

Workshop on Quantum Chaos, Bernoulli Center

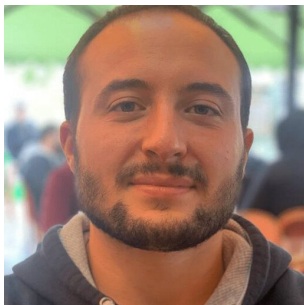
Classical simulation of quantum circuits via Pauli Propagation

October 2nd

Armando Angrisani,
Quantum Information and Computing group

The logo for EPFL (École Polytechnique Fédérale de Lausanne) is displayed in a bold, red, sans-serif font. The letters are thick and blocky, with a distinctive design where the 'E' and 'F' have a small gap at the top.

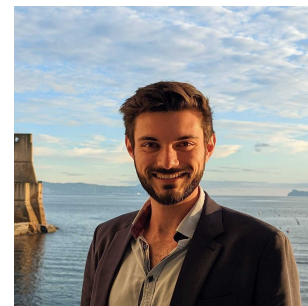
"Classically estimating observables of noiseless quantum circuits." arXiv:2409.01706



Armando Angrisani
EPFL



Alexander Schmidhuber
MIT



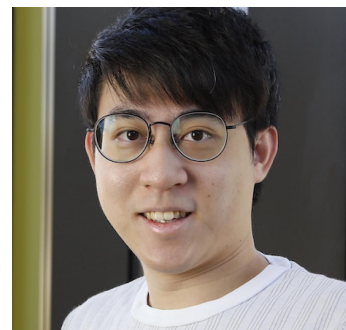
Manuel S. Rudolph
EPFL



Marco Cerezo
LANL



Zoë Holmes
EPFL



Hsin-Yuan Huang
Google, MIT, Caltech

Can chaos enable the classical simulation
of quantum systems?

Can chaos enable classical simulation
of quantum systems?

*it depends on what we mean by “simulation”
and “chaos”*

Classical simulation of quantum circuits

Given a unitary U and an initial state ρ (typically $|0^n\rangle\langle 0^n|$)

- Strong simulation: approximate $U\rho U^\dagger$
- Weak simulation:
 - Sample from the **Born distribution** of $U\rho U^\dagger$
 - Given an observable O , estimate $\text{Tr}[OU\rho U^\dagger]$

Classical simulation of quantum circuits

Given a unitary U and an initial state ρ (typically $|0^n\rangle\langle 0^n|$)

- Strong simulation: approximate $U\rho U^\dagger$ *out of reach for PP*
- Weak simulation:
 - Sample from the **Born distribution** of $U\rho U^\dagger$ *feasible with PP for **locally scrambling noisy** circuits*
 - Given an observable O , estimate $\text{Tr}[OU\rho U^\dagger]$ *feasible with PP for **locally scrambling** circuits*

PP = Pauli Propagation, a family of simulation algorithms

An extreme case: Haar-random state

- Sampling from the Born distribution of a Haar-random state is “classically hard”
Still true for the output of random circuits

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⇒ Experimental demonstration of quantum supremacy

Article | Published: 23 October 2019

Quantum supremacy using a programmable superconducting processor

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- Any expectation value will be close to zero!

An extreme case: “highly random” states

- Sampling from the Born distribution of a Haar-random state is hard
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⇒ Experimental demonstration of quantum supremacy

Article | Published: 23 October 2019

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[Frank Arute](#), [Kunal Arya](#), [Ryan Babbush](#), [David B. Boness](#), [Cristian Burchard](#), [Sergio Boixo](#), [Fernando G. S. L. Brandao](#), [David Buco](#), [Roberto Collins](#), [William Courtney](#), [Andrew D. C. Davis](#), [Edward T. Dickerson](#), [John D. B. Ford](#), [Austin G. Fowler](#), [Lloyd C. Grover](#), [Michael J. Heuley](#), [Srinivas Aravamudan](#), [John M. Martinis](#), [Thomas J. Y. Yang](#), [John A. Martens](#), [Mark A. Naeff](#), [M. P. A. Fisher](#), [Antonio C. G. Azeiteiro](#), [Roberto C. O. Azeiteiro](#), [Marissa Giustina](#), [Rob Graff](#), [Keith Guerin](#), [John M. Martinis](#)

Ordinary computers can beat Google's quantum computer after all

Superfast algorithm put crimp in 2019 claim that Google's machine had achieved “quantum supremacy”

2 AUG 2022 · 5:05 PM · BY ADRIAN CHO

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- Any expectation value will be close to zero!

Chaos and scrambling

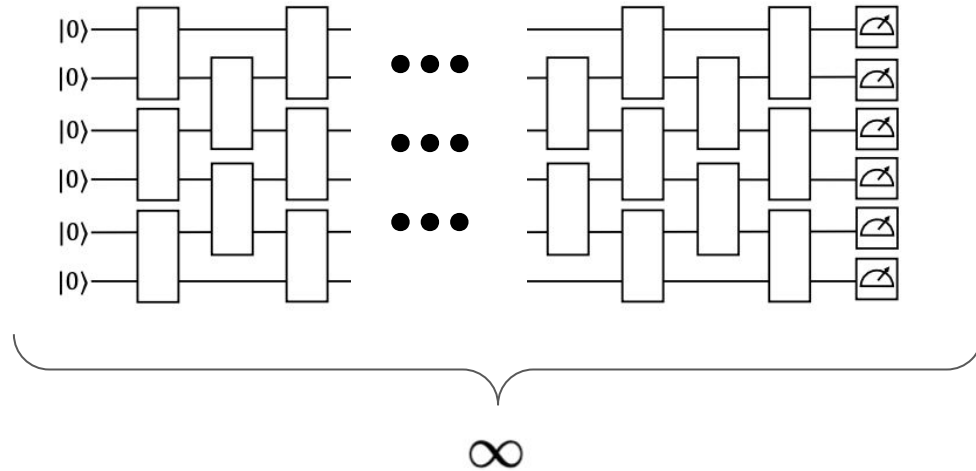
- Quantum chaos “scrambles” local information quickly across an entire system (e.g. black holes)
- The time evolution of the system is a **pseudo-Haar-random** process

$$W(t) = e^{iHt} W e^{-iHt}$$

(H chaotic Hamiltonian, W local operator)

Chaos and Random Quantum Circuits

Infinite-depth random circuits are Haar-random



Probing chaos

H Hamiltonian

V, W local operators

$$W(t) = e^{iHt} W e^{-iHt}$$

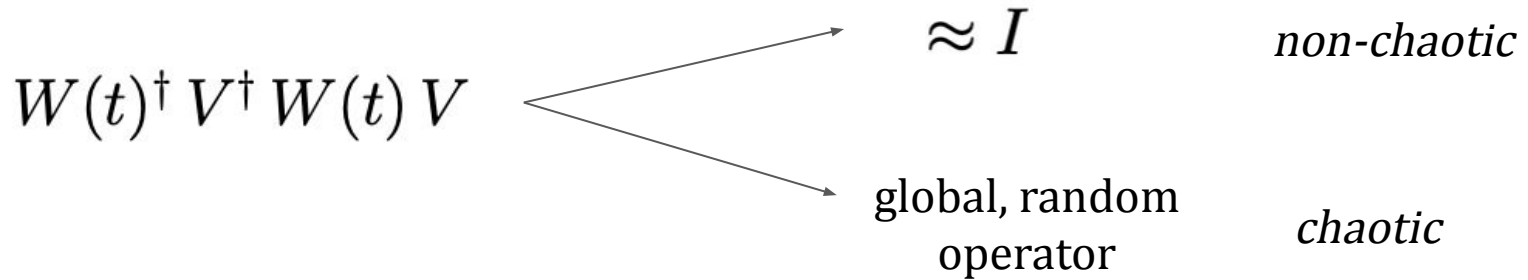
$$W(t)^\dagger V^\dagger W(t) V$$

(group commutator)

Probing chaos

H chaotic Hamiltonian
V, W local operators

$$W(t) = e^{iHt} W e^{-iHt}$$



Probing chaos

H chaotic Hamiltonian

V, W local operators

$$W(t) = e^{iHt} W e^{-iHt}$$

$$W(t)^\dagger V^\dagger W(t) V \longrightarrow \approx I \quad \text{non-chaotic}$$

Probing chaos with **Out-of-Time-Order** 4-point functions

H chaotic Hamiltonian

V, W local operators

$$W(t) = e^{iHt} W e^{-iHt}$$

$$\text{Tr}[W(t)^\dagger V^\dagger W(t) V \rho_H] \begin{cases} \approx 1 & \text{non-chaotic} \\ \approx 0 & \text{chaotic} \end{cases}$$

$$\rho_H = \frac{e^{-\beta H}}{\text{Tr}[e^{-\beta H}]}$$

Probing chaos with **Out-of-Time-Order** 4-point functions

H chaotic Hamiltonian

V, W local operators

$$W(t) = e^{iHt} W e^{-iHt}$$

$$\text{Tr}[W(t)^\dagger V^\dagger W(t) V \rho_H] \begin{cases} \approx 1 & \text{non-chaotic} \\ \approx 0 & \text{chaotic} \end{cases}$$

what's in the middle?

$$\rho_H = \frac{e^{-\beta H}}{\text{Tr}[e^{-\beta H}]}$$

OTO functions and 2-designs

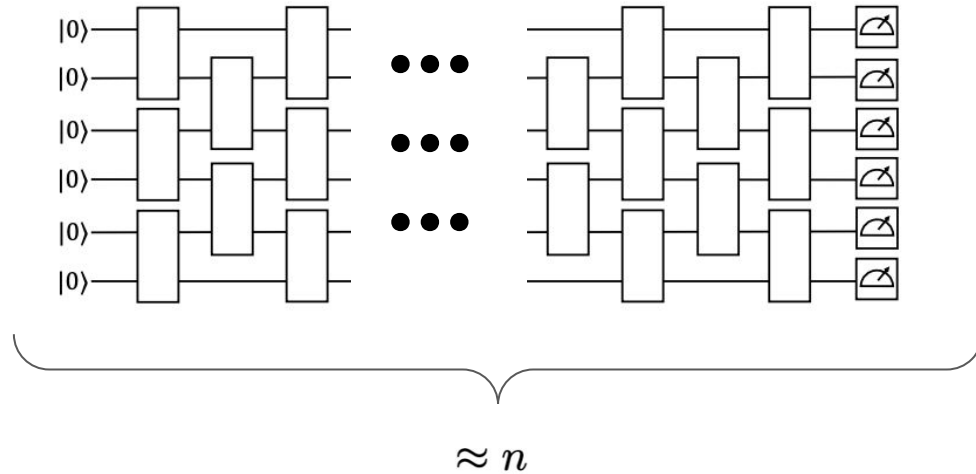
$$\text{Tr}[W(t)^\dagger V^\dagger W(t) V \rho_H]$$

The OTO function contains only second moments of the chaotic Hamiltonian
⇒ We can approximate the chaotic evolution with a unitary 2-design

$$\mathbb{E}_{U \sim \nu} U^{\dagger \otimes 2} O U^{\otimes 2} \approx \mathbb{E}_{U \sim \nu_{\text{Haar}}} U^{\dagger \otimes 2} O U^{\otimes 2}$$

Chaos and Random Quantum Circuits, II

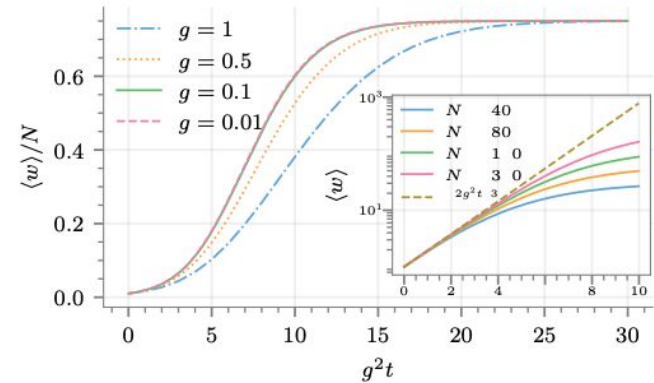
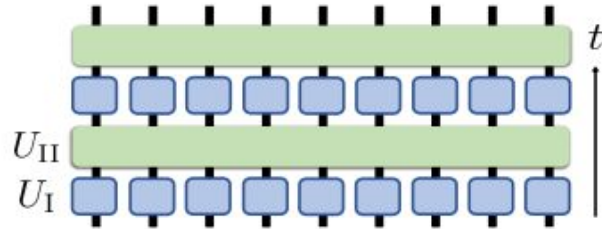
Linear-depth random circuits are approximate 2-designs



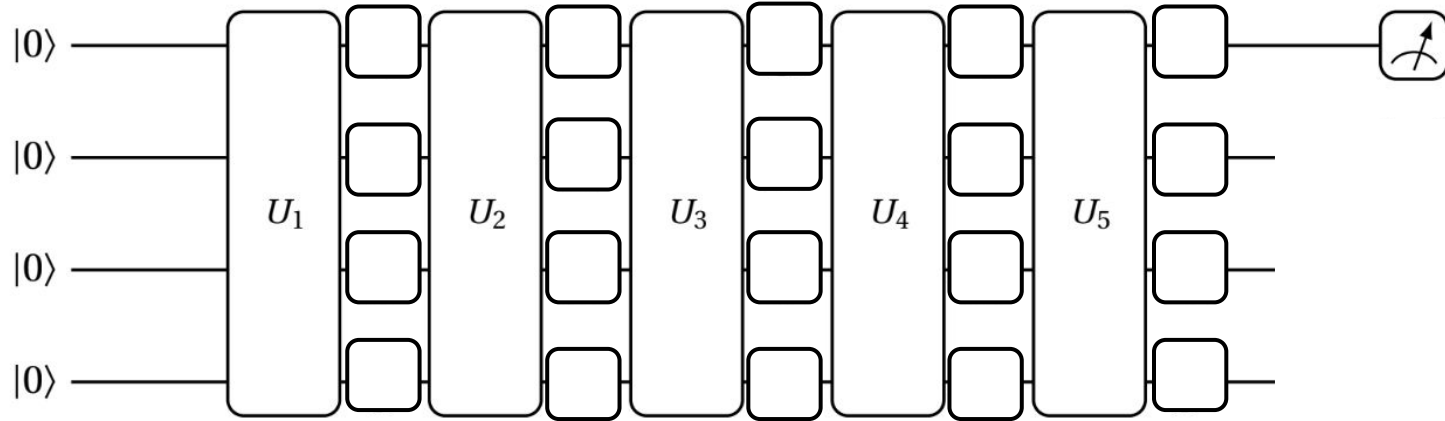
Chaos and Random Quantum Circuits, III

Random single qubit gates suffices for highly chaotic behaviour

$$U_I = \prod_{i=1}^N U_{H,i}, \quad U_{II} = e^{-i \frac{g}{2\sqrt{N}} \sum_{i<j} Z_i Z_j}$$



Locally scrambling circuit



- All single-qubit gates sampled i.i.d. from local 2-designs
- Arbitrary entangling layers U_i 's

The problem

Given a locally scrambling quantum U ,
can we approximate $\text{Tr}[OU\rho U^\dagger]$ with a classical computer?

The problem

Given a locally scrambling circuit U ,
can we approximate $\text{Tr}[OU\rho U^\dagger]$ with a classical computer?

Yes!

It takes $\text{poly}(n)$ time, for any constant precision

Motivations

- Many physical processes are approximately locally scrambling
- Variational quantum circuits are often initialized at random
- Understand why the Pauli Propagation algorithm works so well in practice
 - This is a “simple” model in which we can prove things

The Heisenberg picture

$$\text{Tr}[OU\rho U^\dagger] = \text{Tr}[U^\dagger OU\rho]$$

→ Compute (approximately) $U^\dagger OU$ and project it onto ρ

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→ Compute (approximately) $U^\dagger OU$ and project it onto ρ

$$O = \frac{1}{2^n} \sum_{P \in \{I, X, Y, Z\}^{\otimes n}} \text{Tr}[OP] P$$

$$U = U_L U_{L-1} \dots U_1$$

$$\rho = |0^n\rangle\langle 0^n|$$

$$\left\{ \begin{array}{l} U_j^\dagger P U_j = \sum_{P \in \{I, X, Y, Z\}^{\otimes n}} a_P P \\ \text{with } \sum_{P \in \{I, X, Y, Z\}^{\otimes n}} a_P^2 = 1 \end{array} \right.$$

The Heisenberg picture

$$\text{Tr}[OU\rho U^\dagger] = \text{Tr}[U^\dagger OU\rho]$$

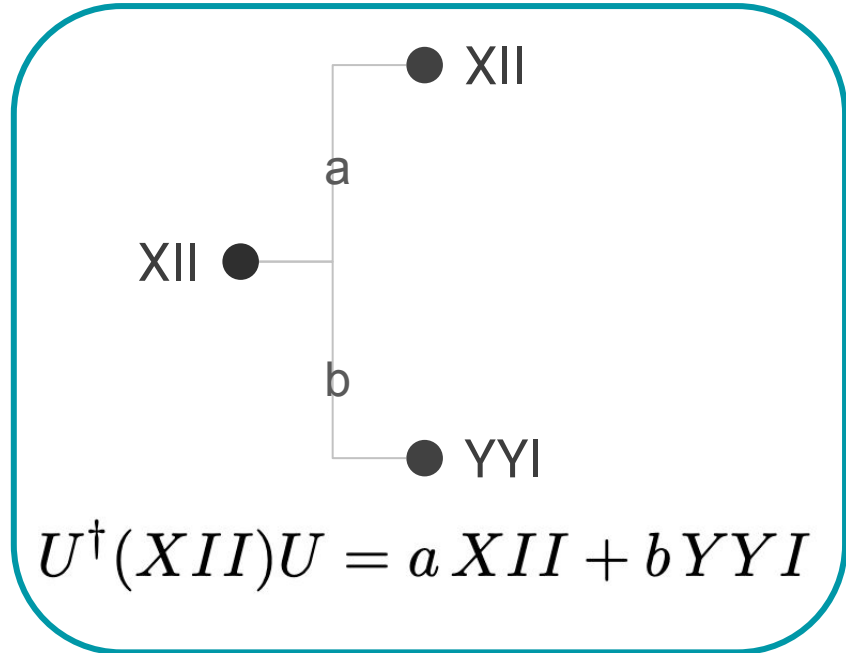
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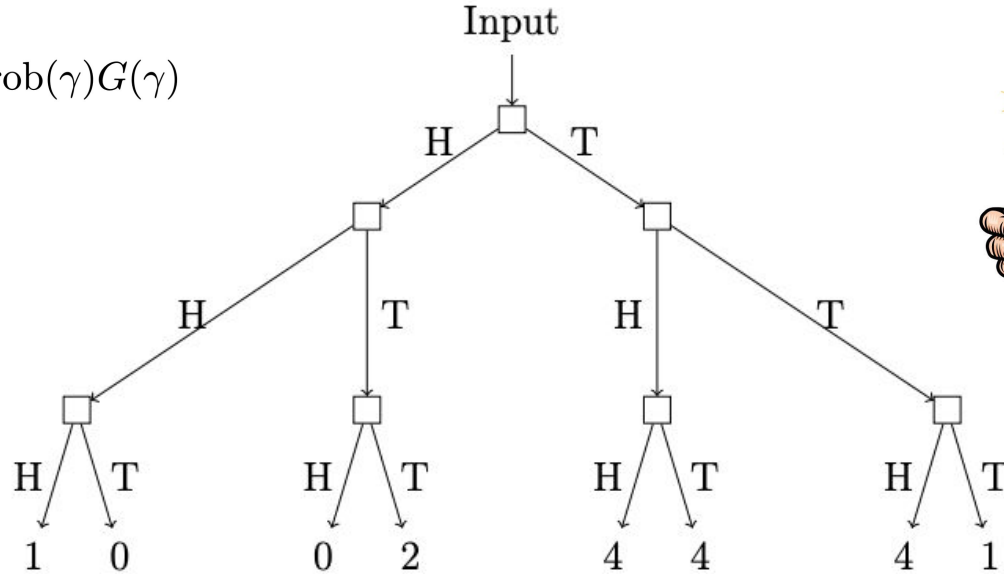
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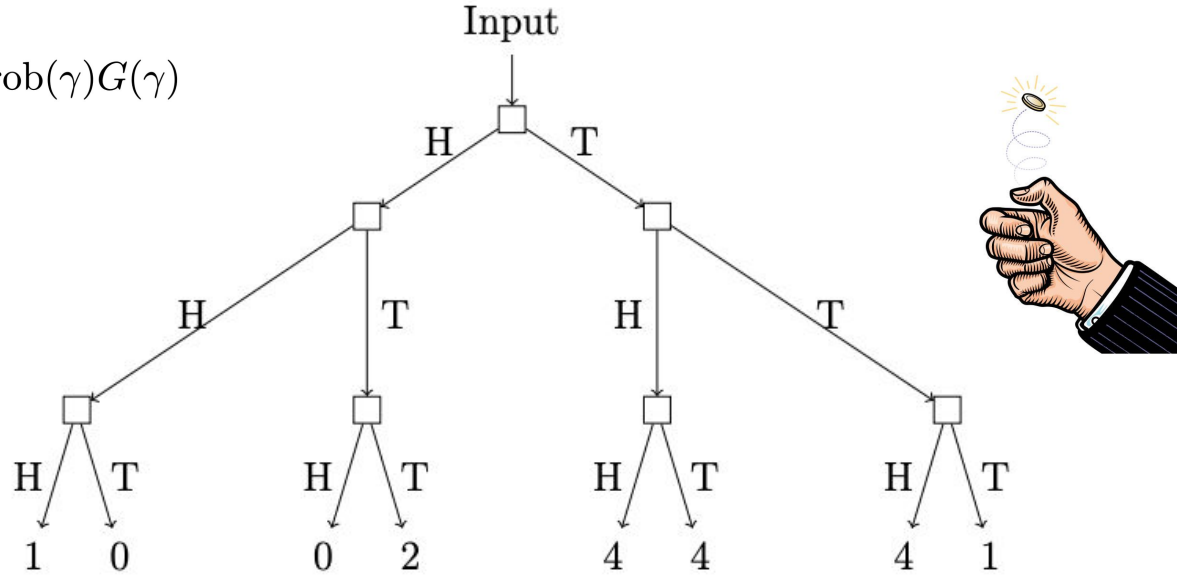
A step back: classical Randomized Computation

$$F(\text{Input}) = \sum_{\gamma \in \text{paths}} \text{Prob}(\gamma) G(\gamma)$$



A step back: classical Randomized Computation

$$F(\text{Input}) = \sum_{\gamma \in \text{paths}} \text{Prob}(\gamma) G(\gamma)$$



Expected values efficiently computable by Monte-Carlo sampling paths!

- Computing exactly or Monte Carlo sampling Pauli paths is classically intractable
- Can we find a sub-tree which carries most of the useful information?

Are all Pauli paths created equal?

Pauli-weight: number of non-identity terms in a Pauli operator

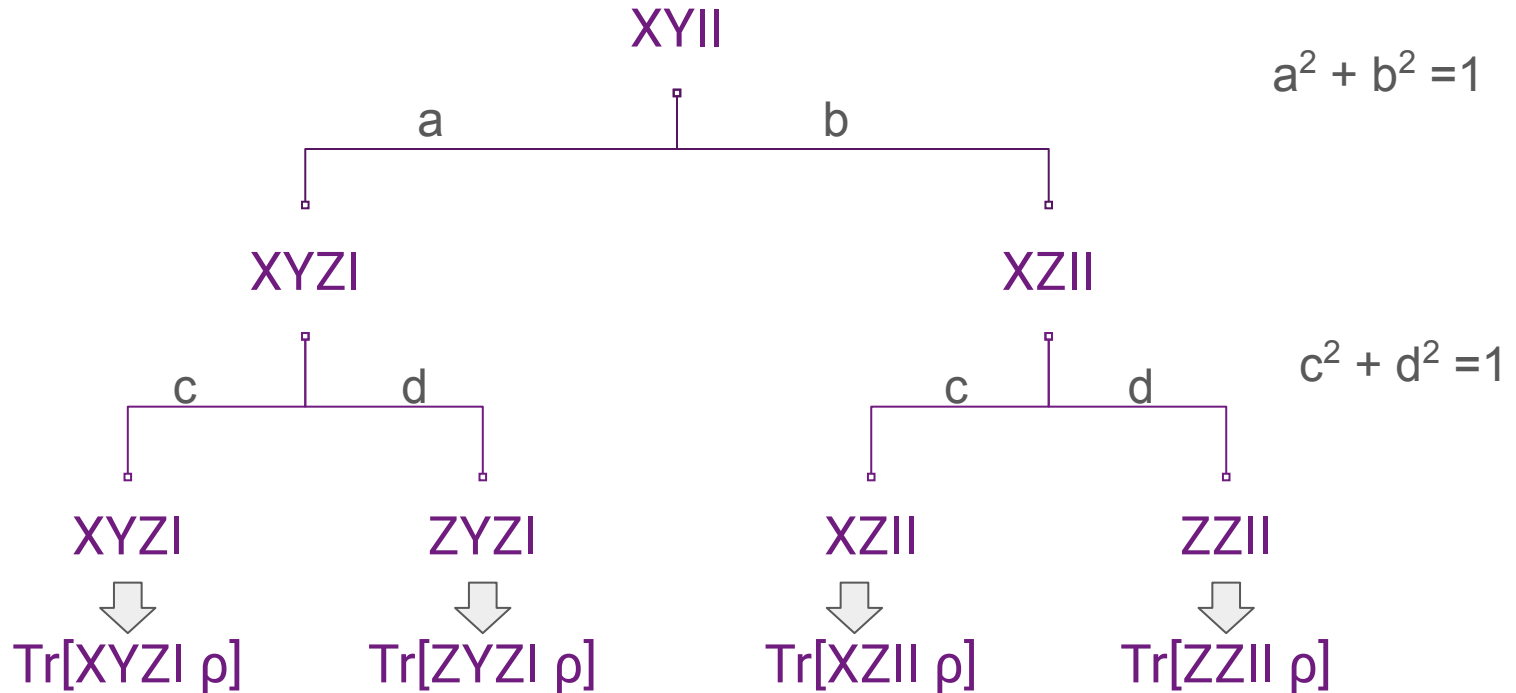
ex. $\text{weight}(\text{IXYIZ}) = 3$

$\text{weight}(\text{IIIII}) = 0$

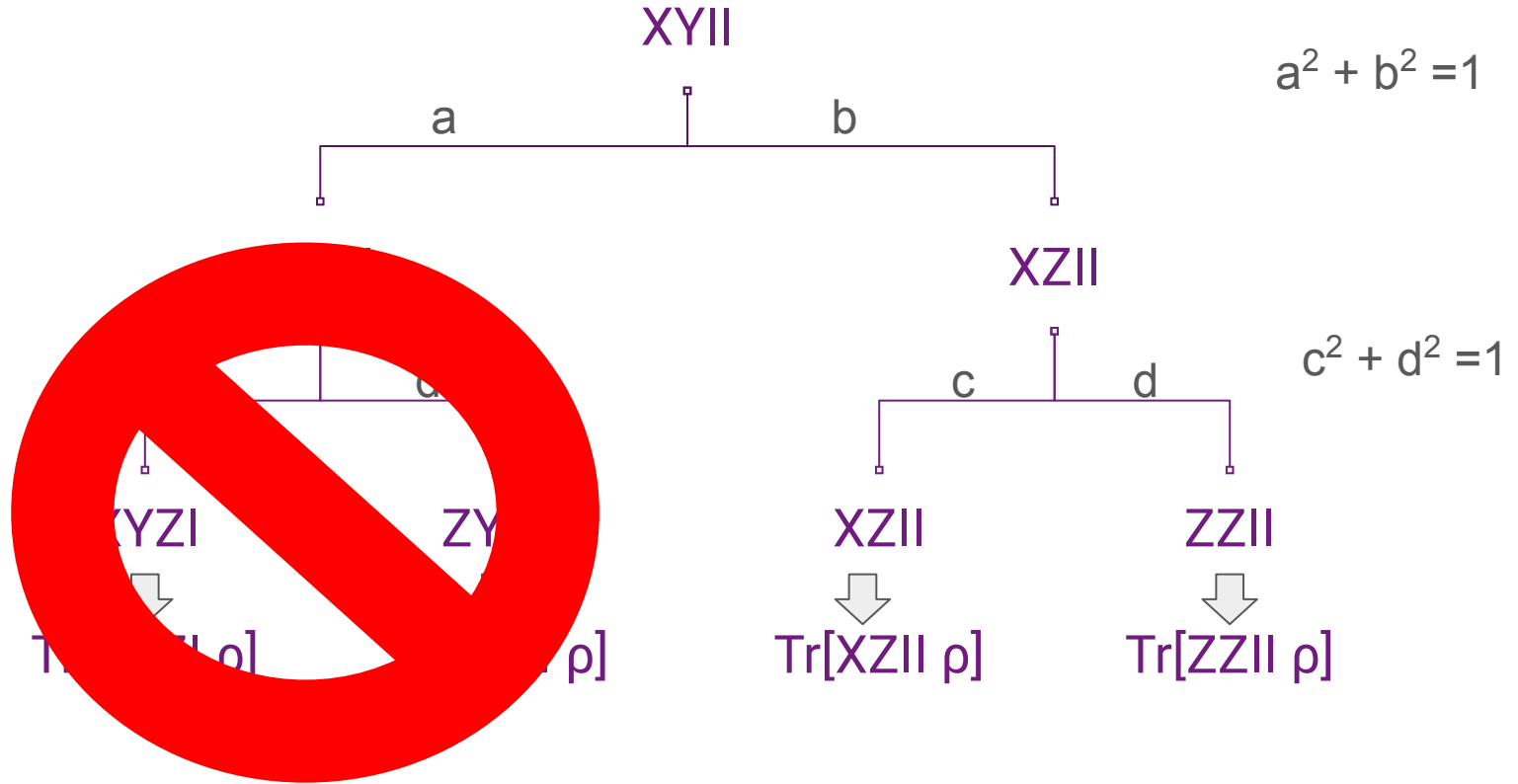
Paths with high-weight Paulis can be truncated if

- the circuit is noisy
- the circuit is locally scrambling,
and we want to estimate expectation values

Tree-like representation of Quantum Computation



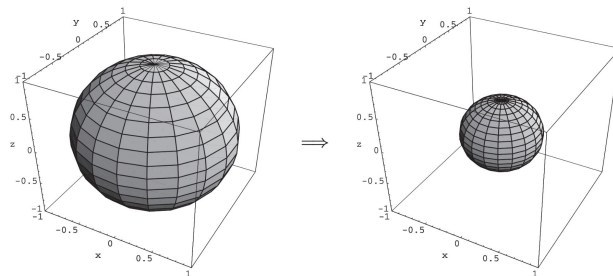
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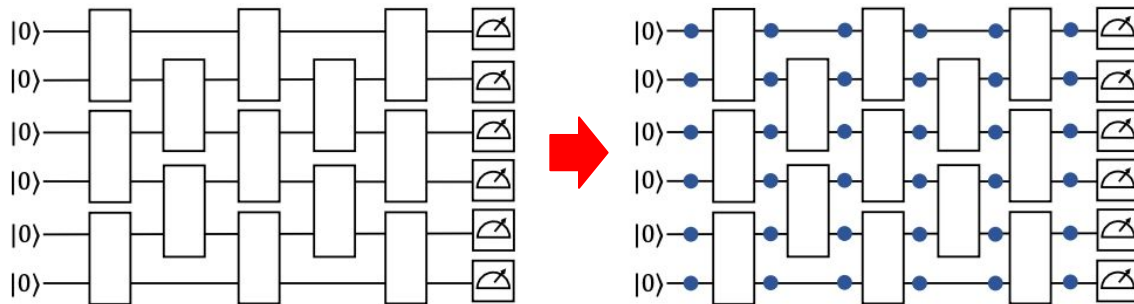
The noisy case

- Single-qubit depolarizing noise

$$\mathcal{D}_p(\rho) = p\frac{I}{2} + (1-p)\rho$$



- Noisy circuits

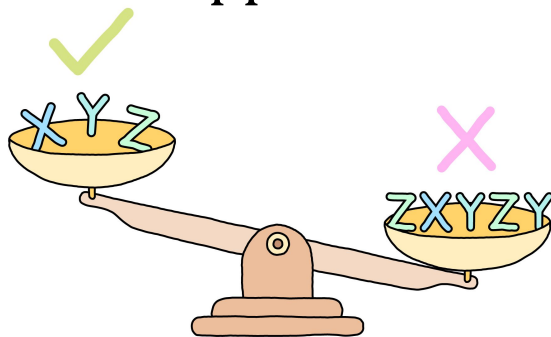


The noisy case

$$\mathcal{D}_p^{\otimes n}(P) = (1 - p)^{\text{weight}(P)} P$$

$$\implies |\text{Tr}[\mathcal{D}_p^{\otimes n}(P)\rho]| \leq (1 - p)^{\text{weight}(P)}$$

Paths containing **high-weight** Paulis are **exponentially** suppressed!



The noisy case

A polynomial-time classical algorithm for noisy random circuit sampling

Dorit Aharonov* Xun Gao† Zeph Landau‡ Yunchao Liu§ Umesh Vazirani¶

‘22

Classical simulations of noisy variational quantum circuits

Enrico Fontana^{1,2,3,*} Manuel S. Rudolph⁴, Ross Duncan², Ivan Rungger³, and Cristina Cirstoiu²

‘23

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locally scrambling circuits

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parametrized circuits

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Thomas Schuster^{1,*} Chao Yin^{2,*} Xun Gao^{2,3} and Norman Y. Yao⁴

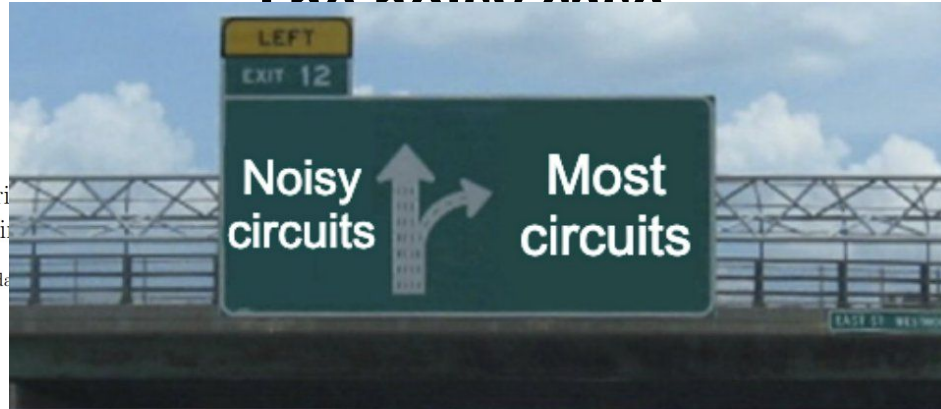
‘24

arbitrary circuits with input sampled from a 1-design

The noisy case

A polynomial-time classical algorithm for sampling from the output distribution of noisy quantum circuits

Dorit Aharonov* Xun Gao† Zeph Landau



Sampling circuits

Classical simulations of noisy quantum circuits

Enrico Fontana^{1,2,3,*} Manuel S. Rudolph⁴, Ross D.



circuits

A polynomial-time classical algorithm for sampling from the output distribution of noisy quantum circuits with input a 1-design

Thomas Schuster,^{1,*} Chao Yin,^{2,*} Xun

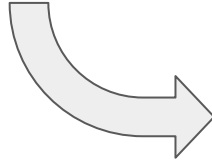
circuits with input a 1-design

However, the same simulation algorithm works well
on “typical” noiseless circuits

Classical surrogate simulation of quantum systems with LOWESA

Manuel S. Rudolph,^{1,2,*} Enrico Fontana,^{3,4,5,6} Zoë Holmes,¹ and Lukasz Cincio²

simulates



Evidence for the utility of quantum computing before fault tolerance

[Youngseok Kim](#) , [Andrew Eddins](#) , [Sajant Anand](#), [Ken Xuan Wei](#), [Ewout van den Berg](#), [Sami Rosenblatt](#), [Hasan Nayfeh](#), [Yantao Wu](#), [Michael Zaletel](#), [Kristan Temme](#) & [Abhinav Kandala](#) 

Nature **618**, 500–505 (2023)

Quantum Convolutional Neural Networks are (Effectively) Classically Simulable

Pablo Bermejo,^{1,2,3} Paolo Braccia,⁴ Manuel S. Rudolph,⁵ Zoë Holmes,⁵ Lukasz Cincio,⁴ and M. Cerezo^{1,*}

However, the



works well

Classical surrogate

Manuel S. Rudolph

with LOWESA

Lukasz Cincio²

simula

quantum computing before

and, Ken Xuan Wei, Ewout van den Berg, Sami
Alel, Kristan Temme & Abhinav Kandala

Quantum Convolution

Pablo Bermejo,^{1,2,3} Paolo Braccia,

Classically Simulable

Manuel S. Rudolph,¹ Zoc Holmes,¹ Lukasz Cincio,⁴ and M. Cerezo^{1,*}

2-designs in the Heisenberg picture

Given a unitary U sampled from a 2-design, we have for any non-identity P

$$\mathbb{E}_U U^{\dagger \otimes 2} P^{\otimes 2} U^{\otimes 2} = \frac{1}{4^n - 1} \sum_{P \neq I} P^{\otimes 2}$$

The single-qubit case

$$\mathbb{E}_U U^{\otimes 2 \dagger} X^{\otimes 2} U^{\otimes 2} = \frac{1}{3} (X^{\otimes 2} + Y^{\otimes 2} + Z^{\otimes 2})$$

Scrambling vs Expectation values

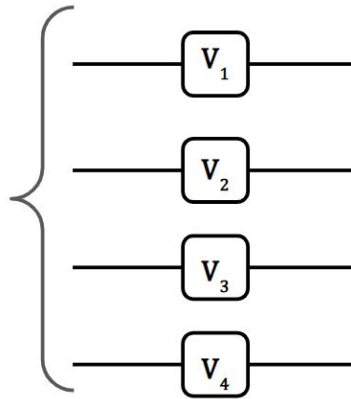
- Without randomness

$$\text{Tr}[X\rho] \in [-1, 1]$$

- With randomness

$$\begin{aligned} \mathbb{E}_U \text{Tr}[XU\rho U^\dagger]^2 &= \frac{1}{3} \left(\text{Tr}[X\rho]^2 + \text{Tr}[Y\rho]^2 + \text{Tr}[Z\rho]^2 \right) \\ &= \frac{2}{3} \cdot \underbrace{\frac{1}{2} \left(\text{Tr}[X\rho]^2 + \text{Tr}[Y\rho]^2 + \text{Tr}[Z\rho]^2 \right)}_{\langle \text{Tr}[\rho^2] \rangle \leq 1} \leq \frac{2}{3} \end{aligned}$$

Locally scrambling layer

$$U = V_1 \otimes V_2 \otimes V_3 \otimes V_4$$


The diagram shows four horizontal lines representing qubits. Each line passes through a rounded rectangular box labeled V_1 , V_2 , V_3 , and V_4 respectively. A large curly brace on the left groups these four qubits together.

$$\mathbb{E}_U \text{Tr} [PU\rho U^\dagger]^2 \leq \binom{2}{3}^{\text{weight}(P)}$$



$$\mathbb{E}_U \text{Tr}[PU\rho U^\dagger]^2 \leq \left(\frac{2}{3}\right)^{\text{weight}(P)}$$

$$|\text{Tr}[\mathcal{D}_p^{\otimes n}(P)\rho]| \leq (1-p)^{\text{weight}(P)}$$

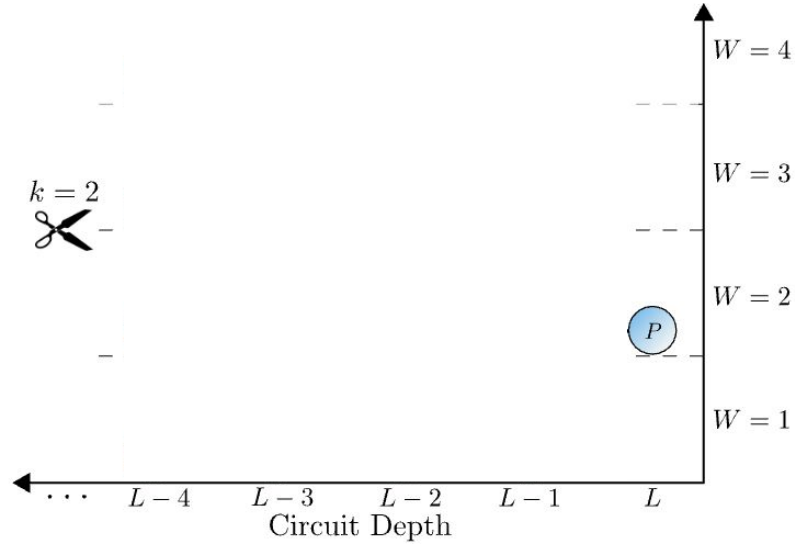
A “weight-truncated” observable

$$O = \sum_P a_P P$$

$$O^{(k)} = \sum_{\substack{P: \\ \text{weight}(P) < k}} a_P P$$

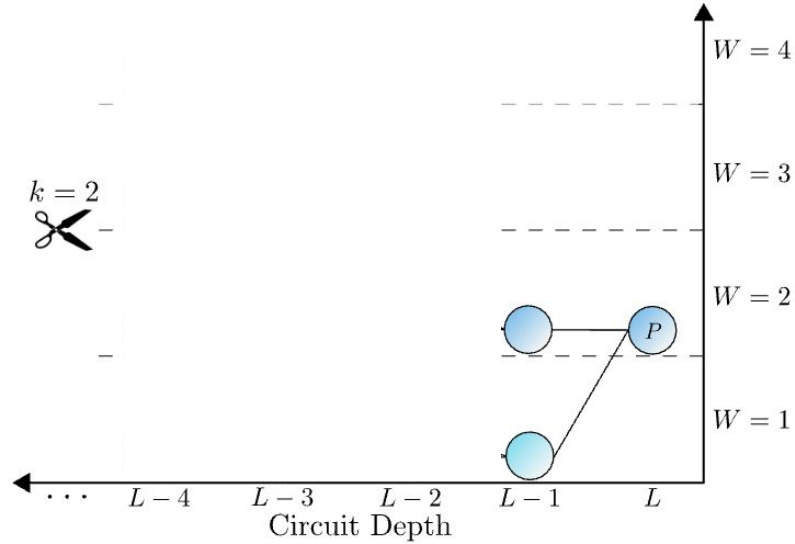
$$\mathbb{E}_U \text{Tr} \left[\left(O - O^{(k)} \right) U \rho U^\dagger \right]^2 \leq \left(\frac{2}{3} \right)^k \|O\|$$

“Weight-truncated” Pauli Propagation



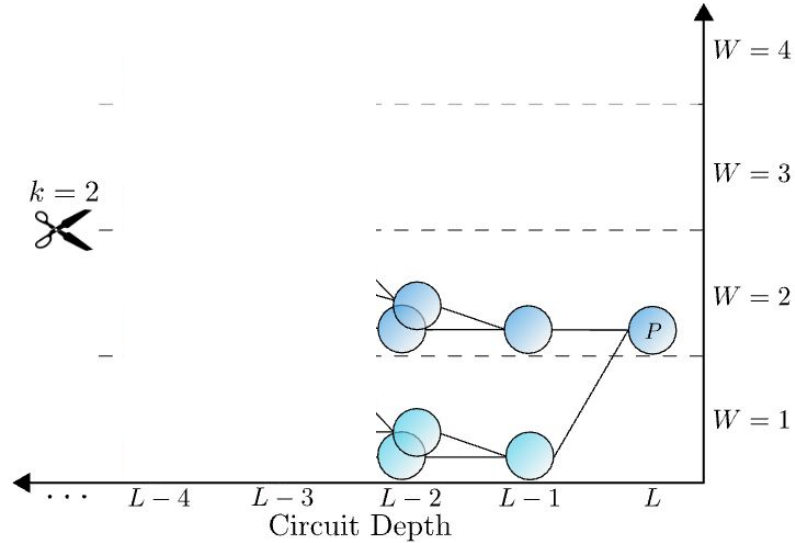
All Paulis with weight $> k$ are truncated

“Weight-truncated” Pauli Propagation



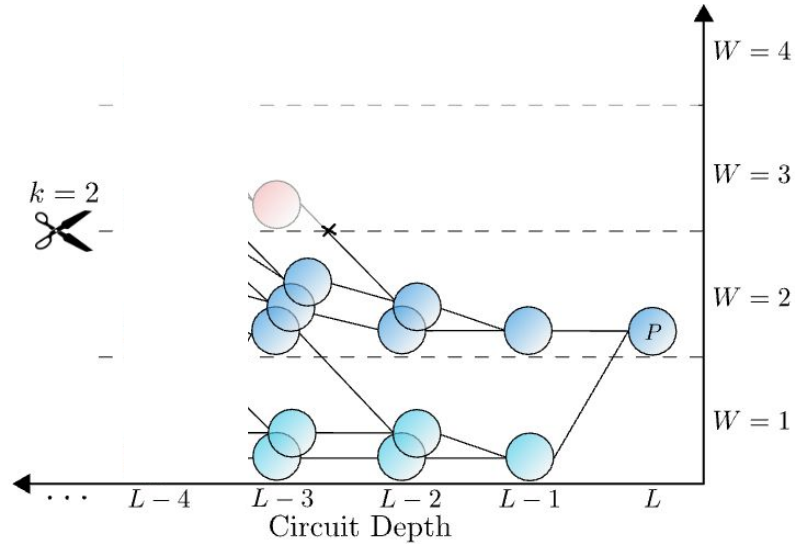
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“Weight-truncated” Pauli Propagation



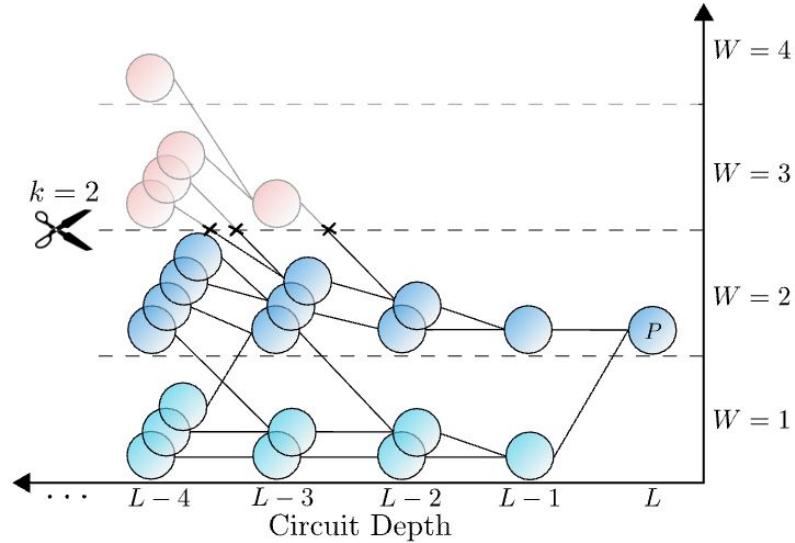
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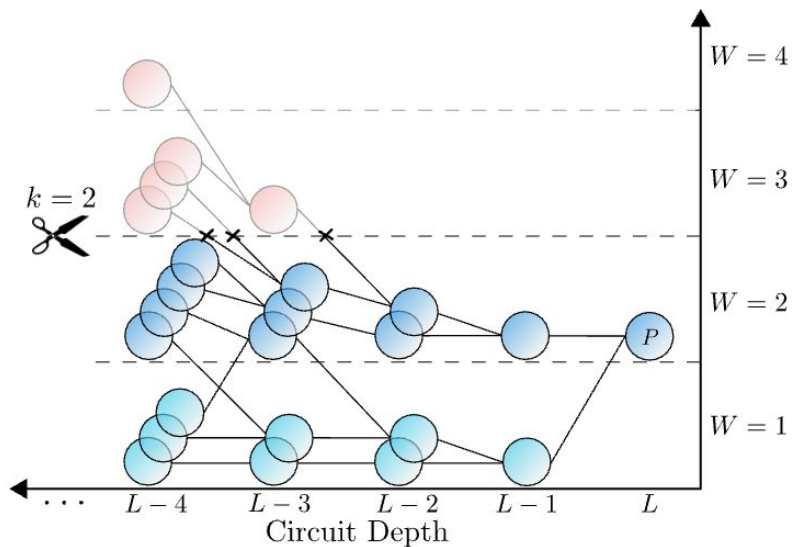
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“Weight-truncated” Pauli Propagation



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“Weight-truncated” Pauli Propagation

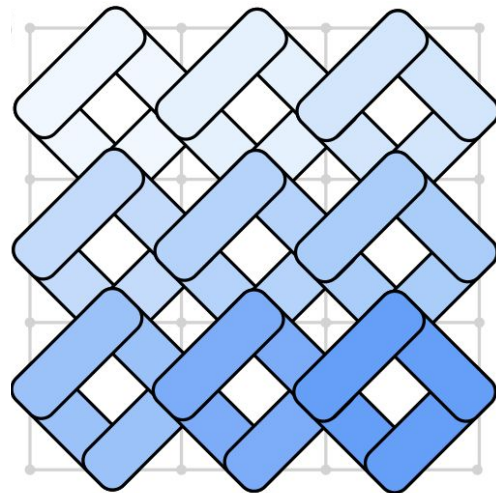
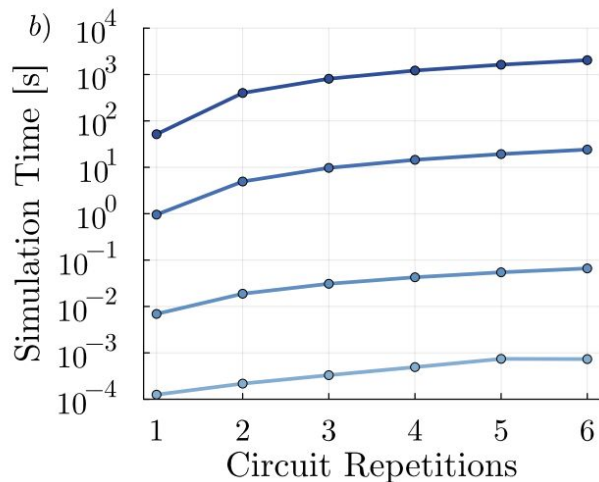
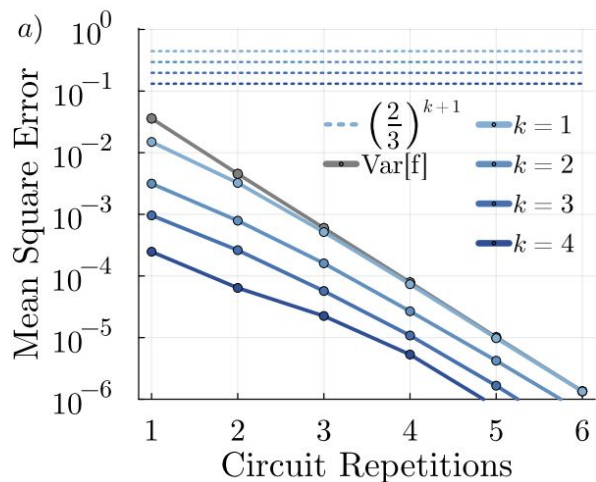


$$\mathbb{E}_U \text{Tr} \left[\left(U^\dagger O U - (U^\dagger O U)^{(\text{apx})} \right) \rho \right]^2 \leq \left(\frac{2}{3} \right)^k \|O\|$$

“Weight-truncated” Pauli Propagation

- Error bound depends only on k
- Polynomial complexity for constant error
- Easy to “certify” numerically
 - Experiments demonstrate faster convergence

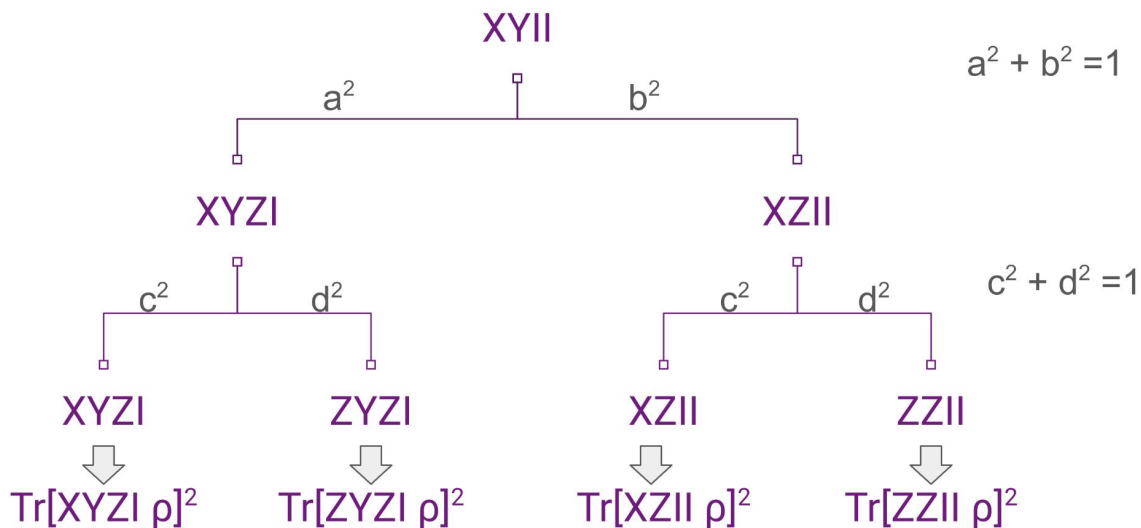
2D Staircase Topology on 64 qubits



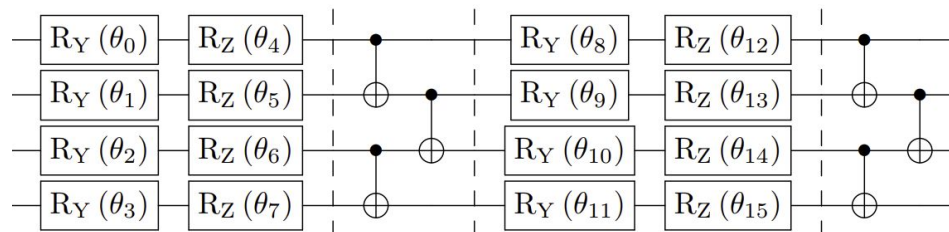
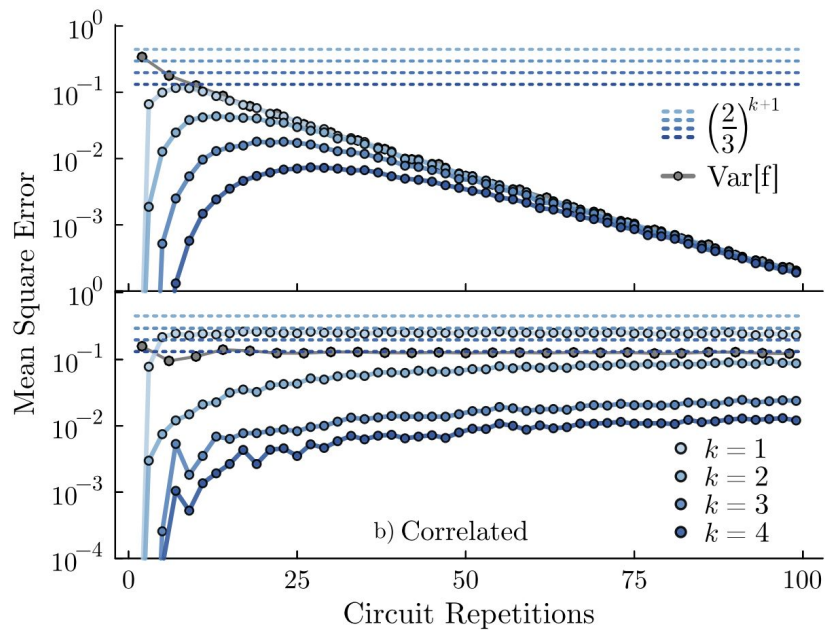
Zhang, Hao-Kai, Shuo Liu, and Shi-Xin Zhang. "Absence of barren plateaus in finite local-depth circuits with long-range entanglement." *Physical Review Letters* 132.15 (2024): 150603.

Numerical certification: intuition

- The average error depends only on the squared amplitudes $a^2, b^2, c^2, d^2 \dots$
- Then we can MC sample the Pauli tree!



Experiments beyond our bounds



Classical Simulation via Pauli Propagation

- Naturally harnesses noise and scrambling
- Not natively hindered by entanglement or circuit connectivity

Caveats

- Theoretical guarantees only in certain settings (uncorrelated gates)
- Not suitable for sampling from noiseless circuits

Related works on quantum chaos

Low-weight truncations strategies have been developed independently by several research communities!

- Von Keyserlingk, Curt, Frank Pollmann, and Tibor Rakovszky. **"Operator backflow and the classical simulation of quantum transport."** Physical Review B 105.24 (2022): 245101.
- Rakovszky, Tibor, C. W. Von Keyserlingk, and Frank Pollmann. **"Dissipation-assisted operator evolution method for capturing hydrodynamic transport."** Physical Review B 105.7 (2022): 075131.
- Ramos-Marimón, Carlos, Stefano Carignano, and Luca Tagliacozzo. **"Pauli weight requirement of the matrix elements in time-evolved local operators: dependence beyond the equilibration temperature."** arXiv:2409.13603 (2024).

Related works on quantum chaos

Low-weight truncations strategies have been developed independently by several research communities!


- Von Keyserlingk, Curt, Frank Pollmann, and Tibor Rakovszky. **"Operator backflow and the classical simulation of quantum transport."** Physical Review B 105.24 (2022): 245101.
- Rakovszky, Tibor, C. W. Von Keyserlingk, and Frank Pollmann. **"Dissipation-assisted operator evolution method for capturing hydrodynamic transport."** Physical Review B 105.7 (2022): 075131.
- Ramos-Marimón, Carlos, Stefano Carignano, and Luca Tagliacozzo. **"Pauli weight requirement of the matrix elements in time-evolved local operators: dependence beyond the equilibration temperature."** arXiv:2409.13603 (2024).

We need a unified framework!



Thanks!
Questions?

An illustration of a balance scale. The left pan is higher and contains the letters 'XYZ' with a green checkmark above it. The right pan is lower and contains the letters 'ZYXZ' with a pink 'X' above it.

 Demo
Classically estimating
expectation values from...

