10^3)$	2996	\pm 7	0.005	1	Z peak cross sections Luminosity measurement
$R_b (\times 10^6)$	216290	\pm 660	0.3	< 60	Ratio of $b\bar{b}$ to hadrons Stat. extrapol. from SLD
$A_{\text{FB},0}^b (\times 10^4)$	992	\pm 16	0.02	1-3	b-quark asymmetry at Z pole From jet charge
$A_{\text{FB}}^{\text{pol},\tau} (\times 10^4)$	1498	\pm 49	0.15	< 2	τ polarization asymmetry τ decay physics
τ lifetime (fs)	290.3	\pm 0.5	0.001	0.04	Radial alignment
τ mass (MeV)	1776.86	\pm 0.12	0.004	0.04	Momentum scale
τ leptonic ($\mu\nu_\mu\nu_\tau$) B.R. (%)	17.38	\pm 0.04	0.0001	0.003	e/μ /hadron separation
m_W (MeV)	80350	\pm 15	0.25	0.3	From WW threshold scan Beam energy calibration
Γ_W (MeV)	2085	\pm 42	1.2	0.3	From WW threshold scan Beam energy calibration

TOP AND FLAVOUR

- Threshold region: most precise measurements of top quark mass and width, FCNC...
- Flavour factory
 - Clean environment, precise momentum of b/c/ τ pairs from Z decays (like in B-factories), ~ 10 times more bb/cc than final Belle-II statistics
- Potential in boosted b/ τ
 - Higher efficiency than at B factories for modes with missing energy (especially multi- ν) and inclusive modes, smaller uncertainties in lepton ID efficiencies.
- Interesting opportunities:
 - e.g. decays of: Rare b-hadron with $\tau\tau$ pairs in the final state and charged-current b-hadrons with $\tau\nu$ in the final state; lepton flavour violating τ decays, or lepton-universality tests in τ decays.

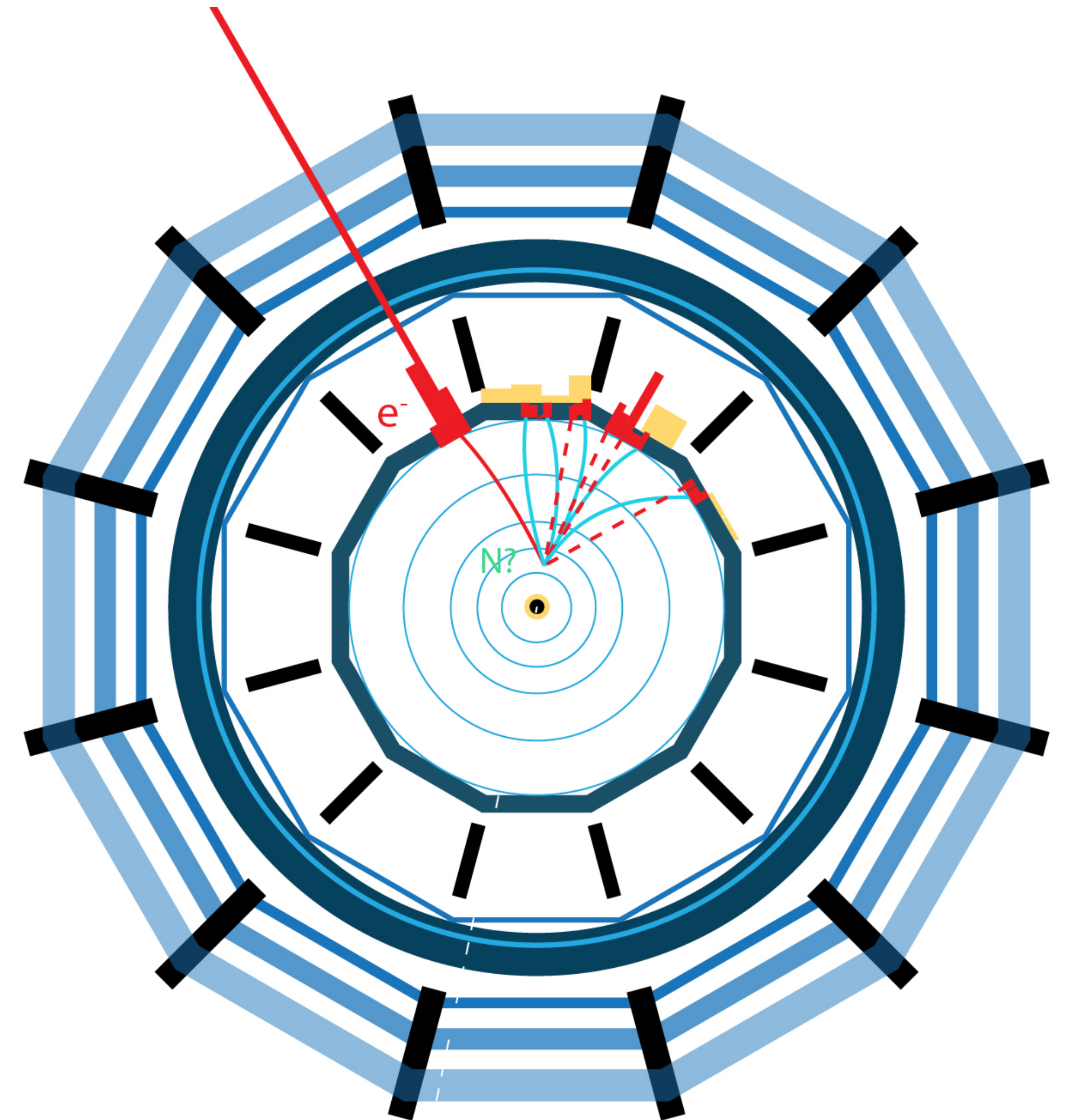


BEYOND THE STANDARD MODEL

- Potential for indirect BSM exploration: SMEFT, and other precision/search cases
- But also direct searches:
 - Clean environment, high luminosity, and large acceptance, direct scrutiny of $O(1-100)$ GeV mass range for new particles
 - Dark/hidden sectors that connect feebly to the SM via mediators (dark photon)
 - Exotic decays of the Z or Higgs boson
 - Specially interesting are signature-driven searches for non-mainstream signals
 - Here come **long-lived particles** again!

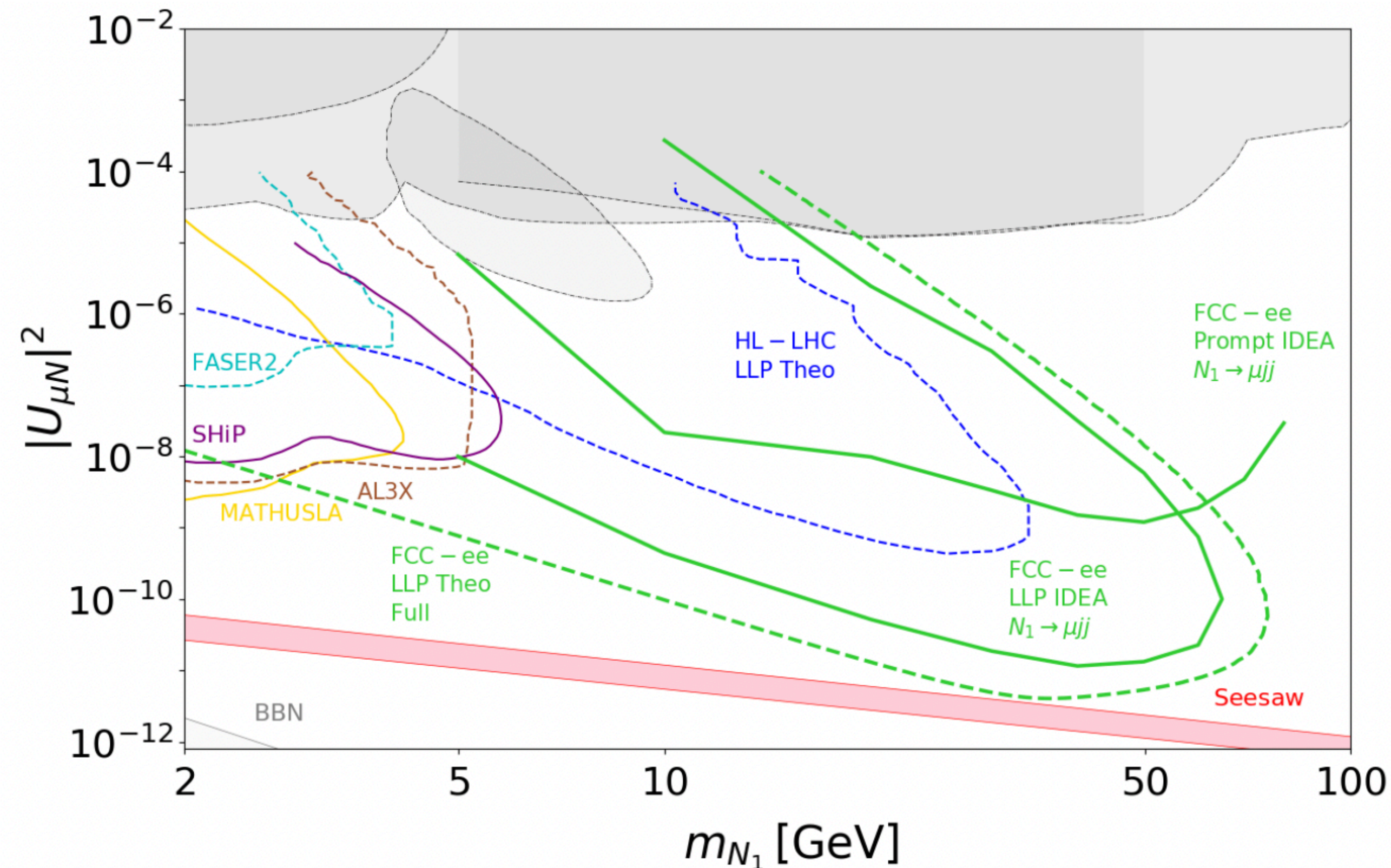
THE FLAGSHIP, HNLs

- Sterile/right-handed neutrinos
 - Could give many answers: neutrino masses, DM, BAU
- Many of the current limits cover large neutrino mixing angles.
 - For small values of the mixing angle → LLP signature (displaced vertex search)
- The FCC-ee will offer a fantastic reach for HNL at the Z-Pole
 - FCC will probe space not constrained by astrophysics or cosmology, complementary to other dedicated experiments



HEAVY NEUTRAL LEPTONS

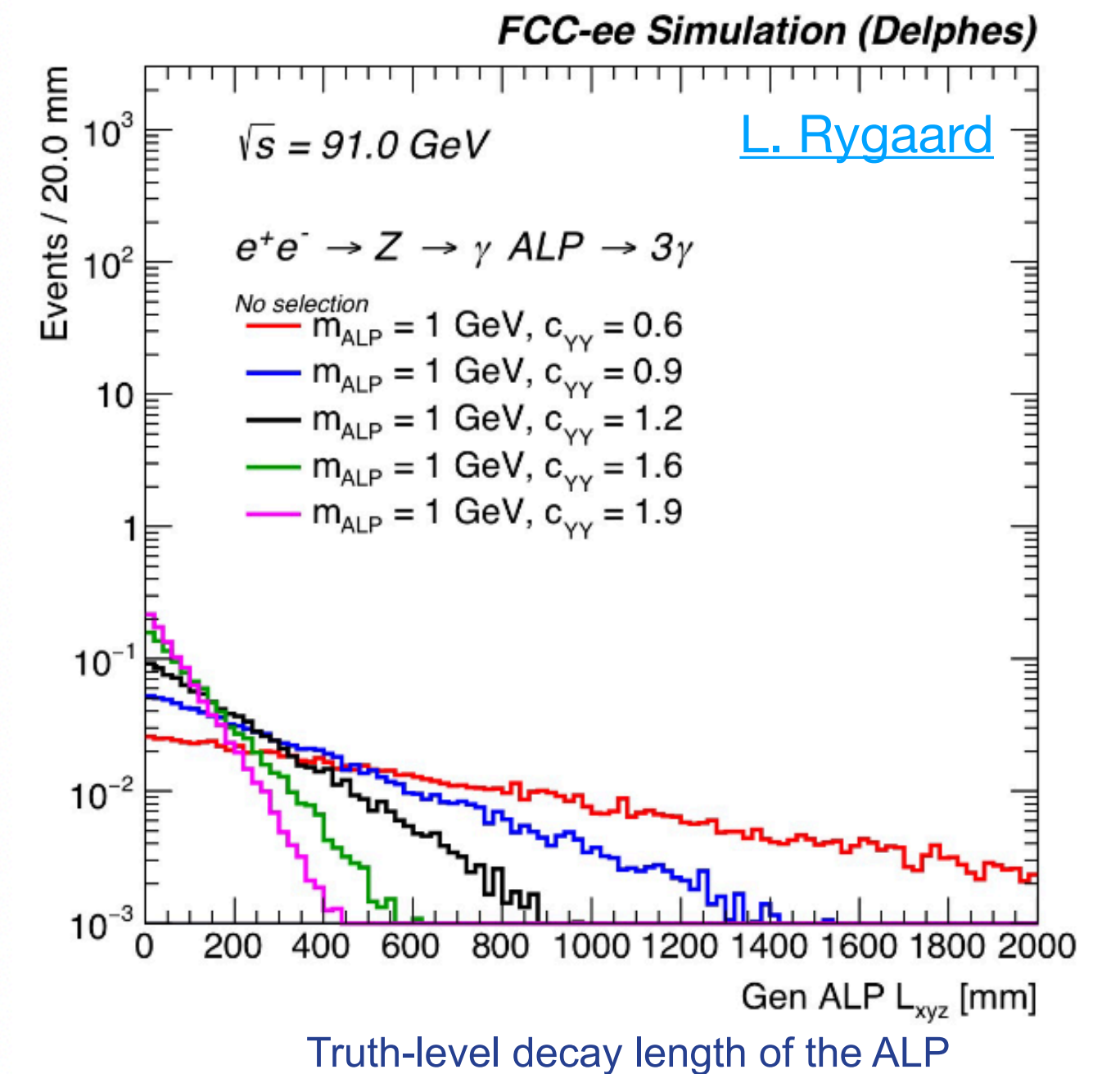
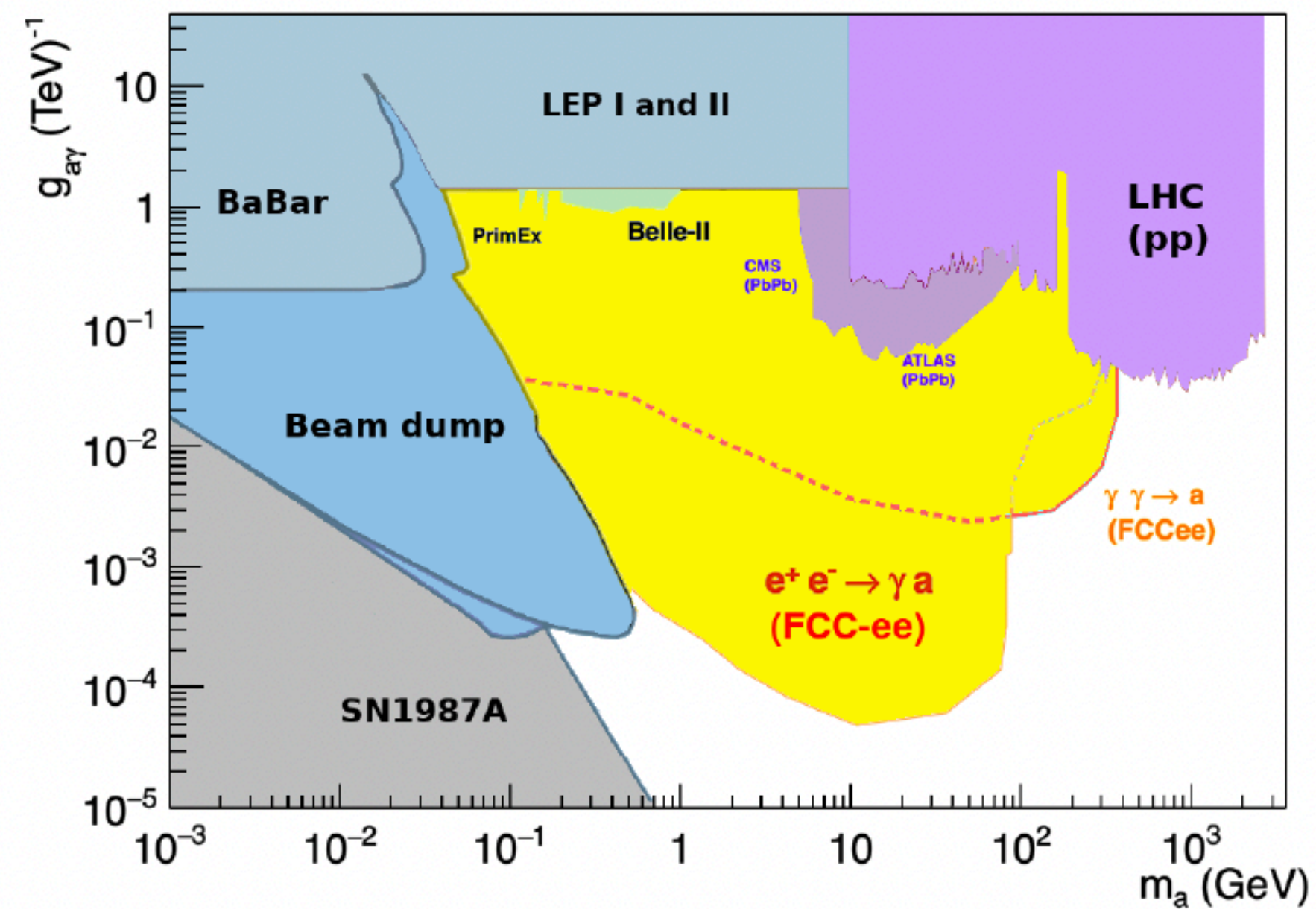
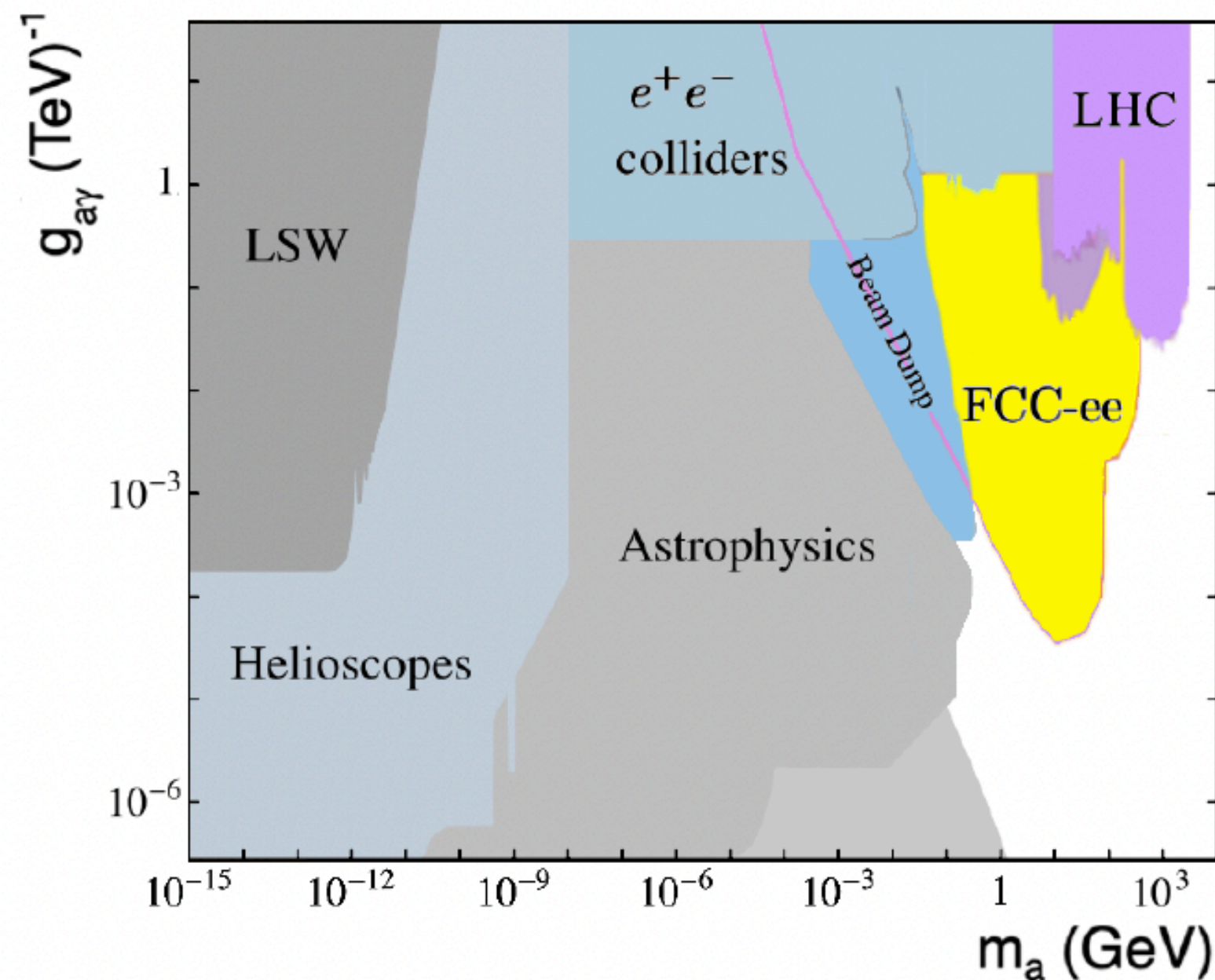
- Working towards a complete sensitivity analysis implemented in FCC software
- A few channels ($e\nu$, μjj , $\mu\mu\nu$) considered prompt and long-lived



AXION-LIKE PARTICLES

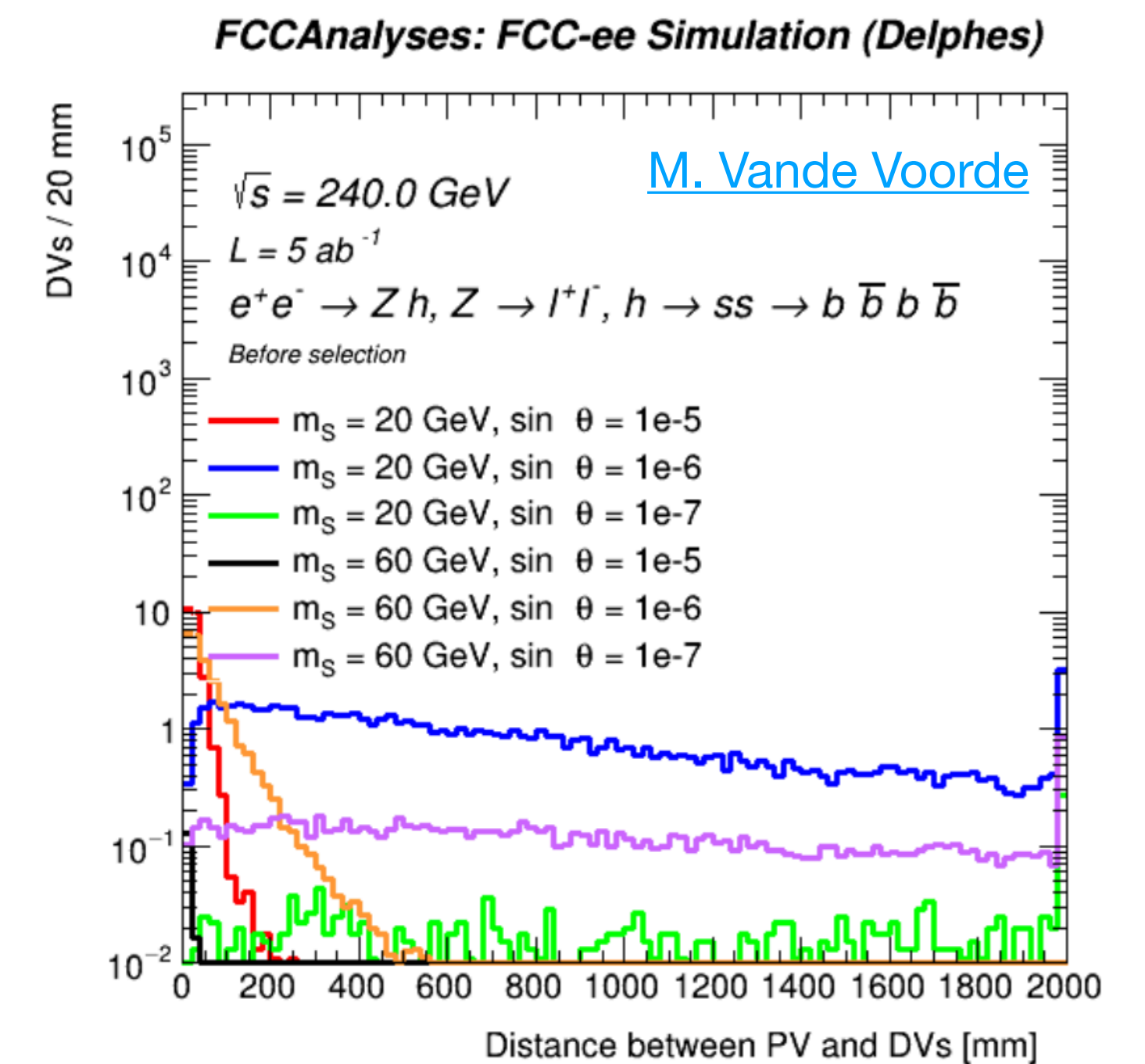
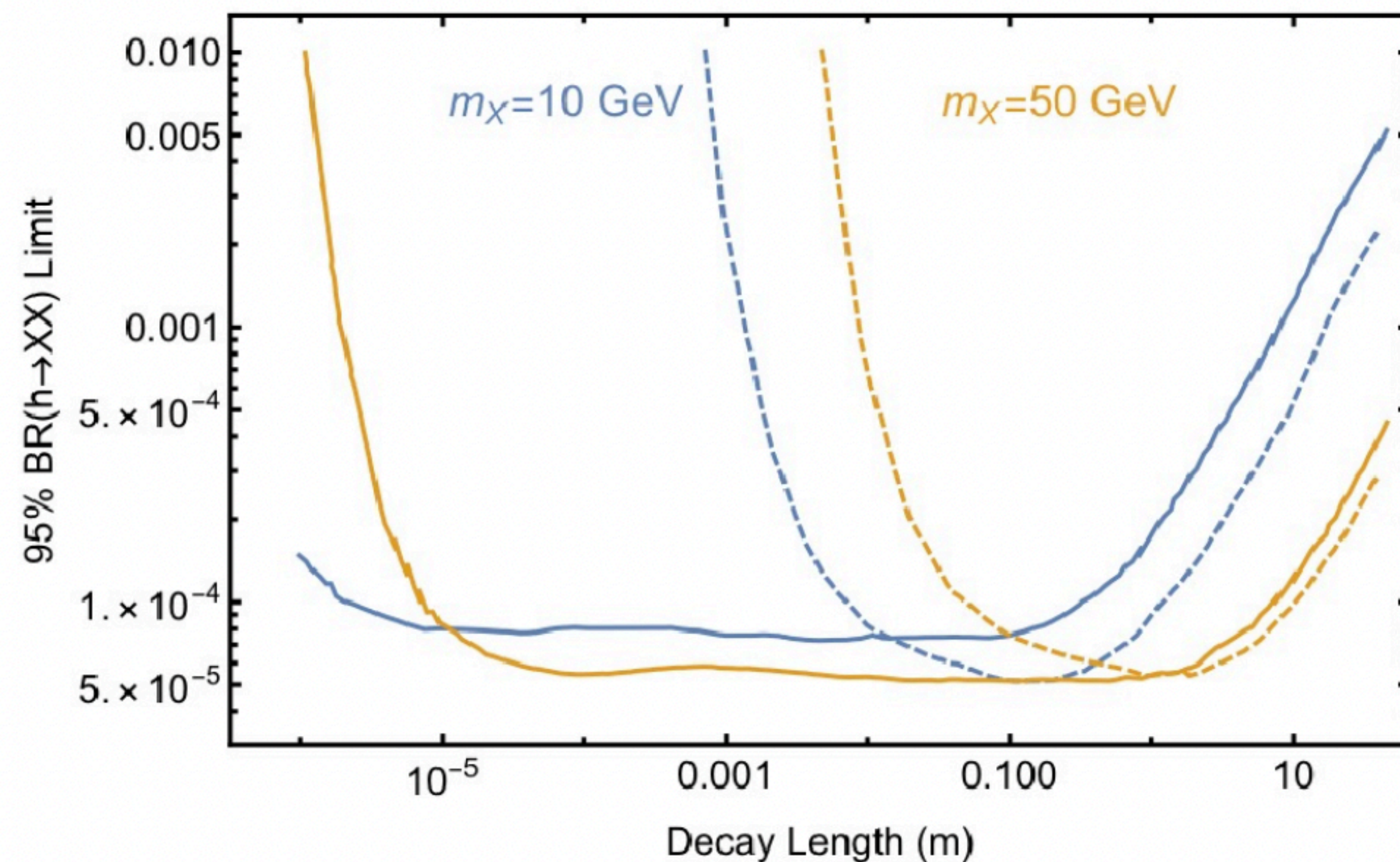
extending current limits > than 3 orders of magnitude

- Unconstrained mass range (0.1 – 100 GeV) accessible via $e^+e^- \rightarrow a\gamma$ and $e^+e^- \rightarrow a, a \rightarrow \gamma\gamma$. Sensitivity to couplings to $< 10^{-4} \text{ TeV}^{-1}$
- At small coupling values \rightarrow macroscopic decay lengths



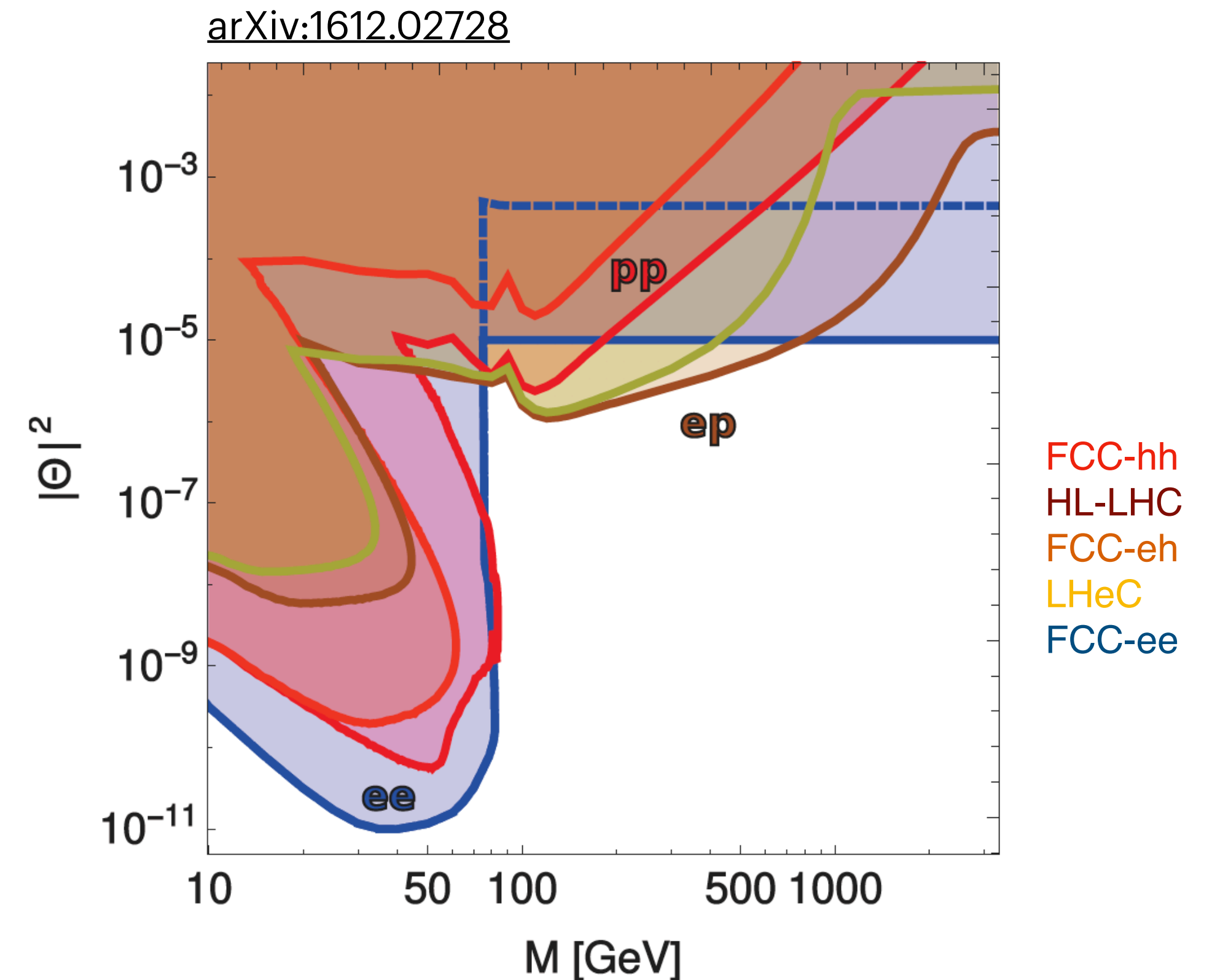
EXOTIC HIGGS BOSON DECAYS

- Possible and present in many models:
 - SM extensions with scalars/fermions/ vectors, MSSM, NMSSM, Hidden Valleys, Twin Higgs (arXiv:1312.4992, arXiv:1812.05588, arXiv:1712.07135)

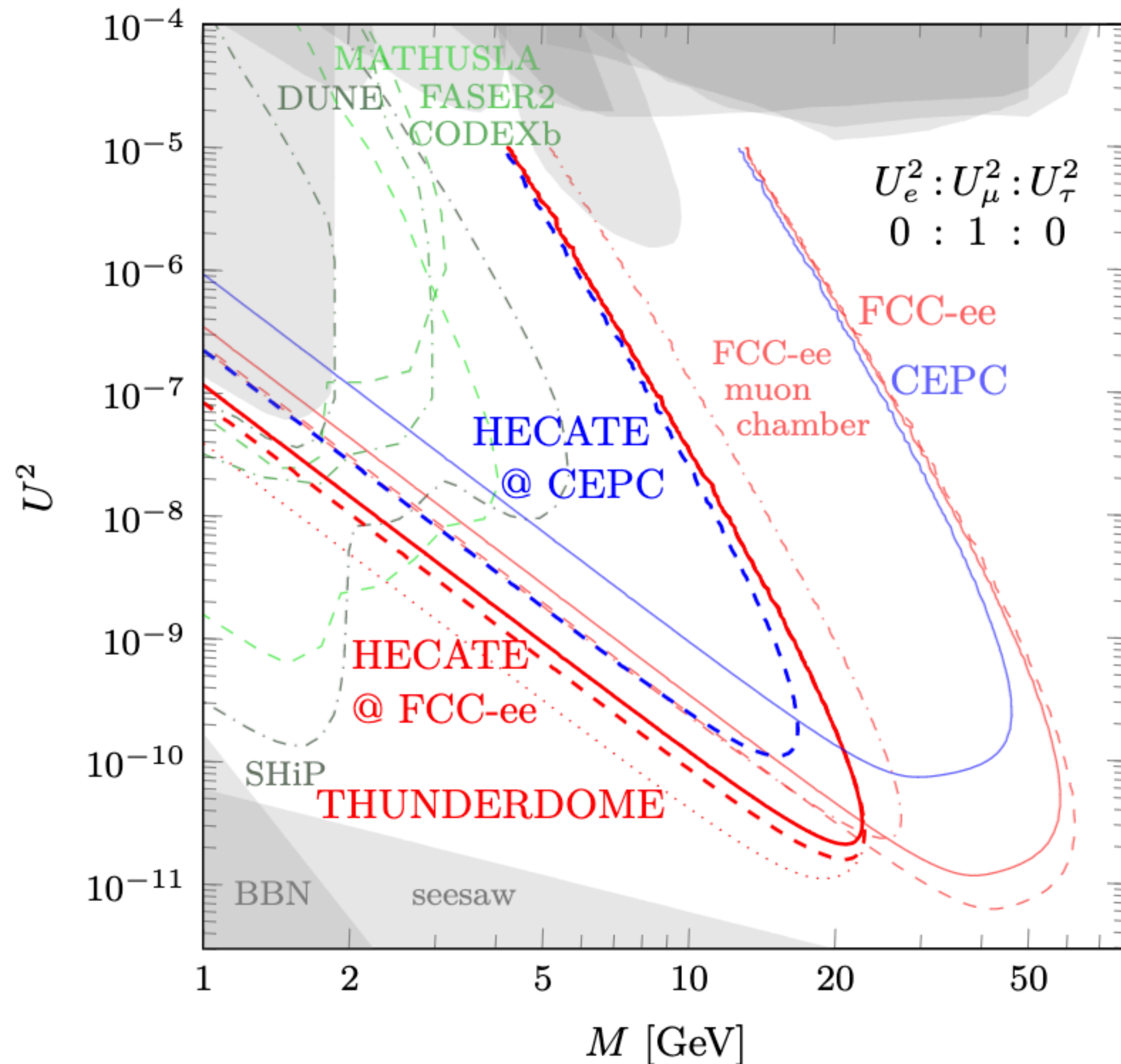


SYNERGIES ACROSS STAGES, RUNNING SCENARIOS

- In FCC-ee: precision measurements in W and Z will have a positive impact in the Higgs couplings measurements in the Higgs runs
- Precision Higgs measurements in FCC-ee will have an impact in the Higgs self-coupling studies in FCC-hh
- Coverage for prompt and long-lived searches will be very different in the different stages



EXTRA DETECTORS!



- Following the plans for different additional LLP experiments at the HL-LHC it is possible to also envision similar concepts at other future colliders
- **HECATE: A long lived particle detector concept for the FCC-ee or CEPC:**
[arXiv:2011.01005](https://arxiv.org/abs/2011.01005)
- The civil engineering of the FCC-ee will have much bigger detector caverns than needed for a lepton collider (to use them further for a future hadron collider)
- We could install extra instrumentation at the cavern walls to search for new long lived particles

FAr Detectors
[arXiv:1911.06576](https://arxiv.org/abs/1911.06576)
 for ALPs at FCC-ee, CepC
[arXiv:2201.08960](https://arxiv.org/abs/2201.08960)

BUT THIS IS JUST ONE OPTION

Timeline for the update of the European Strategy for Particle Physics



LET'S RECAP

WE HAVE EXPLORED

- We had a look at our current best tool for discovery and precision measurements (the LHC)
- We reviewed briefly how the general purpose LHC experiments work
- And how they take data
- We have taken stock on how much data we have available and how much is coming
- We had gone back in time to the early LHC data
- We have looked at the early SM measurements
- We started simple and by now have entered precision regime in many areas
- We mentioned the Higgs boson discovery

BUT TOGETHER WITH THE HIGGS: NO OTHER PARTICLE

- The LHC is made for searches and we certainly are searching
- We started by the easier cases and built to more and more complicated ones
- We are setting limits and cornering many new physics models
- We are on the brink of the high-luminosity upgrade of the LHC
- We will get enough lumi to see Higgs bosons produced in pairs
- There are exciting options for searches too
- However, the HL-LHC will be limited
- And so we are planning on the next collider



It could be YOUR
collider, get involved!

CREDITS

Future colliders:

<https://ecfa.web.cern.ch/>

<https://fcc-ped.web.cern.ch/>

<https://clicdp.web.cern.ch/>

<https://linearcollider.org/>

<http://cepc.ihep.ac.cn/>

<https://muoncollider.web.cern.ch/>

Most of the pictures in this lecture belong to CERN and can be found in the document server:

<https://cds.cern.ch/>

<https://home.cern/>

Physics results from the main LHC experiments can be access in their websites:

<https://atlas.cern/>

<https://cms.cern/>

<http://lhcb.web.cern.ch/>

<https://alice.cern/>

And the public repositories of results:

<https://arxiv.org/>

<https://inspirehep.net/>

I have re-used some of my material from other lectures (at schools and in Uppsala University) and slides presented at seminars/workshops/conferences as well as outreach talks.

Credits/Gratitude to (in no particular order): Katharina Anthony, Giulia Ripellino, Anna Sfyrla, Christophe Grojean, Patrick Janot, Ayres Freitas, Christoph Paus, Roberto Tenchini, Patrizia Azzi, Fabiola Gianotti, Sarah Williams, Juliette Alimena, Frank Zimmermann, Michele Selvaggi, Matthew McCullough and many others!