

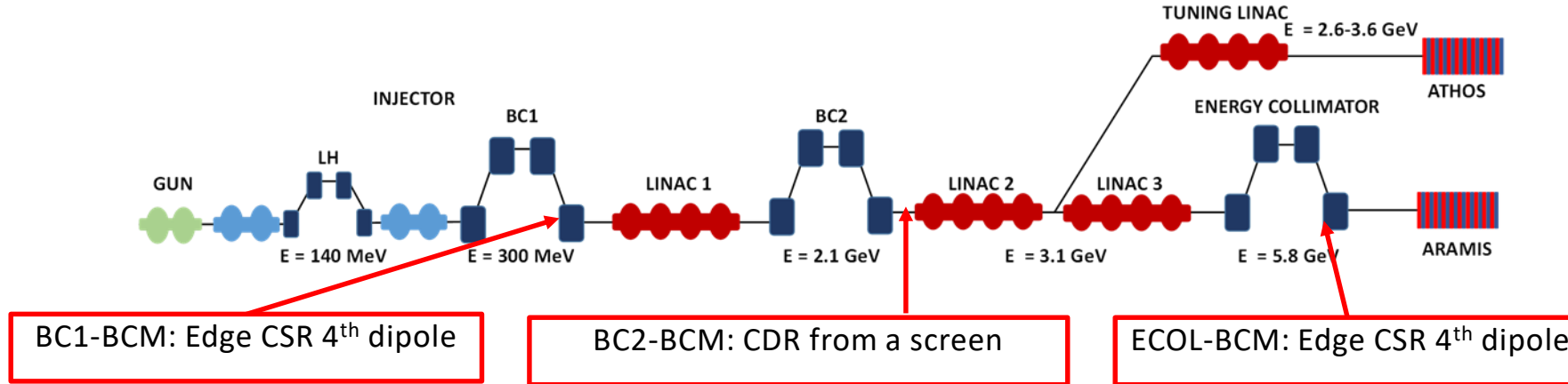


Sub-fs bunch length measurements with ECOL-BCM during 10 pC operations

Gian Luca Orlandi
PSI, 18.06.2024

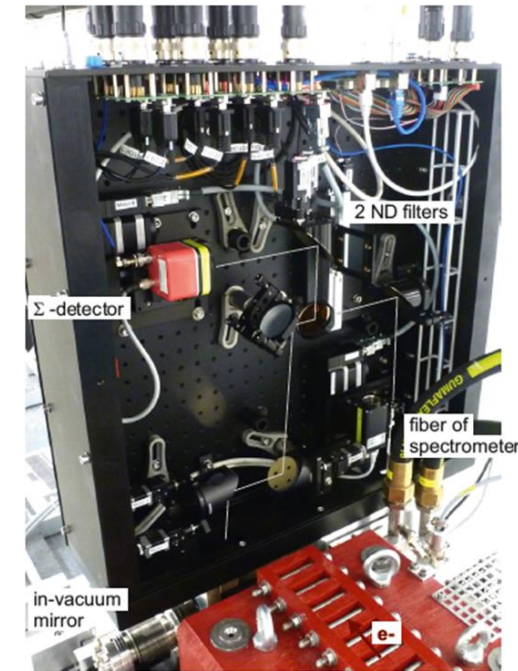
- ✓ ECOL-BCM: the experimental set-up
- ✓ 10 pC and sub-fs bunch-length SwissFEL set-up for Cristallina
- ✓ Method for absolute determination of bunch-length from “two-detector” BCM
- ✓ Preliminary experimental results from data acquired during the Cristallina set-up in March 2024
- ✓ Conclusions and Outlook

ECOL-BCM: the experimental set-up



ECOL-BCM (5.8GeV, Edge SR):

- Vacuum chamber cutoff of the radiation ($<4\mu\text{m}$)
- Design optimized for bunch length 0.7-3.0 fs (10 pC),
- Detectors:
 - Pyroelectric sensor (0.9-4.0) μm
 - Spectrometer (Ocean Optics, 0.9-2.5 μm)
- Output Signals:
 - Pyrodetector: voltage signal, radiation energy spectrum spectral fully integrated over the detector wavelength band
 - Spectrometer: 512-pixel array spectrogram








- Cristallina set-up 05.03.2024 (Prat, Reiche, Dijkstal et al.): 10 pC, 8.36 keV, 3-stage compression
- Bunch length measurements before ECOL with TDS-Cband in linac 3 : $\sigma(\text{rms}) \sim 15$ fs
- Bunch-length prediction downstream the ECOL: $\sigma(\text{rms}) \sim 0.56$ fs

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Rapid Communications

Single- and two-color attosecond hard x-ray free-electron laser pulses with nonlinear compression

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$$\sigma_{z,f} = c \frac{E_i}{E_f} R_{56(3)} \sigma_\delta, \quad (1)$$

where σ_δ is the uncorrelated energy spread before BC1, $E_{i,f}$ are the energies at BC1 and BC3, respectively, and $c = c_1 c_2$ is the accumulated compression before BC3, where $c_n = (1 + R_{56(n)} h_n)^{-1}$ are compression factors in the first and second stages ($c_{1,2}^{-1} \neq 0$).

Method for absolute determination of bunch-length from “two-detector” BCM



$$\frac{dI^{Ne}(\omega)}{d\omega} \simeq N(N-1)F(\omega)\frac{dI^e(\omega)}{d\omega} \simeq N^2F(\omega)\frac{dI^e(\omega)}{d\omega}$$

Radiation energy spectrum by a Ne⁻ bunch at the threshold of the temporal coherence enhancement upon integration over the solid angle of acceptance of the detector

$$F(\omega) = \left| \int_{-\infty}^{+\infty} e^{j\omega z/c} \rho_z(z) dz \right|^2 = e^{-(\frac{\omega\sigma}{c})^2}$$

Gaussian bunch form-factor (longitudinal), highly collimated beams as usual in a FEL

$$I = \int_{\omega_{min}}^{\omega_{max}} d\omega \left(\frac{dI^{Ne}(\omega)}{d\omega} / \frac{dI^e(\omega)}{d\omega} \right)$$

Radiation energy spectrum of the Ne⁻ is normalized w.r.t. the single particle energy spectrum $dI^e/d\omega$ and integrated over the frequency band of acceptance $\Delta\omega$ of the detector. Assumption: $dI^e/d\omega$ either constant or weakly varying over the wavelength band of the detector

$$\ln(I) = 2 \ln(N) + \ln\left(\int_{\omega_{min}}^{\omega_{max}} d\omega F(\omega) \right)$$

Calculation of the natural logarithm upon single particle normalization and integration of the spectrum over $\Delta\omega$

Method for absolute determination of bunch-length from “two-detector” BCM (*)



$$\frac{\Delta I}{I} = 2 \frac{\Delta N}{N} + \frac{\int_{\omega_{min}}^{\omega_{max}} d\omega [F^*(\omega) - F(\omega)]}{\int_{\omega_{min}}^{\omega_{max}} d\omega F(\omega)}$$

$$F^*(\omega) - F(\omega) = e^{-\left(\frac{\omega\sigma}{c}\right)^2 (1 + \frac{\Delta\sigma}{\sigma})^2} - e^{-\left(\frac{\omega\sigma}{c}\right)^2} \simeq -2 \frac{\Delta\sigma}{\sigma} \left(\frac{\omega\sigma}{c}\right)^2 e^{-\left(\frac{\omega\sigma}{c}\right)^2} = -2 \frac{\Delta\sigma}{\sigma} \left(\frac{\omega\sigma}{c}\right)^2 F(\omega)$$

$$G(\sigma, \Delta\omega) = \left\{ \frac{2\sigma}{\sqrt{\pi}c} \frac{\left[e^{-\left(\frac{\omega\sigma}{c}\right)^2} \omega \right]_{\omega_{min}}^{\omega_{max}}}{\left[erf(\omega\sigma/c) \right]_{\omega_{min}}^{\omega_{max}}} - 1 \right\}$$

$$\frac{\Delta I}{I} = 2 \frac{\Delta N}{N} + \frac{\Delta\sigma}{\sigma} G(\sigma, \Delta\omega)$$

$$\frac{\Delta I}{I} = \frac{I^* - I}{I}$$

$$\frac{\Delta N}{N} = \frac{N^* - N}{N}$$

$\Delta I/I$ and $\Delta N/N$: shot-to-shot relative signal variations from the one shot to another or relative statistical fluctuations of the signals w.r.t. a reference value: for instance, the mean value over a sequence of signal readouts acquired under a machine steady state regime.

(*) Orlandi, G.L. Absolute and non-invasive determination of the electron bunch length in a free electron laser using a bunch compressor monitor. Sci Rep 14, 6319 (2024). <https://doi.org/10.1038/s41598-024-56586-1>

Method for absolute determination of bunch-length from “two-detector” BCM



Two-detector BCM: absolute determination of bunch length σ vs $(\Delta I/I)_{j=1,2}$ and $\Delta N/N$

$$\left(\frac{\Delta I}{I}\right)_{j=1,2} = 2\frac{\Delta N}{N} + \frac{\Delta\sigma}{\sigma} G(\sigma, (\Delta\omega)_{j=1,2})$$

For a BCM equipped with two detectors simultaneously illuminated and with different wavelength band of acceptance $(\Delta\omega)_{j=1,2}$

$$\left[\left(\frac{\Delta I}{I}\right)_2 - 2\frac{\Delta N}{N}\right] = \frac{G(\sigma, (\Delta\omega)_2)}{[G(\sigma, (\Delta\omega)_2) - G(\sigma, (\Delta\omega)_1)]} \times \left[\left(\frac{\Delta I}{I}\right)_2 - \left(\frac{\Delta I}{I}\right)_1\right]$$

Formula for the shot-to-shot tracking of the absolute value σ of the bunch length

$$\backslash std \left(\left[\left(\frac{\Delta I}{I} \right)_2 - 2 \frac{\Delta N}{N} \right] \right) - \backslash abs \left(\frac{G(\sigma, (\Delta\omega)_2)}{[G(\sigma, (\Delta\omega)_2) - G(\sigma, (\Delta\omega)_1)]} \right) \times \backslash std \left(\left[\left(\frac{\Delta I}{I} \right)_2 - \left(\frac{\Delta I}{I} \right)_1 \right] \right) = 0$$

Formula for determining the absolute value the bunch length (σ) from the analysis of a sequence of data acquired under a steady state regime of the machine: I and N are the mean values of the BCM and charge readout of the sequence of acquired data.

The variable σ runs into a test interval until the value corresponding to the “zero” of the equation is found

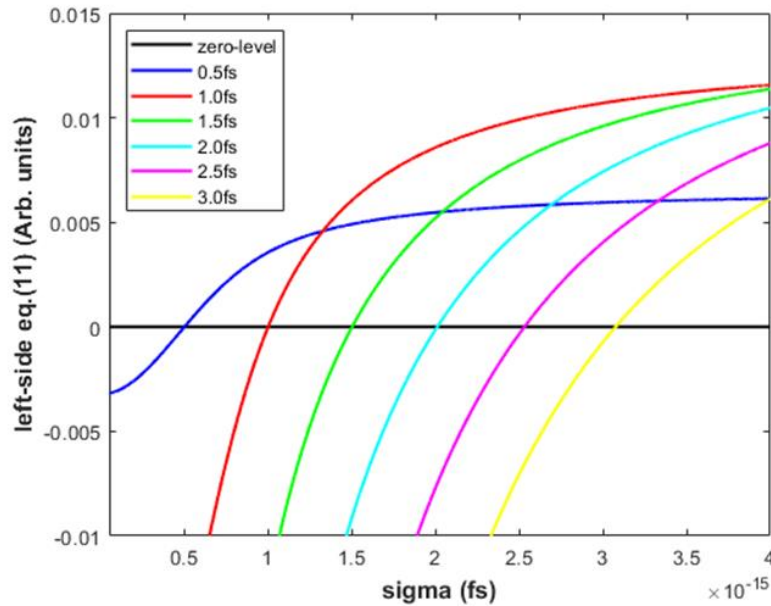
Method for absolute determination of bunch-length from “two-detector” BCM



Results from a numerical simulation based on ECOL-BCM (5.8GeV):

- Design optimized for bunch length 0.7-3 fs (10 pC),
- Detectors: Pyroelectric sensor (0.9-4.0)μm
Optical fiber spectrometer (Ocean Optics, 0.9-2.5 μm)

$$\backslash std \left(\left[\left(\frac{\Delta I}{I} \right)_2 - 2 \frac{\Delta N}{N} \right] \right) - \backslash abs \left(\frac{G(\sigma, (\Delta\omega)_2)}{[G(\sigma, (\Delta\omega)_2) - G(\sigma, (\Delta\omega)_1)]} \right) \times \backslash std \left(\left[\left(\frac{\Delta I}{I} \right)_2 - \left(\frac{\Delta I}{I} \right)_1 \right] \right) = 0$$

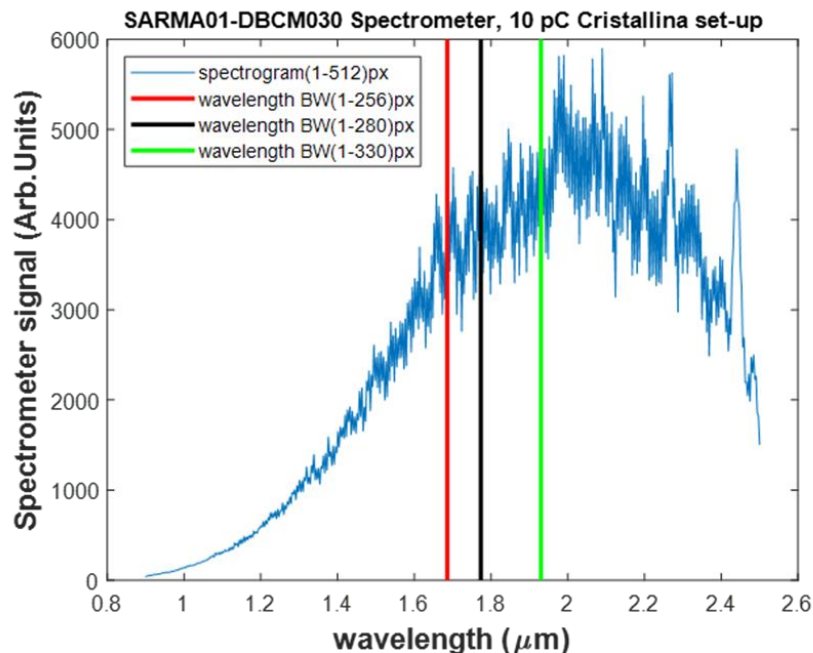


ECOL-BCM:

- numerical prediction of the absolute bunch length σ determination in case of Gaussian Form-Factor for different σ under a steady state machine regime.
- I and N are mean values over a sequence of acquired shots

the abscissae of the curve intercepts with the "zeros-level" line are the model-estimated absolute values of σ . Simulation settings: Gaussian Form Factor with $\sigma = 0.5-3.0$ fs; BCM signals (pyro & spectrometer) simulated with rms deviation of 1 % for $\Delta\sigma/\sigma$ and $\Delta N/N$

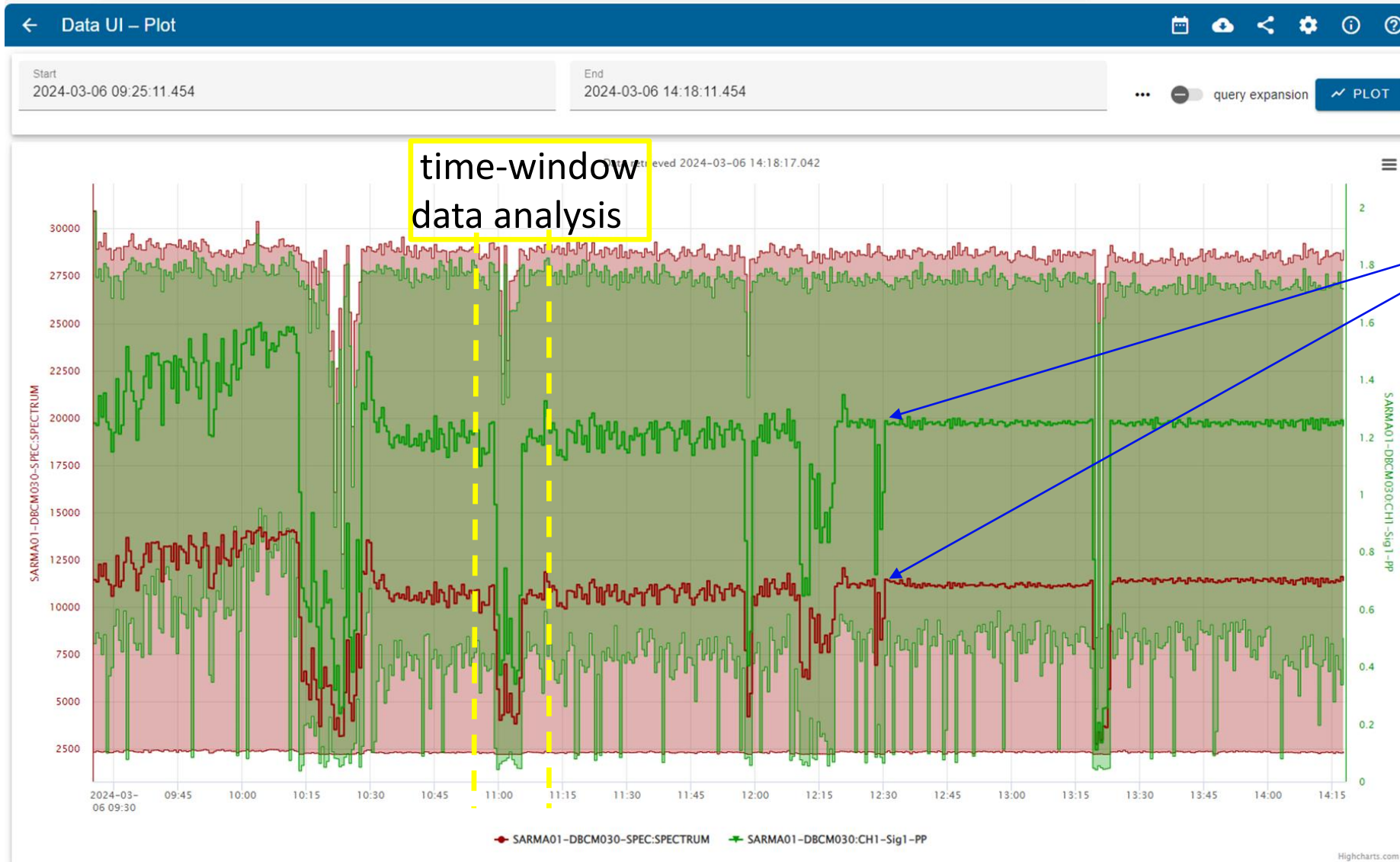
- Optical fiber spectrometer (Ocean Optics, 0.9-2.5 μm):
 - form-factor signature in the spectrogram
 - bunch length (σ) estimate by fitting the spectrogram with a Gaussian form-factor
 - main issue, the diffractive low frequency suppression of the spectrogram: non-univocal results of the Gaussian fit vs the frequency cut-off
 - Alternatively, “two-detector” BCM method



“Two-detector” BCM method, analysis of relative variations of:

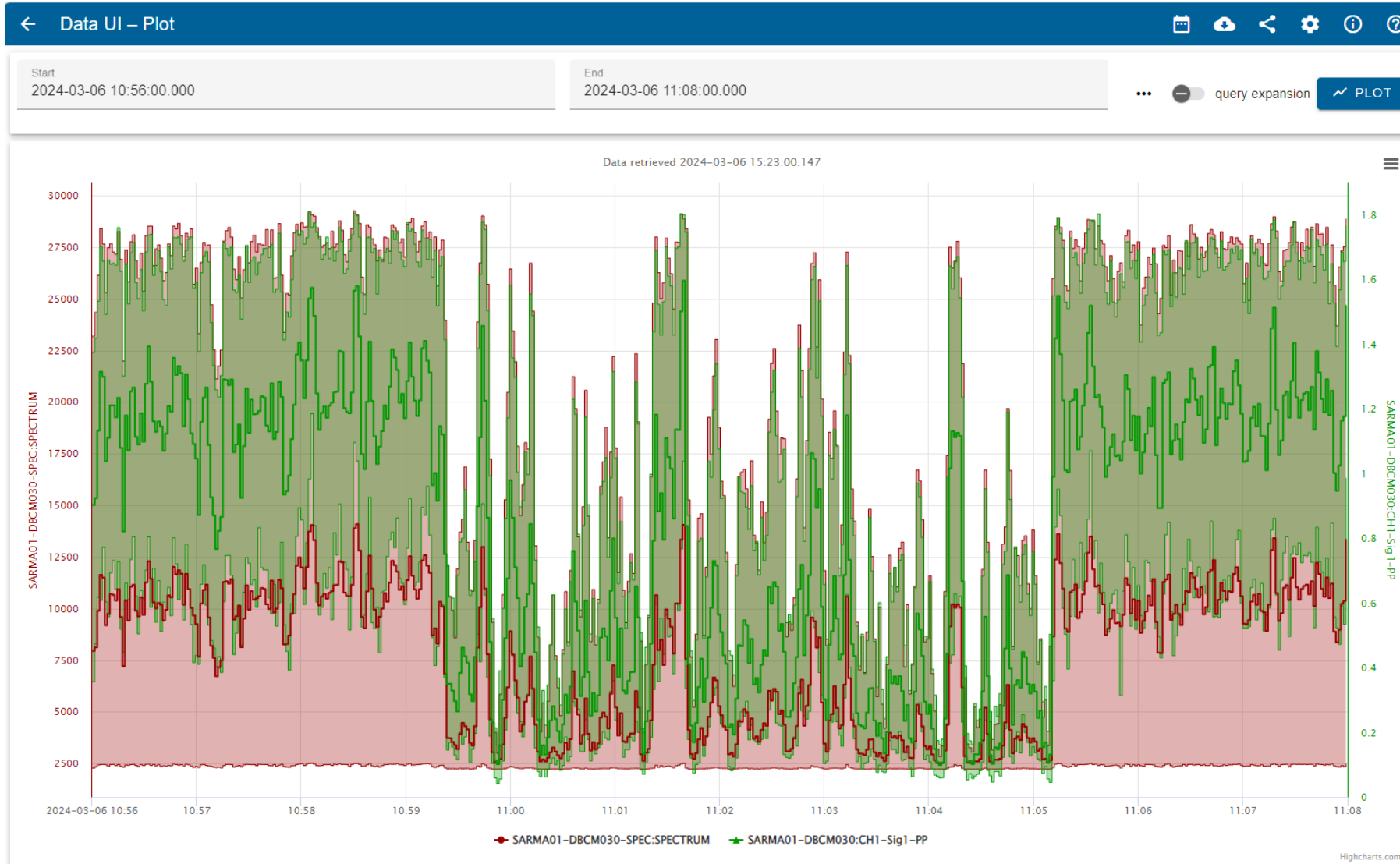
- pyro-detector signal $(\Delta I/I)_{j=1}$
- total spectrogram energy (spectrogram integral) $(\Delta I/I)_{j=2}$
- charge monitor signal $\Delta N/N$

10pC Aramis operations: from the set-up to the compression feedback ON

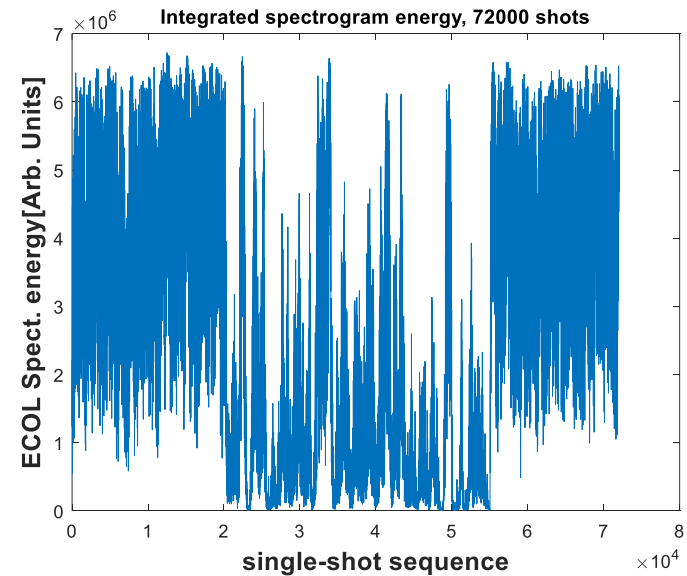
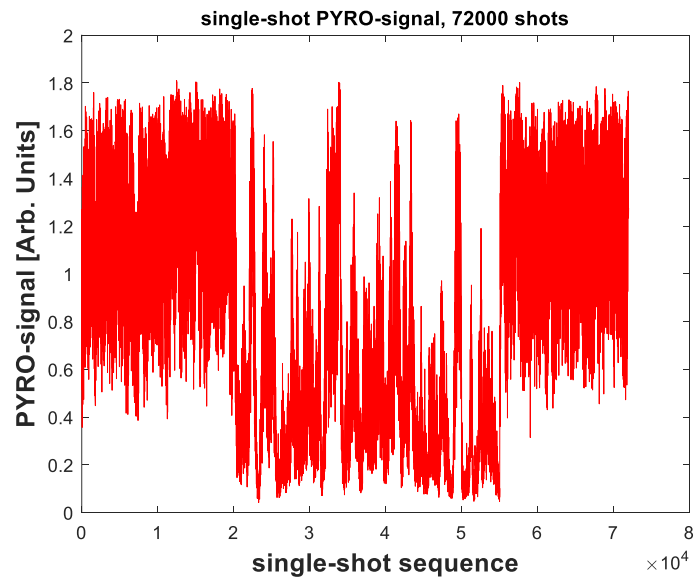
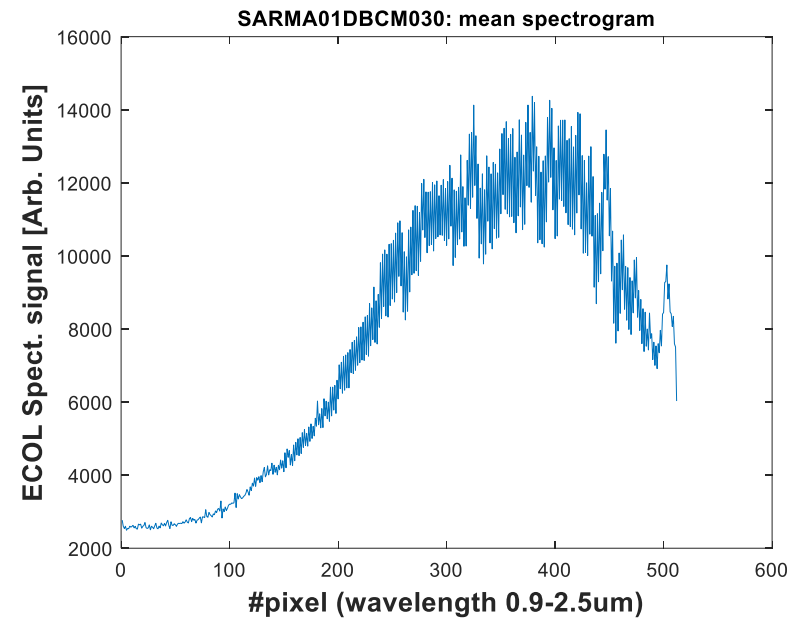
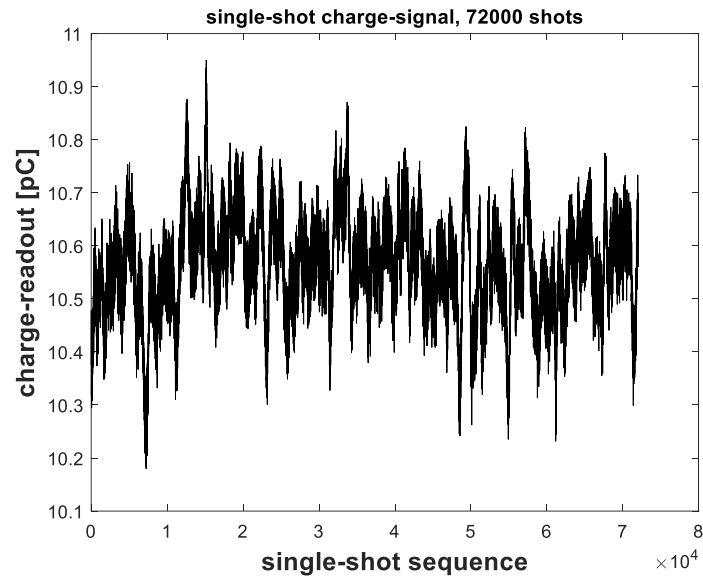


Optimal machine setting for Cristallina achieved and compression feedback locked onto BC1-BCM and ECOL-BCM signals

10pC Aramis operations: time-window of data analysis



Charge, pyrodetector and integrated spectrogram signals: 72000 shots

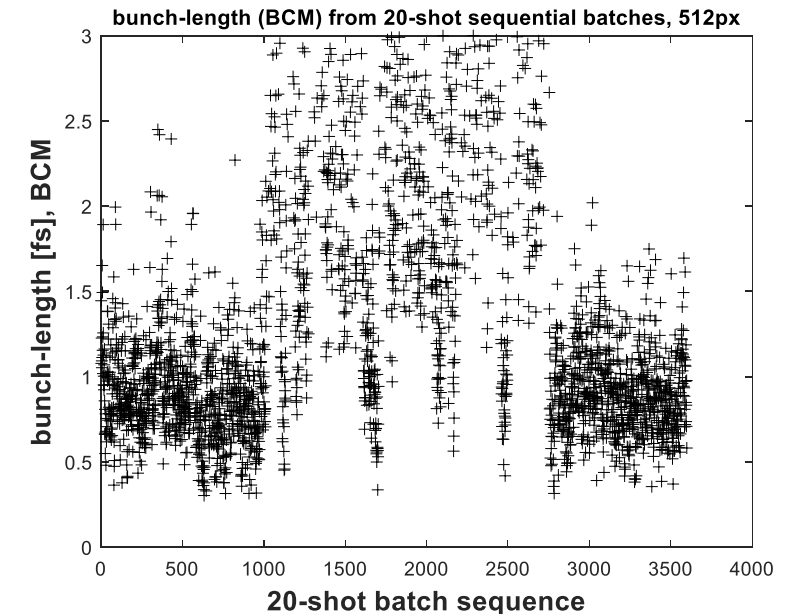
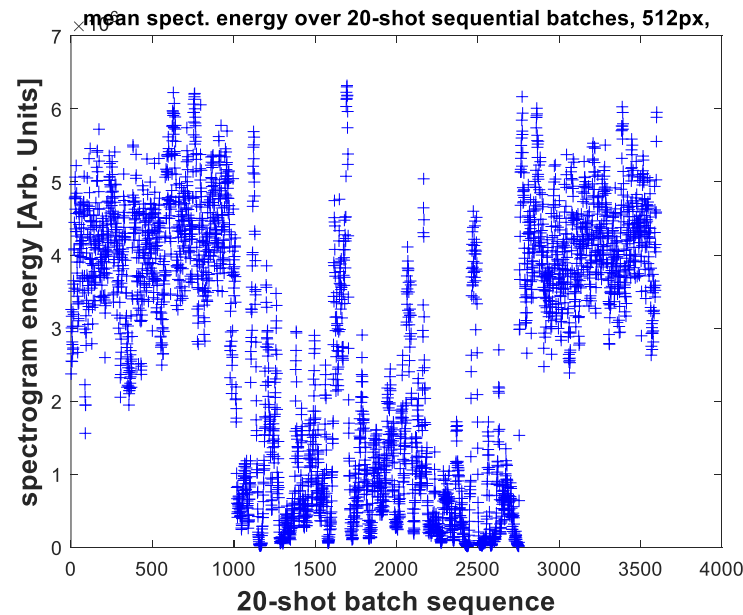
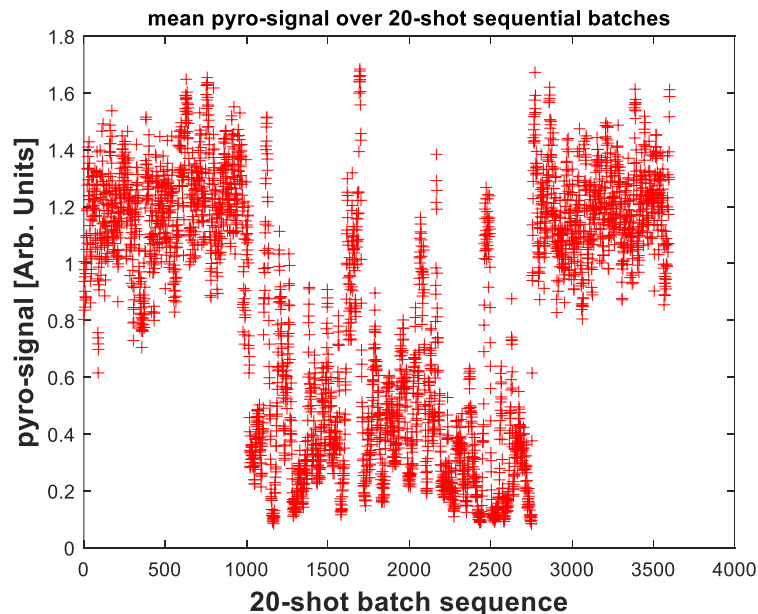


bunch-length characterization with “two-detector” BCM method

Data Analysis processing:

- divide the 72000 data array into 20-shot sequential batches;
- for each batch, calculate the mean values of the signals [charge, pyrodetector, spectrogram-energy-(512px)]
- for each batch, calculate $(\Delta I/I)_{j=1,2}$ and $\Delta N/N$ w.r.t. the resulting mean values
- for each batch, find the “zero” of the equation below as a function of a test σ ranging from 0.3 to 3.0 fs (small step)
- save and plot the test value of σ which equalizes to zero the right side of the equation

$$\backslash std \left(\left[\left(\frac{\Delta I}{I} \right)_2 - 2 \frac{\Delta N}{N} \right] \right) - \backslash abs \left(\frac{G(\sigma, (\Delta\omega)_2)}{[G(\sigma, (\Delta\omega)_2) - G(\sigma, (\Delta\omega)_1)]} \right) \times \backslash std \left(\left[\left(\frac{\Delta I}{I} \right)_2 - \left(\frac{\Delta I}{I} \right)_1 \right] \right) = 0$$

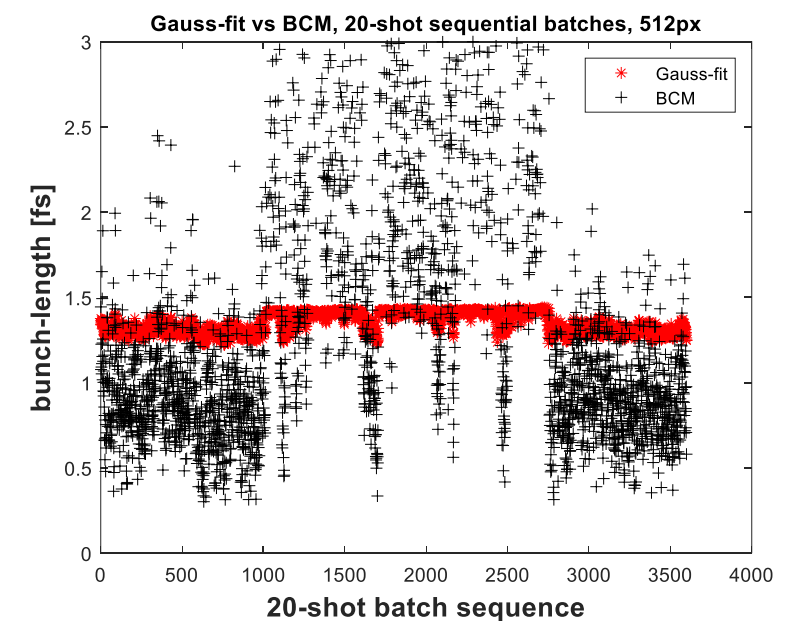
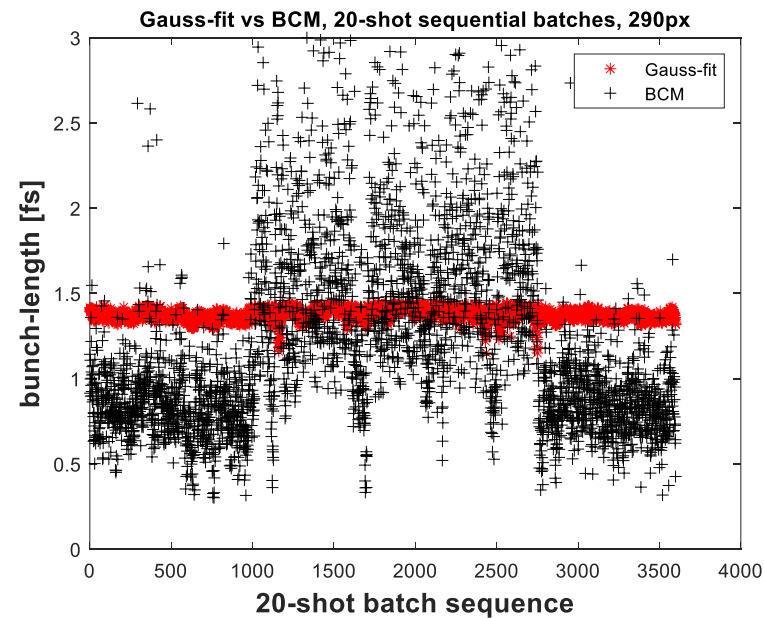
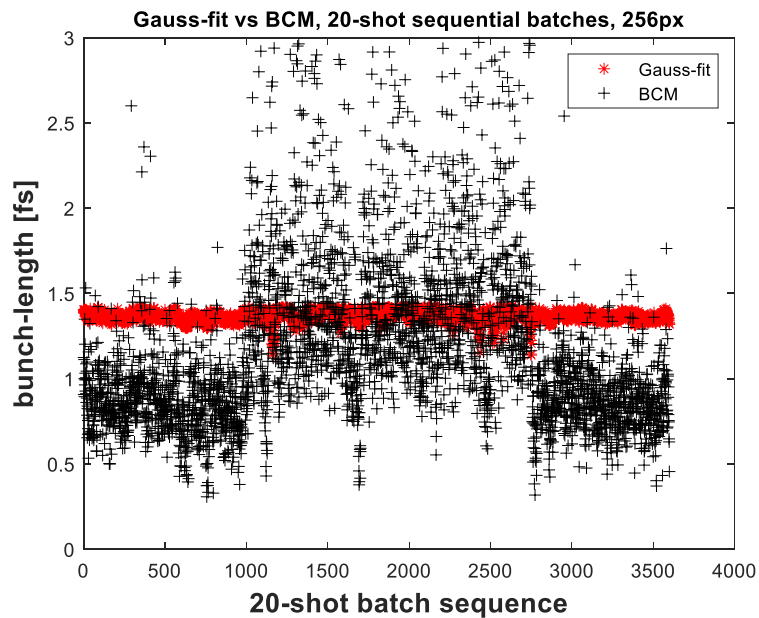


bunch-length characterization with “two-detector” BCM method



“two-detector” BCM method vs Gauss fit of the spectrogram:

- consider spectrogram in different wavelength band (low-frequency cut-off):
 - (1-256)px \leftrightarrow (0.9-1.7) μ m; (1-290)px \leftrightarrow (0.9-1.8) μ m; (1-512)px \leftrightarrow (0.9-2.5) μ m
- divide the 72000 data array into 20-shots sequential batches and calculate $(\Delta I/I)_{j=1,2}$ and $\Delta N/N$ as usual
- apply the «two-detector» BCM method to the processed data
- apply a Gaussian fit to the mean spectrogram of each 20-shot batch for the 3-wavelength band



Conclusions and Outlook



- the processing of the BCM – and charge signal- according to the “two-detector” method allows for an absolute characterization of the bunch-length
- In a FEL/SwissFEL, compression stabilized by feeding back the RF settings with the BCM signals (pyro-detector signal)
- “two-detector” BCM method allows for direct tuning of the compression feedback with the expected absolute value of the bunch-length instead of the bunch-length dependent signal of a BCM
- The method of processing the relative fluctuations of the BCM detector signals $(\Delta I/I)_{j=1,2}$ reduces the sensitivity of the analysis on systematic effect such as non-uniformity of the single particle spectrum or diffractive low-frequency cut-off of the radiation energy spectrum
- To be noticed: absolute value of a physical quantity (bunch-length) obtained from the analysis of the bunch-length induced relative statistical fluctuations of a detector signal