

Finite element based computation of beam coupling impedances



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Wakefield computation in the time domain

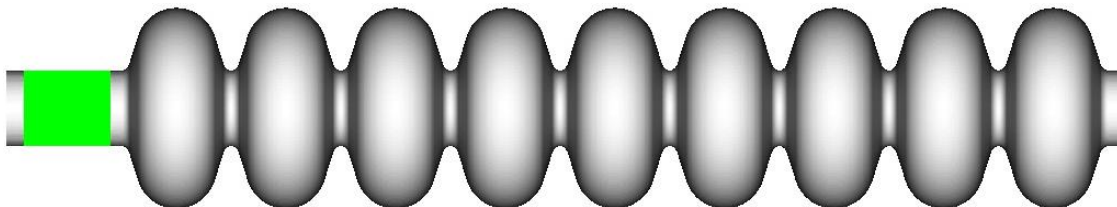
- The wakefield problem in the time domain

$$W_{\parallel}(r, s) = \frac{1}{Q} \int dz E_z(r, z, t(s, z)) \quad t(s, z) = \frac{s + z}{c}$$

- Solve time-dependent Maxwell's eqs. with beam current excitation
- Get impedance by Fourier transform

$$Z_{\parallel}(r, \omega) = -\frac{1}{c} \frac{1}{\tilde{\lambda}(\omega)} \int ds W_{\parallel}(r, s) \exp\left(-i \frac{\omega}{c} s\right)$$

- For short-range wakes: moving window / dispersion-free computation

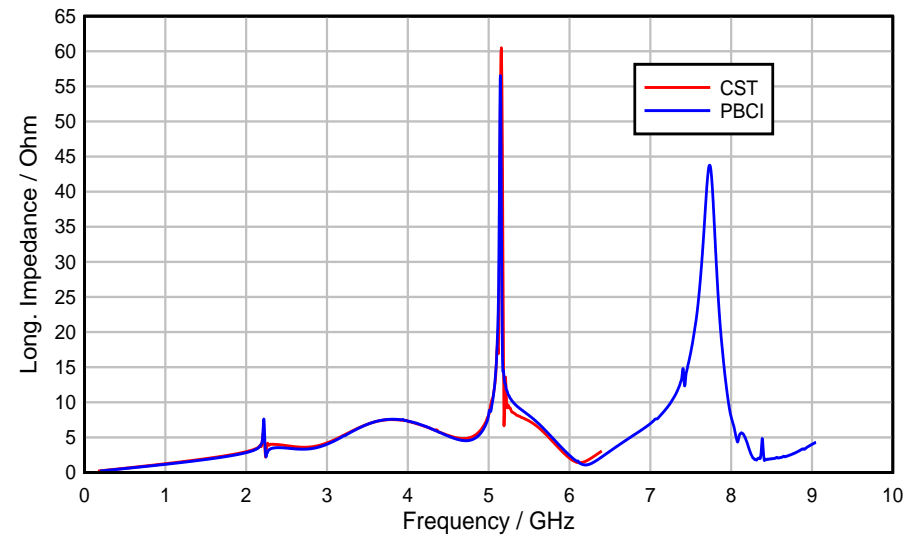
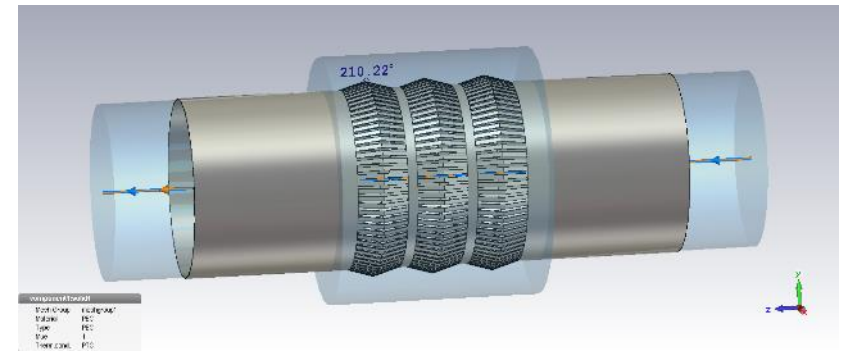
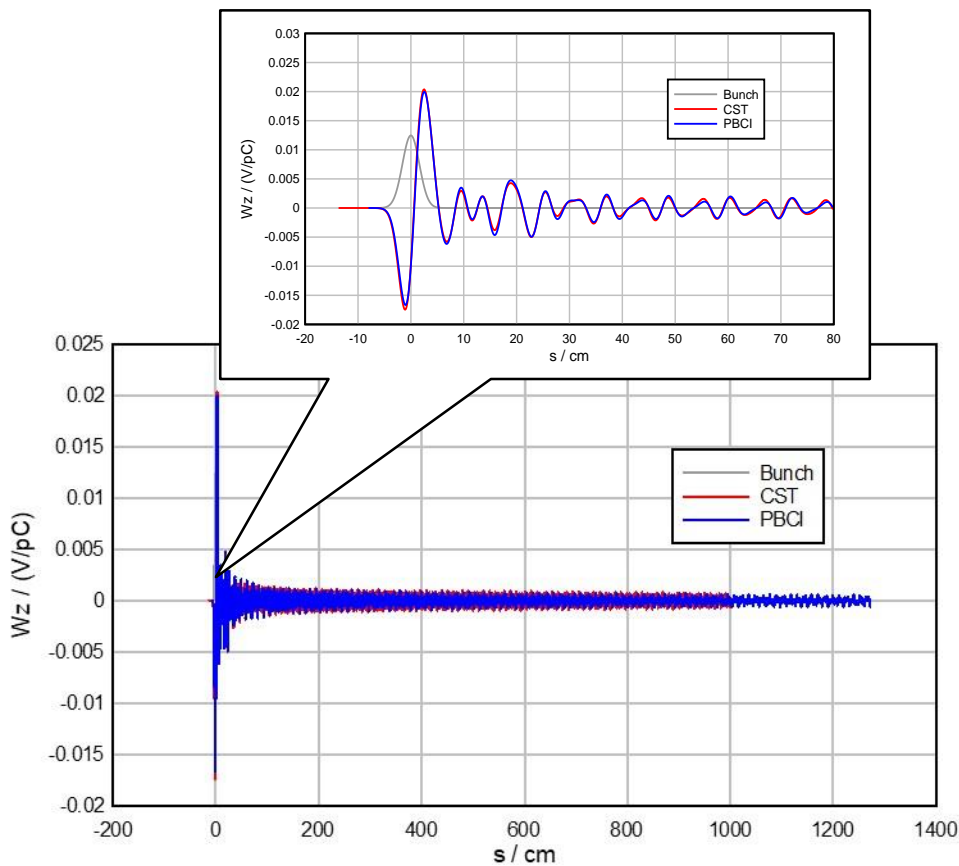


PBCI simulation: wakefield
in a TESLA cavity

E. Gjonaj et al., ICAP'06, MOM2IS02

Wakefield computation in the time domain

▪ LHC RF-Fingers



Contents

- A high order FE method for beam impedance computations
- Shielded CSR wakes
- Generalized S-Matrix formulation
- Lumped model cavity optimization
- Including external loads
- Domain decomposition approach

FE method for beam impedance

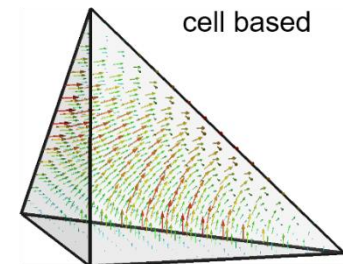
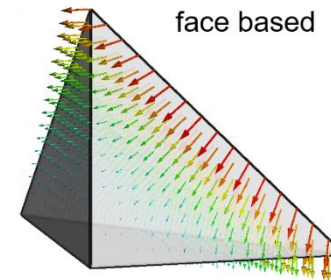
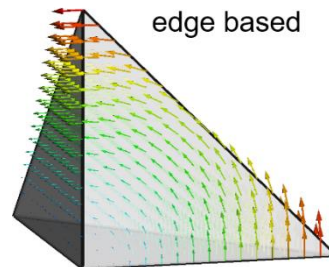
- The frequency domain problem

$$\nabla \times \mu^{-1} \nabla \times E - k_0^2 \varepsilon E = -jk_0 Z_0 J_s \quad J_s(x, y, z, \omega) = \delta(x - x_0) \delta(y - y_0) e^{-i\frac{\omega}{v}z}$$

- Weak formulation: find $E \in H(\text{curl})^*$ such that:

$$\int dV \mu^{-1} \nabla \times E \cdot \nabla \times v_h - k_0^2 \int dV \varepsilon E \cdot v_h = -jk_0 Z_0 \int dV J_s \cdot v_h$$

$$+ \underbrace{\oint_S dS \, n \cdot [v_h \times \mu^{-1} \nabla \times E]}_{\text{boundary term}}$$



FE method for beam impedance

- Treatment of boundary surfaces

$$\int dV \mu^{-1} \nabla \times E \cdot \nabla \times v_h - k_0^2 \int dV \varepsilon E \cdot v_h =$$
$$-jk_0 Z_0 \int dV J_s \cdot v_h + \underbrace{\int_{S_{SIBC}} dS \mathbf{n} \cdot [v_h \times \mu^{-1} \nabla \times E]}_{\text{resistive wall}} + \underbrace{\int_{S_{SWG}} dS \mathbf{n} \cdot [v_h \times \mu^{-1} \nabla \times E]}_{\text{in \& outgoing pipes}}$$

- Surface impedance

$$\oint_{S_{SIBC}} dS \mathbf{n} \cdot [v_h \times \mu^{-1} \nabla \times E] = \dots = j\omega \mathbf{Y}_S(\omega) \oint_{S_{SIBC}} dS v_h \cdot [n \times n \times E]$$

- Simple modification of the system matrix on SIBC surfaces
- No fitting of the surface impedance function or ADE/convolution is needed

FE method for beam impedance

- Treatment of boundary surfaces

$$\int dV \mu^{-1} \nabla \times E \cdot \nabla \times v_h - k_0^2 \int dV \varepsilon E \cdot v_h =$$

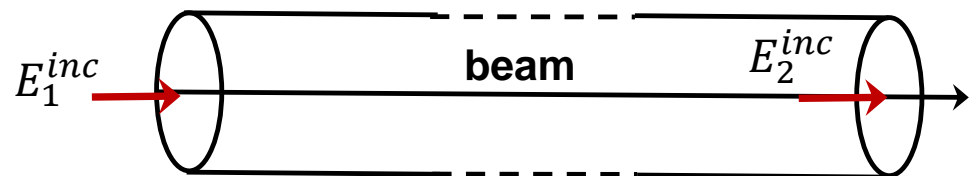
$$-jk_0 Z_0 \int dV J_s \cdot v_h + \underbrace{\int_{S_{SIBC}} dS n \cdot [v_h \times \mu^{-1} \nabla \times E]}_{\text{resistive wall}} + \underbrace{\int_{S_{SWG}} dS n \cdot [v_h \times \mu^{-1} \nabla \times E]}_{\text{in \& outgoing pipes}}$$

- Beam pipes

$$n \times \nabla \times E = n \times \nabla \times E^{inc} + \sum_m a_m^{TE} \gamma_m^{TE} e_m^{TE} + \sum_m a_m^{TM} \frac{-k_0^2}{\gamma_m^{TM}} e_m^{TM}$$

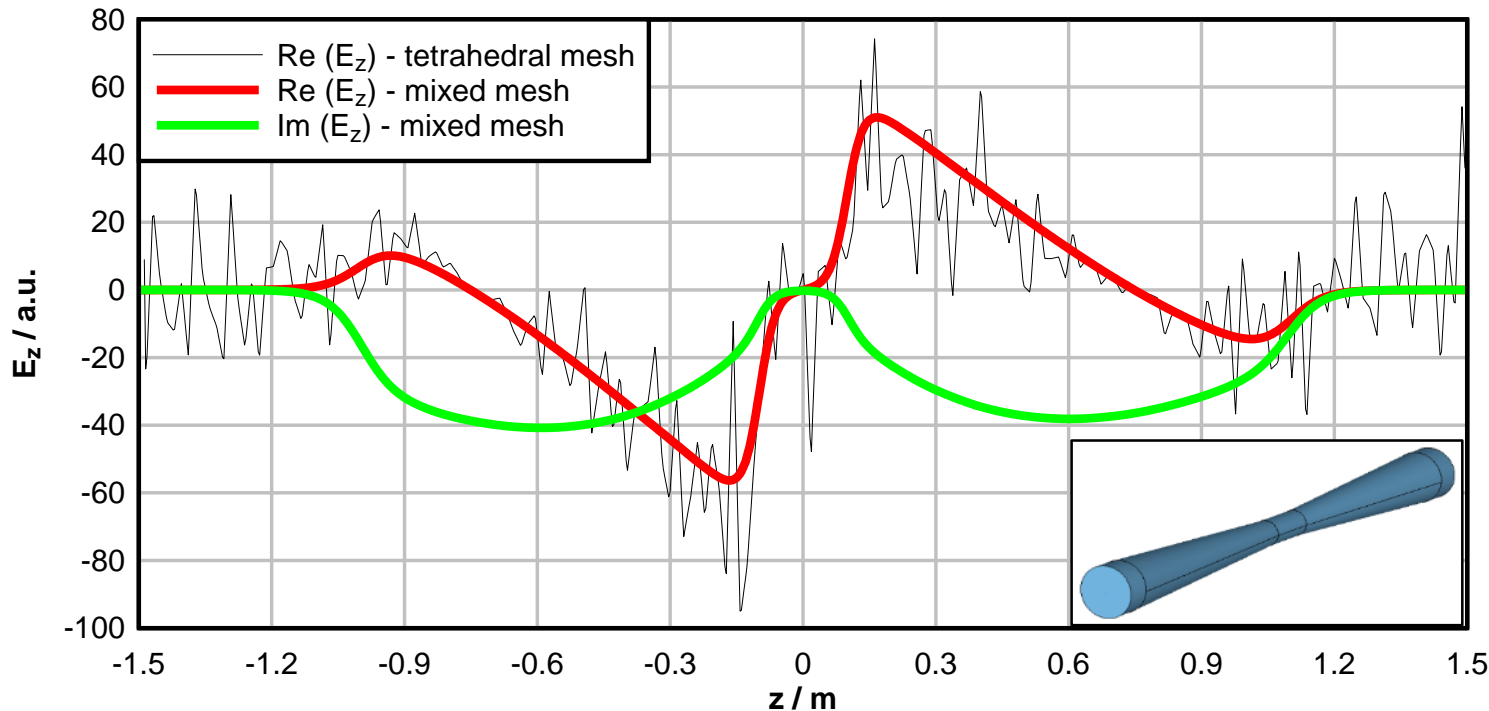
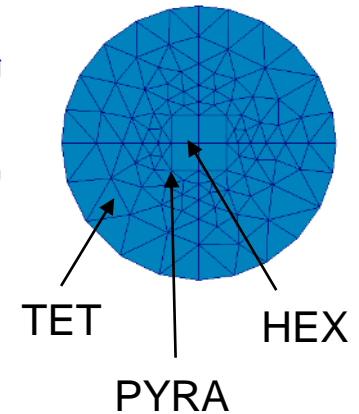
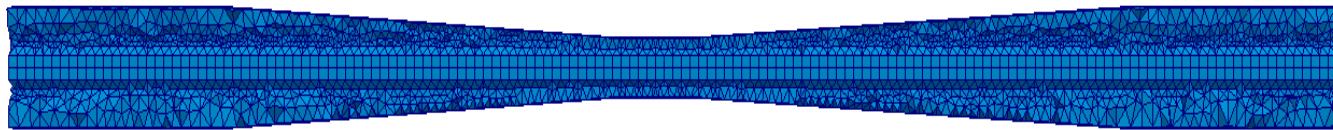
$$a_m^{TE} = \int_{S_{SWG}} dS e_m^{TE} \cdot [E - E^{inc}]$$

$$a_m^{TM} = \int_{S_{SWG}} dS e_m^{TM} \cdot [E - E^{inc}]$$



FE method for beam impedance

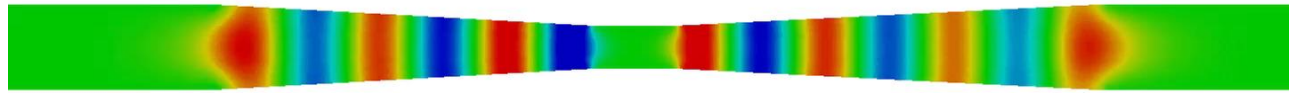
Hybrid meshes – collimator example



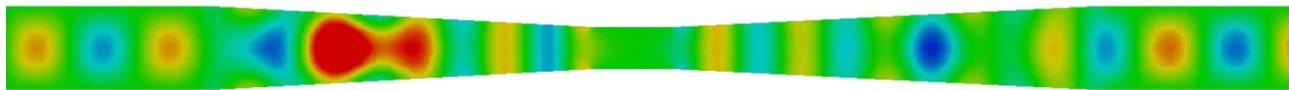
Longitudinal
wakefield on axis
at 100MHz

FE method for beam impedance

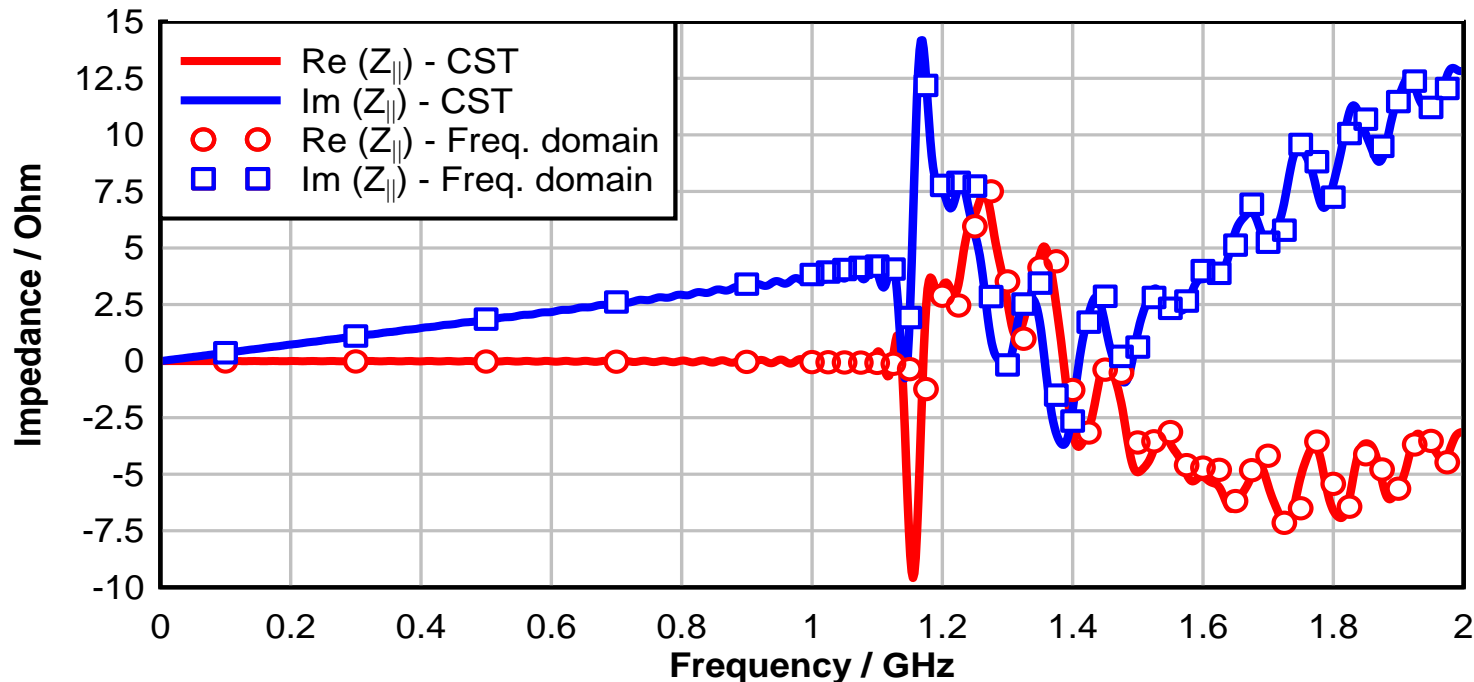
- Hybrid meshes – collimator example



E_z – 1GHz



E_z – 1.5GHz

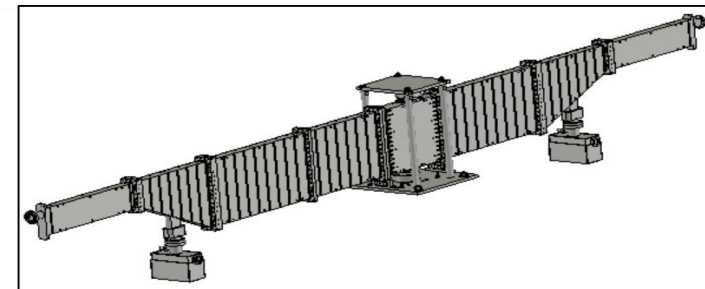
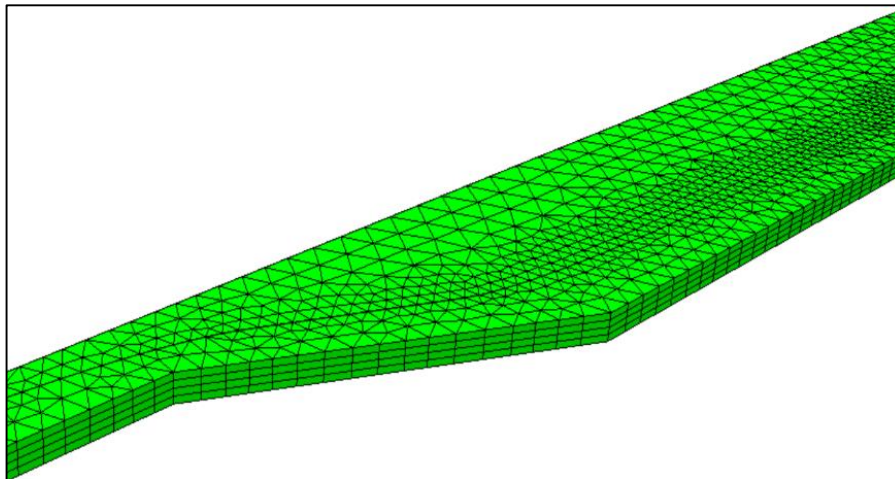
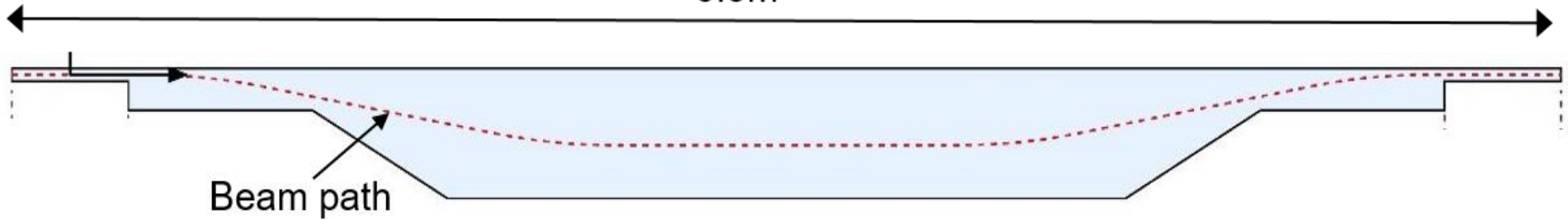


Comparison with
CST

Shielded CSR wakes

- DESY-EuXFEL bunch compressor (BC1)

~6.5m



Prismatic element mesh:

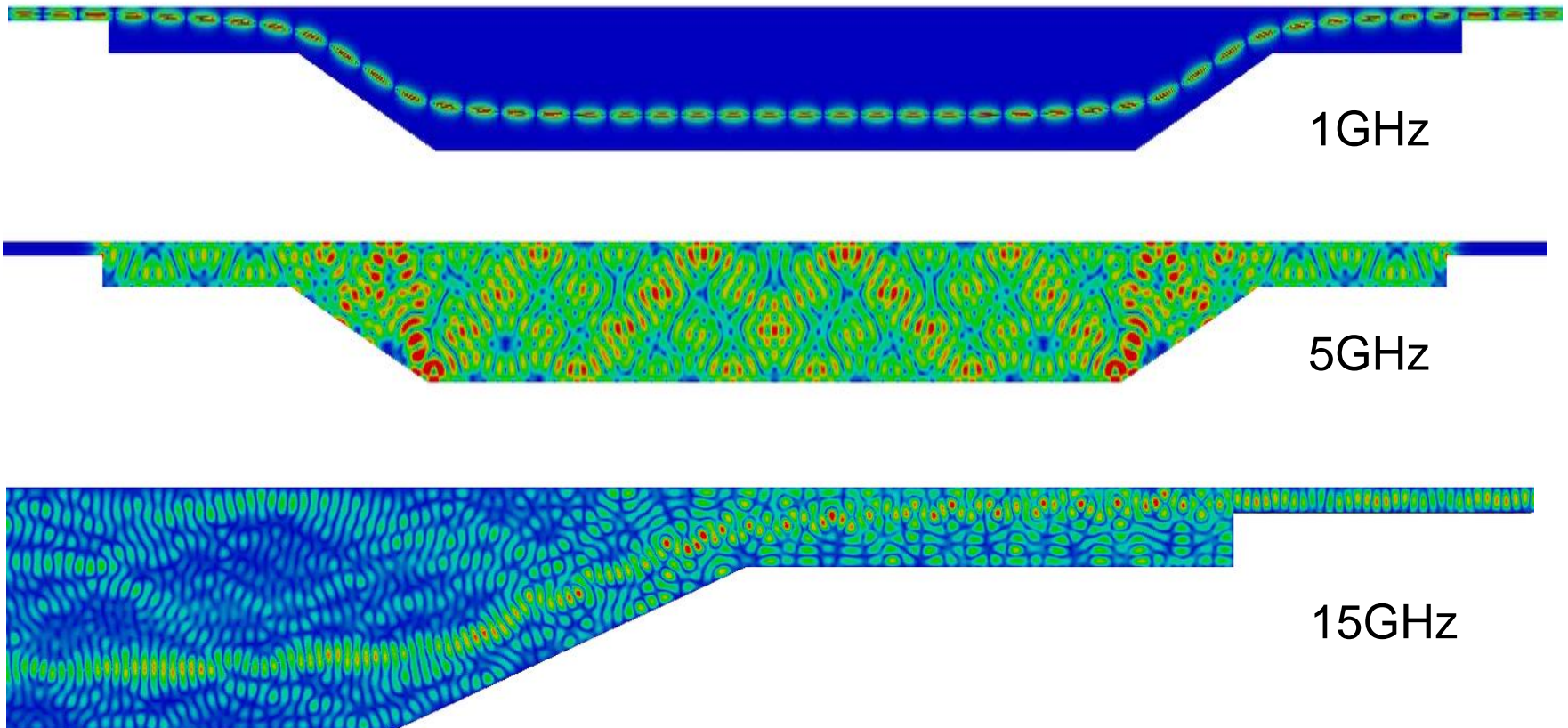
$\Delta \approx 5\text{mm}$

600k cells

4th order FEM

Shielded CSR wakes

- DESY-EuXFEL bunch compressor (BC1)

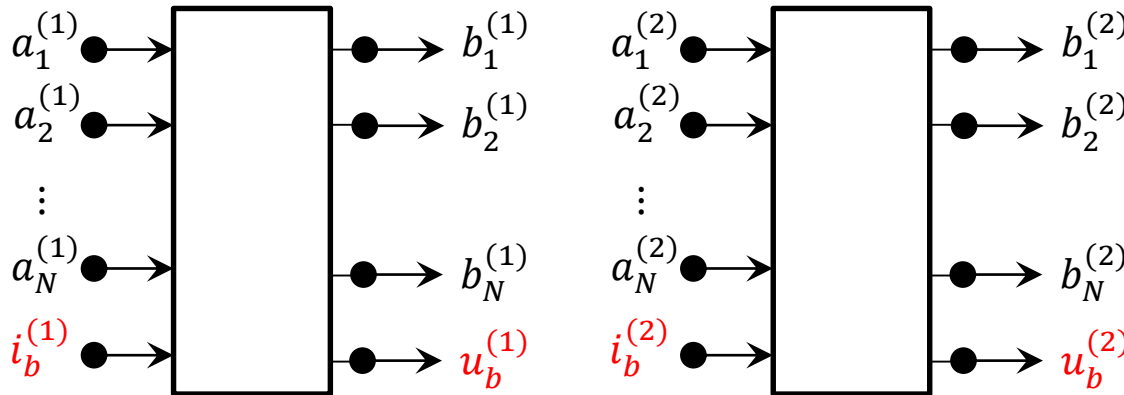


Generalized S-Matrix

- Scattering matrix with beam:

$$\begin{pmatrix} S_{11} & S_{12} \\ S_{21} & S_{22} \end{pmatrix} \begin{pmatrix} a_1 \\ a_2 \end{pmatrix} = \begin{pmatrix} b_1 \\ b_2 \end{pmatrix} \Rightarrow \boxed{\begin{pmatrix} \tilde{S} & \tilde{k} \\ \tilde{h} & Z_b \end{pmatrix} \begin{pmatrix} a_m \\ i_b \end{pmatrix} = \begin{pmatrix} b_m \\ u_b \end{pmatrix}}$$

- Concatenation of cascaded structures:



Matching conditions:

$$b_i^{(n)} = a_i^{(n+1)}$$

$$i_b^{(n)} = i_b^{(n-1)} e^{ik_0 L_{n-1}}$$

$$\sum_n u_b^{(n)} = u_b^{tot}$$

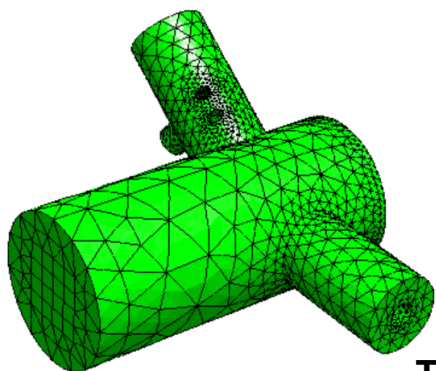
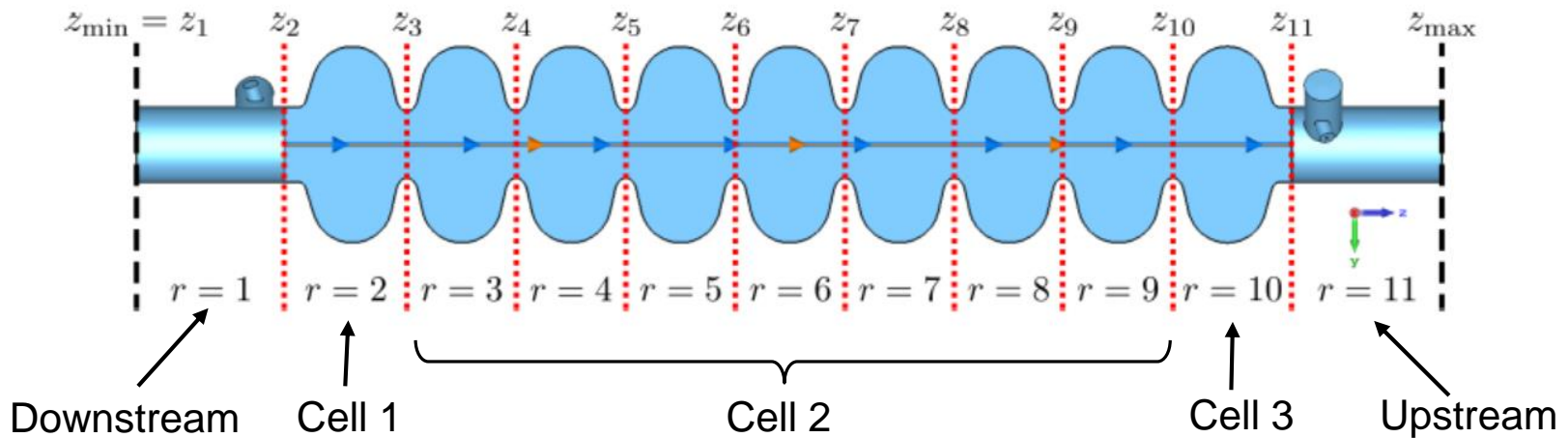
- Total scattering matrix:

$$\boxed{\begin{pmatrix} \tilde{S}^{tot} & \tilde{k} \\ \tilde{h} & Z_b^{tot} \end{pmatrix} \begin{pmatrix} a_m \\ i_b \end{pmatrix} = \begin{pmatrix} b_m \\ u_b^{tot} \end{pmatrix}}$$

Flisgen; Gjonaj, et al.:
PRAB (2020)

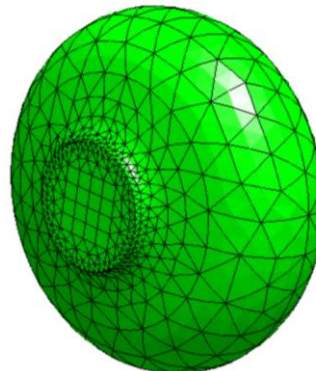
Generalized S-Matrix

- Tesla 1.3GHz cavity

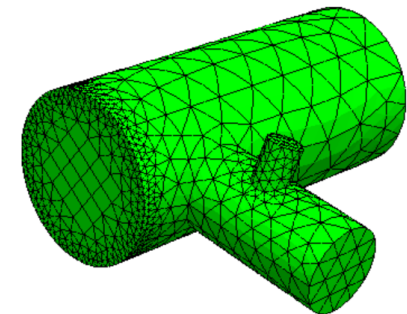


15 TE-Modes
($\omega_{max} = 8.2\text{GHz}$)
15 TM-Modes
($\omega_{max} = 10.6\text{GHz}$)

TEM, ...



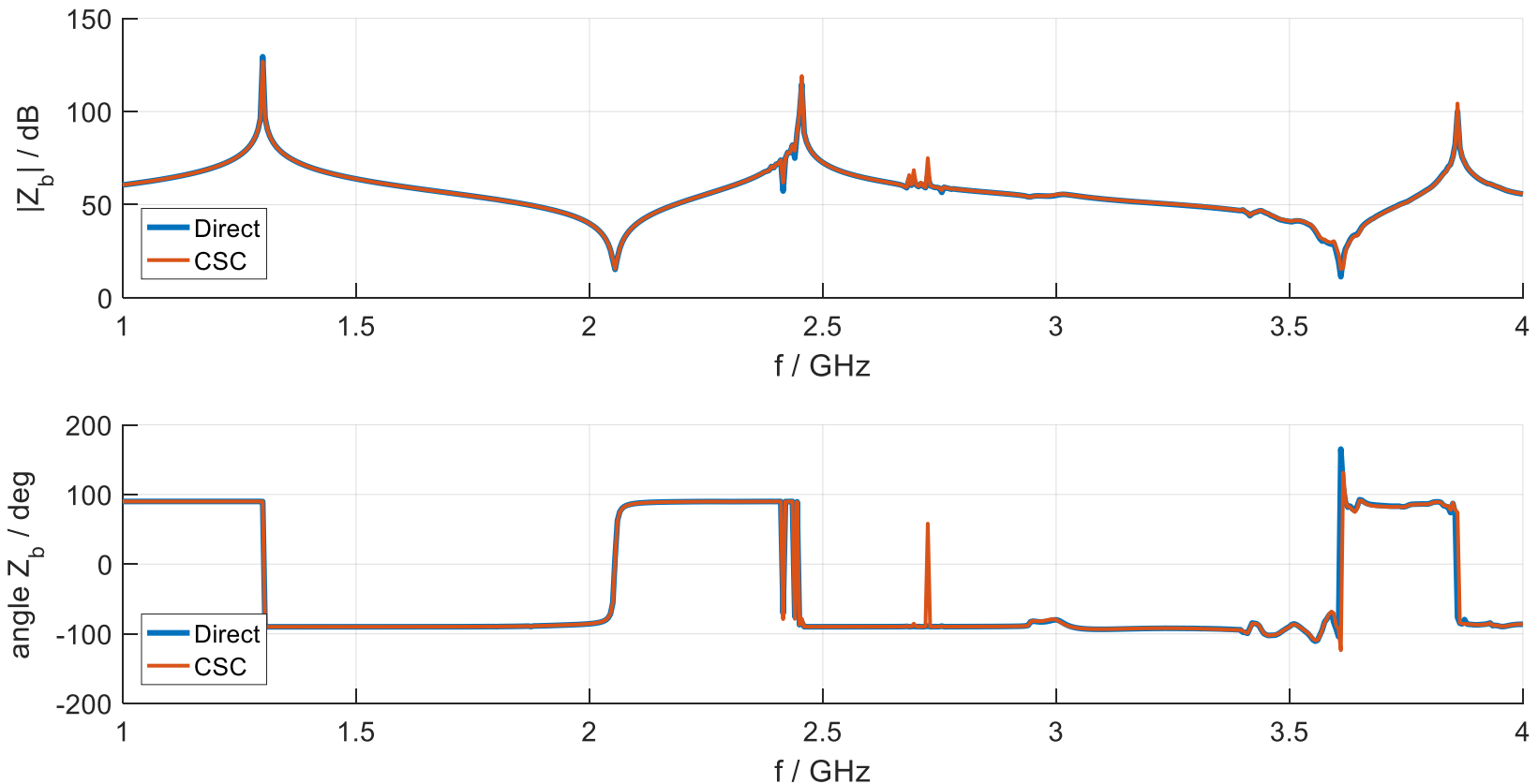
15 TE-Modes
($\omega_{max} = 8.2\text{GHz}$)
15 TM-Modes
($\omega_{max} = 10.6\text{GHz}$)



Generalized S-Matrix

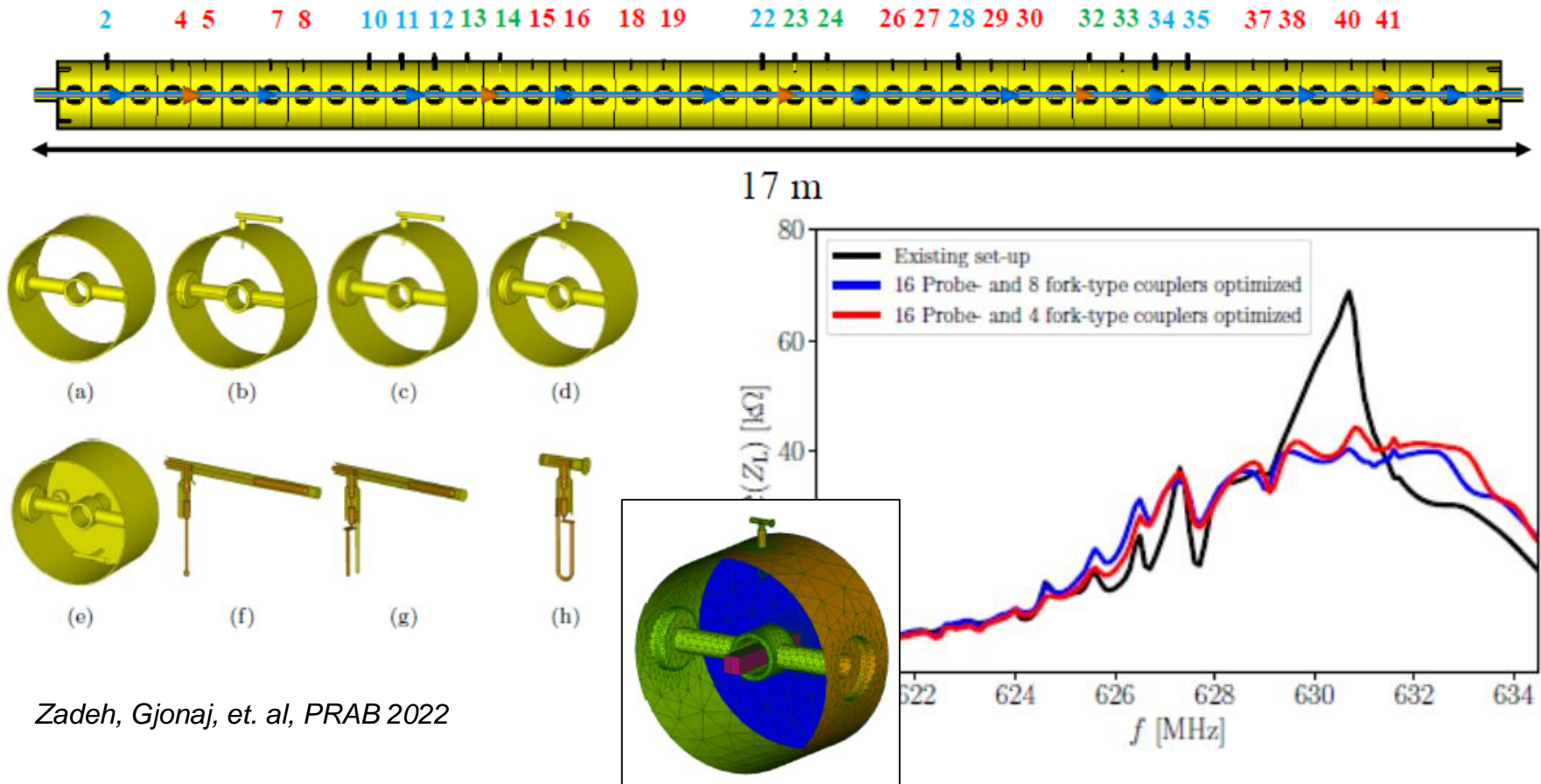
- Tesla 1.3GHz cavity

30 Modes



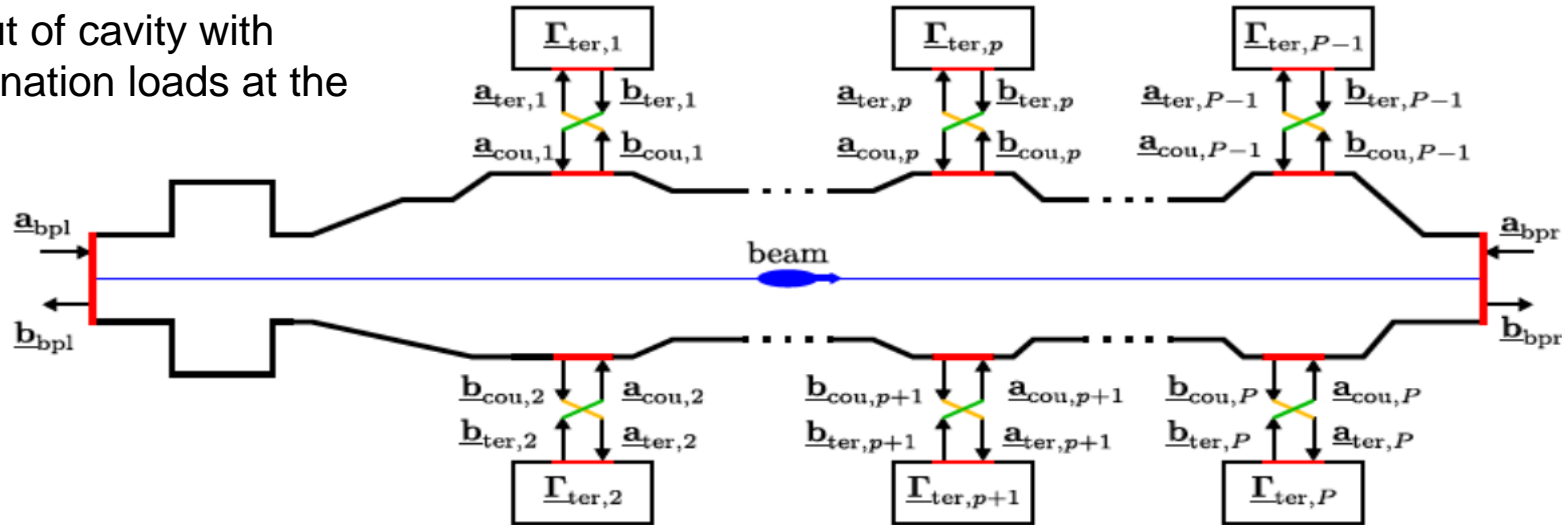
Lumped model cavity optimization

- SPS 44-cell 200MHz TW cavities



Including termination loads

General layout of cavity with external termination loads at the coupler ports



$$\left\{ \begin{array}{l} \begin{pmatrix} \tilde{S} & \tilde{k} \\ \tilde{h} & Z_b \end{pmatrix} \begin{pmatrix} a_m \\ i_b \end{pmatrix} = \begin{pmatrix} b_m \\ u_b \end{pmatrix} \\ a_{cou} = b_{ter} \\ a_{ter} = b_{cou} \\ Ra_{ter} = b_{ter} \end{array} \right.$$

$R = \text{diag}\{\Gamma_{ter,1}, \Gamma_{ter,2}, \dots\}$ termination reflection matrix

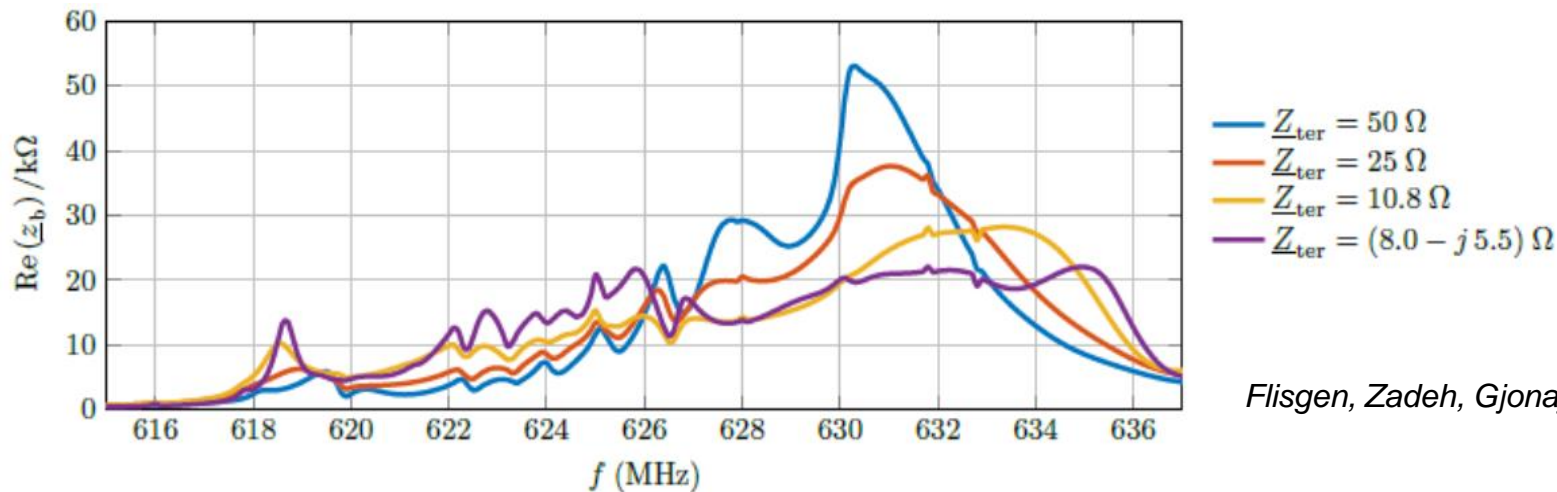
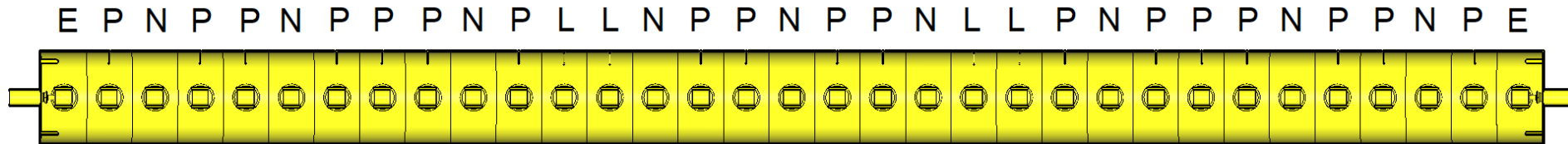
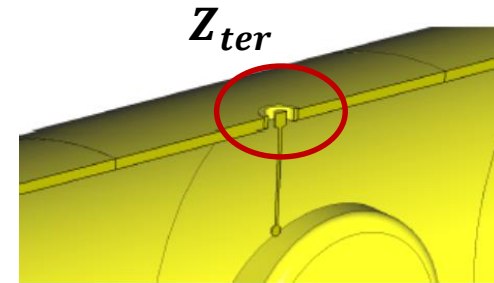
$$\tilde{G}_{mod} \begin{pmatrix} a_m \\ i_b \end{pmatrix} = \begin{pmatrix} b_m \\ u_b \end{pmatrix}$$

scattering matrix with
termination + beam impedance

Flisgen, Zadeh, Gjonaj, PRAB 2023

Including termination loads

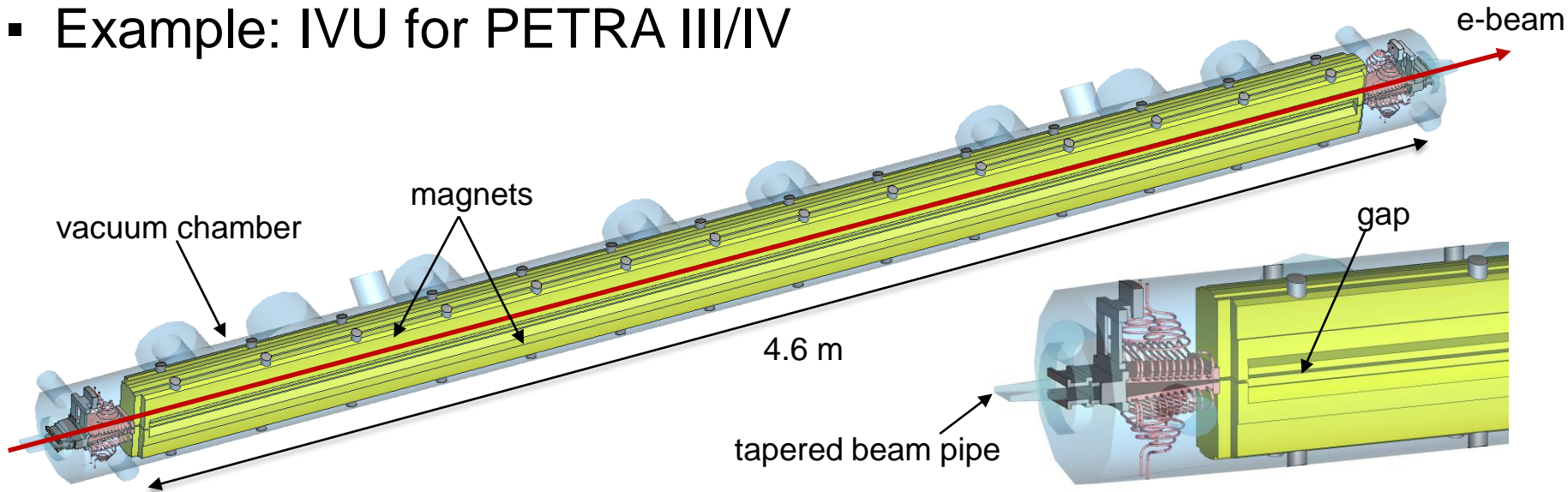
- SPS 33-cell 200 MHz TW-structure



Flisgen, Zadeh, Gjonaj, PRAB 2023

Domain decomposition

▪ Example: IVU for PETRA III/IV

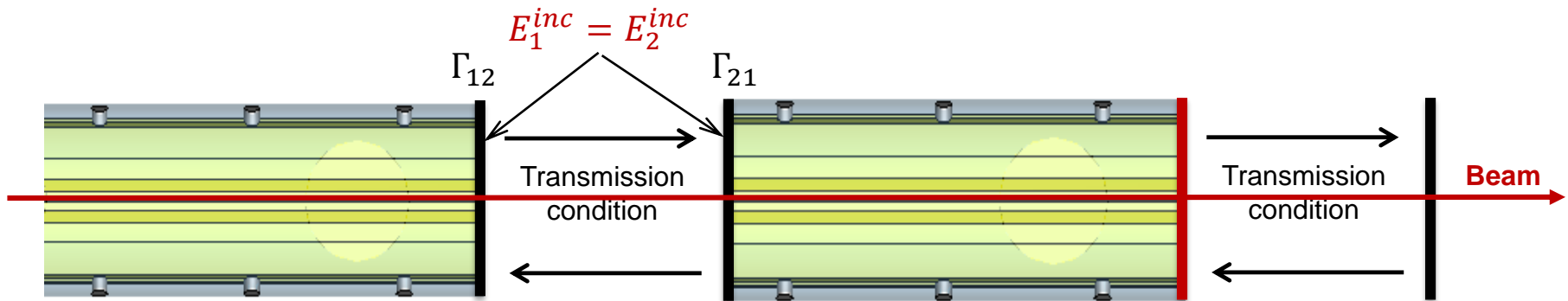


Order	# DoFs	Order	# DoFs
0	0.9e6	3	19e6
1	2e6	4	38e6
2	7e6	5	

- Multiple trapped modes at 0.1-1GHz
- High-Q dominated by wall losses
- Local heat spots / surface fields
- ~1500GB RAM
- ~ weeks of CPU-time

Domain decomposition

- Decompose geometry by suitable boundary conditions



$$F_{12}^{out} = \underbrace{\mathcal{S}(n_1 \times n_1 \times E_1')}_{\text{absorbing operator}} + n_1 \times (\nabla \times E_1') = \underbrace{\mathcal{S}(n_2 \times n_2 \times E_2') - n_2 \times (\nabla \times E_2')}_{\text{one-way wave operator}} = F_{21}^{in} \quad \text{on } \Gamma_{12}$$

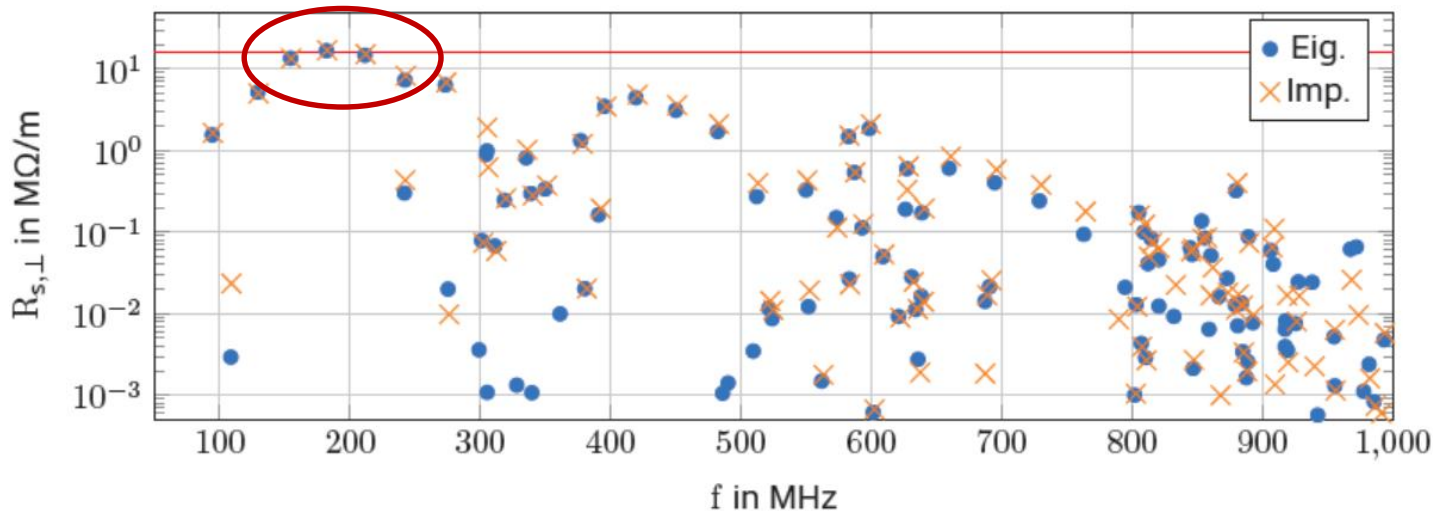
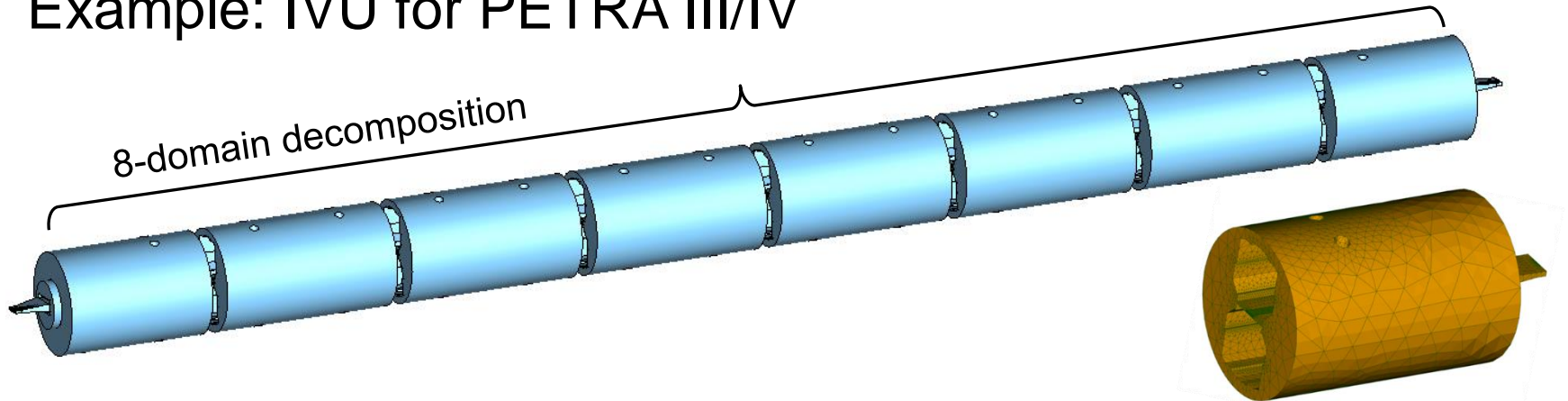
with iteration prescription: $\begin{cases} F_{21}^{in}(n+1) = F_{12}^{out}(n) \\ F_{12}^{in}(n+1) = F_{21}^{out}(n) \end{cases}$ (using fixed-point or a Krylov iteration)

Apply TC using waveguide port operator*: $F_{12}^{out} = P(E_1 - E_1^{inc}) + n_1 \times (\nabla \times (E_1 - E_1^{inc}))$

*Gjonaj, et al.: IPAC'24, THPC62

Domain decomposition

- Example: IVU for PETRA III/IV



Transverse shunt
impedances

Summary & Conclusions

- The Finite-Element-Frequency-Domain approach
 - Fills the gap for some important wakefield / impedance problems
 - Complicated chamber geometries
 - Resonant structures → long-range wakefields
 - Resistive, rough surfaces, dispersive materials, waveguide openings
 - Curved beam trajectories and CSR
 - Allows lumped parameter modeling
 - Concatenation of large structures by generalized S-Matrix
 - Fast cavity impedance optimization including termination loads
 - Quasi-periodic THz structures
 - Limitation: huge size of 3D discrete problems for ultra-high frequencies
 - Parallel iterative solvers by non-overlapping DDM

