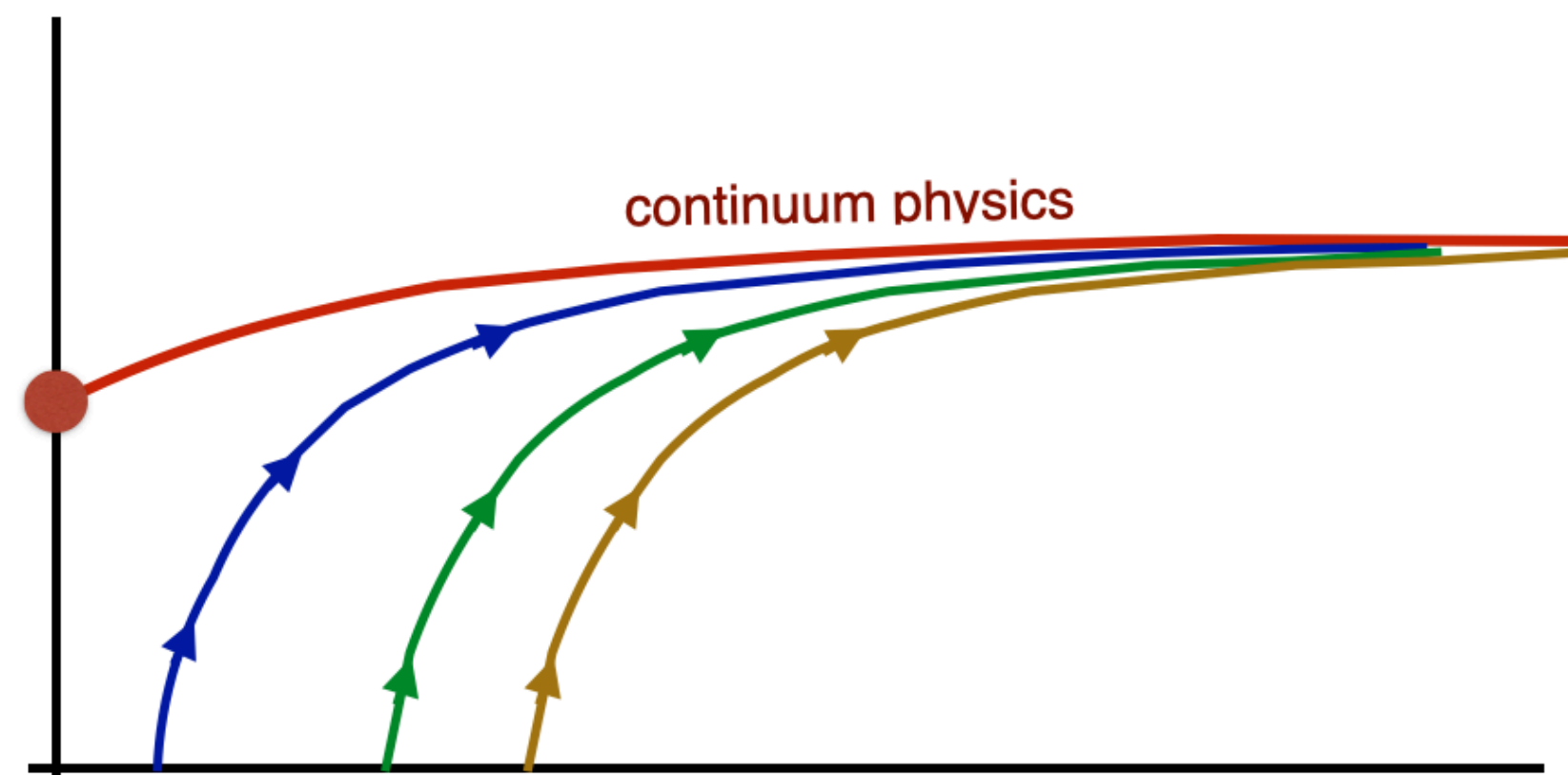


Gradient flow and the (Wilsonian) renormalization group

Anna Hasenfratz
University of Colorado Boulder

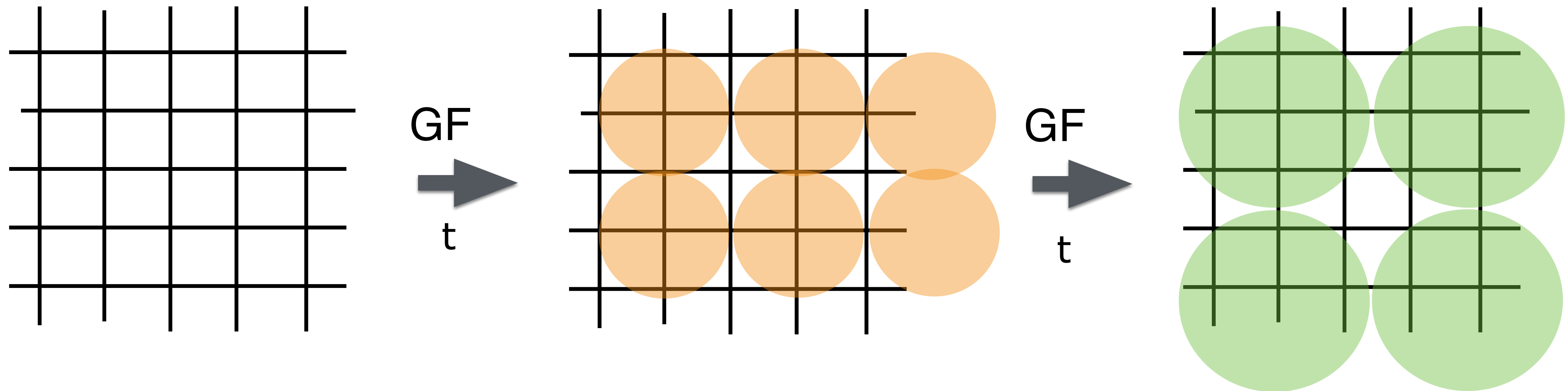
*Gradient Flow Workshop, Zurich
Feb 14, 2025*



Gradient flow: non-perturbative interpretation

GF is a continuous smearing transformation of radius $\rho \propto \sqrt{8t}$:
➔ defines “block spins” or “block links”

A. Carosso, AH, E. Neil,
PRL 121,201601 (2018)
Sonoda, H., Suzuki,
H. PTEP,023B05 (2021)



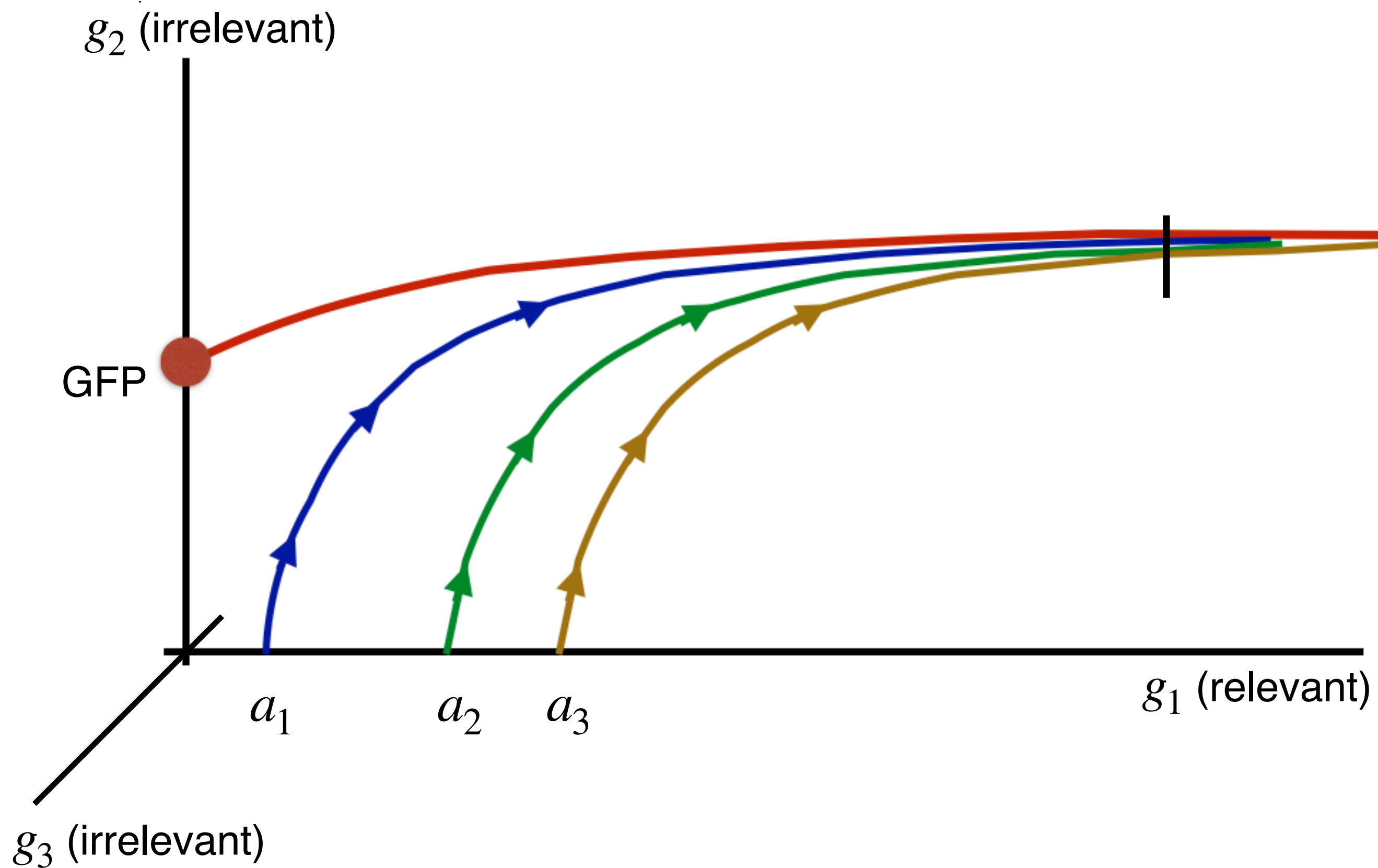
GF does not have coarse graining, not an RG transformation;
BUT it can be *interpreted* as Wilsonian RG with $\mu \propto 1/\sqrt{8t}$ for *local operators*

GF from lattice to continuum

Bare action parameter space:

$$\text{GF/RG: } \begin{array}{l} t = 0 : S(g_i, m_i; U) \longrightarrow t/a^2 : S(g_i(t), m_i(t); U(t)) \\ \mathcal{O}(g_i) \equiv \mathcal{O}(U) \qquad \qquad \qquad \mathcal{O}(g_i(t)) \equiv \mathcal{O}(U(t)) \end{array}$$

Urs Wegner's talk

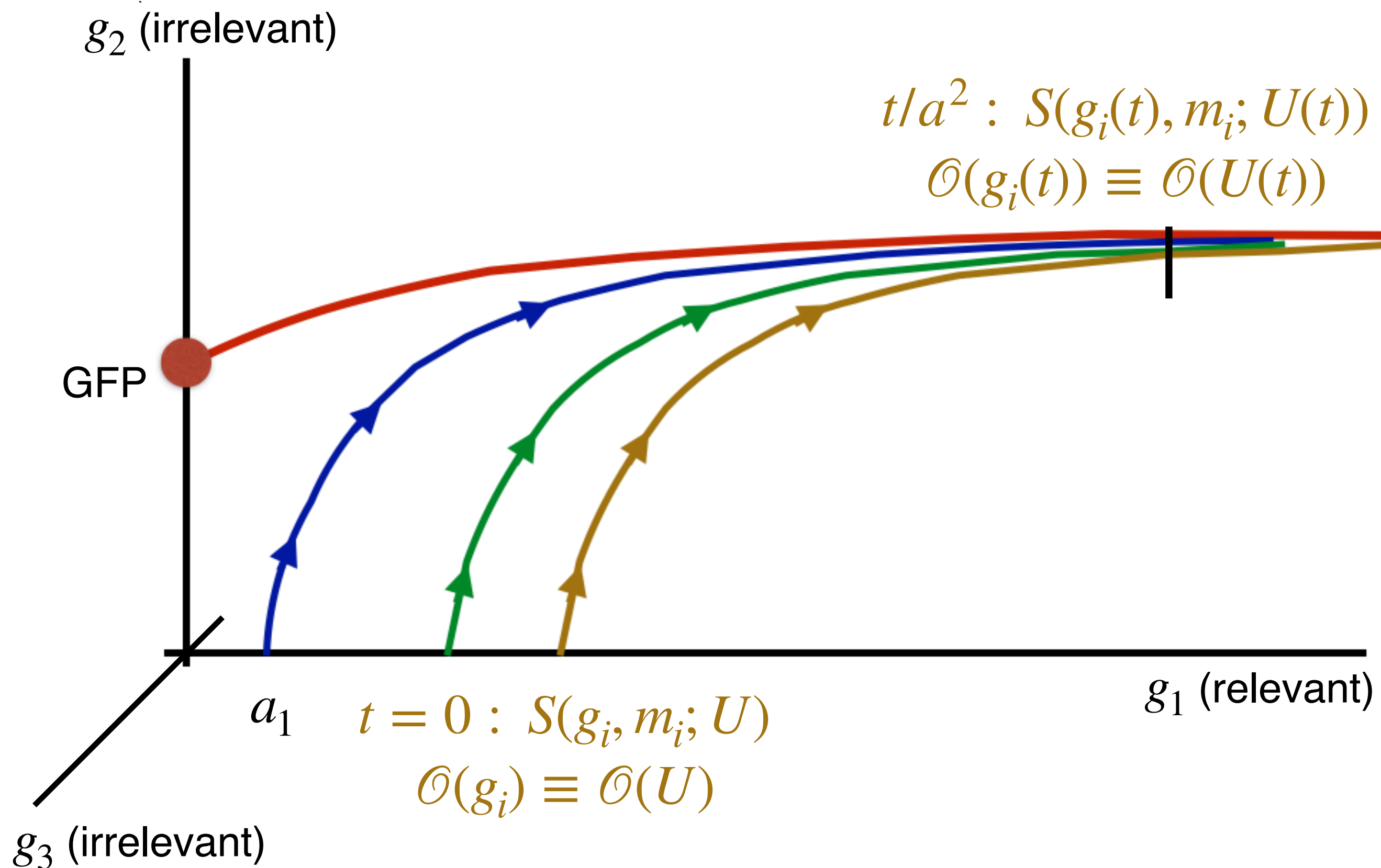


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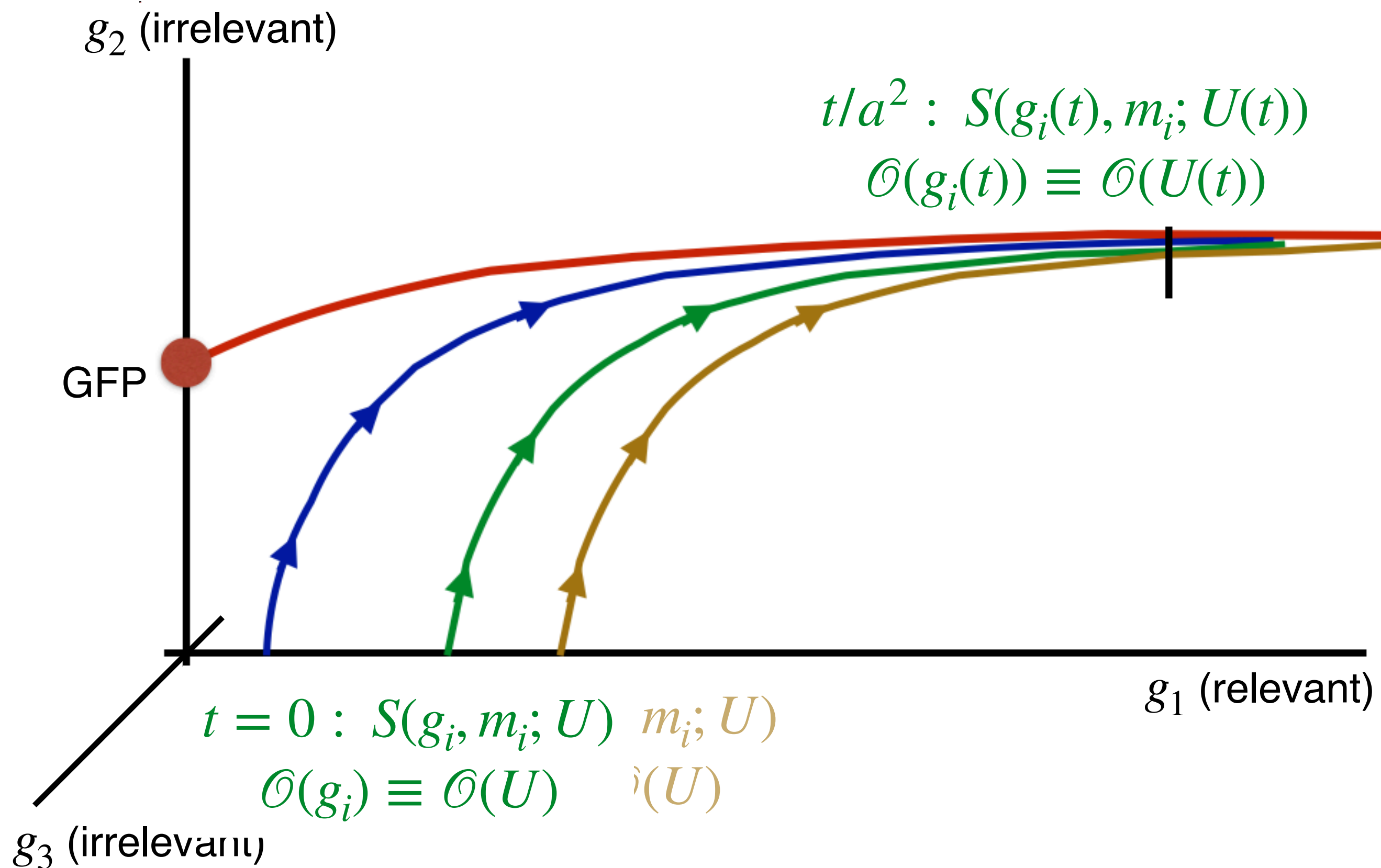
Urs Wegner's talk



GF from lattice to continuum

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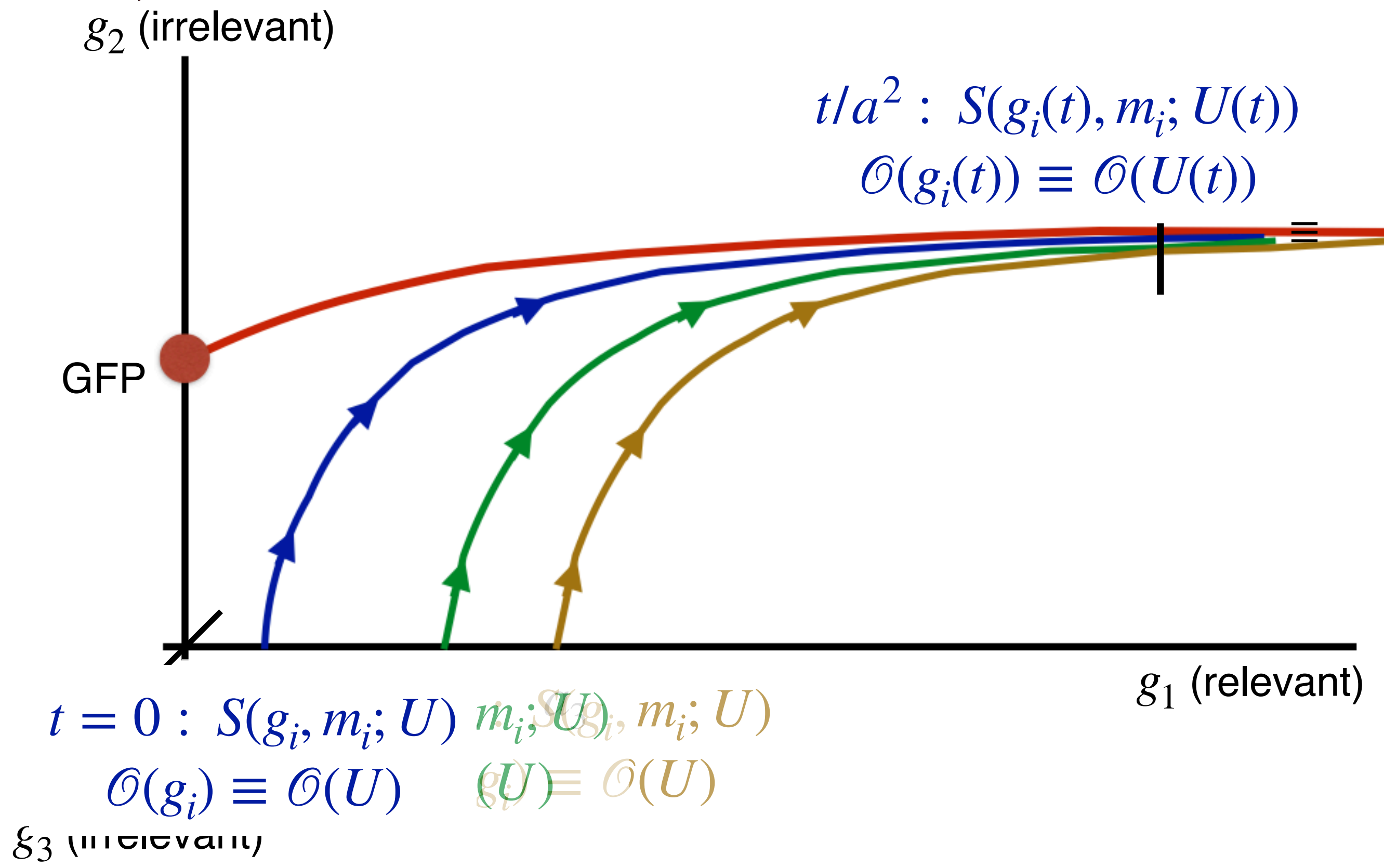
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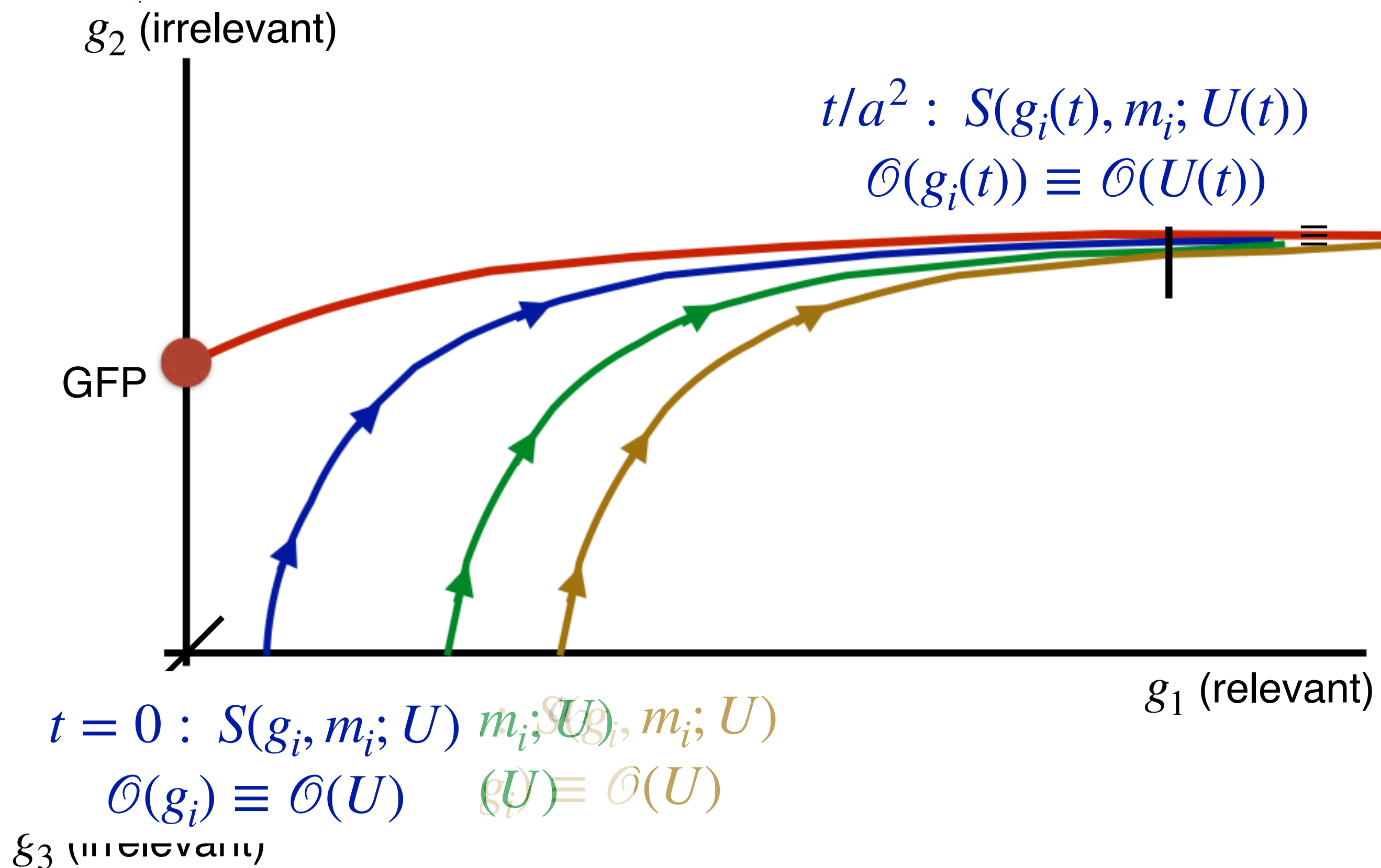


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Flowed actions, operators are (nearly) identical: renormalized



$$\begin{array}{l} t/a^2 : S(g_i(t), m_i(t); U(t)) \\ \mathcal{O}(g_i(t)) \equiv \mathcal{O}(U(t)) \end{array}$$

|||

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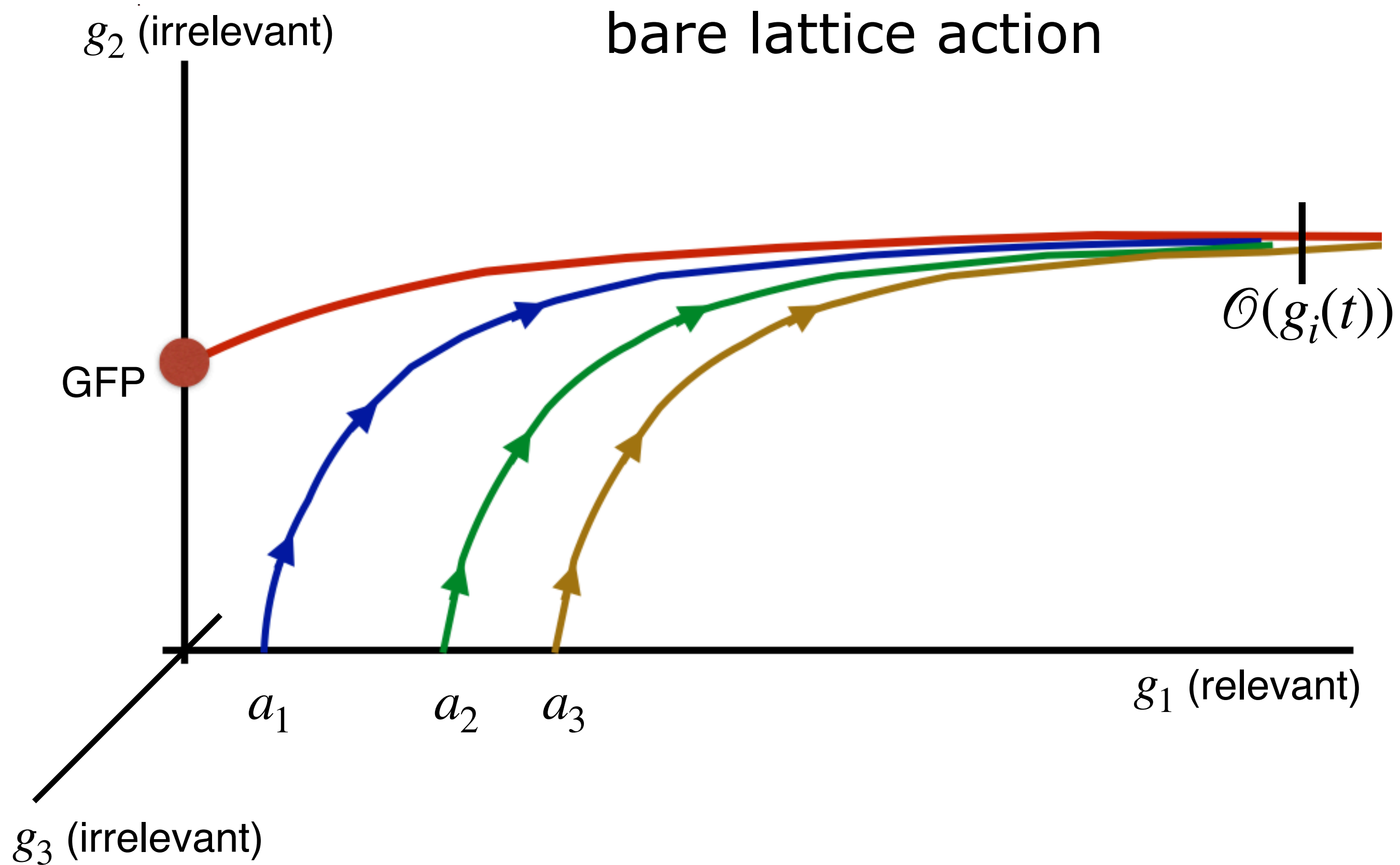
|||

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From lattice to continuum

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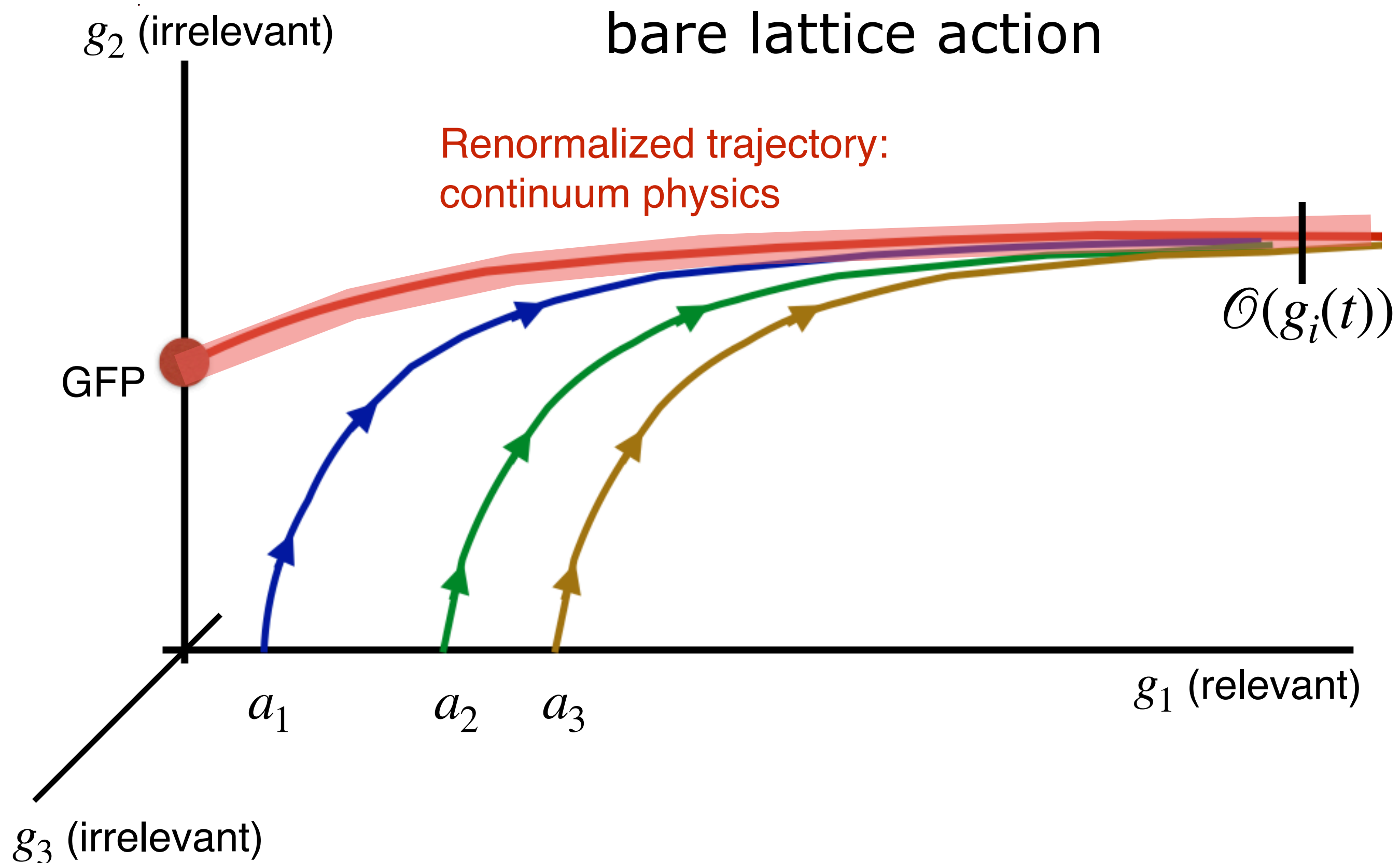
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From lattice to continuum

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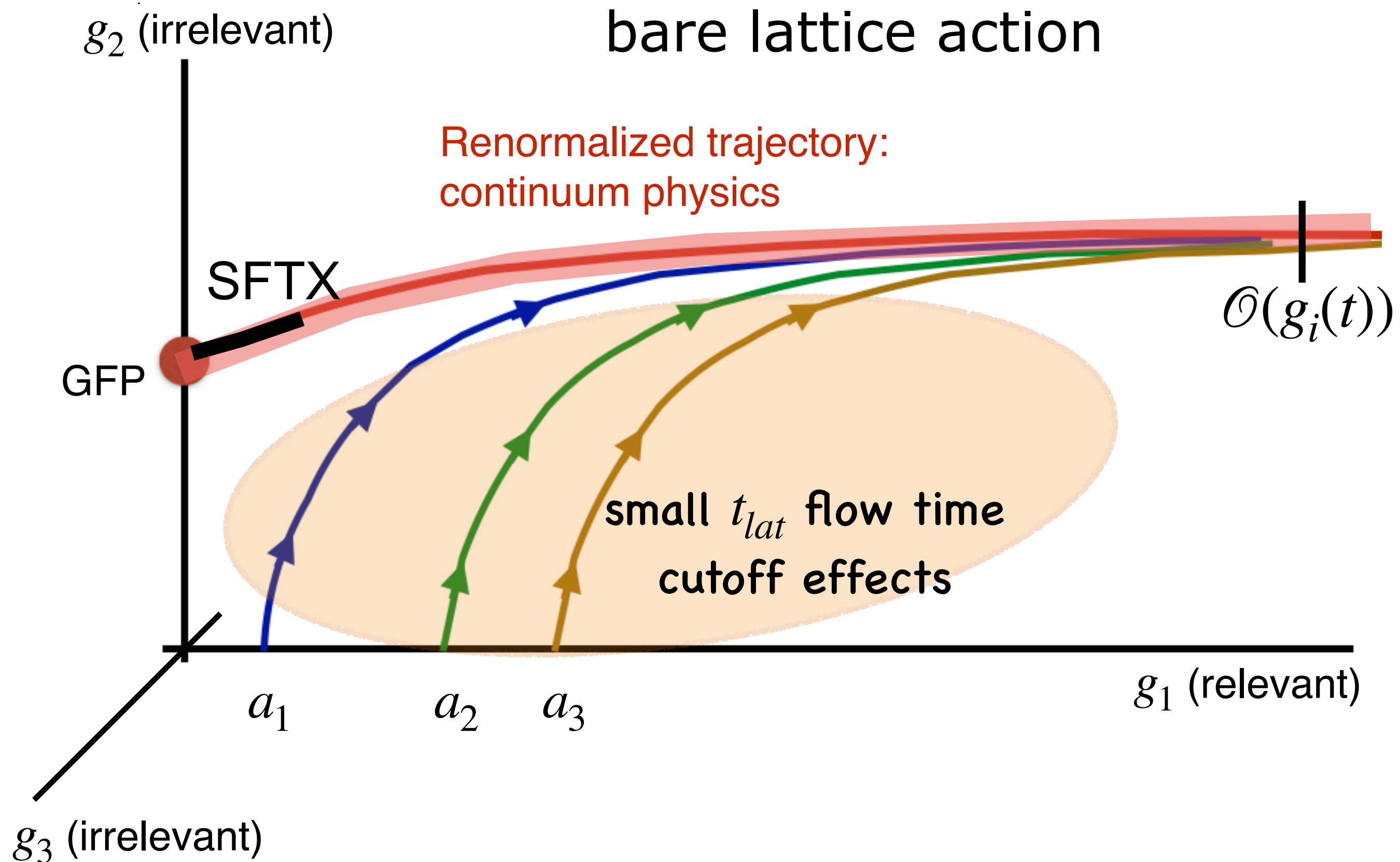
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Vary flow time to map RT and $\mathcal{O}(t)$

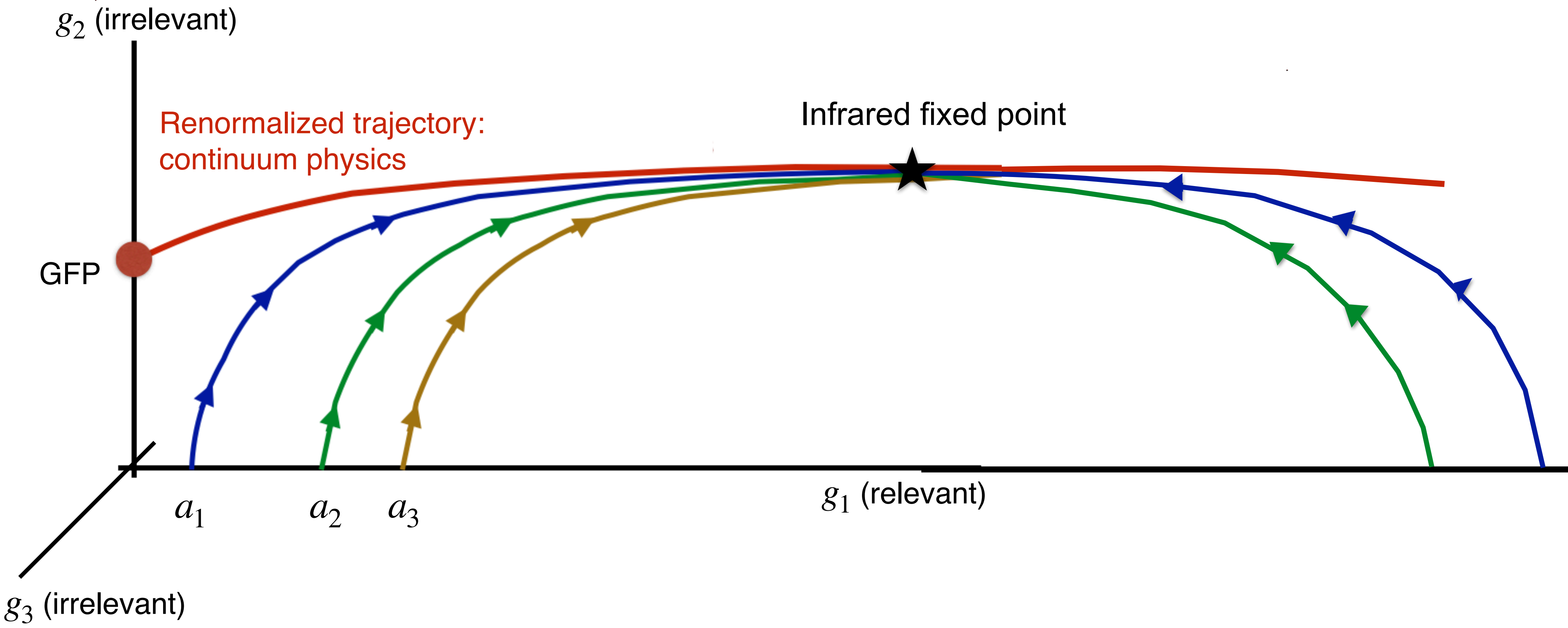
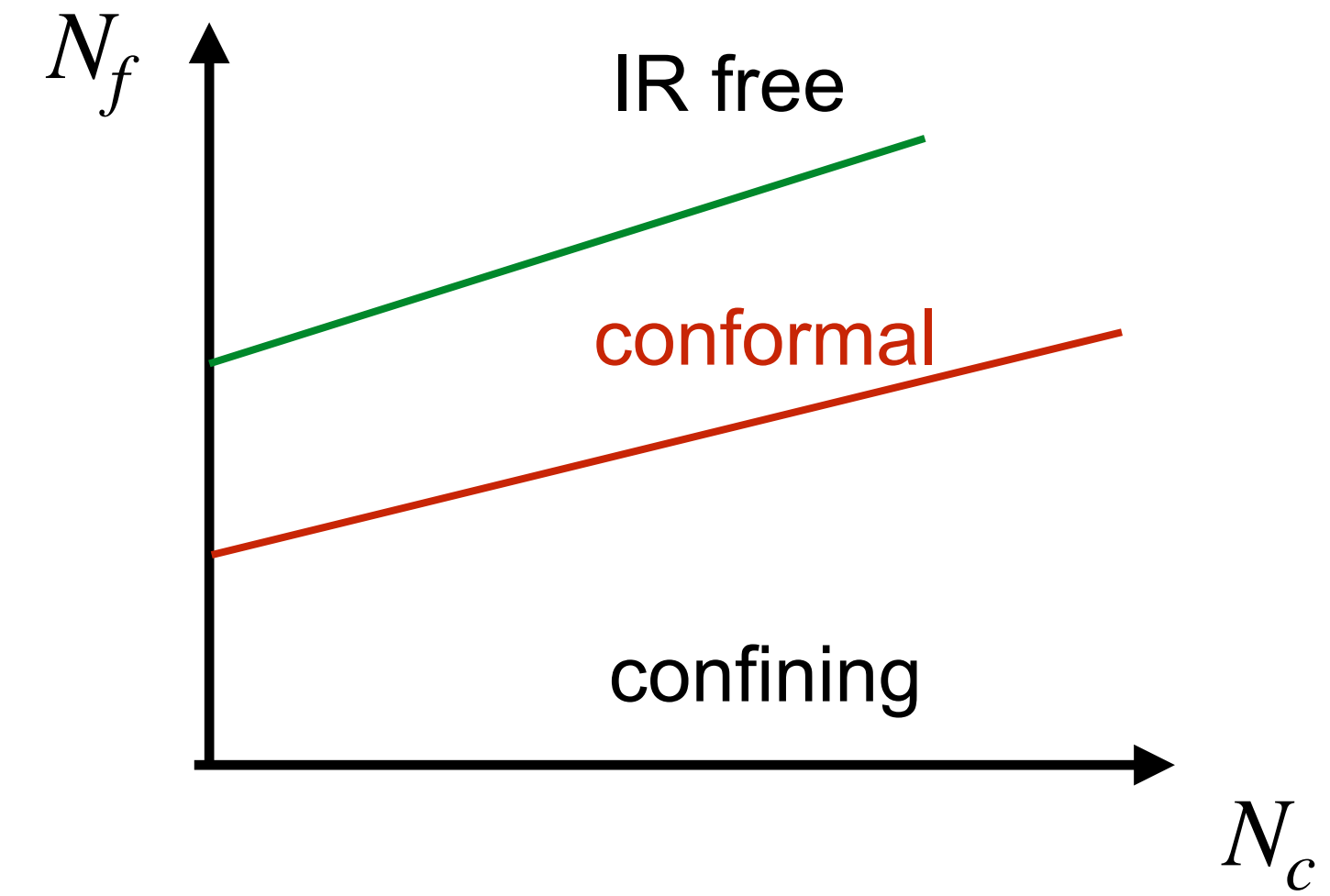
- small t_{lat} flow time: cutoff effects
- short flow time expansion is continuum, on RT

Lattice and SFTX are both on the RT
They should match!

In the conformal window

$SU(N_c)$ gauge + N_f fermions within the conformal window

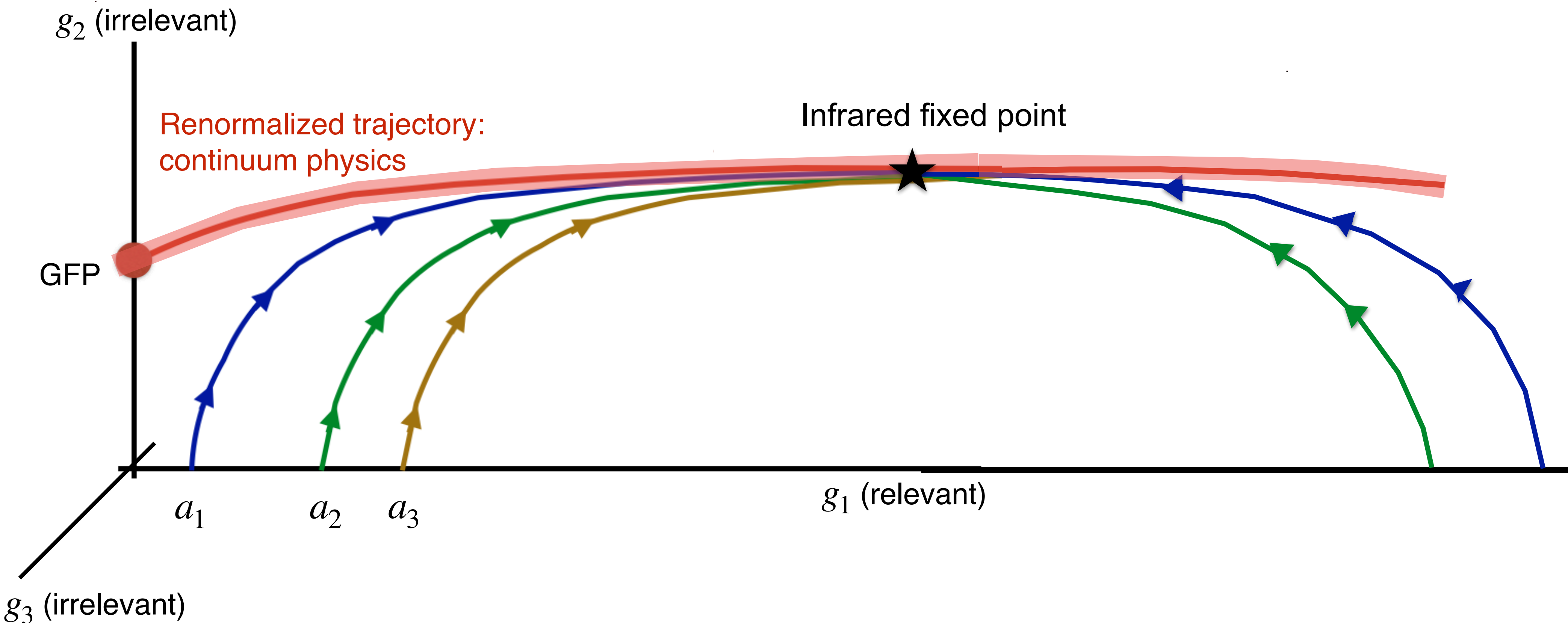
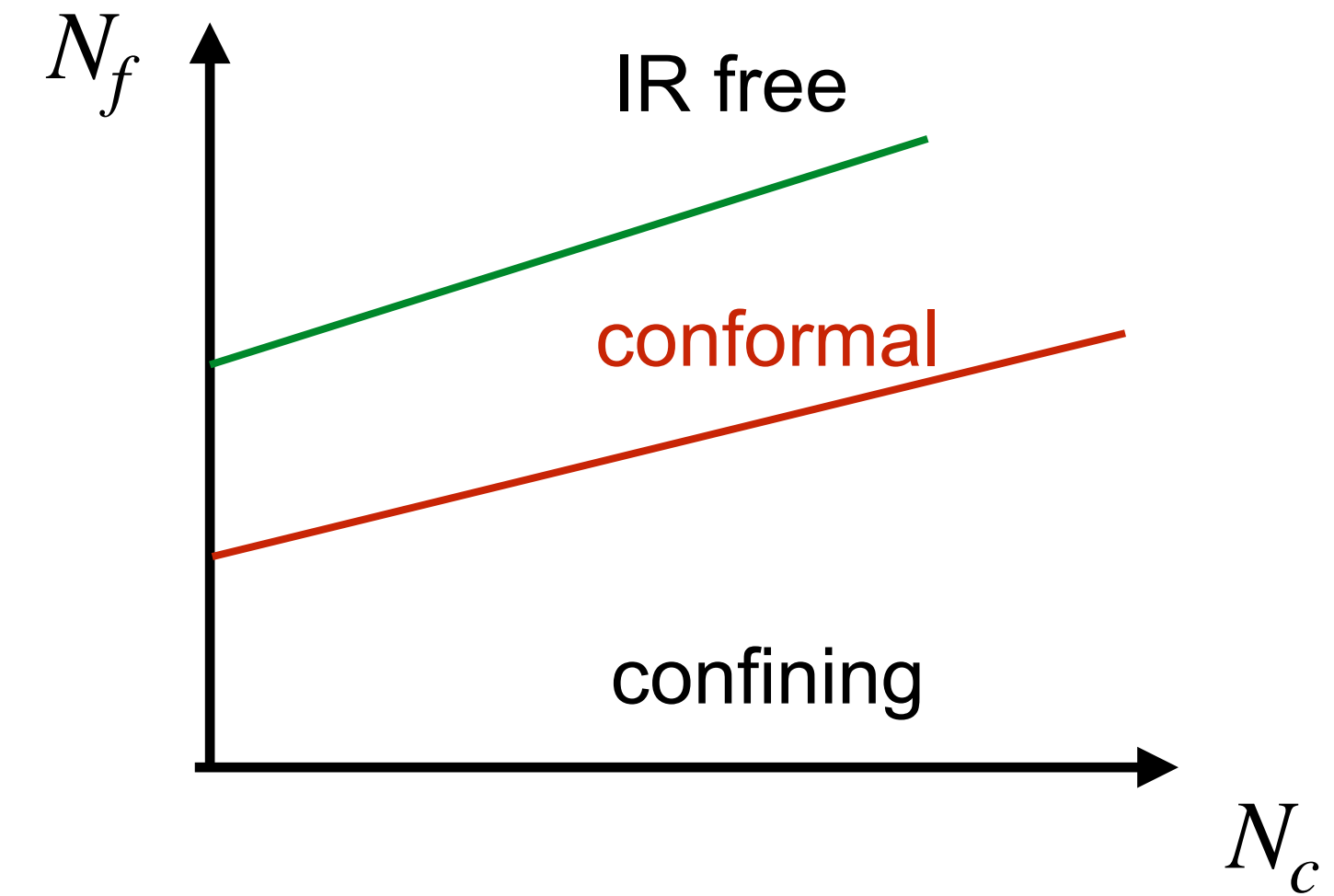
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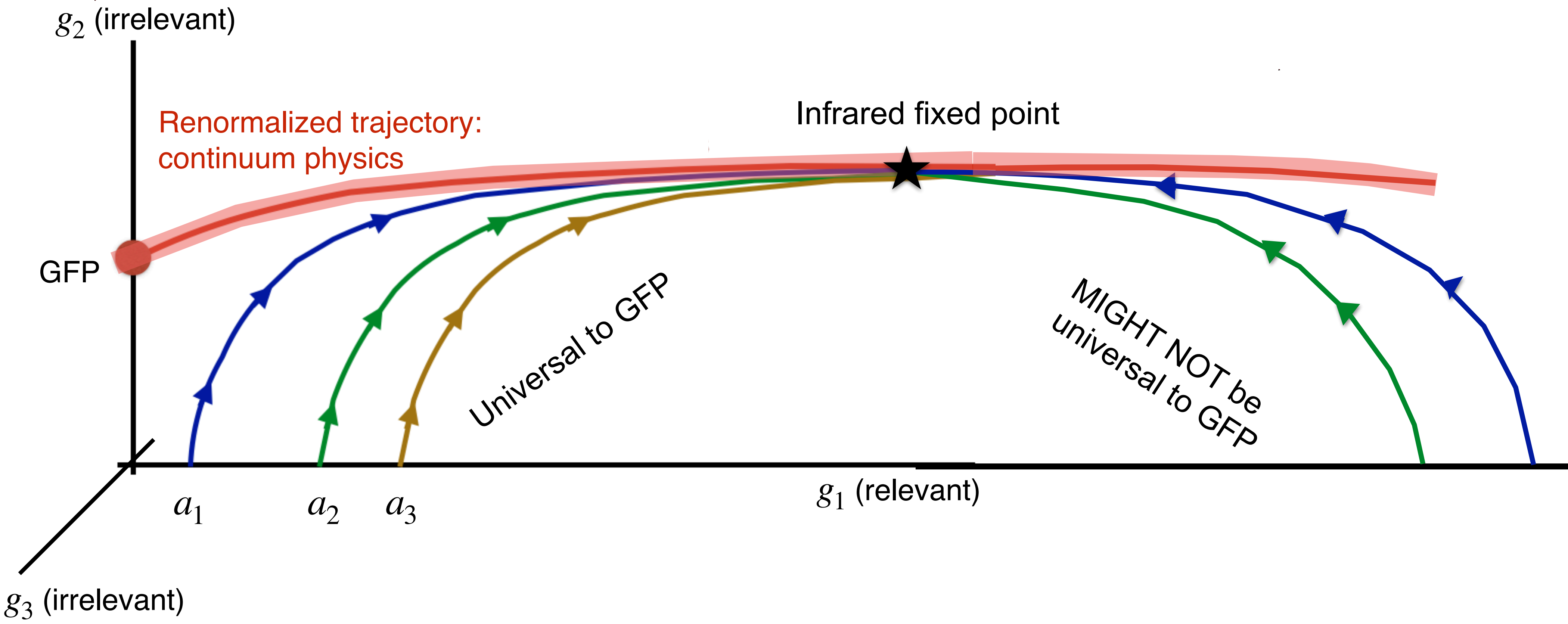
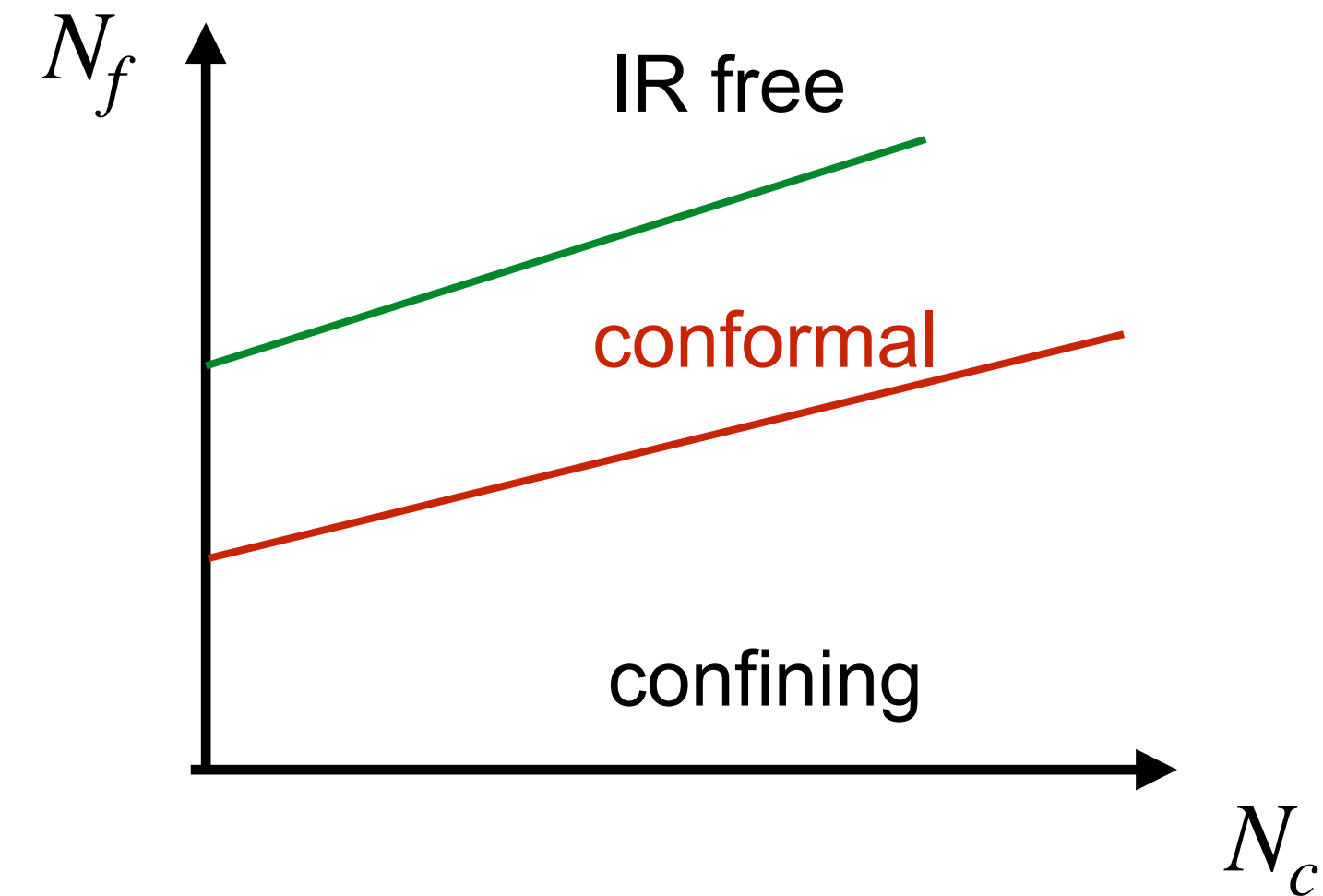
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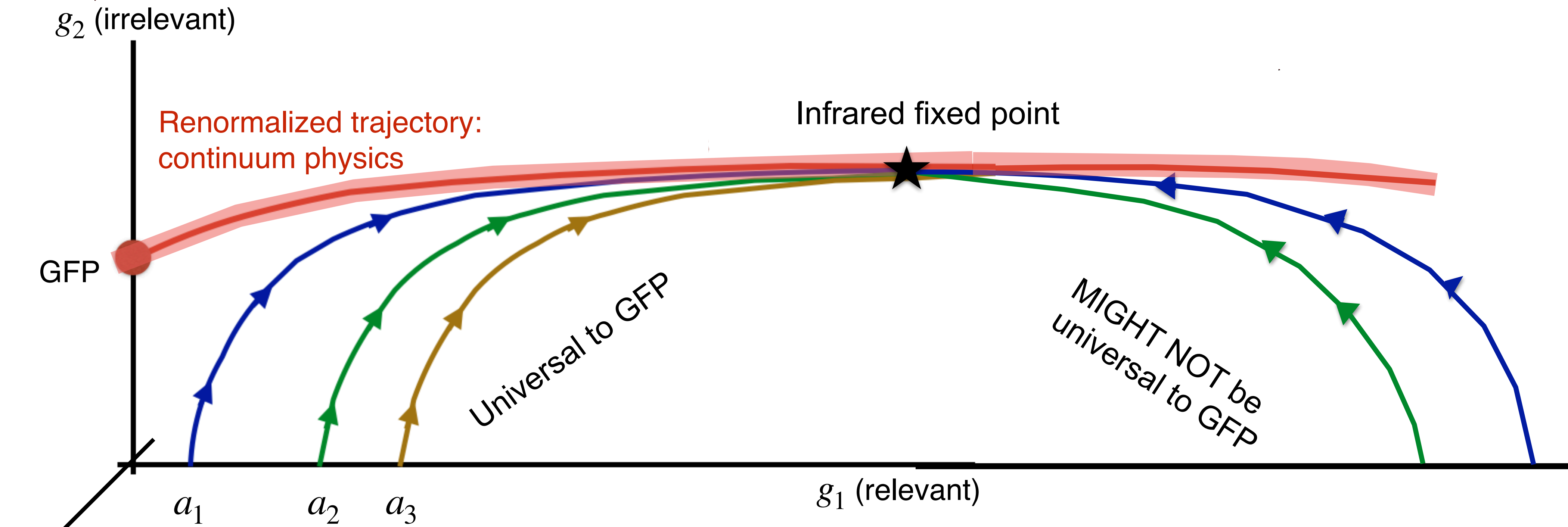
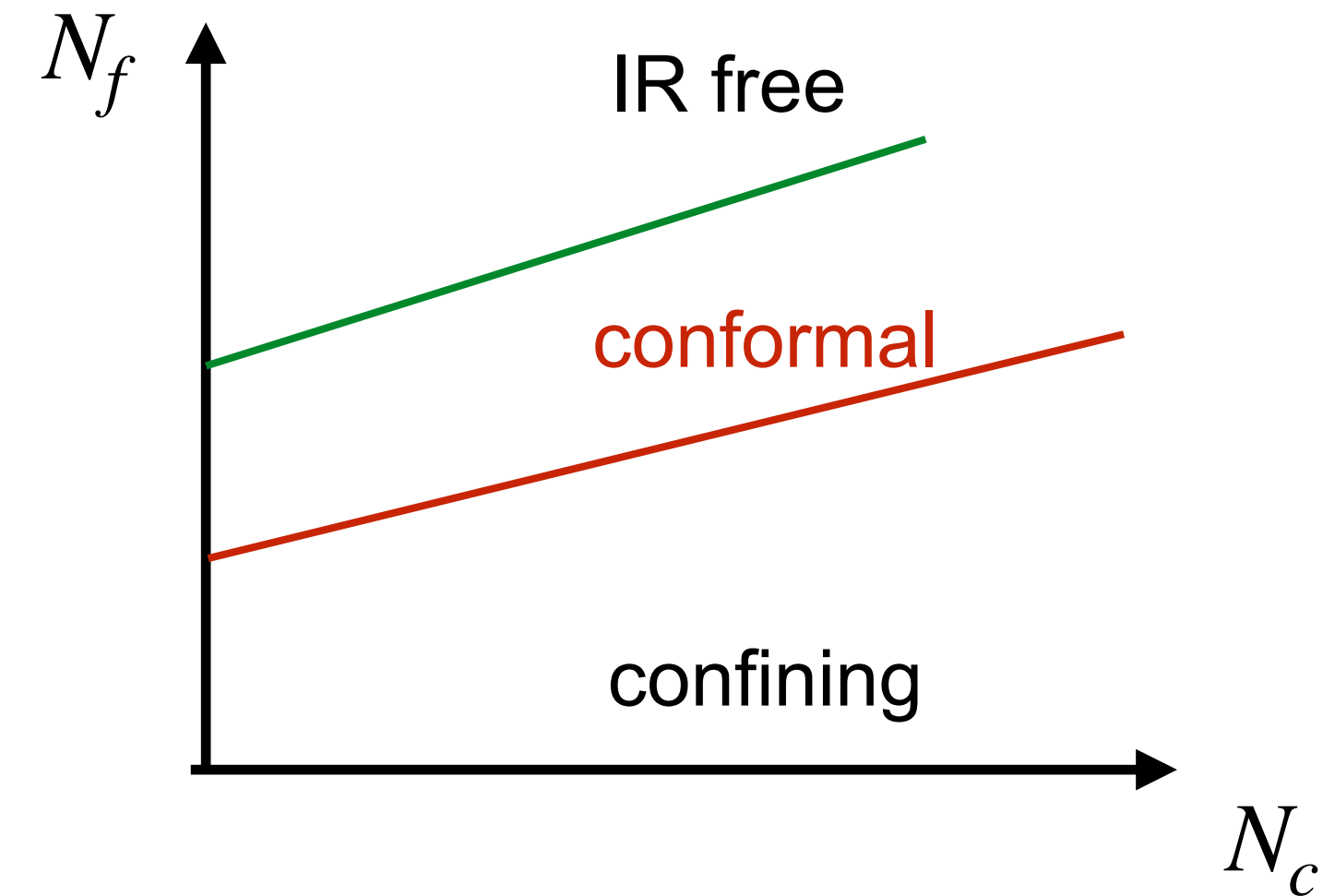
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The nonperturbative RG interpretation is still valid!

Operators:

Gauge flow:

$$\mathcal{O} \equiv g_{GF}^2 = \mathcal{N}t^2 \langle E(t) \rangle$$

→ $g_{GF}^2(t)$ is dimensionless, $\gamma_{g^2} = 0$

- along the RT it measures the flow
- renormalized coupling

$$g_{\overline{MS}}^2 = g_{GF}^2 + c g_{GF}^4 + \dots$$

→ RG β function: $\beta(g^2) = -t \frac{d g_{GF}^2}{d t}$

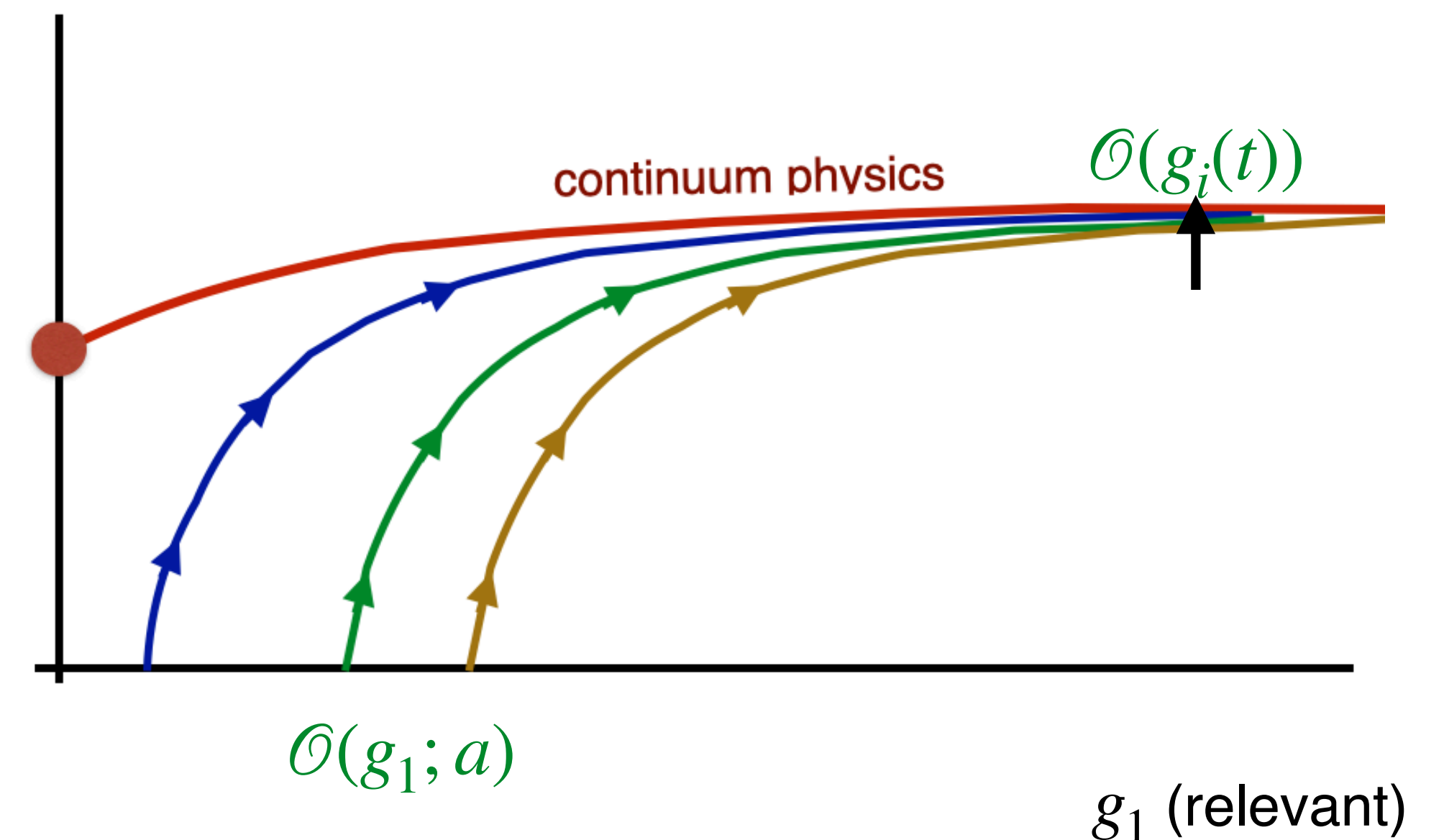
▸ Applications:

▸ QCD-like: $\Lambda_{GF}, \Lambda_{\overline{MS}} \rightarrow \alpha_{strong}$

▸ conformal: $\beta(g_{IRFP}^2) = 0 \rightarrow \text{IRFP}$

(continuous β fn vs step scaling)

Luscher *JHEP* 08 (2010) 071



AH, O. Witzel, *Phys.Rev.D* 101 (2020) 3

Fodor et al, *EPJWeb Conf.* 175, 08027 (2018)

Operators:

Fermion flow:

Luscher *JHEP* 04 (2013) 123

▸ bare operator: $\mathcal{O}(a, t = 0) = \bar{\psi}(x)\Gamma\psi(x)$

➔ flowed operator is renormalized

▸ running anomalous dimension $2t \frac{d\mathcal{O}(t)}{dt} \sim \gamma_{\mathcal{O}}(t)$

▸ match to SFTX

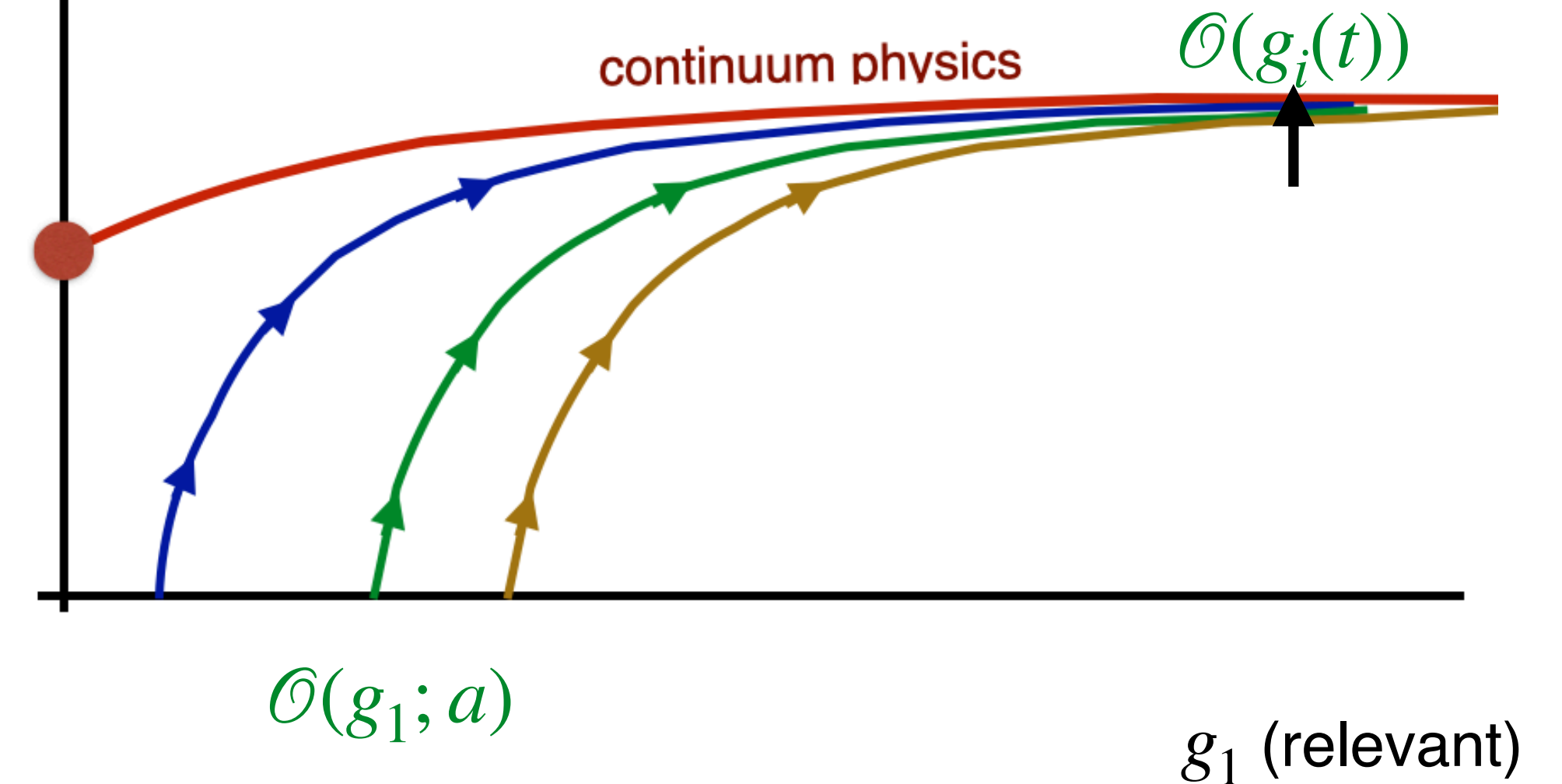
▸ RG scheme: matching factor

$$Z_{\mathcal{O}}^{-1}(a, t) \mathcal{O}(g_i(t), t) = \mathcal{O}(g_i(t), t)^0 \quad \text{tree level}$$

▸ Applications:

▸ QCD-like: running $\gamma_{\mathcal{O}}$, RG scheme

▸ conformal: predict $\gamma_{\mathcal{O}}^*$ at IRFP



Hasenfratz.,Neil, Shamir, Svetitsky, Witzel, *Phys.Rev.D* 108 (2023) 7
 Hasenfratz, Monahan, Schindler, Rizik, Witzel, in prep.

Fermion Operators:

A. Carosso, AH, E. Neil,
PRL 121,201601 (2018)

Often $\langle \mathcal{O}_\Gamma(t) \rangle = 0$; consider a GF two-point function

$$G_{\mathcal{O}}(x_4, t) = \int d^3x d^3x' \langle \mathcal{O}(\mathbf{x}, x_4; t) \mathcal{O}(\mathbf{x}', 0; t = 0) \rangle$$

- Only one operator is flowed, GF-RG equivalence (coarse graining) is OK

- Also need $x_4 \gg \sqrt{8t}$

- The scaling dimension of $G_{\mathcal{O}}(x_4, t)$ is $\Delta_{\mathcal{O}} = d_{\mathcal{O}} + \gamma_{\mathcal{O}} + \eta$
 - canonical dimension \nearrow $d_{\mathcal{O}}$
 - anomalous dimension \uparrow $\gamma_{\mathcal{O}}$
 - wave function renormalization $Z_\chi = e^{t\eta/2}$ \nwarrow η

The vector and axial charge operators have no anomalous dimension

→ in the ratio $\mathcal{R}_{\mathcal{O}}(x_4; t) = \frac{G_{\mathcal{O}}(x_4; t)}{G_V(x_4; t)}$ both Z_χ and $d_{\mathcal{O}}$ cancel

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Anomalous dimension and RG

$\gamma_{\mathcal{O}}$ is the logarithmic derivative

$$\gamma_{\mathcal{O}}(a; t) = 2t \frac{d \log \mathcal{R}_{\mathcal{O}}(a; x_4, t)}{dt} \quad ;$$

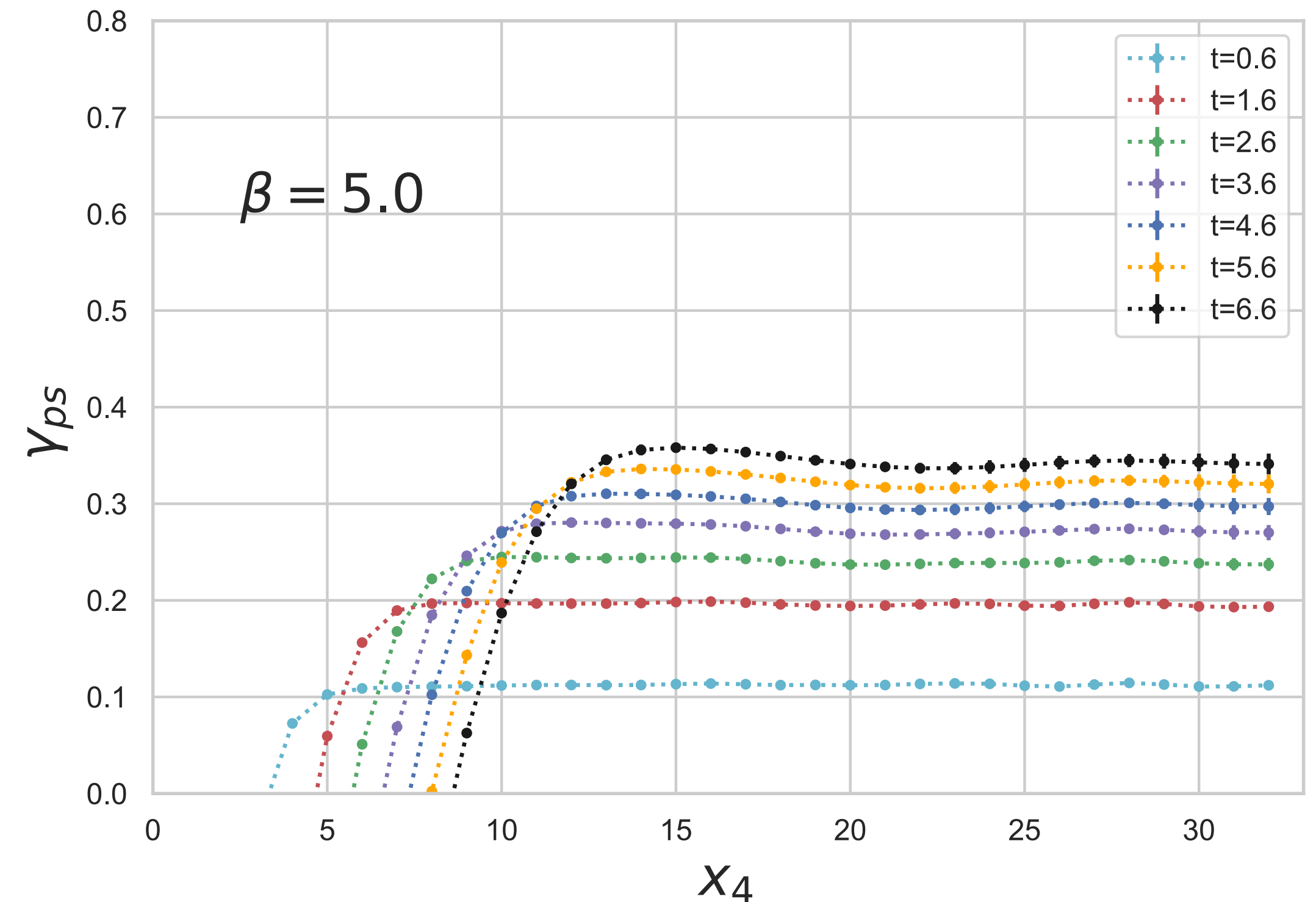
$$\mathcal{R}_{\mathcal{O}}(x_4; t) = \frac{G_{\mathcal{O}}(x_4; t)}{G_V(x_4; t)}$$

Typical correlator when $x_4 \gg \sqrt{8t}$

$$G_{\mathcal{O}}(t) = A_1(t)e^{-m_1 x_4} + A_2(t)e^{-m_2 x_4} + \dots$$

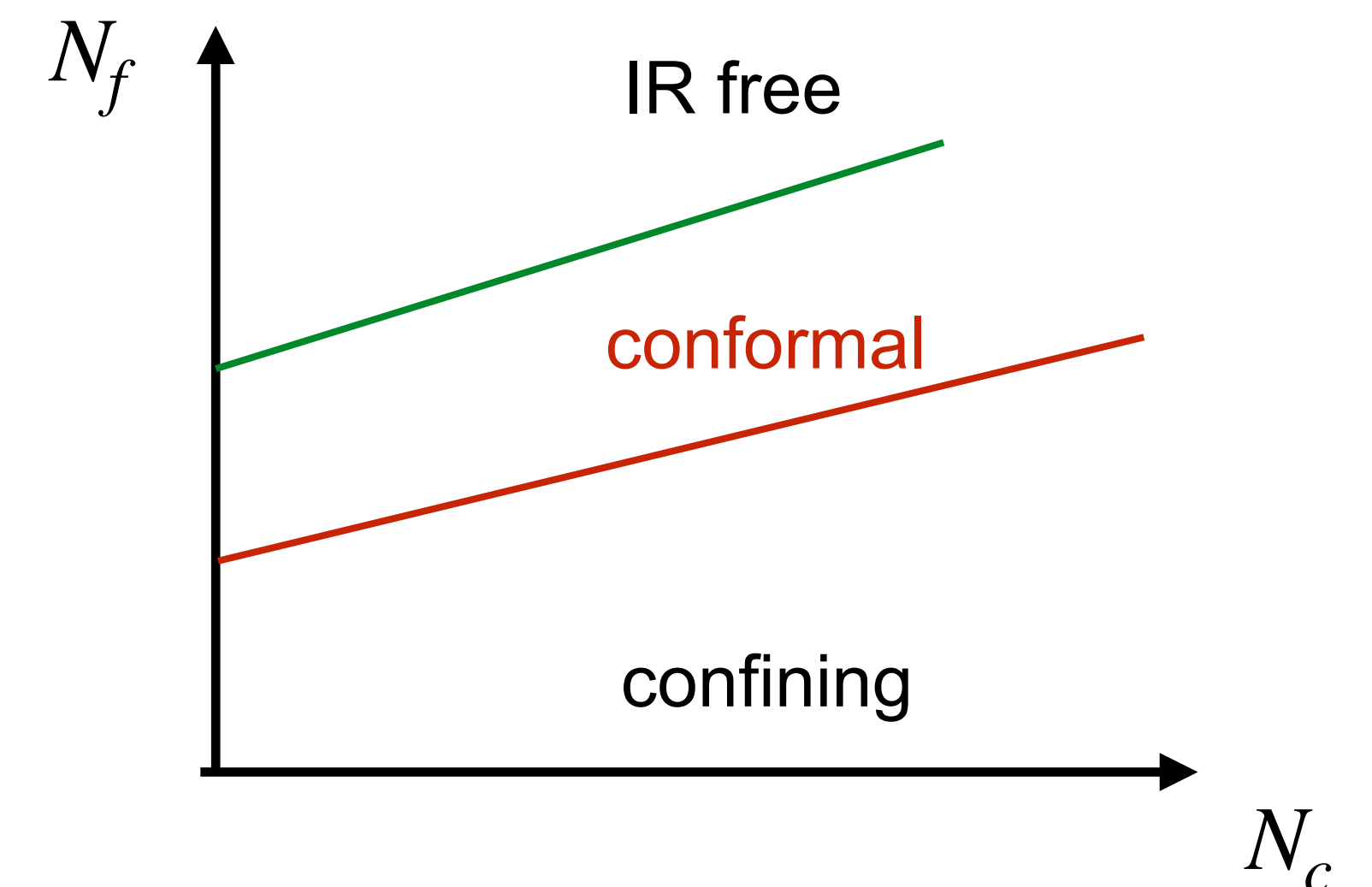
$$2t \frac{d \log G_{\mathcal{O}}(t)}{dt} = \frac{d \log A_1(t)}{dt} + \mathcal{O}(e^{-(m_2 - m_1)x_4})$$

- ▶ $\gamma_{\mathcal{O}}(t)$ is independent of x_4 if $x_4 \gg \sqrt{8t}$
- ▶ $\gamma_{\mathcal{O}}(t)$ corresponds to the lightest state; all others die out



Some recent applications:

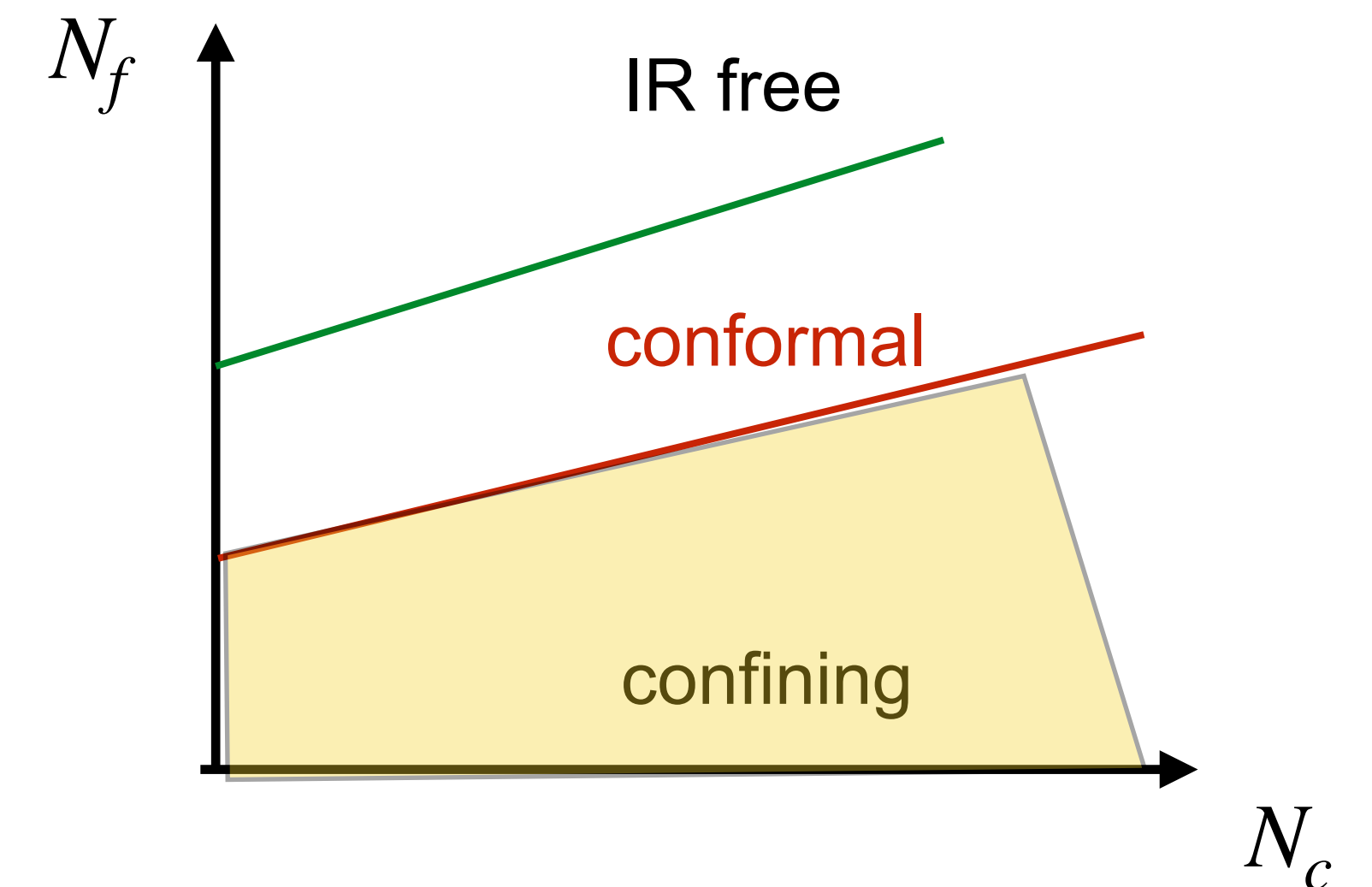
- SU(3) gauge, no fermions
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- SU(3) gauge + $N_f = 10$ fundamental fermions; Wilson lattice fermions
- SU(4) gauge + $N_f = 4 + 4$ fundamental+sextet fermions; Wilson lattice fermions
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Some recent applications:

QCD-like: scale setting, Λ_{QCD} , etc

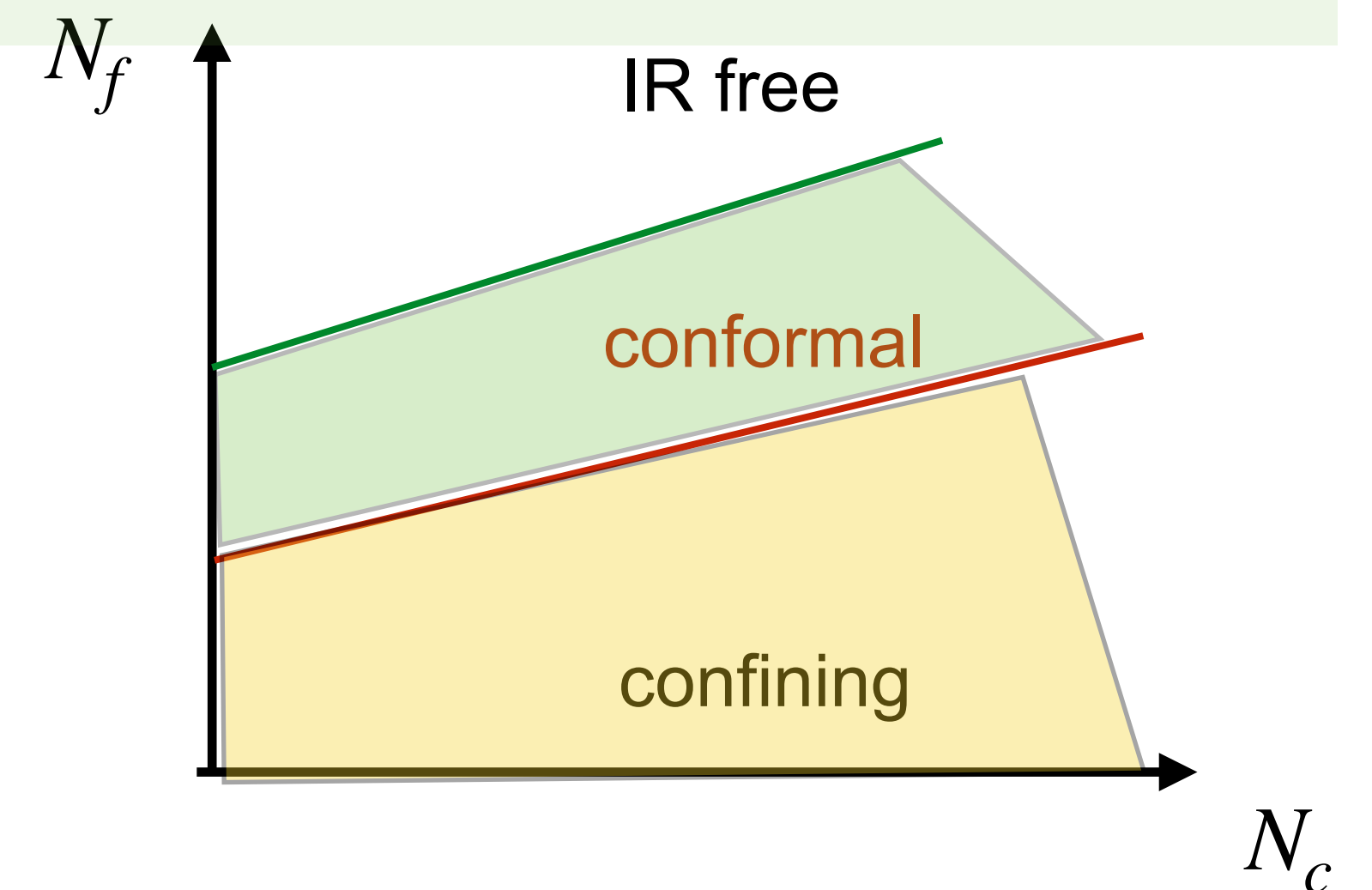
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Details, details,
(lattice)

Limits

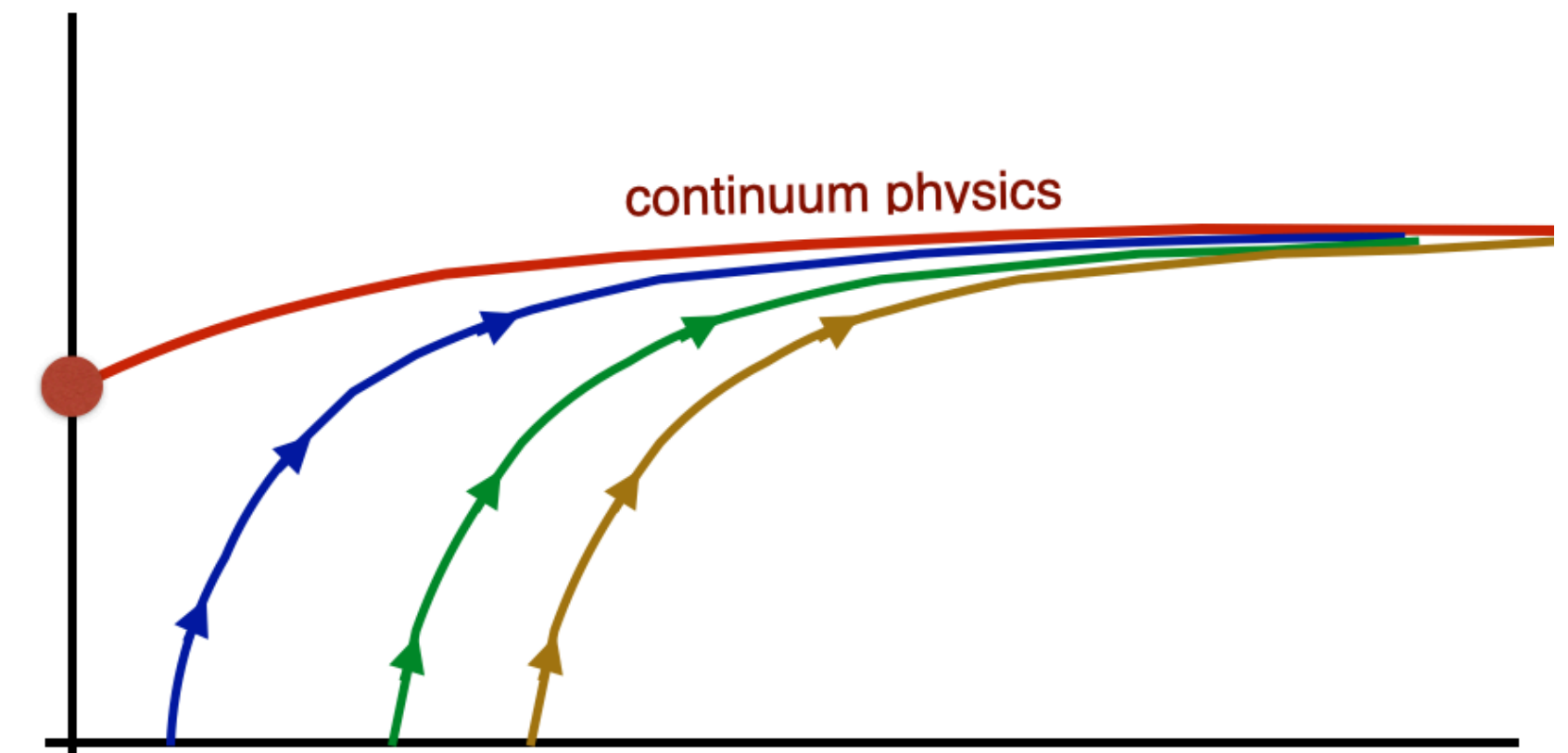
The RG picture I discussed is valid in infinite volume:

- take the $1/L \rightarrow 0$ limit ($g^2(L = \infty) = g^2(L) + \frac{c}{L^4} + \dots$)

Cancel small-flow time cutoff effects:

- take the $a^2/t \rightarrow 0$ continuum limit
(forces the bare coupling g_0^2 UVFP)

Can be delicate but straightforward



Improved actions

Improved actions reduce cutoff effects:

- perturbative : Symanzik program - useful in the weak coupling
- non-perturbative: (empirical)
 - cutoff effects can trigger unphysical bulk phase transitions :

➔ **Pauli-Villars improvement:**

AH, Shamir, Svetitsky, PRD104, 074509 (2021)

Add *heavy* PV bosons

- same interaction as fermions but with bosonic statistics
- $S_{eff} < 0$, $\beta = 2N_c/g_0^2$ increases : UV fluctuations decrease
 - in the IR the heavy flavors decouple, do not change physics
 - equivalently: range of effective gauge action is $\sim \exp(-2am_{PV})$

This only modifies the gauge action

Add many PV bosons reduce the lattice fluctuations

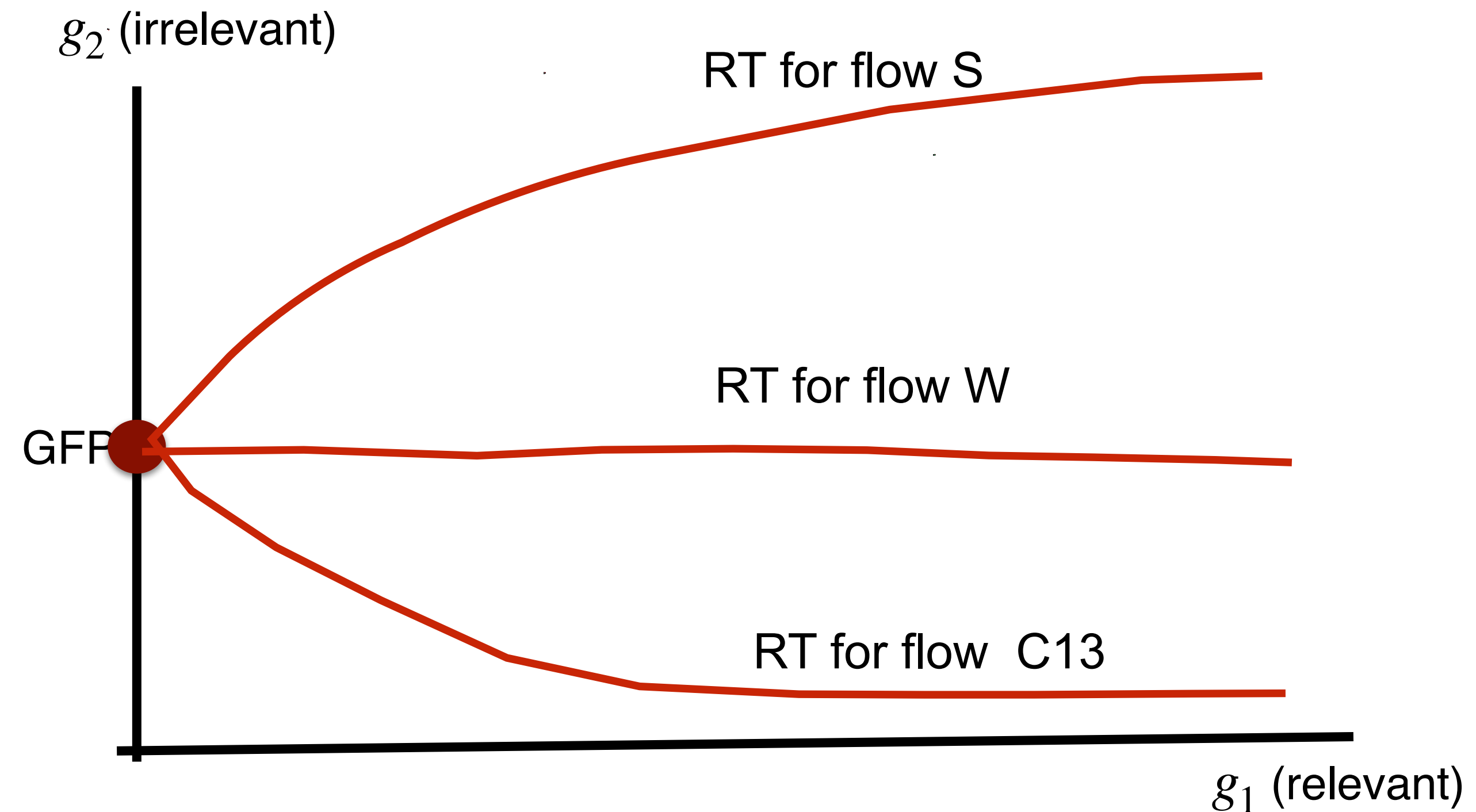
PV improvement was essential in the β function results with $N_f = 12, 10, 8$

Improved flows

When the RT is too far from the action, lattice artifacts of the flow are severe

➔ improve the flow

- perturbative : Zeuthen flow
- empirical non-perturbative:
Symanzik-like flow with different coefficients c_p, c_r
constraint: $c_p + 8c_r = 1$
 - $c_p = 1$: Wilson flow
 - $c_p = 5/3$: Symanzik flow
 - $c_p = 1/3$: C13 flow
 - etc but $c_p > 0$

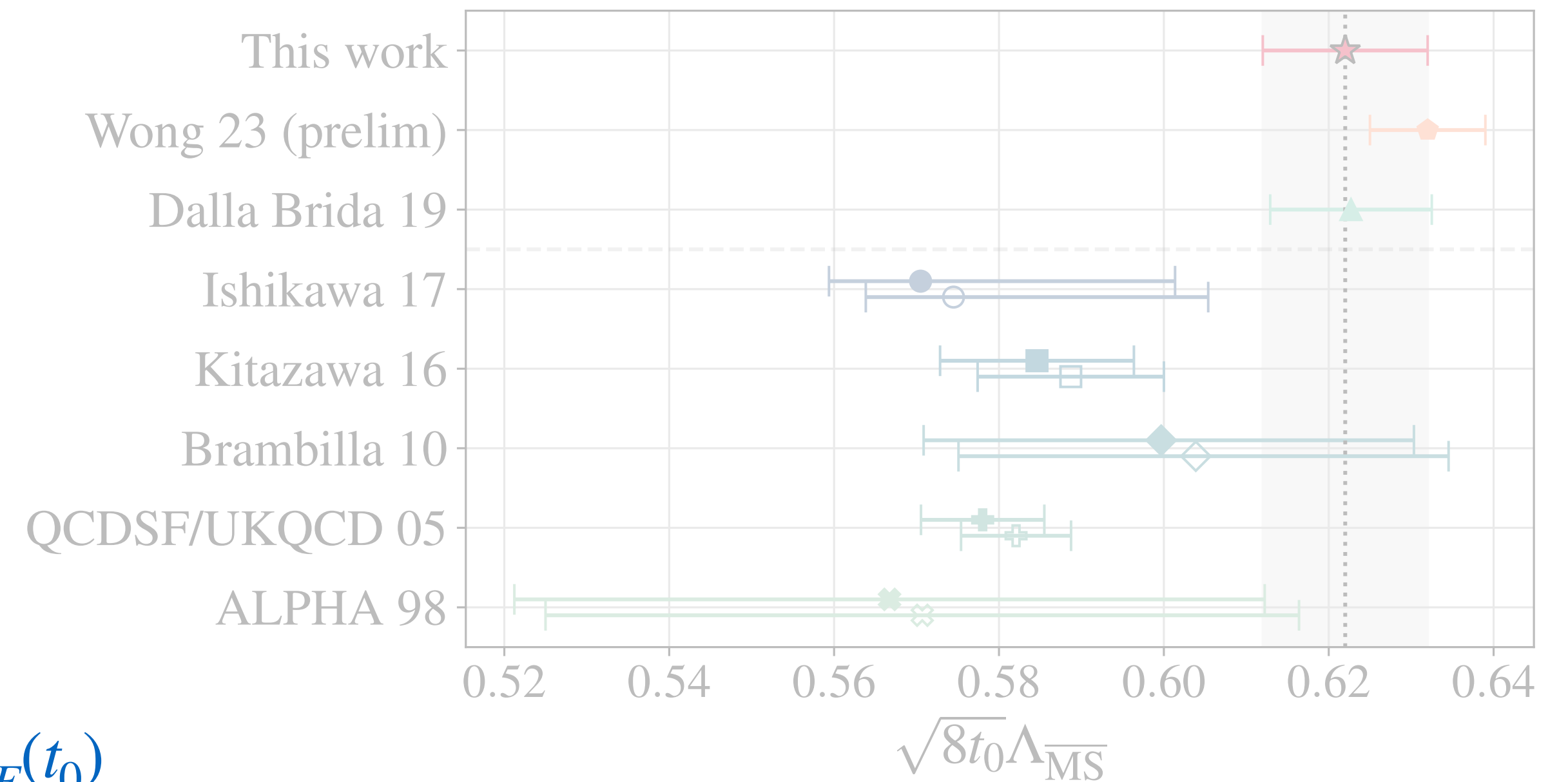
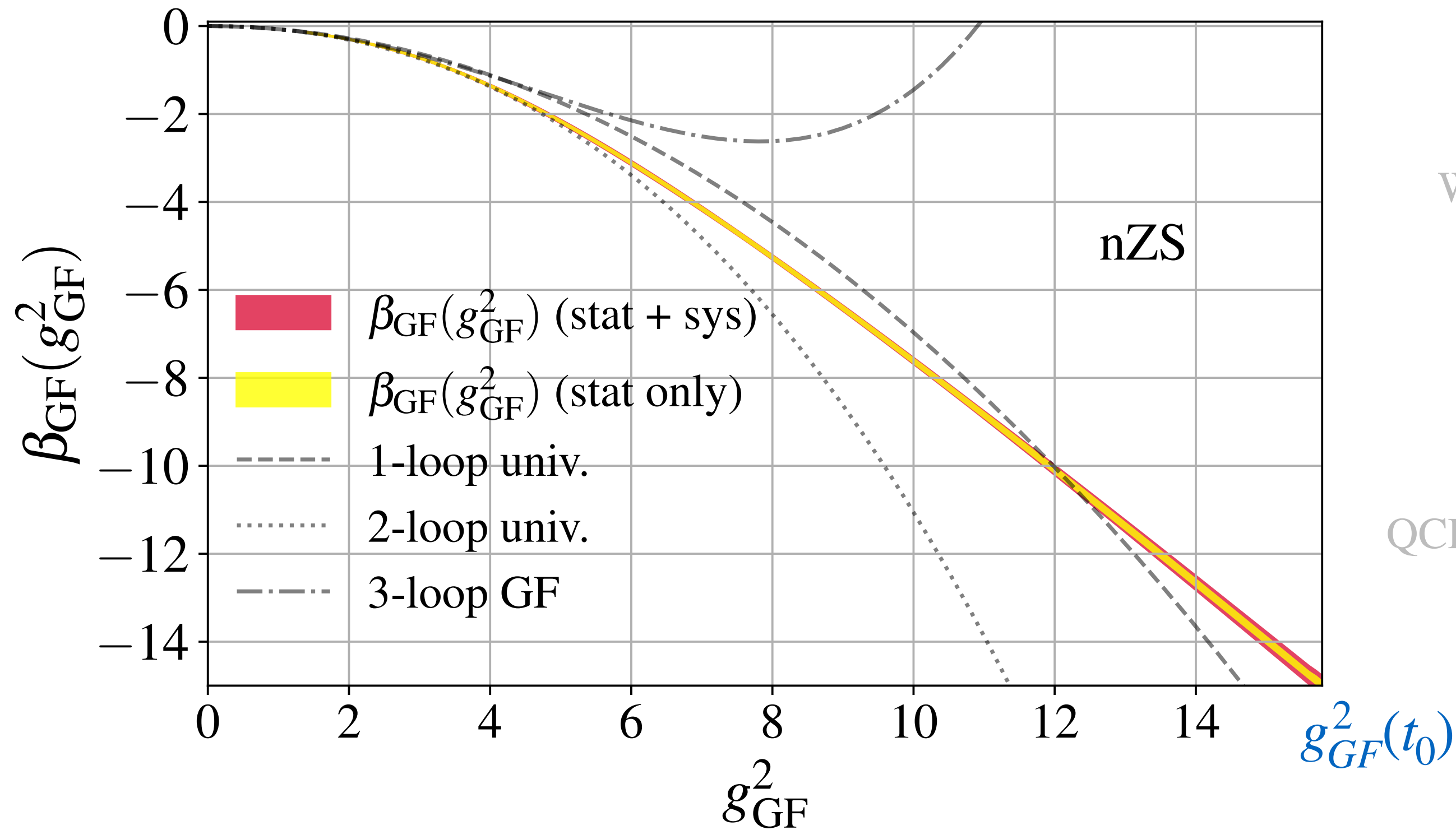


A few recent results

SU(3) Yang-Mills

(AH,C.Peterson,J.VanSickle,O.Witzel, arXiv: [2303.00704](https://arxiv.org/abs/2303.00704))

Wong et al - next talk



$\Lambda_{\overline{MS}}$ consistent with other GF results
 in tension with other (older) results

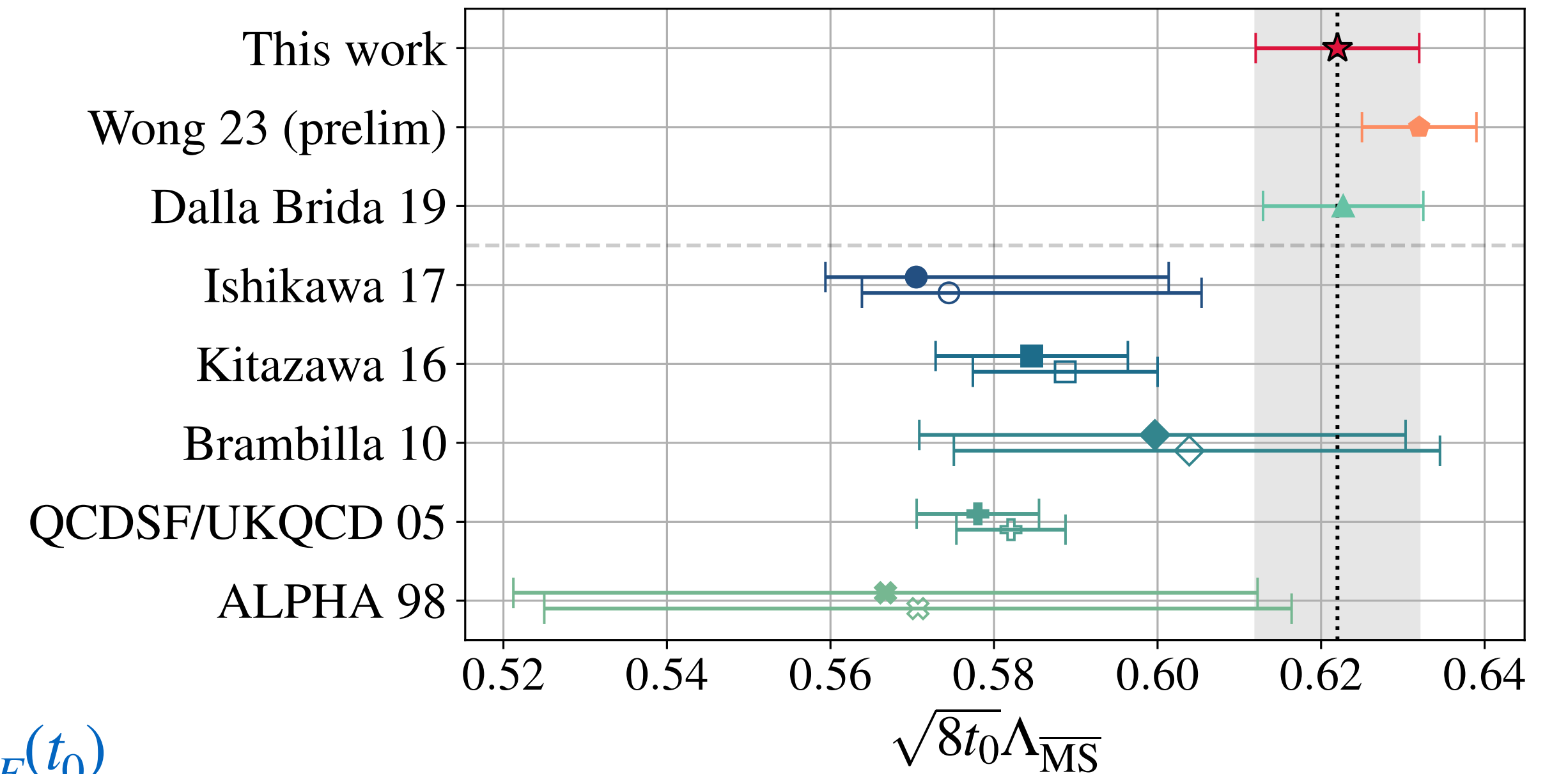
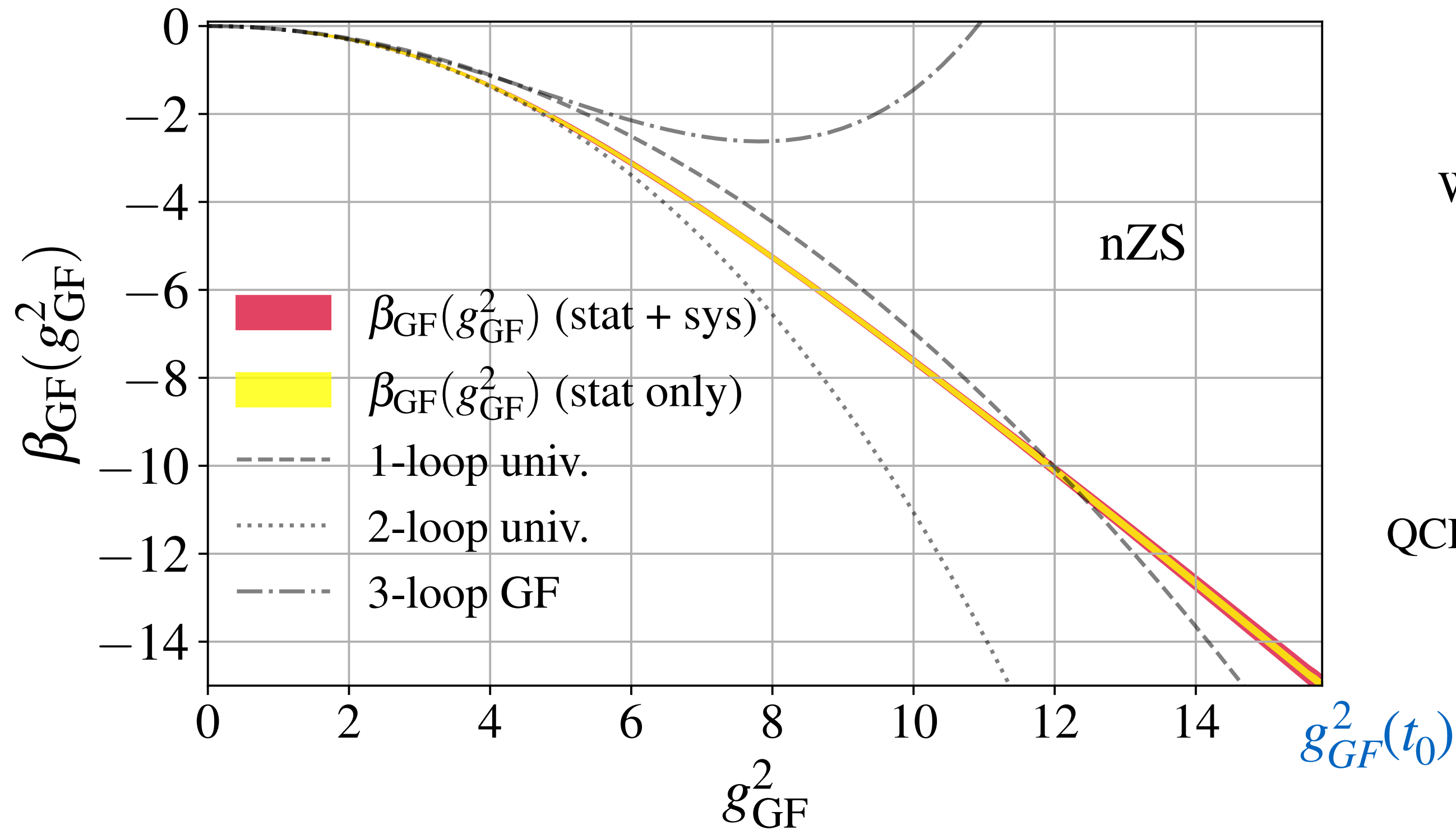
RG β function in the gradient flow scheme

- maps into perturbative curves at weak coupling
- very linear in the strong coupling (why?)

SU(3) Yang-Mills

(AH, C. Peterson, J. VanSickle, O. Witzel, arXiv: [2303.00704](https://arxiv.org/abs/2303.00704))

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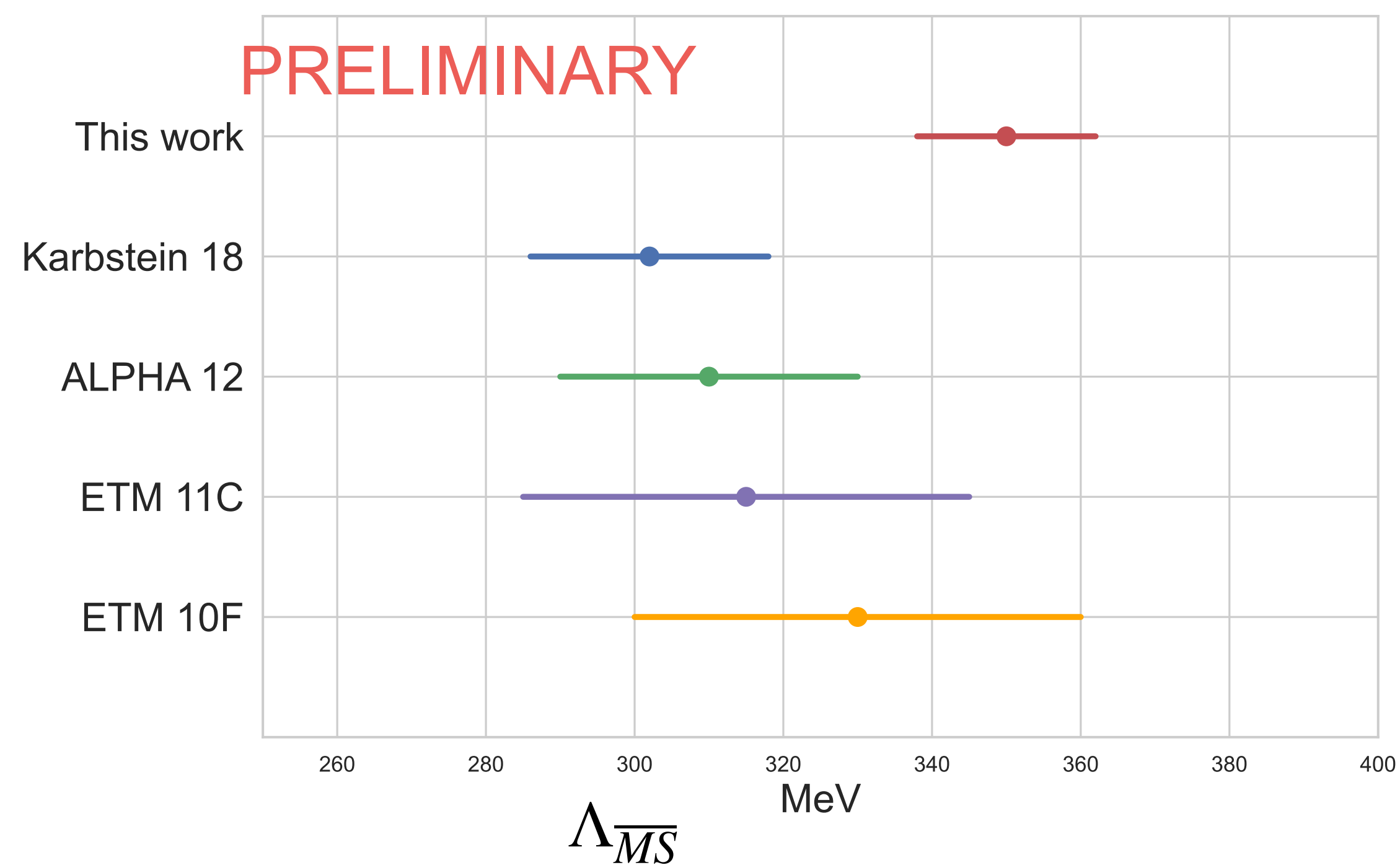
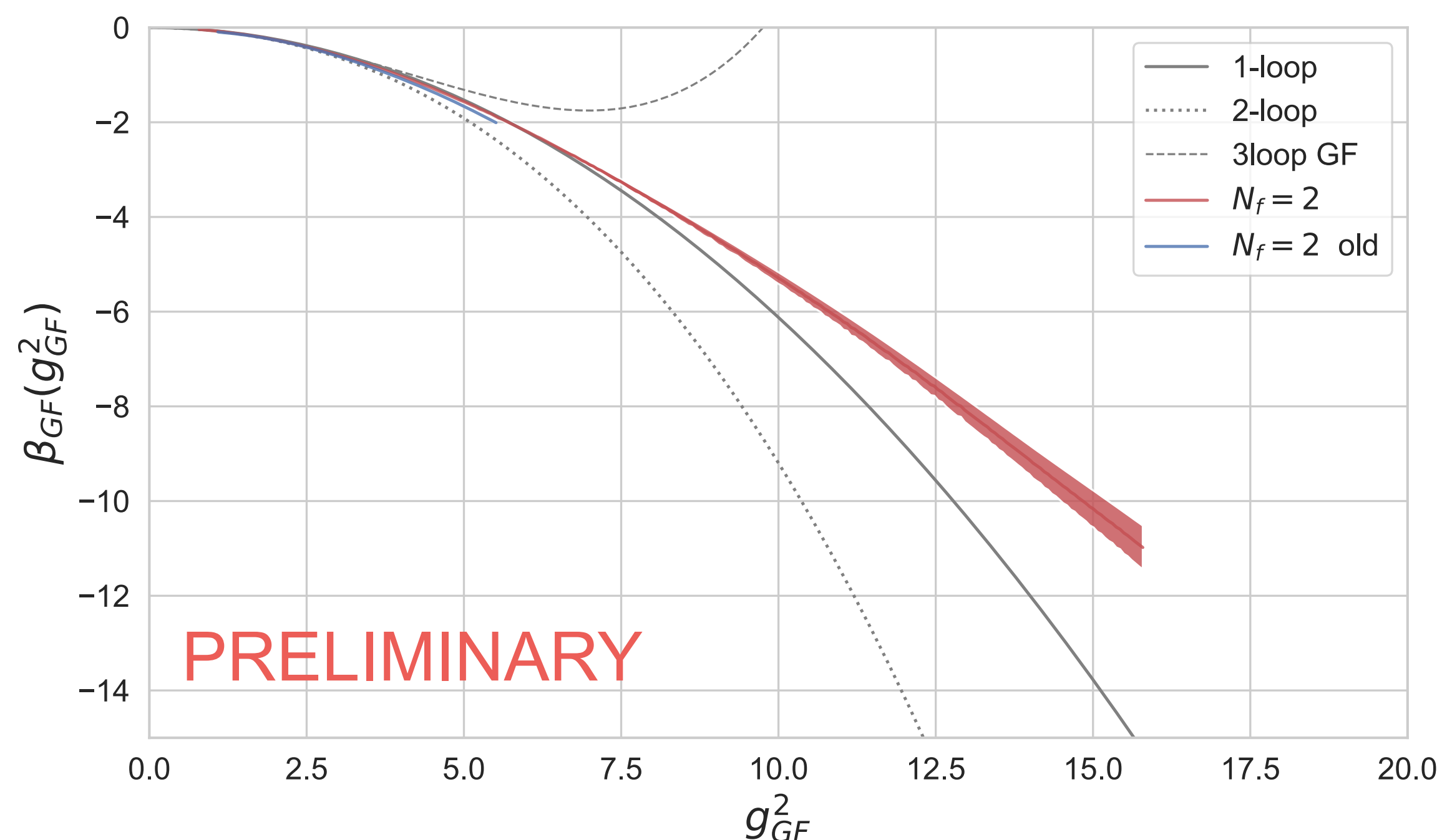
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$$SU(3) + N_f = 2$$

A.H., C. Monahan, M. Rizik,
A. Shindler and O. Witzel
Lattice'21 (arXiv:2201:09740)
+ in preparation

Pilot study with domain wall fermions



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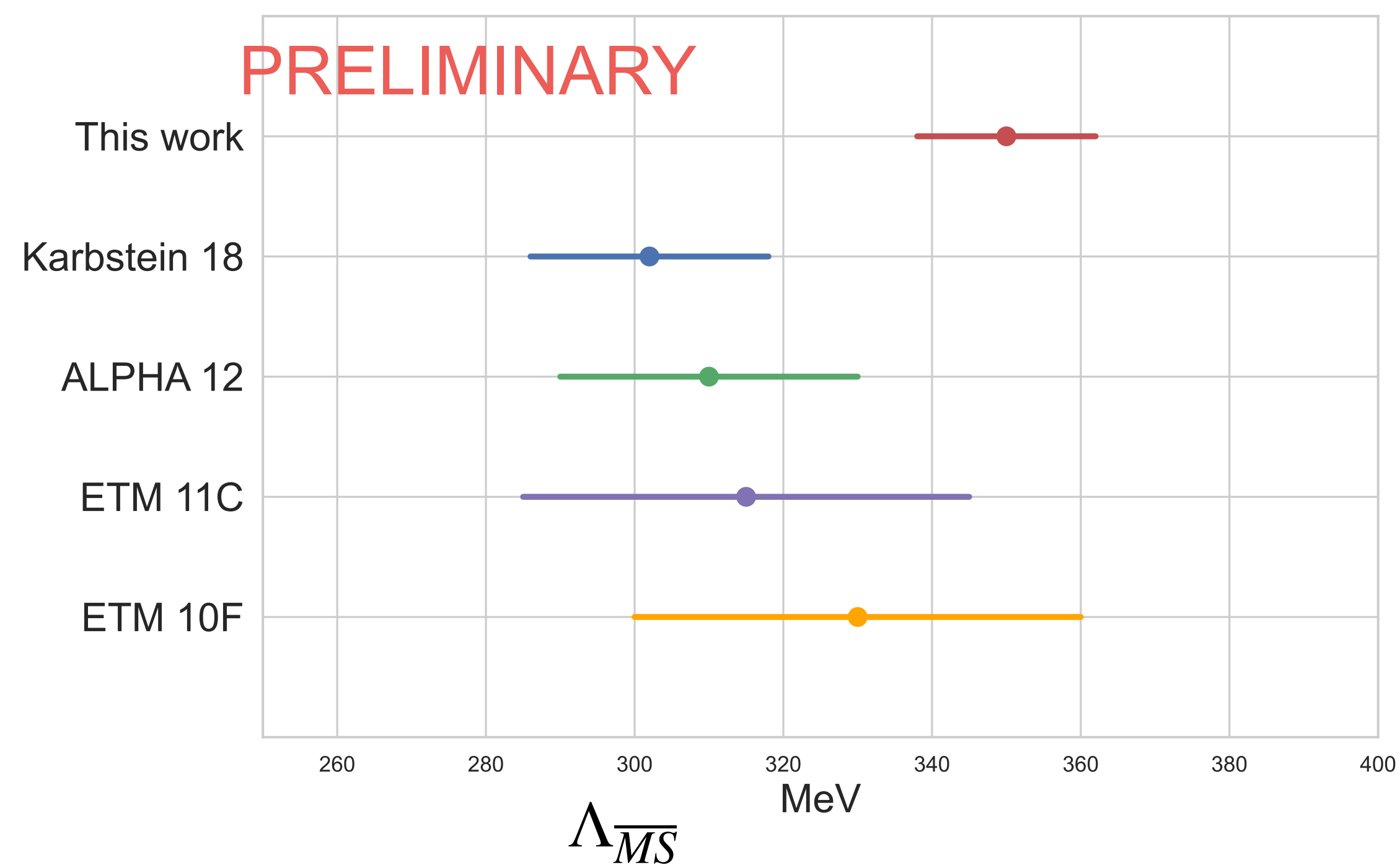
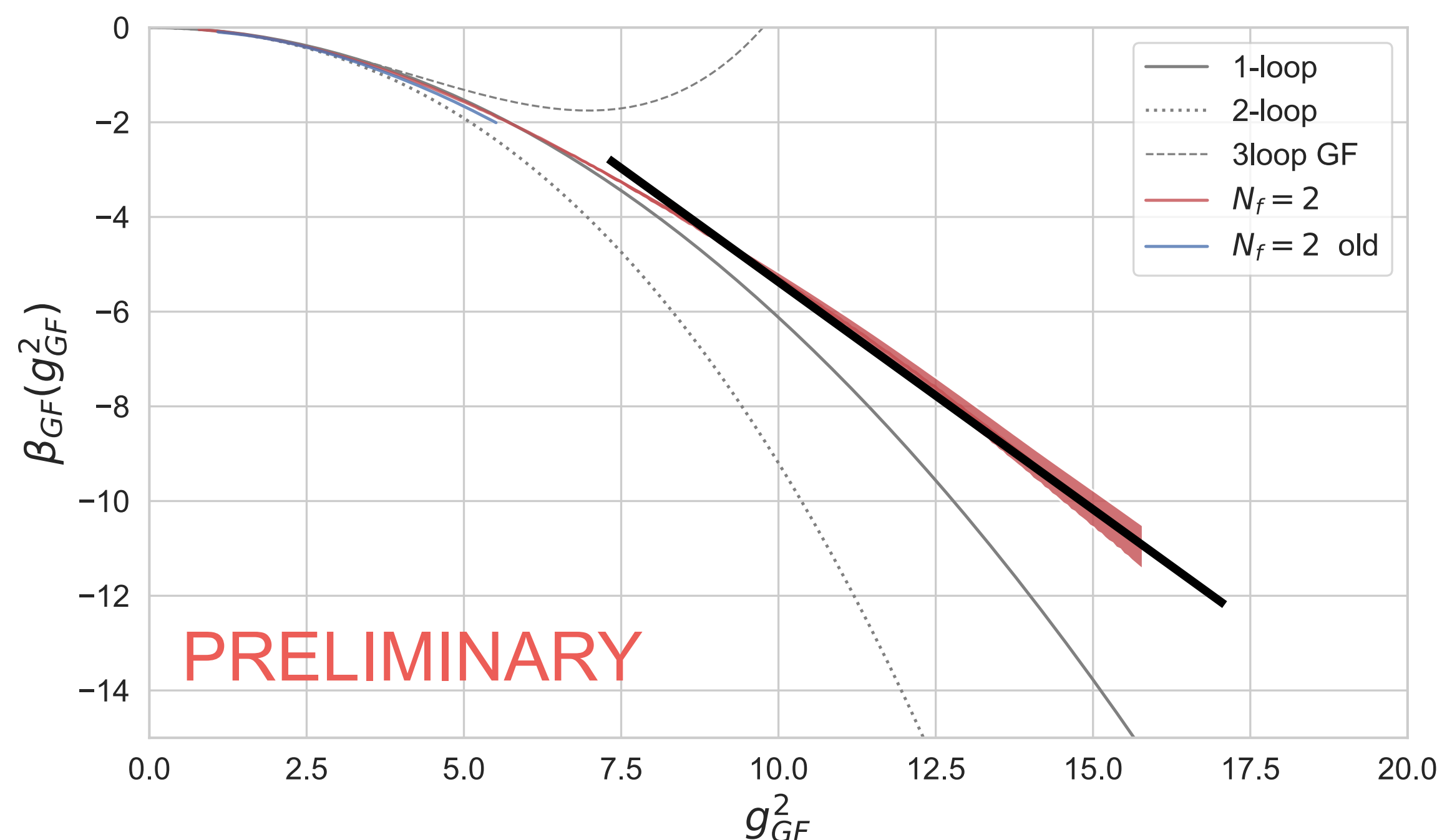
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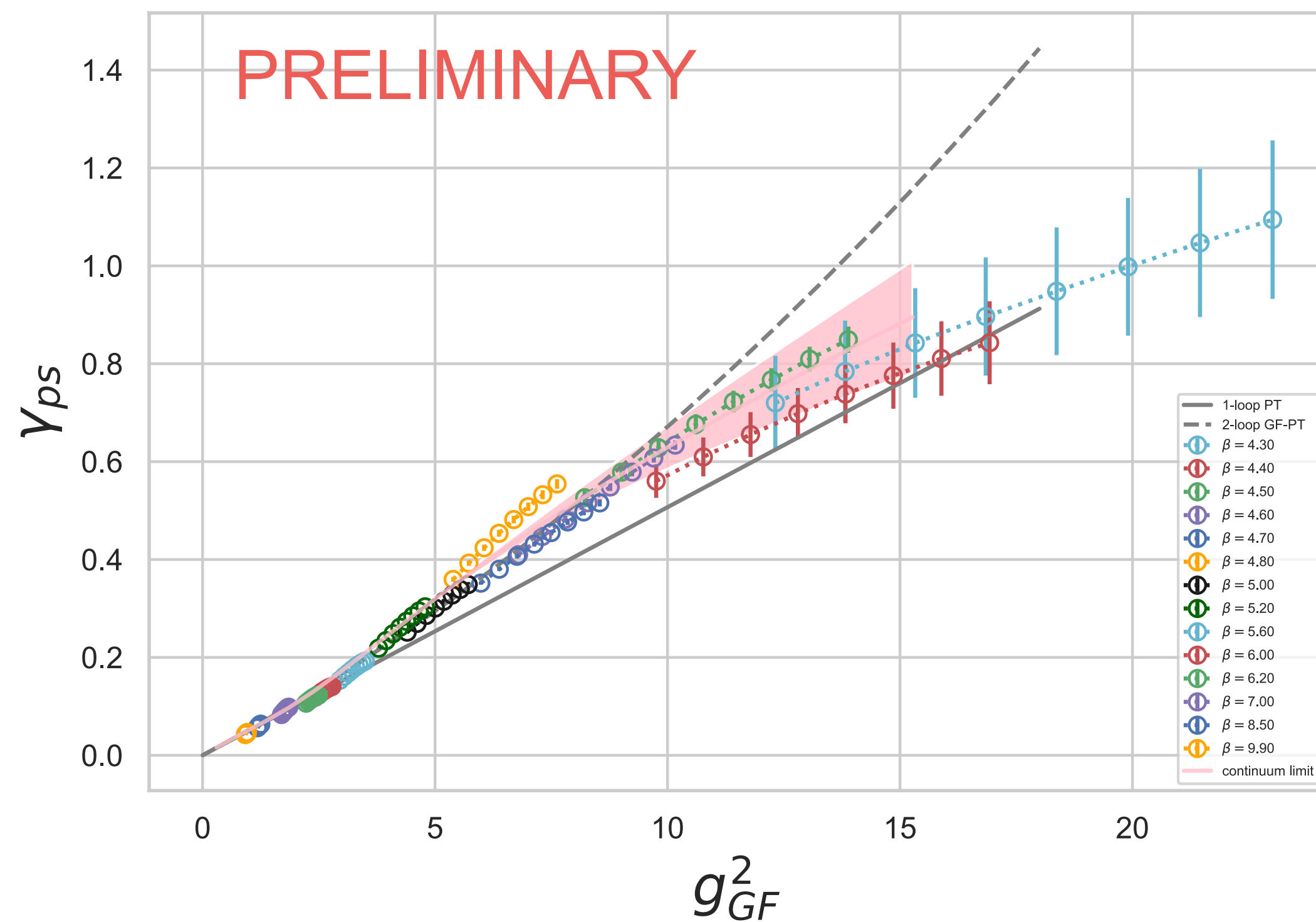
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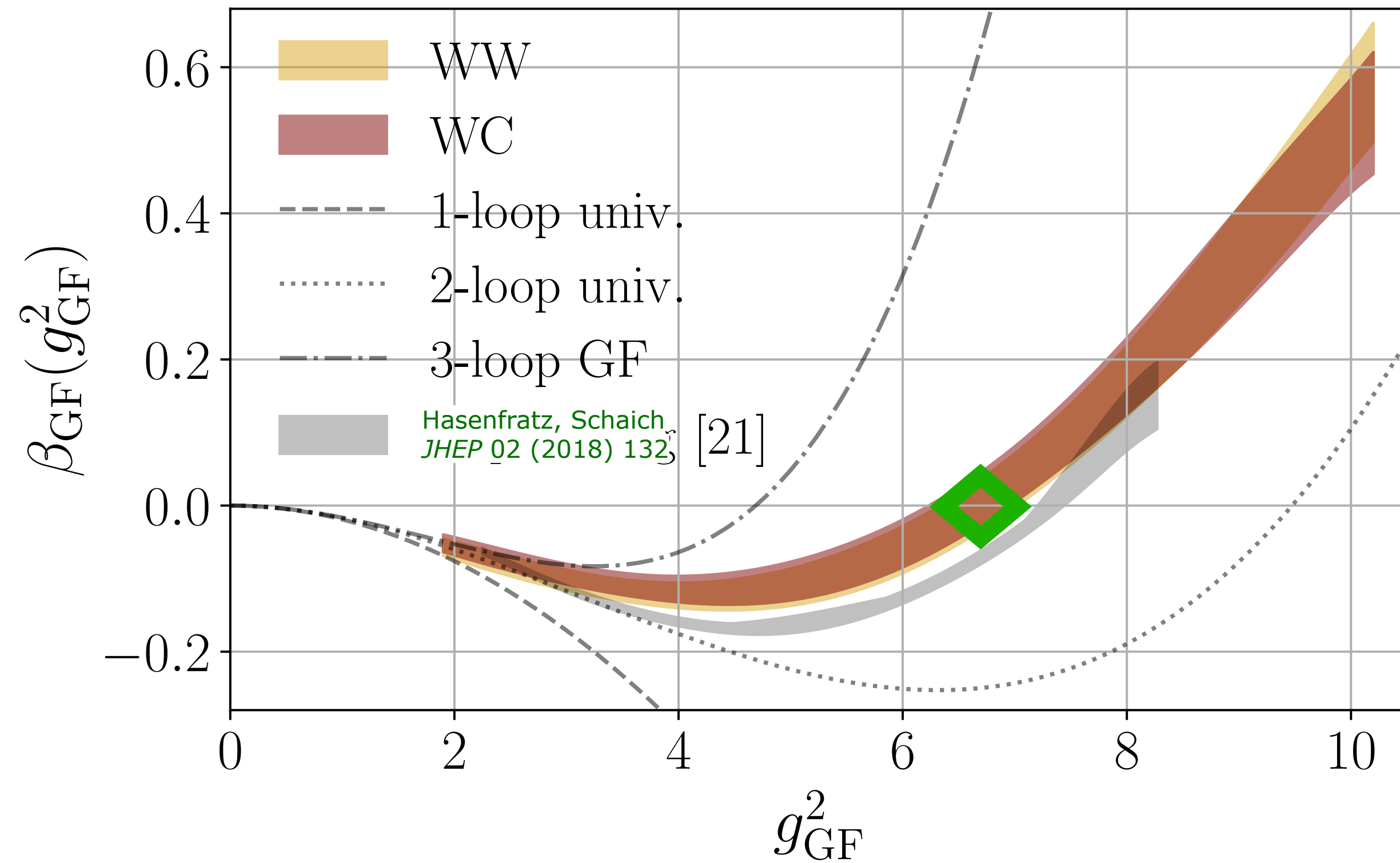
colored points show raw data:
at fixed bare coupling, changing flow time

Anomalous dimension of $\bar{\psi}\gamma_5\psi$

- small cutoff effects
- close to 1-loop

$SU(3)+N_f = 12$, staggered

Peterson, Hasenfratz, 2402.1803



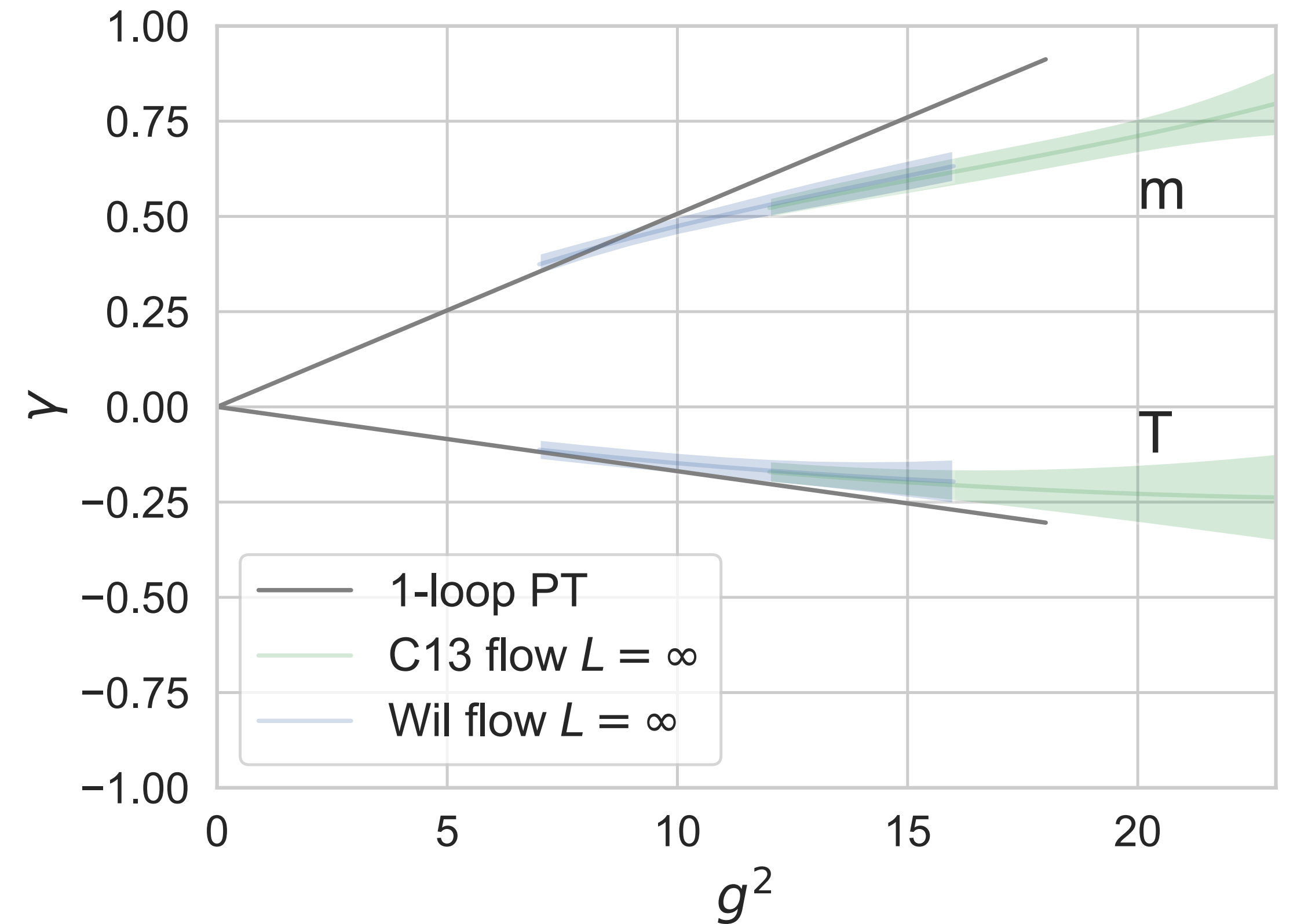
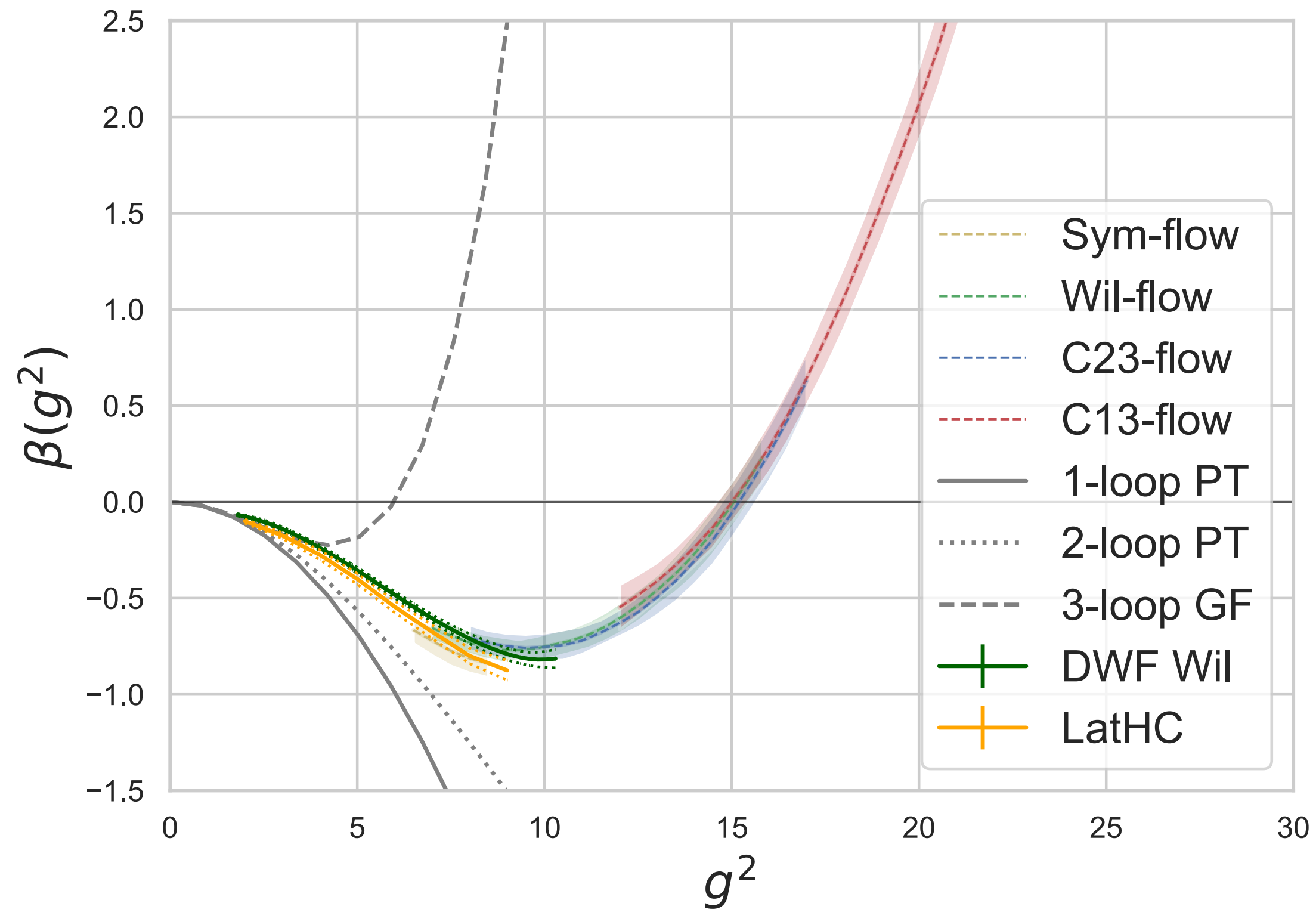
RG β function in the gradient flow scheme exhibits an IRFP

$$g_{GF}^2 \approx 6.8$$

SU(3)+ $N_f = 10$, Wilson

Hasenfratz.,Neil, Shamir, Svetitsky, Witzel,
Phys.Rev.D 108 (2023) 7

PV improved action allowed to reach $g_{GF}^2 \gtrsim 20$

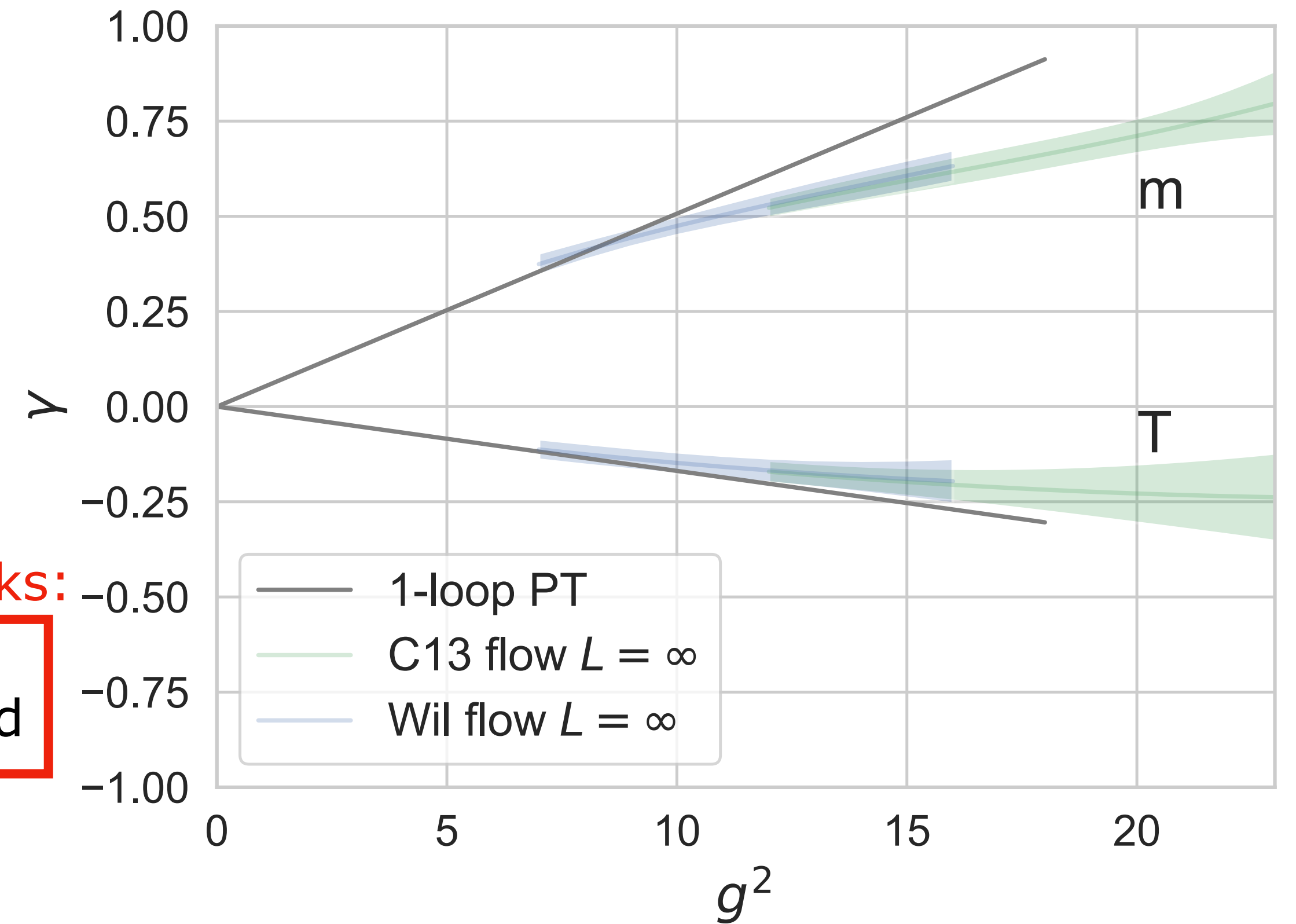
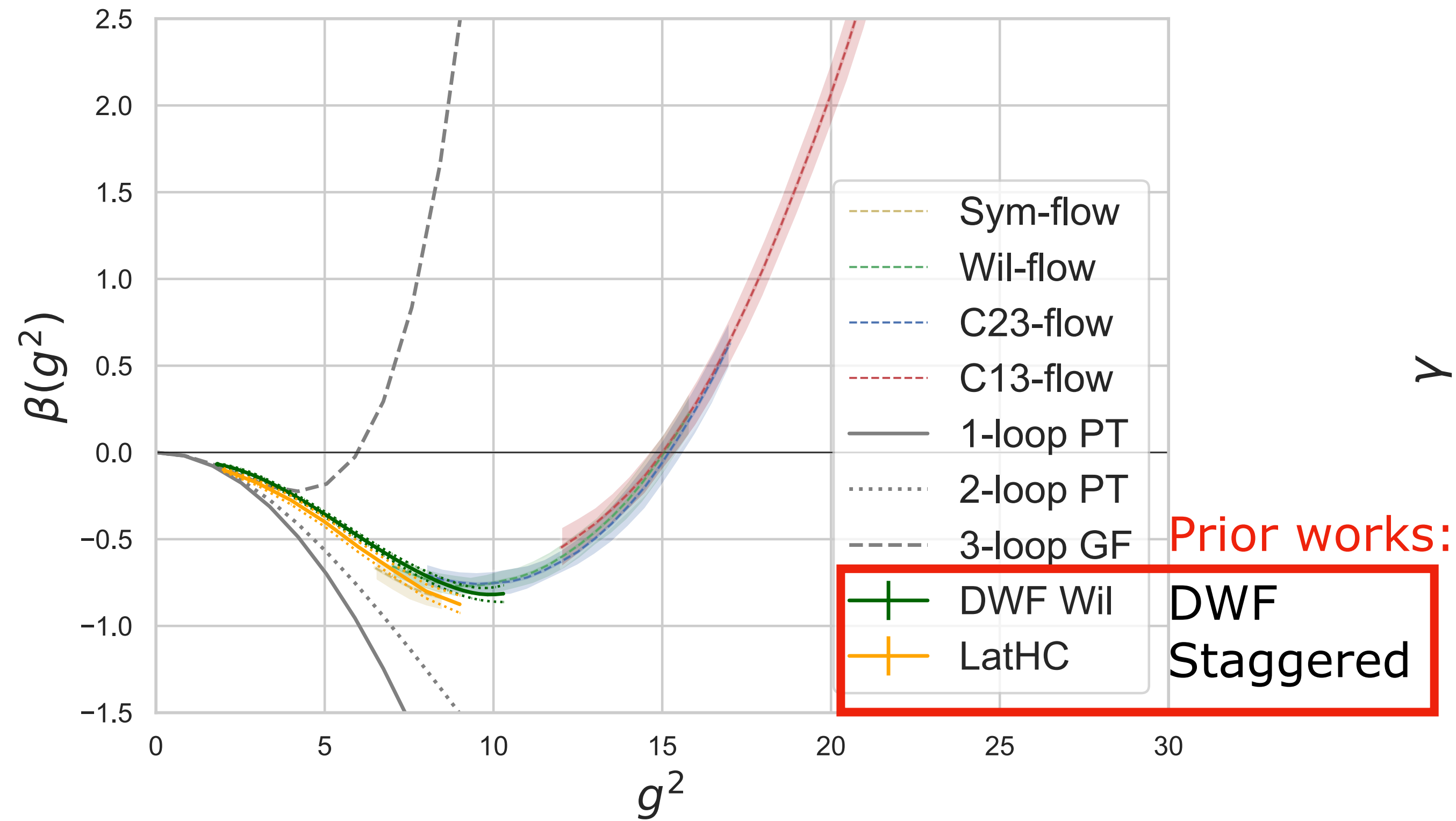


RG β function in the gradient flow scheme exhibits an IRFP at $g_{GF}^2 \approx 15$, $\gamma_m \approx 0.6$

SU(3)+ $N_f = 10$, Wilson

Hasenfratz.,Neil, Shamir, Svetitsky, Witzel,
Phys.Rev.D 108 (2023) 7

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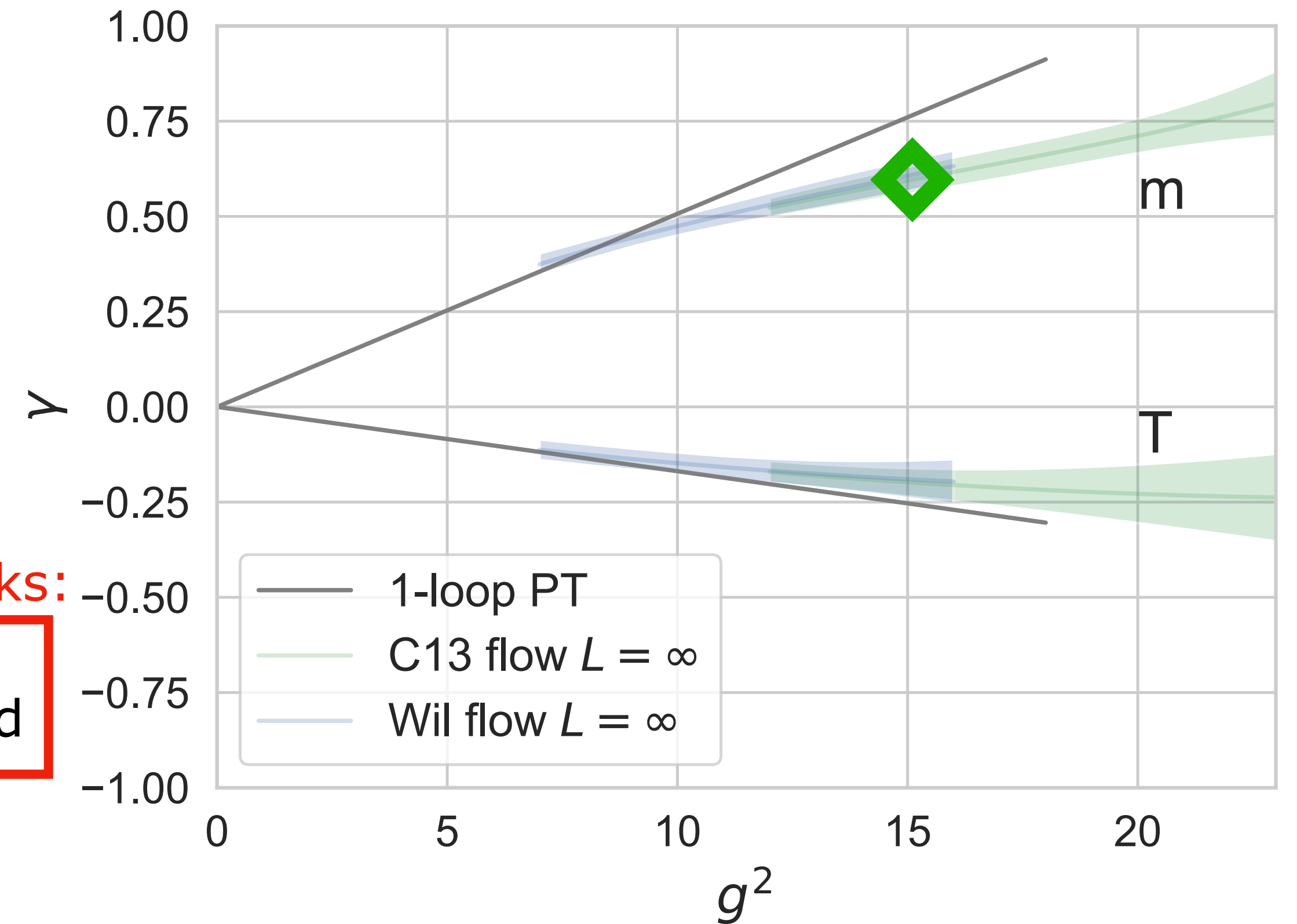
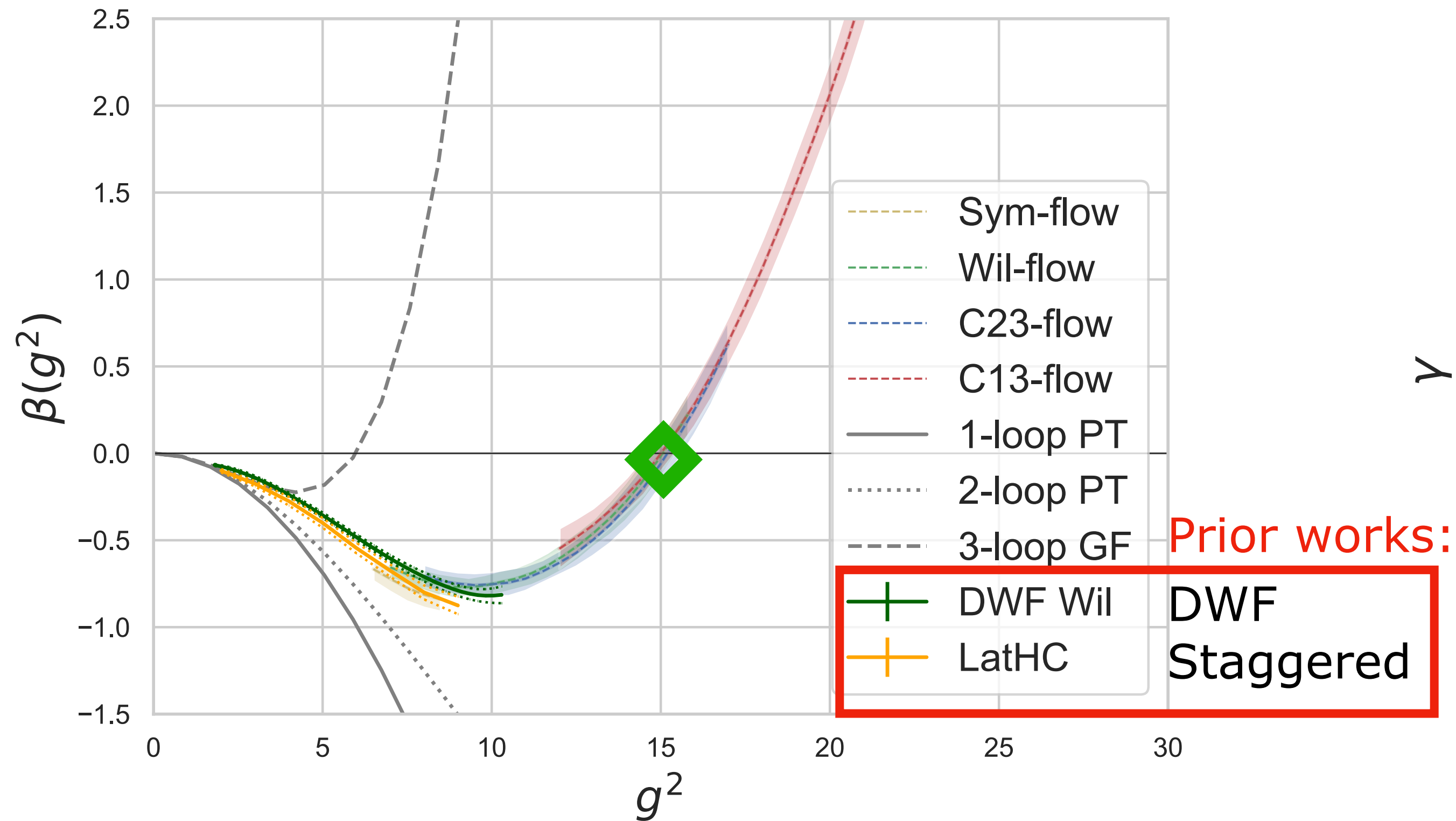


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Hasenfratz.,Neil, Shamir, Svetitsky, Witzel,
Phys.Rev.D 108 (2023) 7

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$SU(3)+N_f = 8$, staggered

“But even Wikipedia says $N_f = 8$ is chirally broken!”

While it has been long favored as composite Higgs model with slowly walking coupling and a light scalar (Higgs), this is not the case

$SU(3)+N_f = 8$, staggered

LSD Collaboration 2306.06095,
Phys.Rev.D 108 (2023) 9

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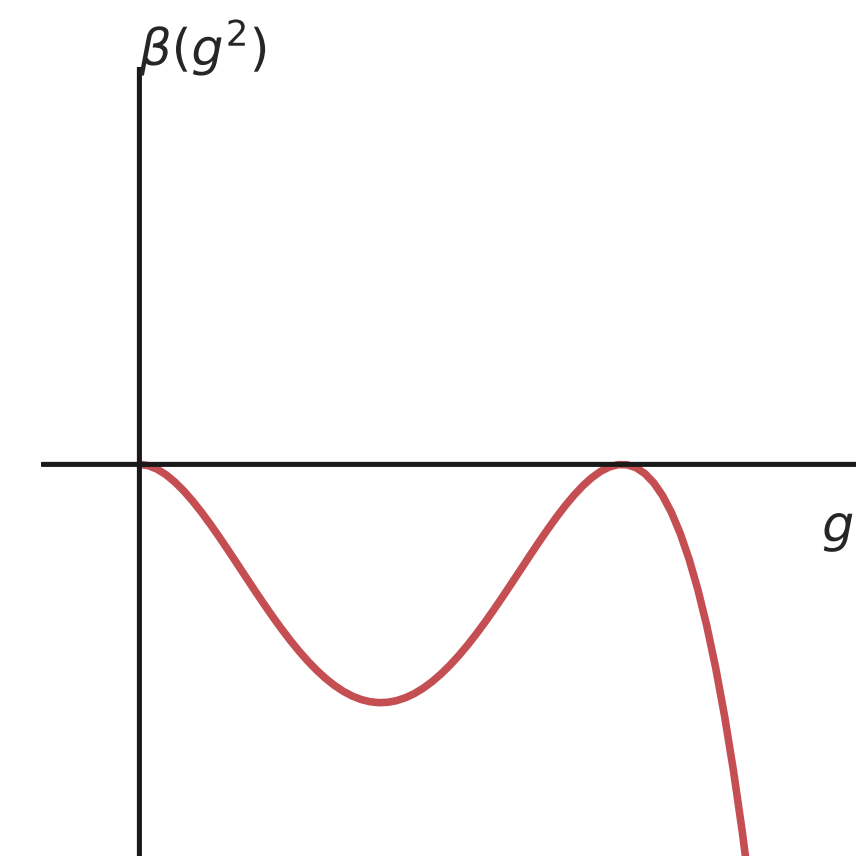
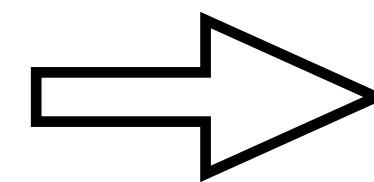
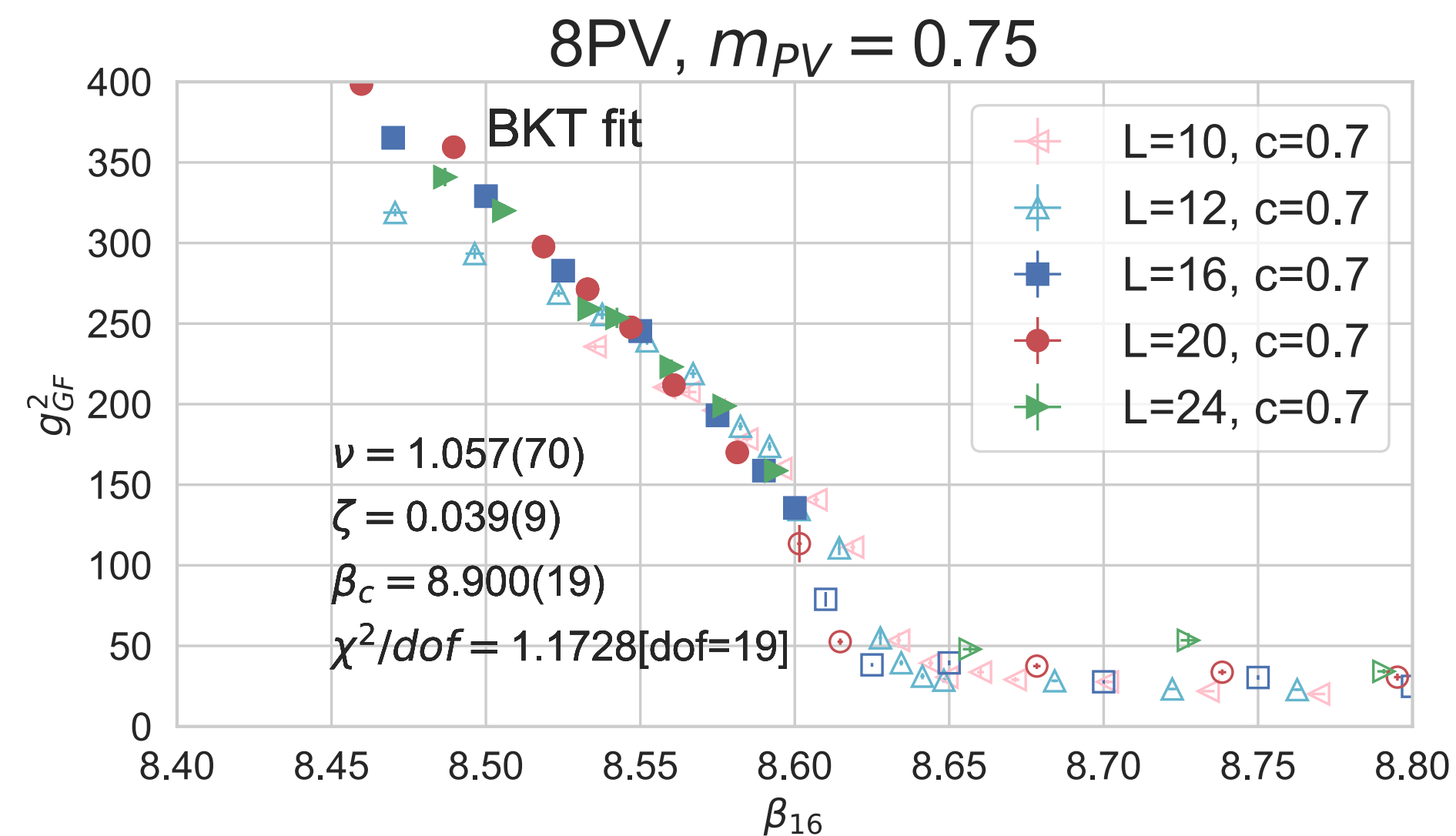
PV improved action reveals a different picture:

$SU(3)+N_f = 8$ could be the opening of the conformal window

Finite size scaling analysis using finite volume g_{GF}^2 shows

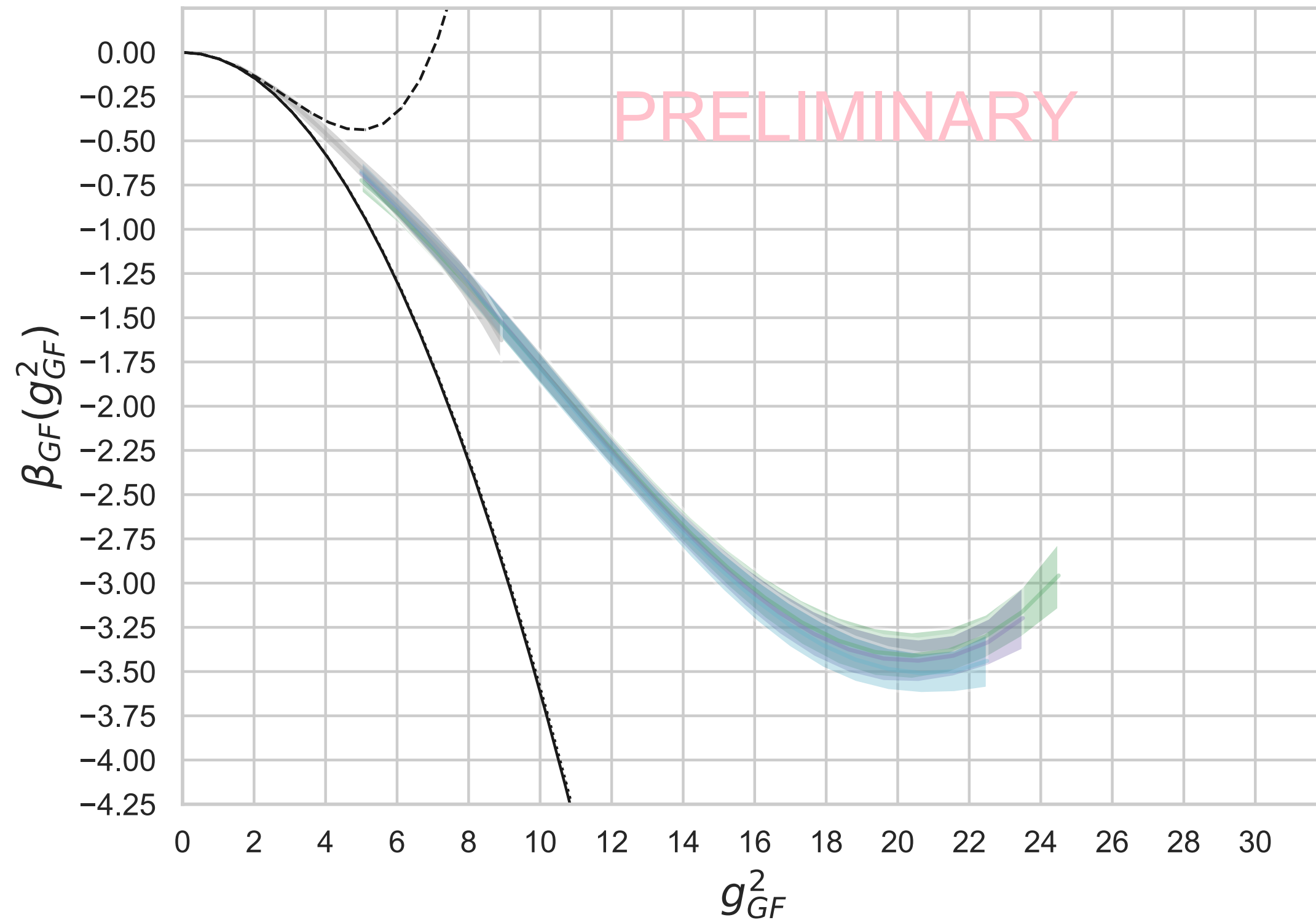
- continuous, likely BKT phase transition: $\xi \propto e^{-\zeta(\beta-\beta_c)^{-\nu}}$
- Symmetric Mass Generation phase

A.H. PRD 106 (2022) 014513
A.H, O. Witzel 2412.10322 (LSD)



$SU(3)+N_f = 8$, staggered

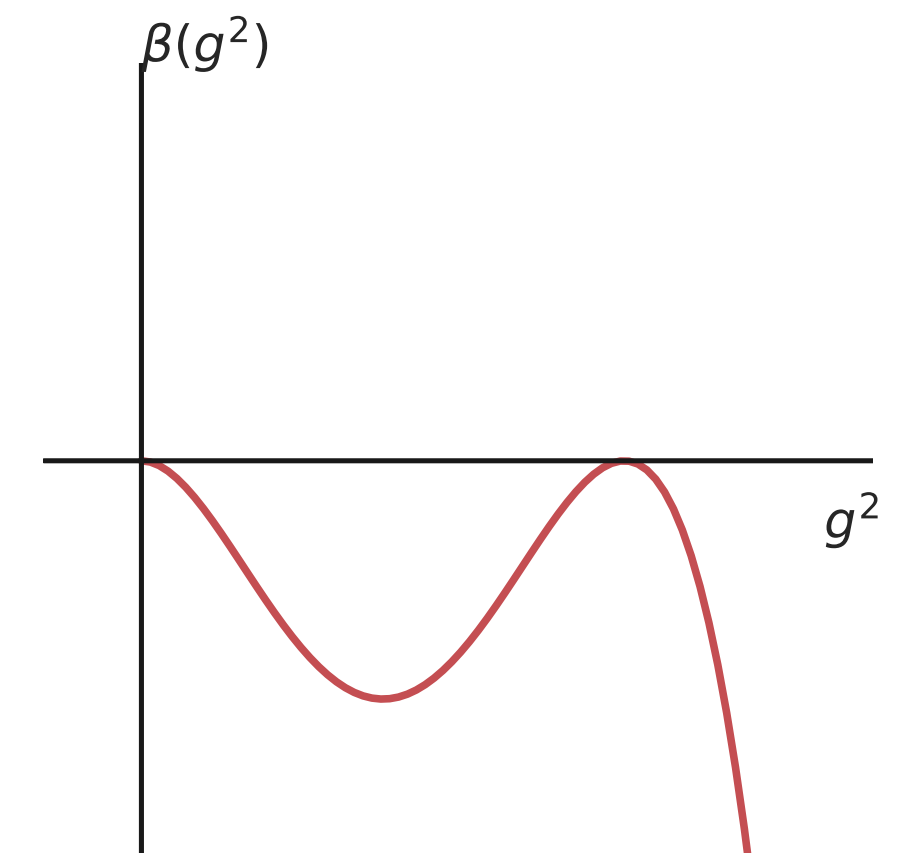
Hasenfratz, Peterson, in prep.



In practice, nonperturbative cutoff effects due to instantons make the predicted $|\beta(g^2)|$ too large

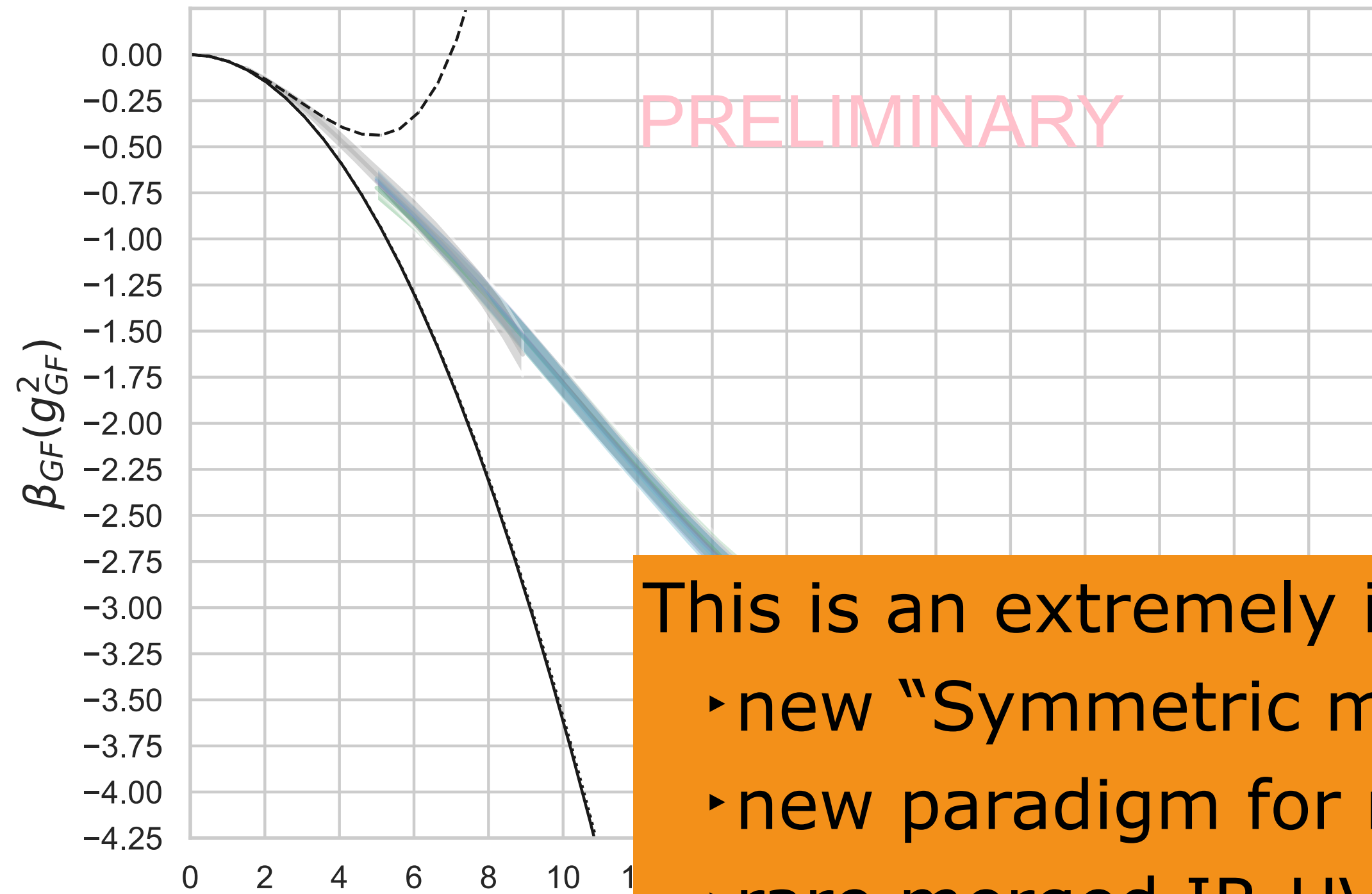
RG β function in the gradient flow scheme shows an uptick, but not a fixed point

We determine $\beta(g_{GF}^2)$ while tuning $g_0^2 \rightarrow 0$
A BKT transition would require infinite flow time



$SU(3)+N_f = 8$, staggered

Hasenfratz, Peterson, in prep.



RG β function in the gradient flow scheme shows an uptick, but not a fixed point

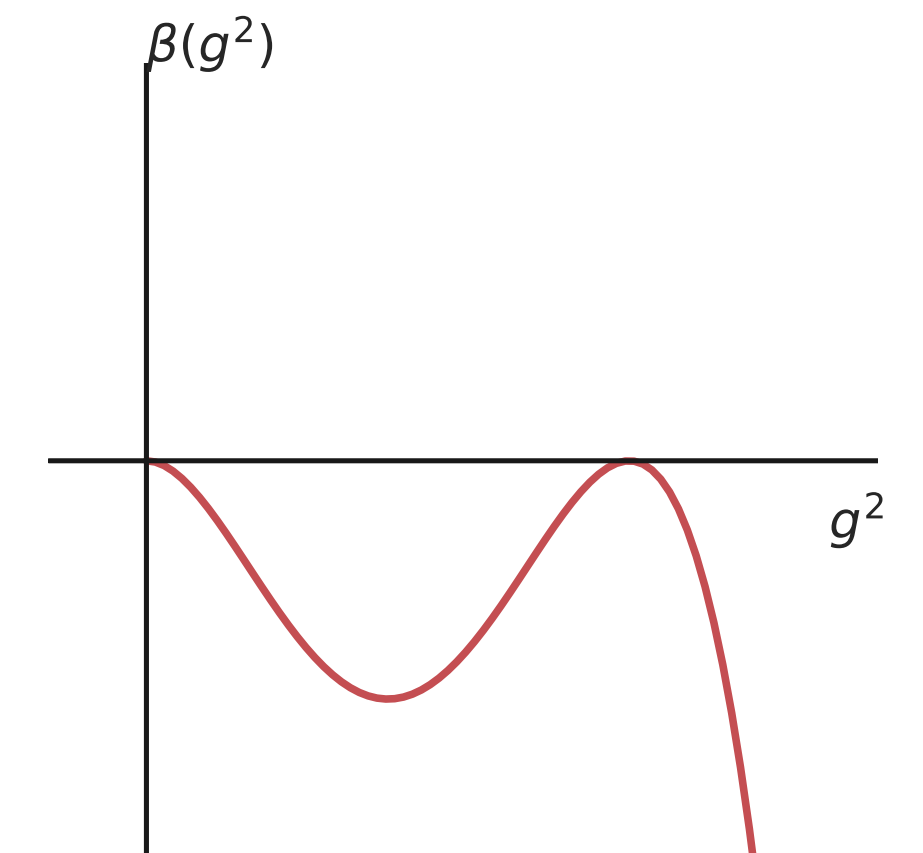
We determine $\beta(g_{GF}^2)$ while tuning $g_0^2 \rightarrow 0$

A BKT transition would require infinite flow

This is an extremely interesting system with a

- new “Symmetric mass generation” phase
- new paradigm for mass generation
- rare merged IR-UV fixed point

In practice, nonperturbative cutoff effects due to instantons make the predicted $|\beta(g^2)|$ too large



Summary and outlook

When gradient flow is interpreted as a continuous RG transformation, a plethora of non-perturbative methods open up:

- β and γ functions in the chirally broken regime: $\longrightarrow \Lambda_{QCD}$ and renormalization factors
- β and γ functions in conformal systems: \longrightarrow IRFP and universal γ^*
- g_{GF}^2 in finite volume can be used in finite size scaling \longrightarrow novel phase?

Future directions:

- Flowed observables can give new, fully non-perturbative renormalization scheme
- QCD α_{strong} could be determined with high accuracy
- Many applications in non-QCD-like systems