

APRENDE

Cross section measurements in the resolved resonance range

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E. Mendoza, CIEMAT



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Introduction

The topic of this talk is *cross section measurements in the RRR*.

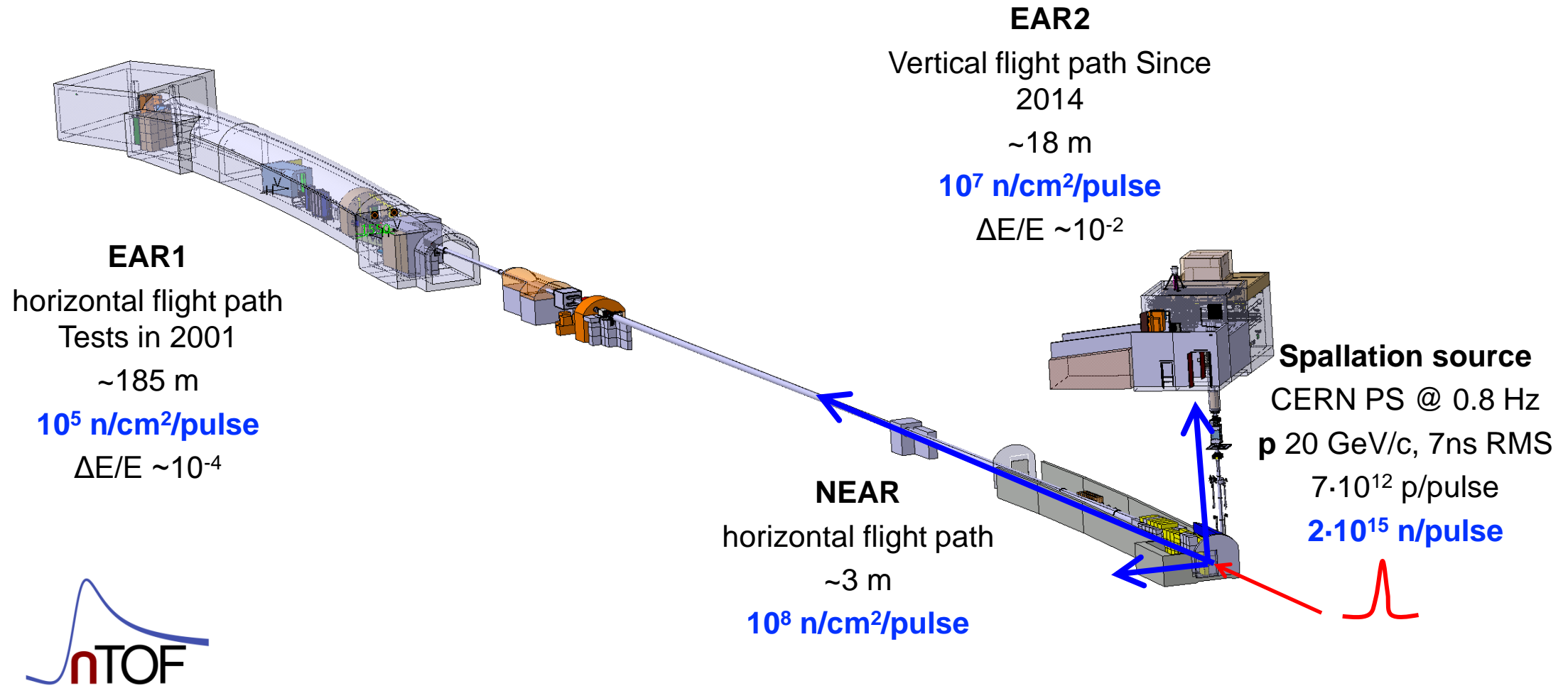
My expertise is quite focused in neutron capture (n_TOF):

- E. Mendoza et al., *Measurement and analysis of the ^{243}Am neutron capture cross section at the n_TOF facility at CERN*, Phys. Rev. C 90, 034608 (2014).
- E. Mendoza et al., *Measurement and analysis of the ^{241}Am neutron capture cross section at the n_TOF facility at CERN*, Phys. Rev. C 97, 054616 (2018).
- J. Balibrea, E. Mendoza et al., *Measurement of the alpha ratio and (n, γ) cross section of ^{235}U from 0.2 to 200 eV at n_TOF*, Phys. Rev. C 102, 044615 (2020).
- V. Alcayne, E. Mendoza et al., *Measurement of the ^{244}Cm capture cross sections at both CERN n_TOF experimental areas*, EPJ Web of Conf. 239, 01034 (2020).
- V. Alcayne, A. Kimura, E. Mendoza et al., *Measurement and analysis of the ^{246}Cm and ^{248}Cm neutron capture cross-sections at the EAR2 of the n_TOF facility at CERN*, Eur. Phys. Jour. 60, 246 (2024).
- A. Sánchez-Caballero, ..., E. Mendoza et al., *Experimental setup of the ^{239}Pu neutron capture and fission cross-section measurements at n_TOF, CERN*, EPJ Web of Conf. 294, 01003 (2024).

In this talk I will try to make an overview of the experimental procedures performed in a TOF experiment, focusing in neutron capture measurements (at n_TOF). I will also try to present possible biases in the analysis procedure.



The n_TOF facility at CERN



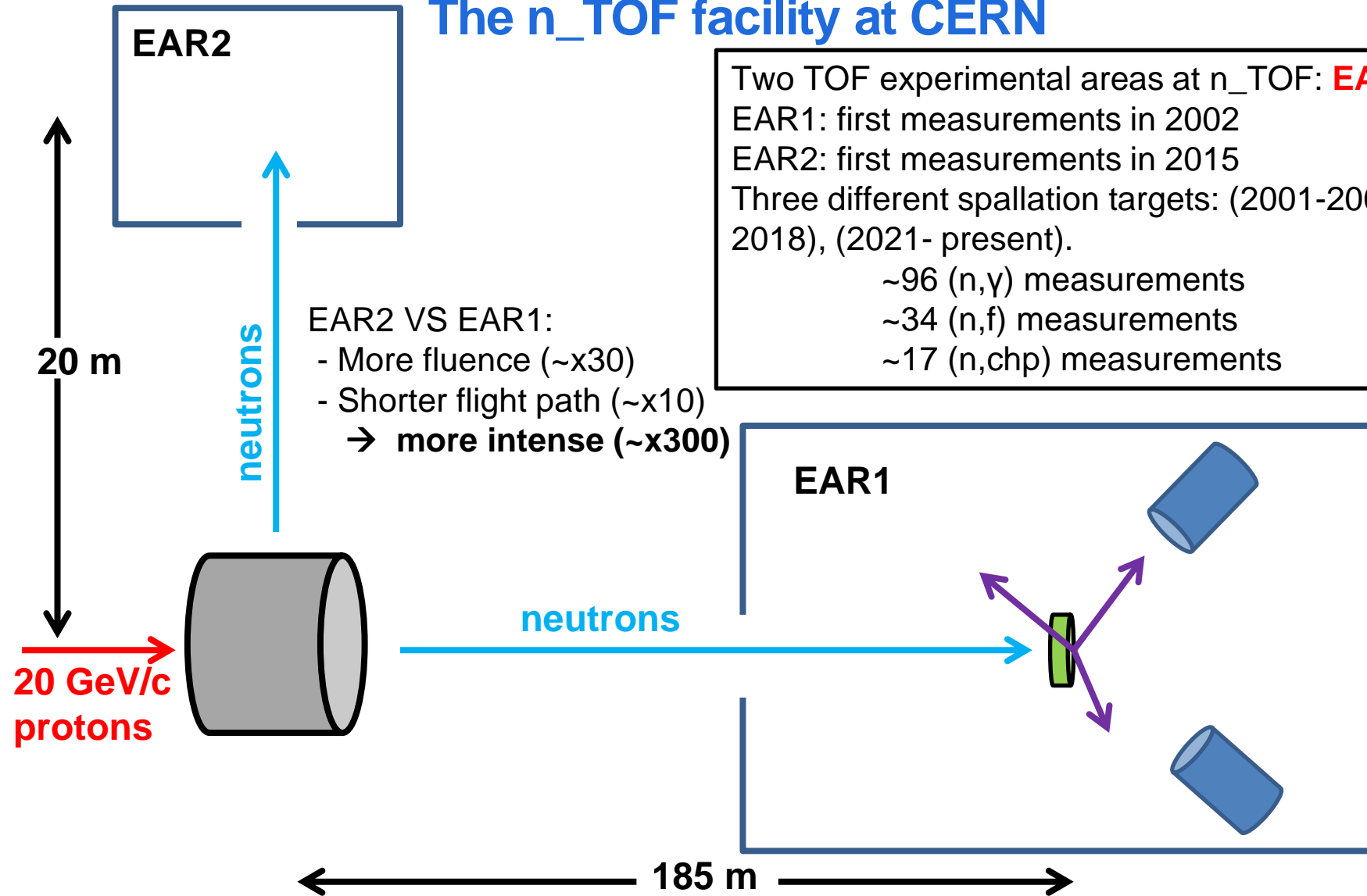
N. Patronis et al. EPJ Techniques and Instrumentation 10, 13 (2023)



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The n_TOF facility at CERN



Two TOF experimental areas at n_TOF: **EAR1** and **EAR2**
EAR1: first measurements in 2002
EAR2: first measurements in 2015
Three different spallation targets: (2001-2004), (2009-2018), (2021- present).
~96 (n, γ) measurements
~34 (n,f) measurements
~17 (n,chp) measurements

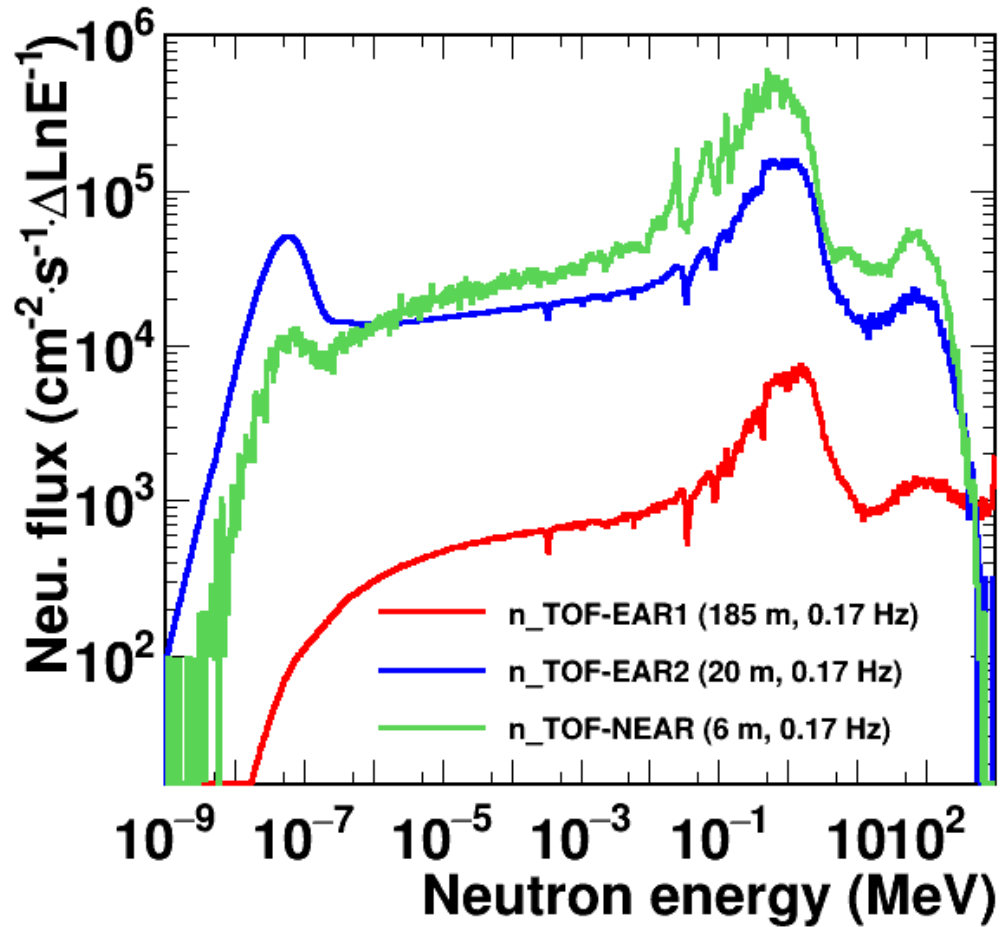
EAR2 VS EAR1:
- More fluence (~x30)
- Shorter flight path (~x10)
→ more intense (~x300)



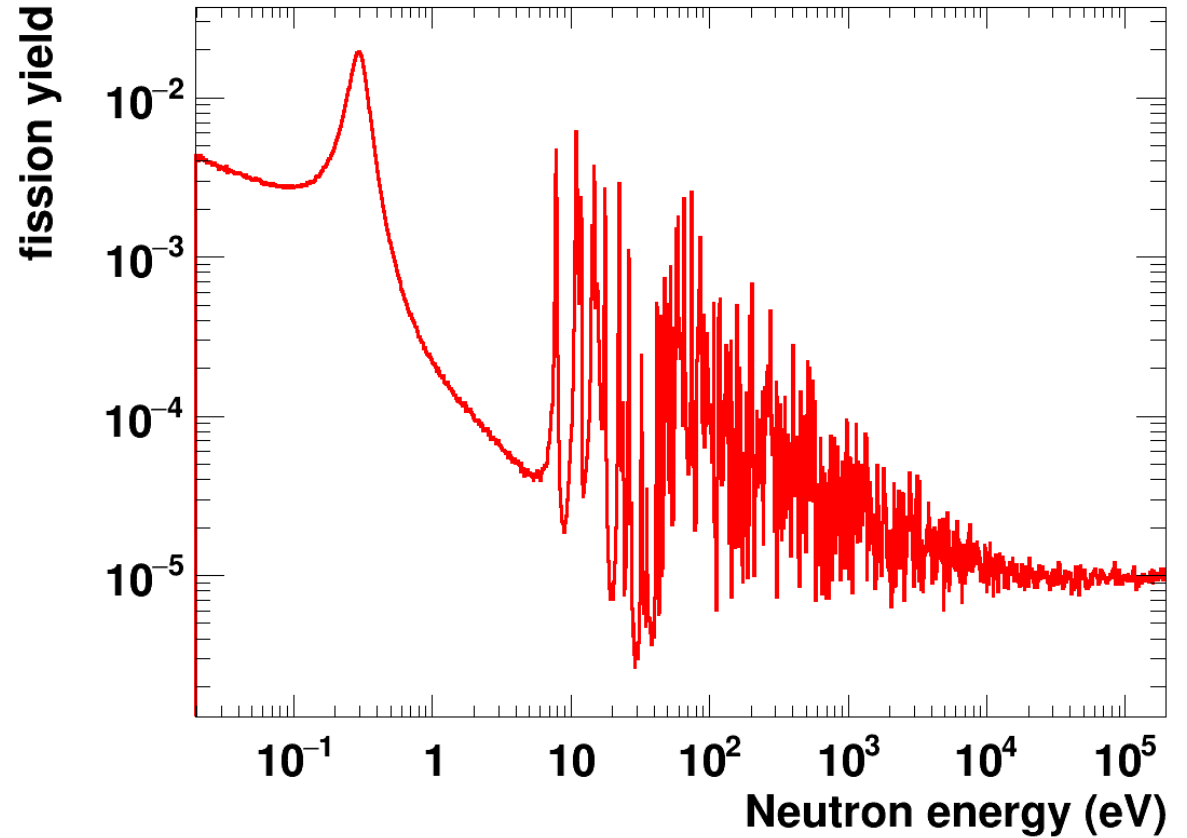
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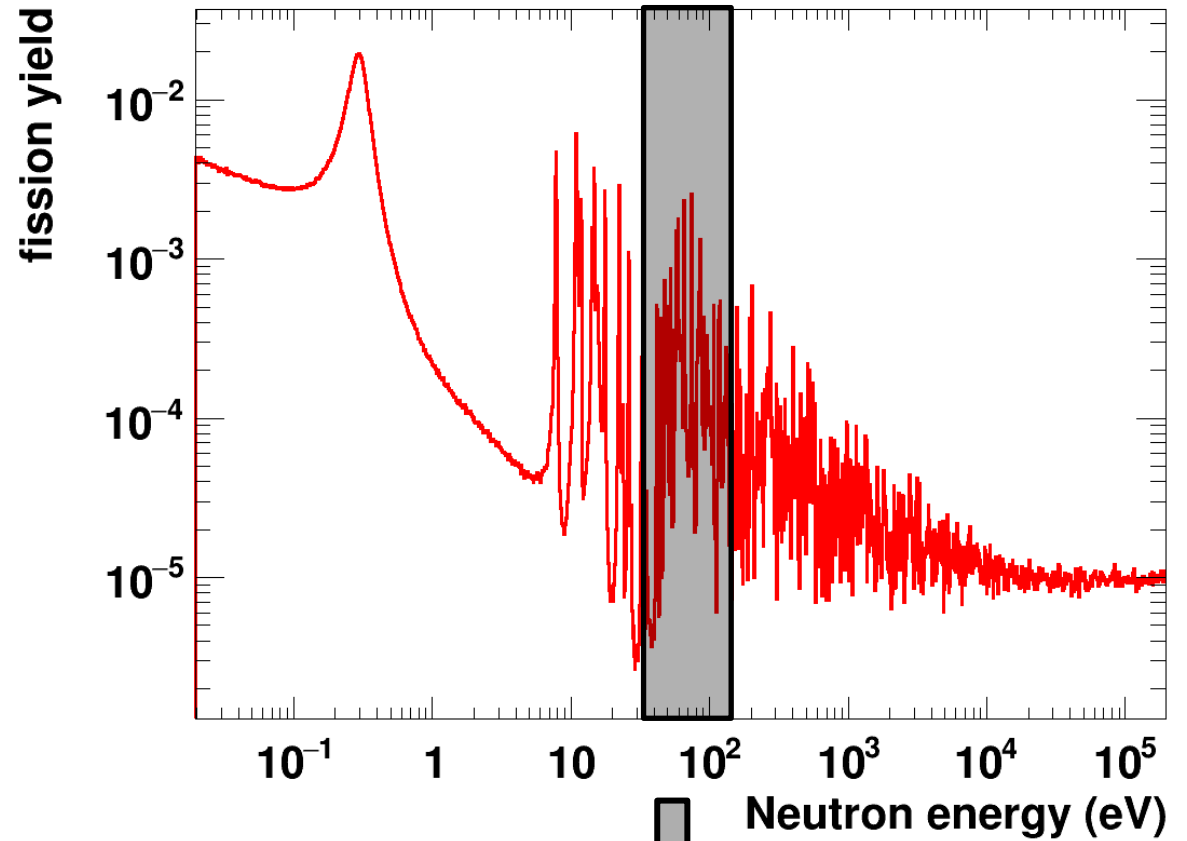
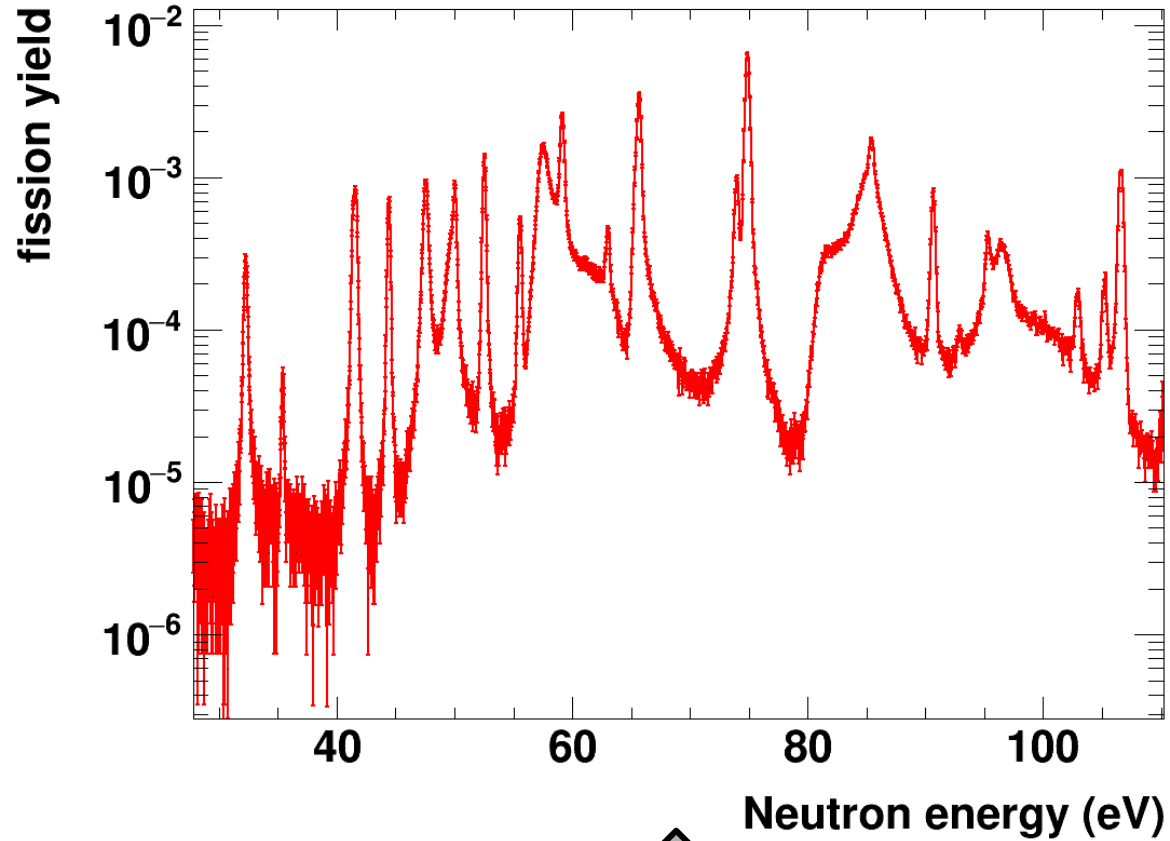


$^{239}\text{Pu}(n,f)$ measurement @ n_TOF-EAR1 (2022)



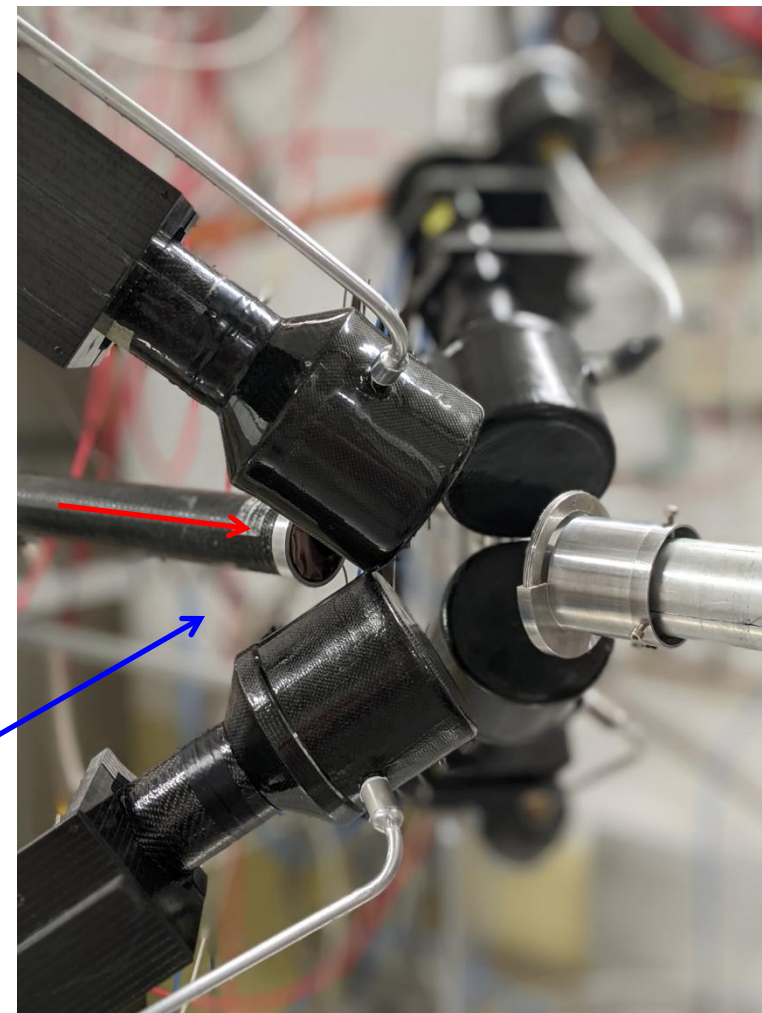
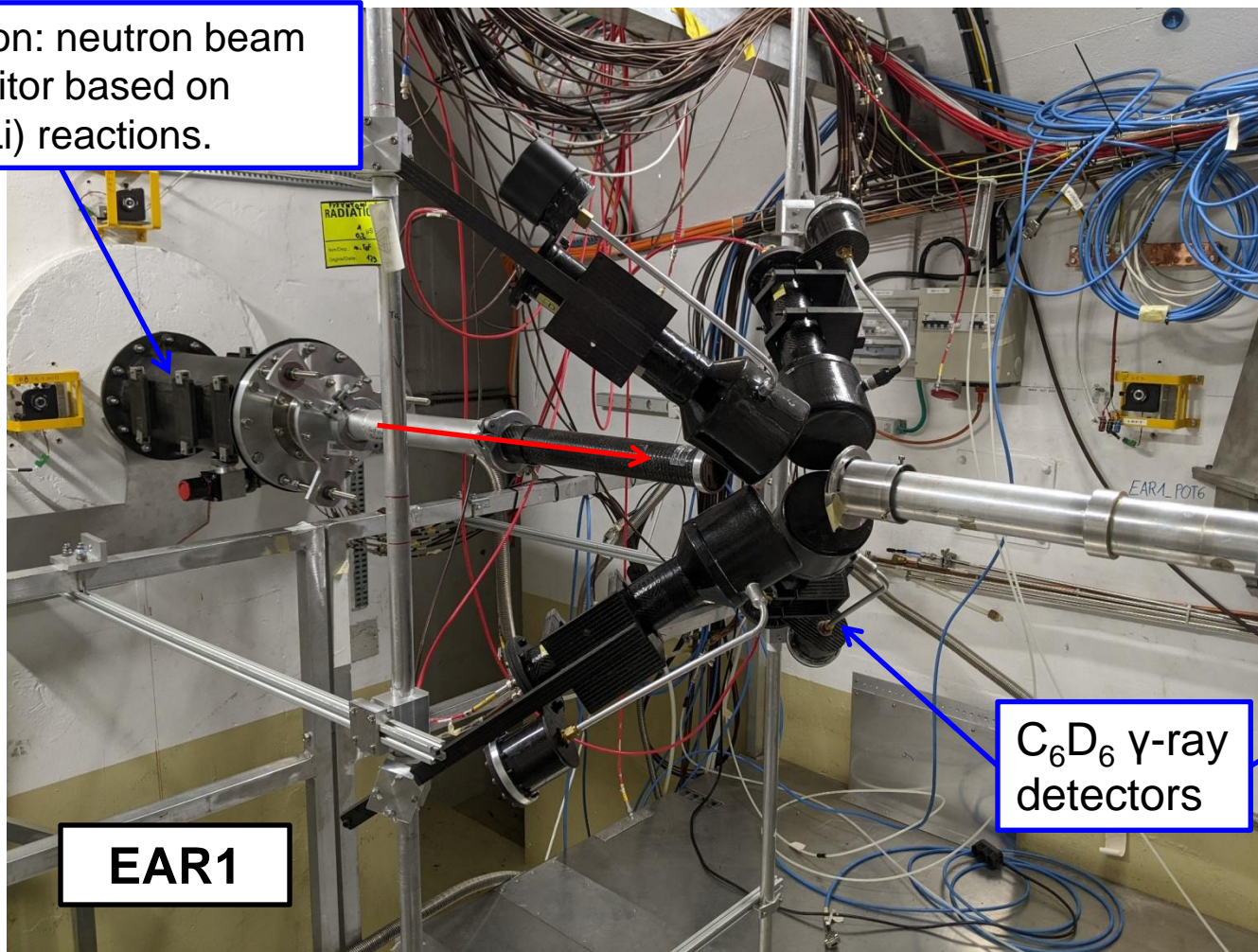
The n_TOF facility at CERN

$^{239}\text{Pu}(n,f)$ measurement @ n_TOF-EAR1 (2022)



Detection systems

SiMon: neutron beam monitor based on $(n, {}^6\text{Li})$ reactions.

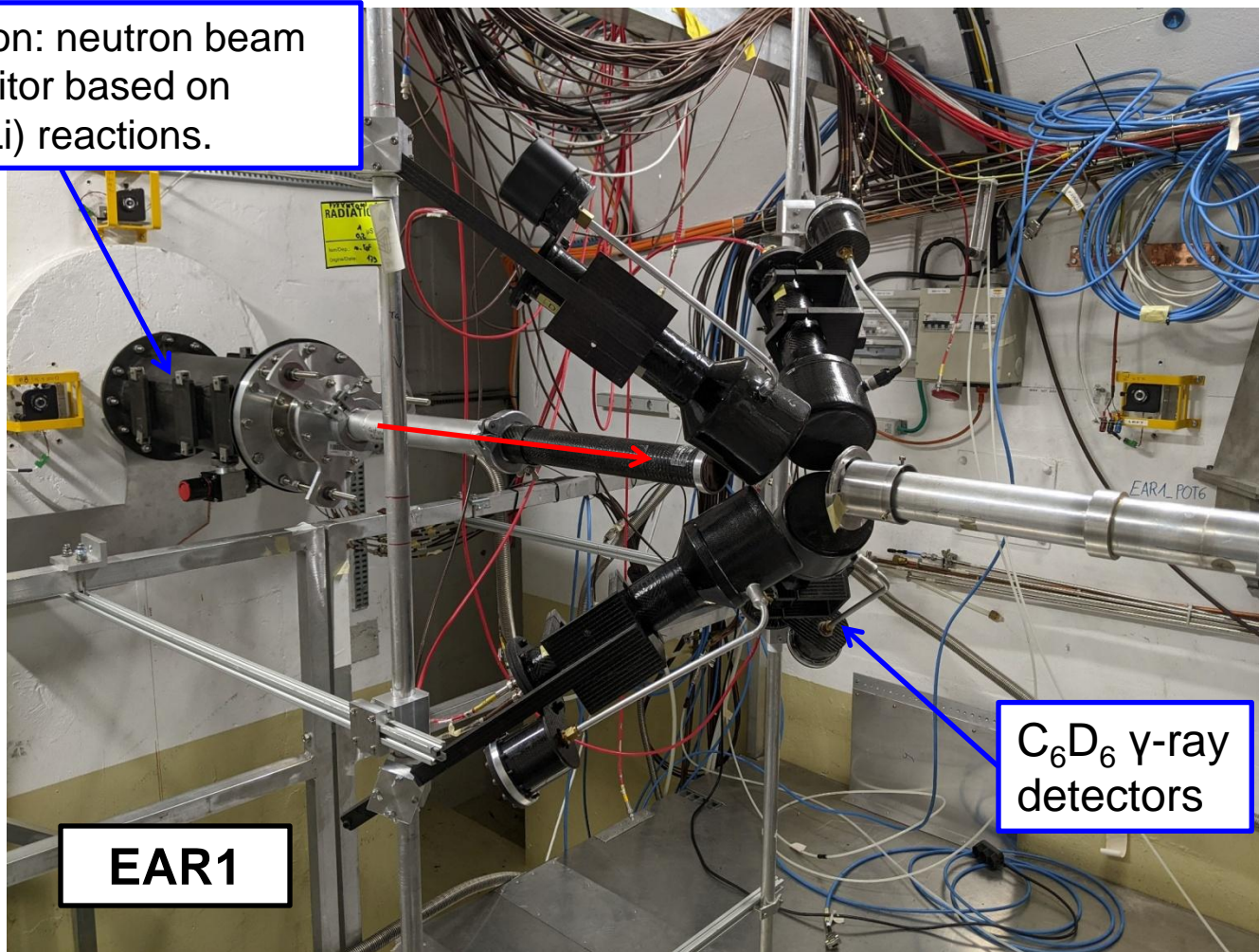


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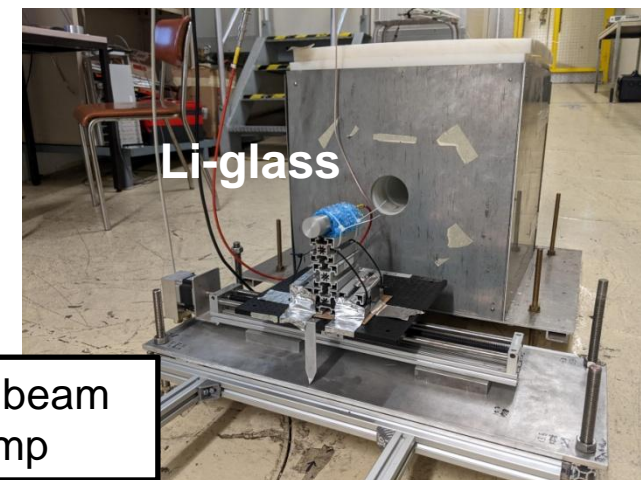
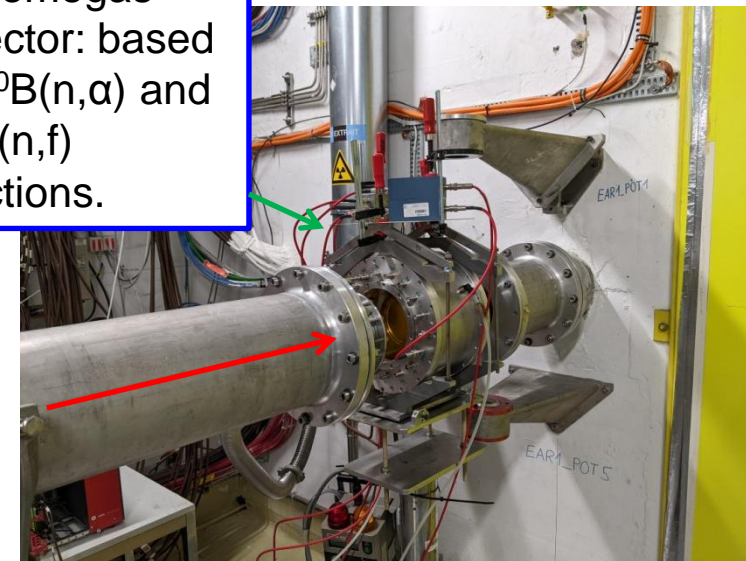
SiMon: neutron beam monitor based on $(n, {}^6\text{Li})$ reactions.



C_6D_6 γ -ray detectors

EAR1

Micromegas detector: based on ${}^{10}\text{B}(n, \alpha)$ and ${}^{235}\text{U}(n, f)$ reactions.



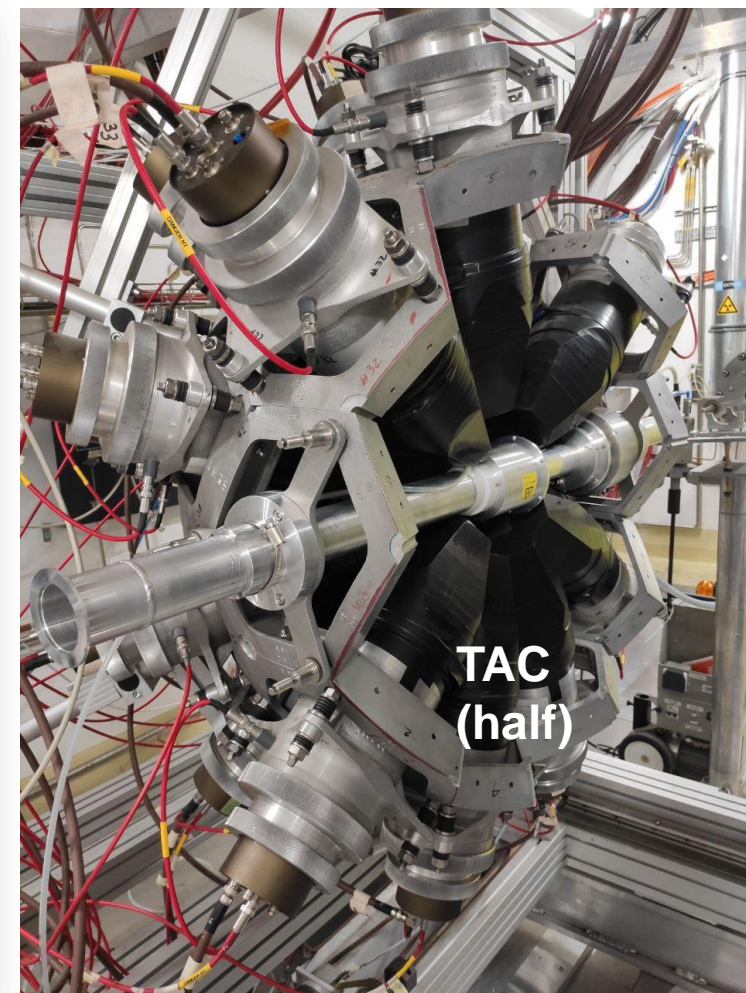
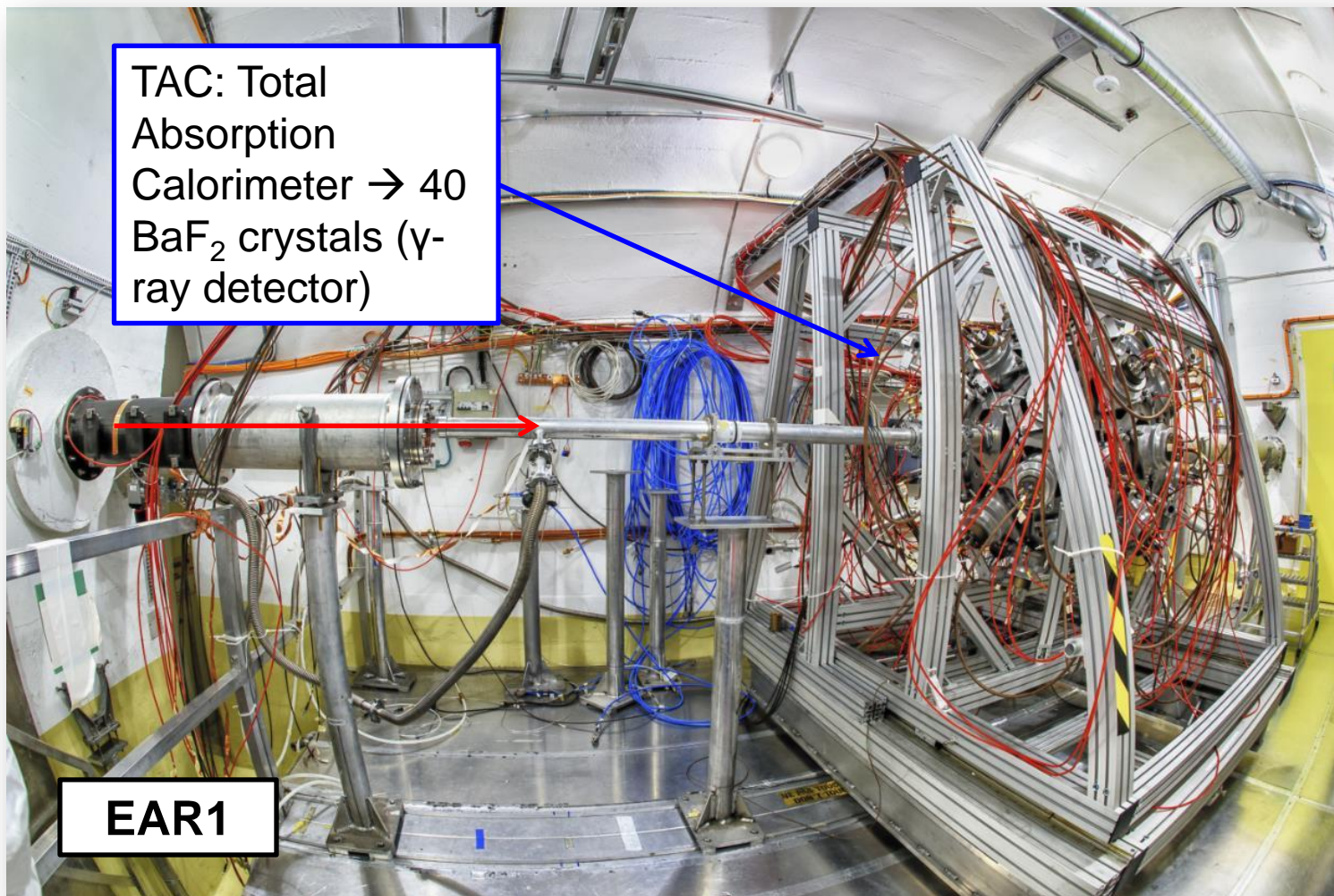
EAR1 beam dump



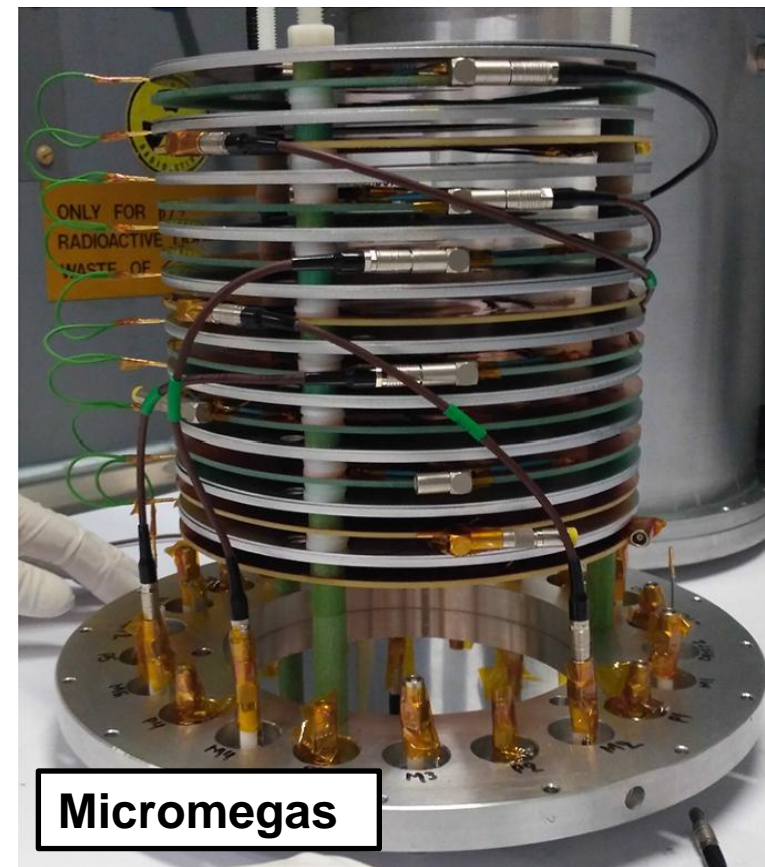
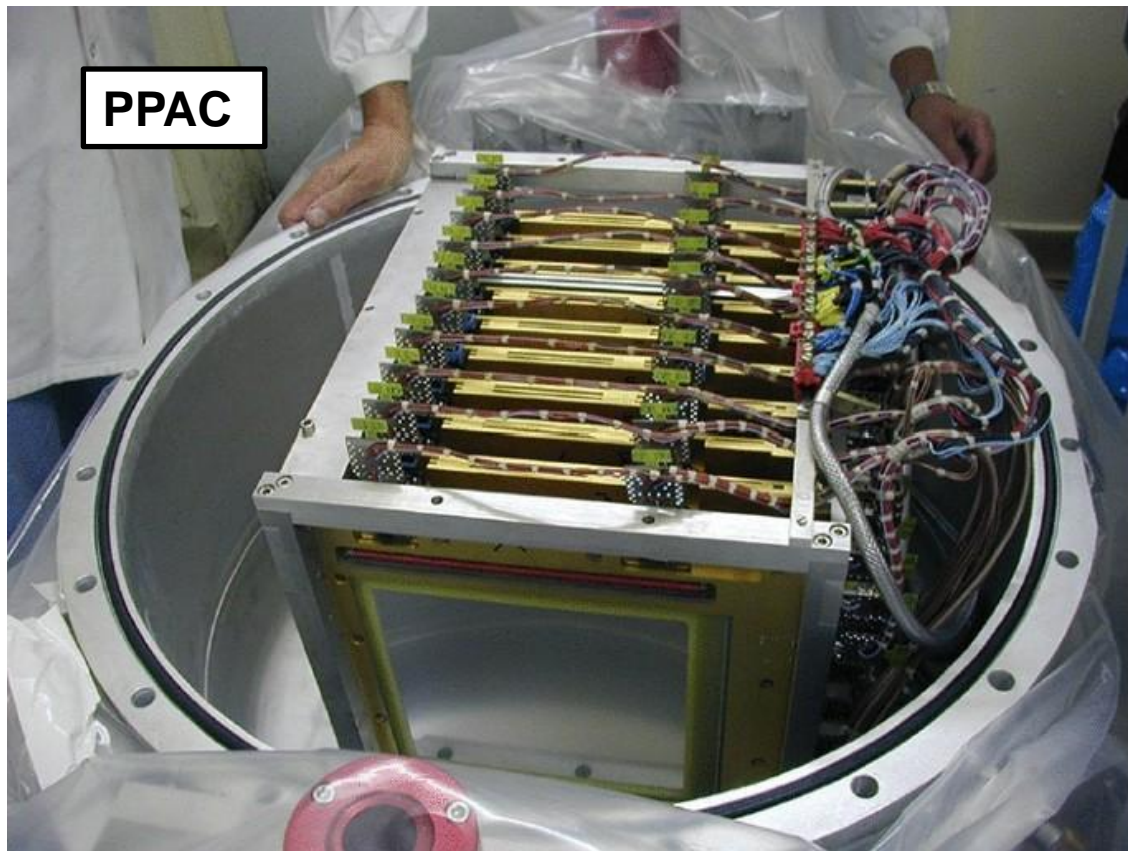
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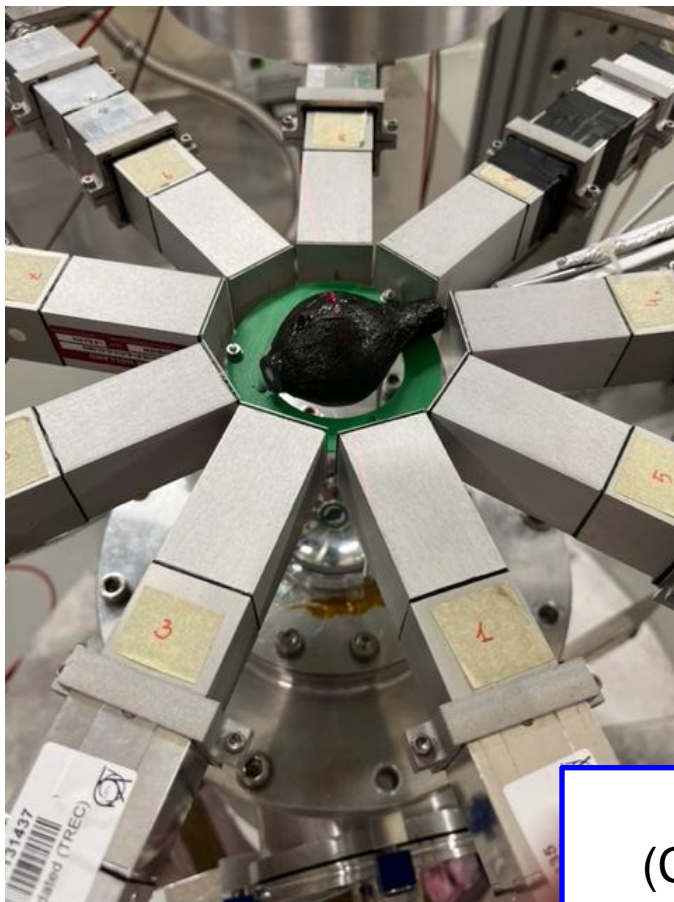
Detection systems



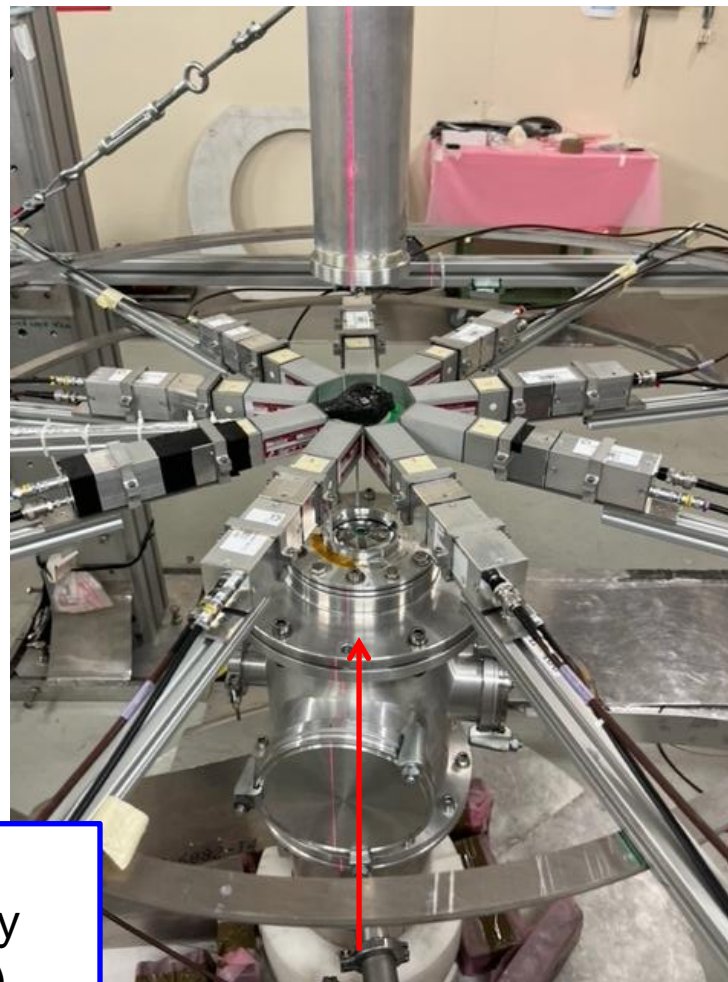
Detection systems



Detection systems



EAR2



sTED
(C_6D_6 γ -ray
detectors)



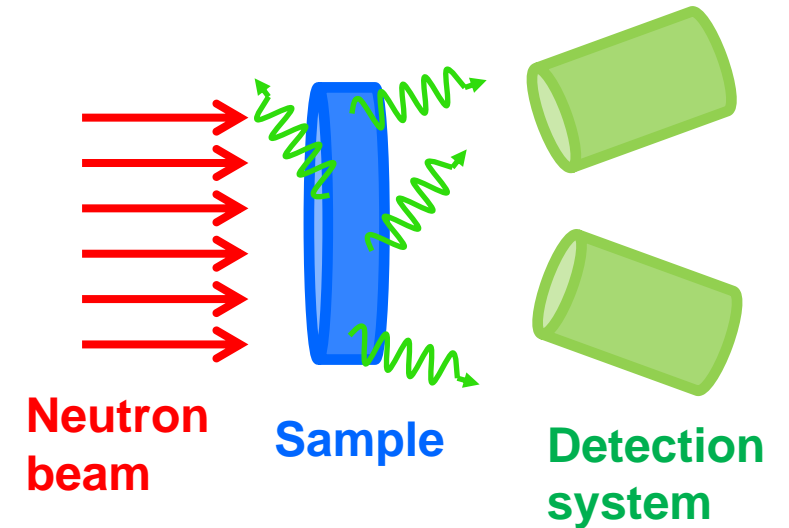
Neutron capture measurements

The experimental capture yield is usually calculated from an expression similar to this:

$$Y_{(n,\gamma)}(E_n) = \frac{C(E_n) - B(E_n)}{\varepsilon(E_n) \cdot \phi(E_n)}$$

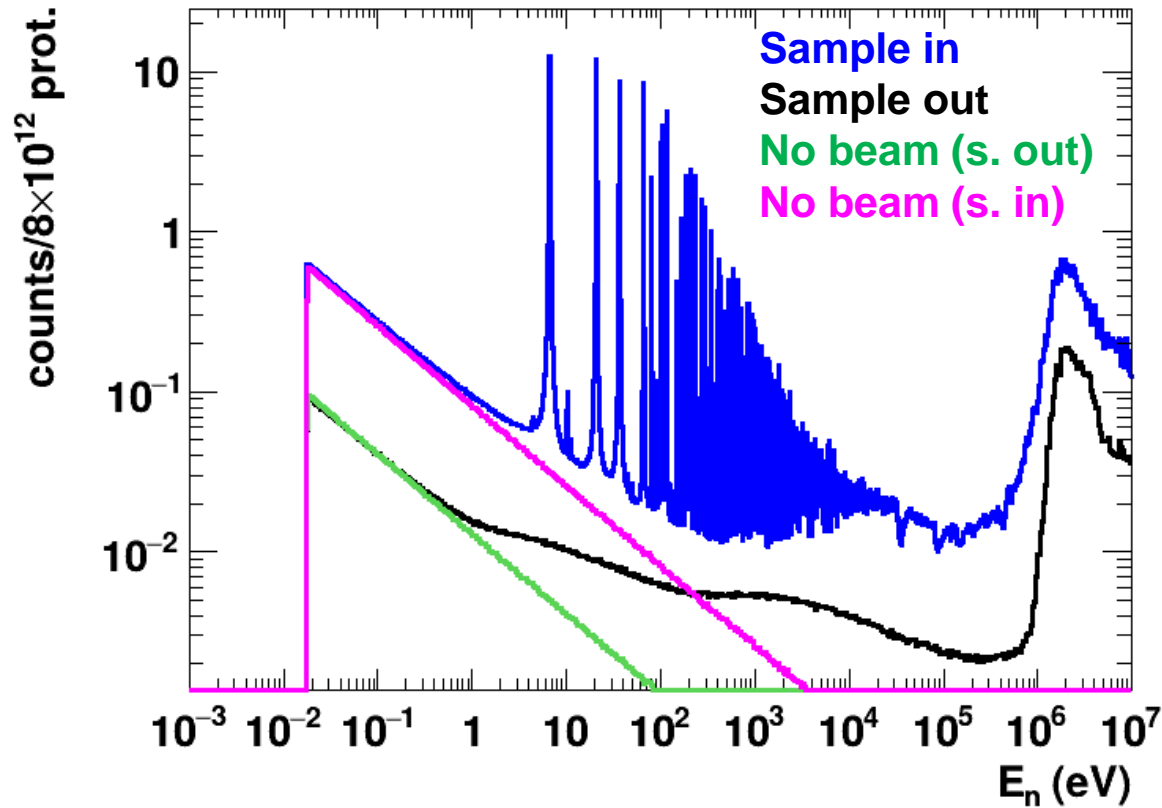
where:

- $Y_{(n,\gamma)}$ is the capture yield.
- E_n is the neutron energy.
- $C(E_n)$ are the number of counts in the detector [per unit something].
- $B(E_n)$ is the background.
- $\varepsilon(E_n)$ is the detection efficiency.
- $\phi(E_n)$ are the number of neutrons impinging in the sample [per unit something].

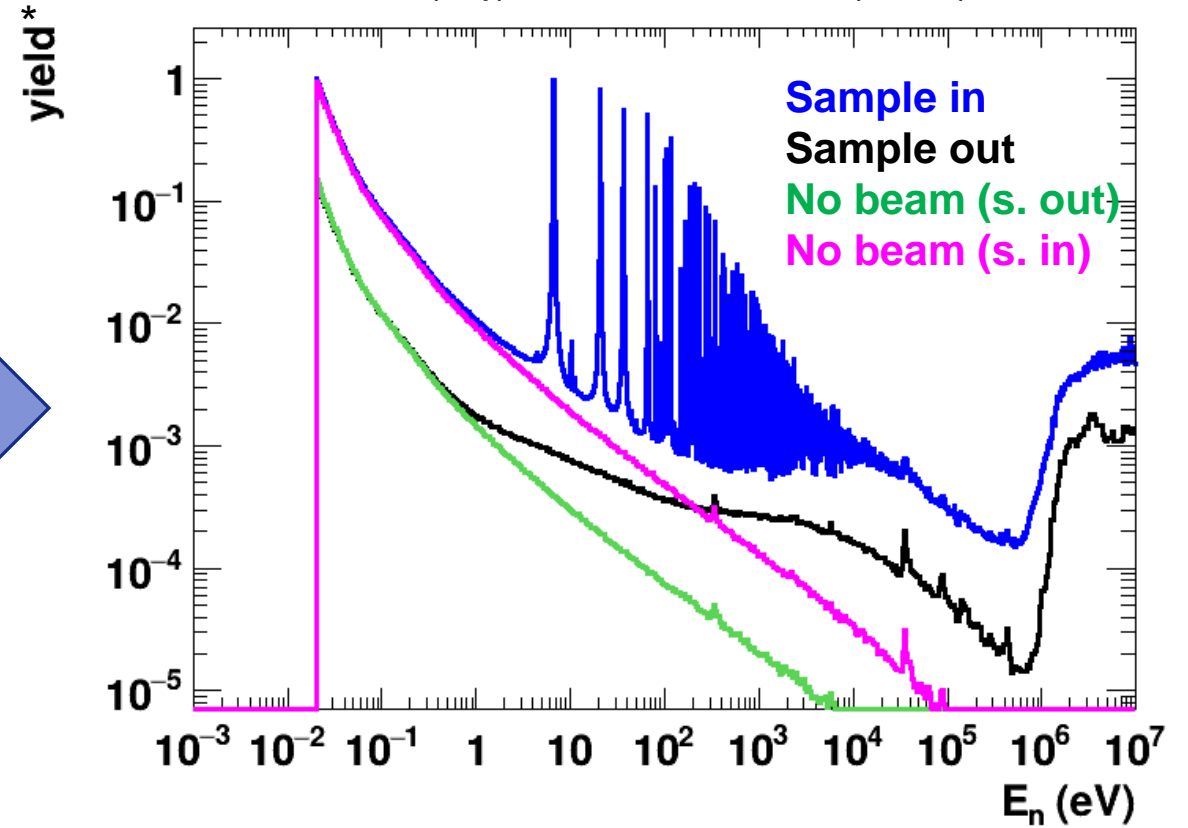


Neutron capture measurements

$^{238}\text{U}(n,\gamma)$ @ n_TOF-EAR1 (2024)



$^{238}\text{U}(n,\gamma)$ @ n_TOF-EAR1 (2024)



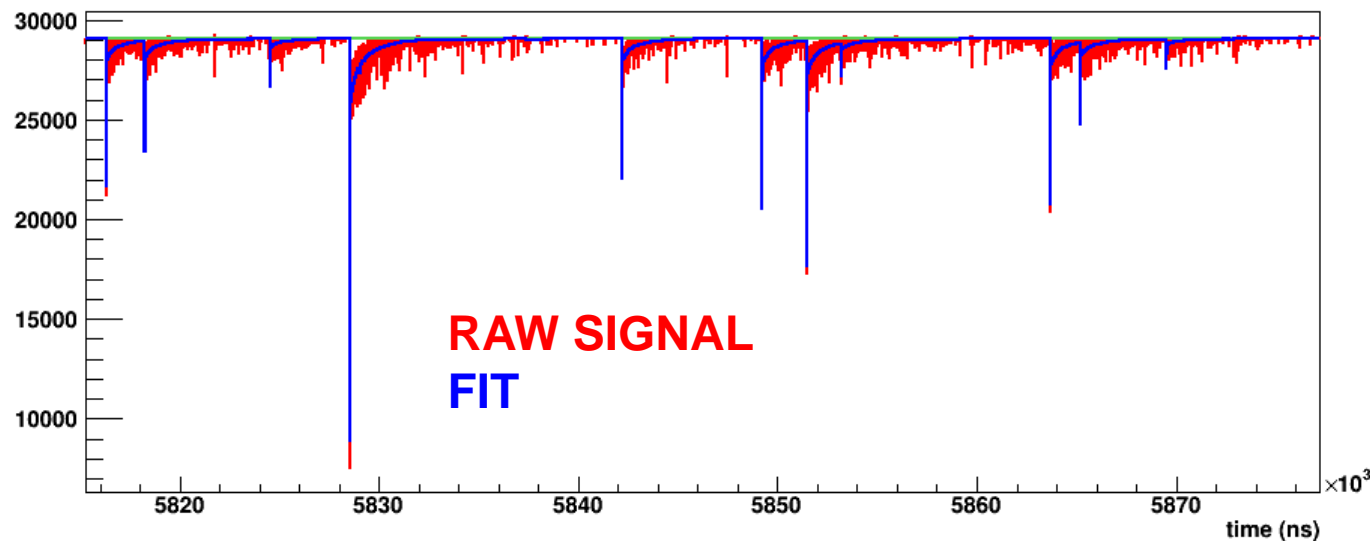
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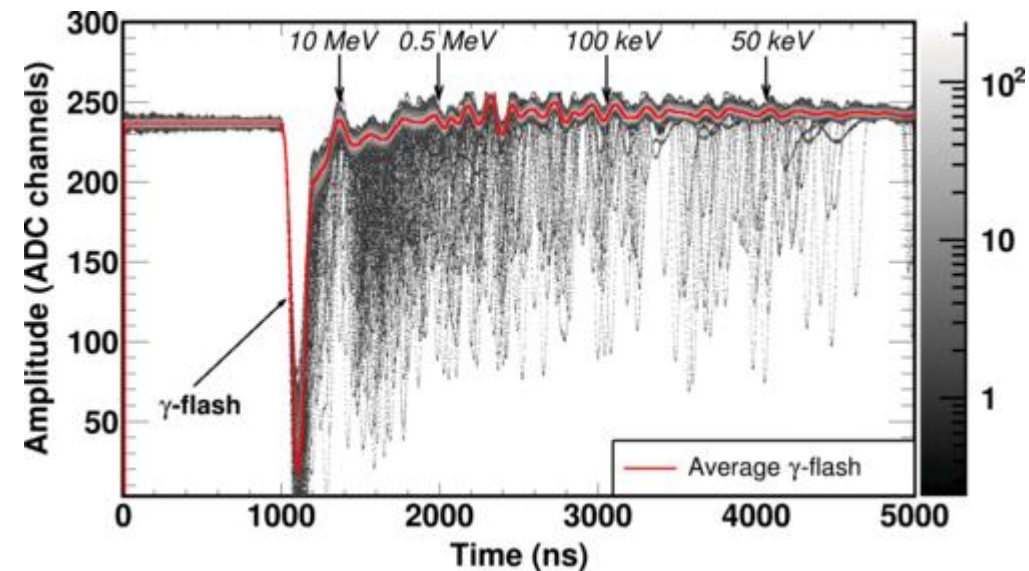
Signal processing

Signals from all detectors are digitized and stored for later analysis by dedicated pulse analysis routines. In most cases there is not much of a problem with this, but there are some exceptions:

- Pulse pile-up
- Signals close to the γ -flash



Signals from the a BaF₂ module of the TAC



Signals close to the γ -flash in a $^{240}\text{Pu}(n,f)$ measurement. Figure from A. Stamatopoulos et al., PRC 102, 014616 (2020).



Pulse pile-up

Pile-up effects appear when the counting rate is large enough so the probability that two signals overlap is not small.

Very often pile-up effects are small (or negligible) and can be corrected without any problem (C_6D_6).

Pile-up occurs when counting rates are high, i.e.:

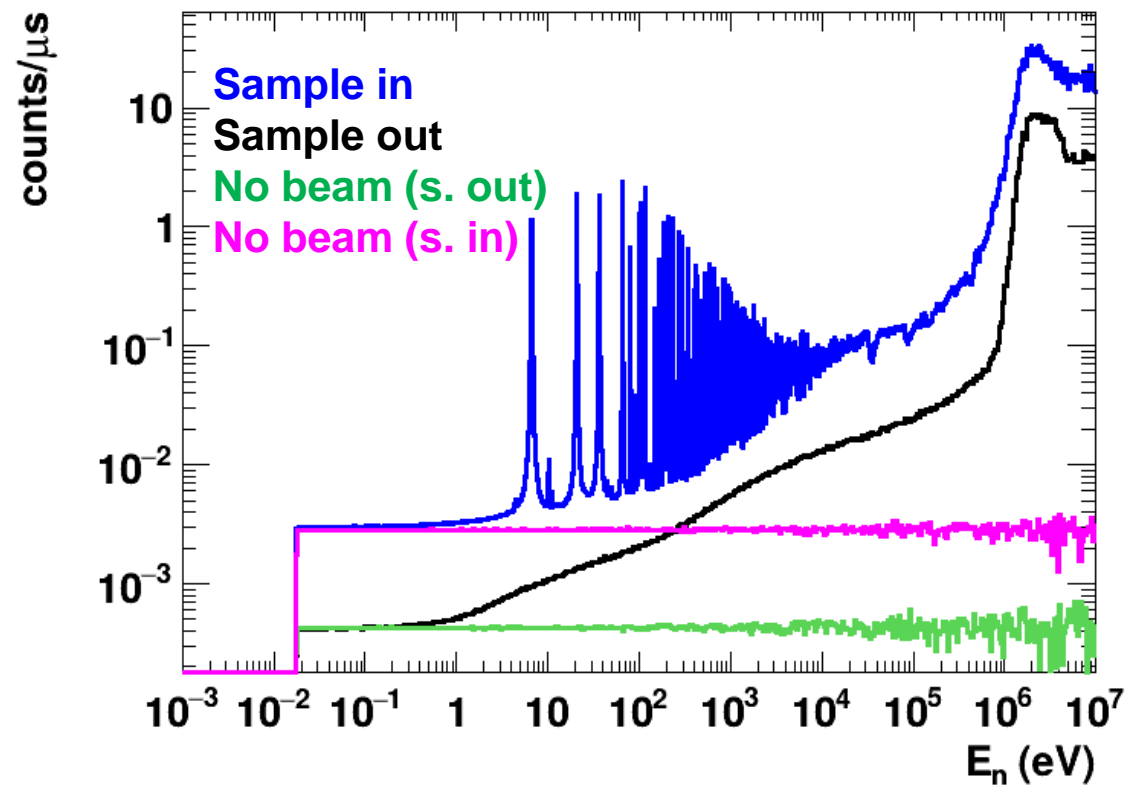
- Strongest resonances.
- At high neutron energies.
- At all neutron energies (high radioactive samples).

BaF_2 signals are much wider than C_6D_6 → sizeable pile-up.

Some works regarding pile-up corrections:

- *E. Mendoza et al., NIMA 768, 55 (2014).*
- *C. Guerrero et al., NIMA 777, 63 (2015).*

$^{238}U(n,\gamma)$ @ n_TOF-EAR1 (2024)



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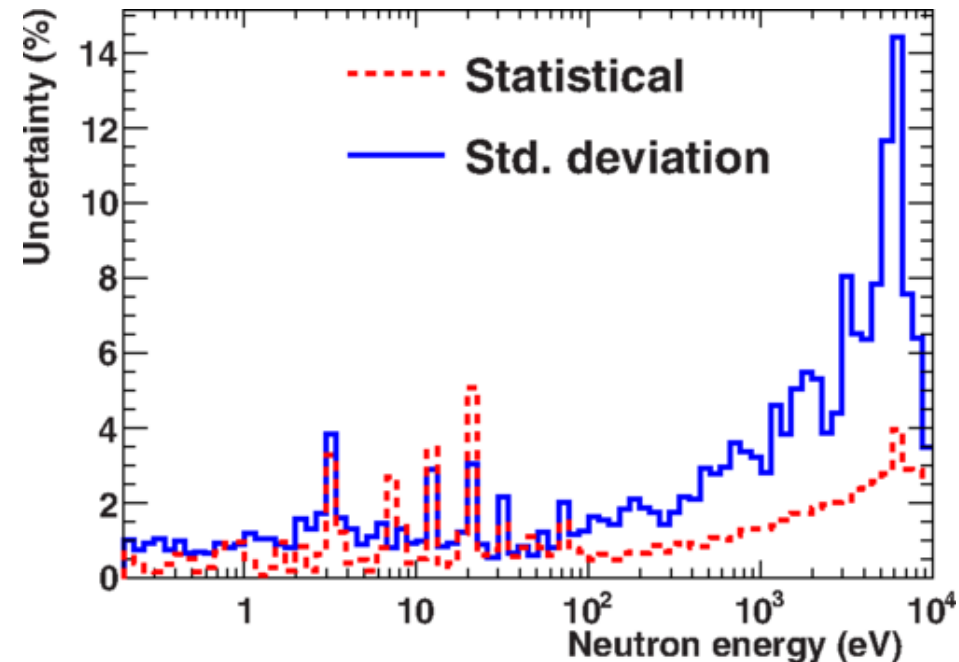
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Possible biases:

- In the strongest resonances.
- Increasing with the neutron energy.



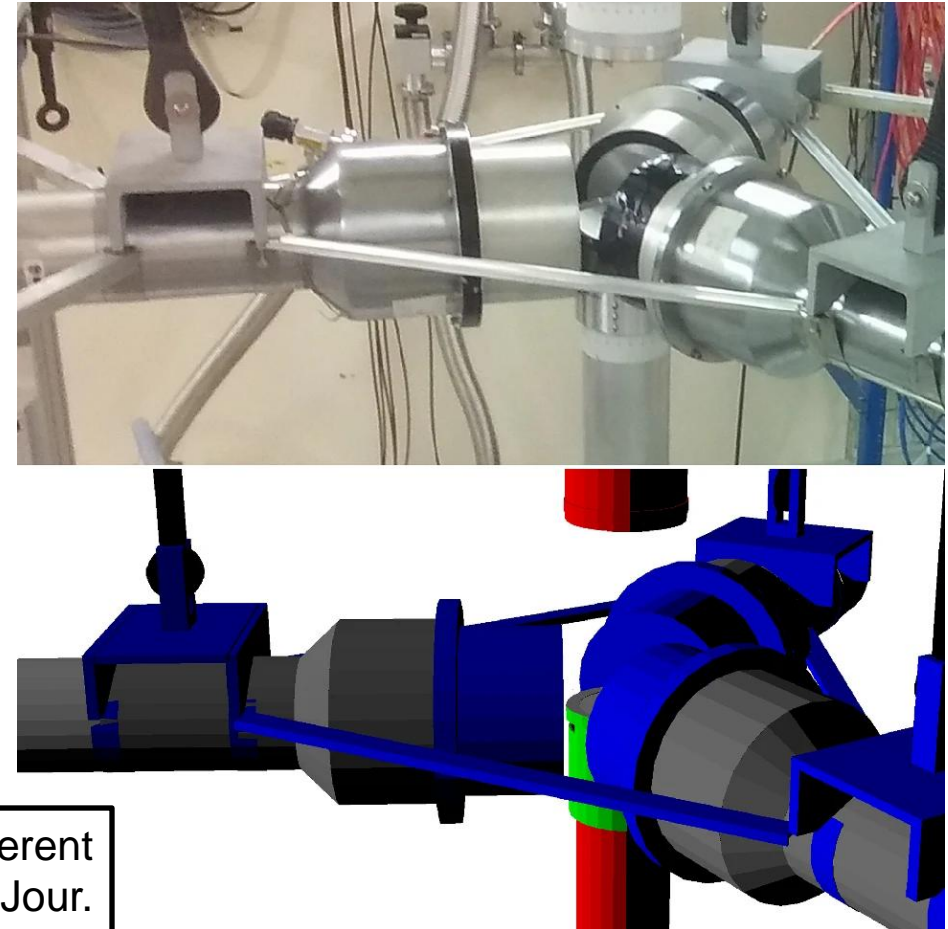
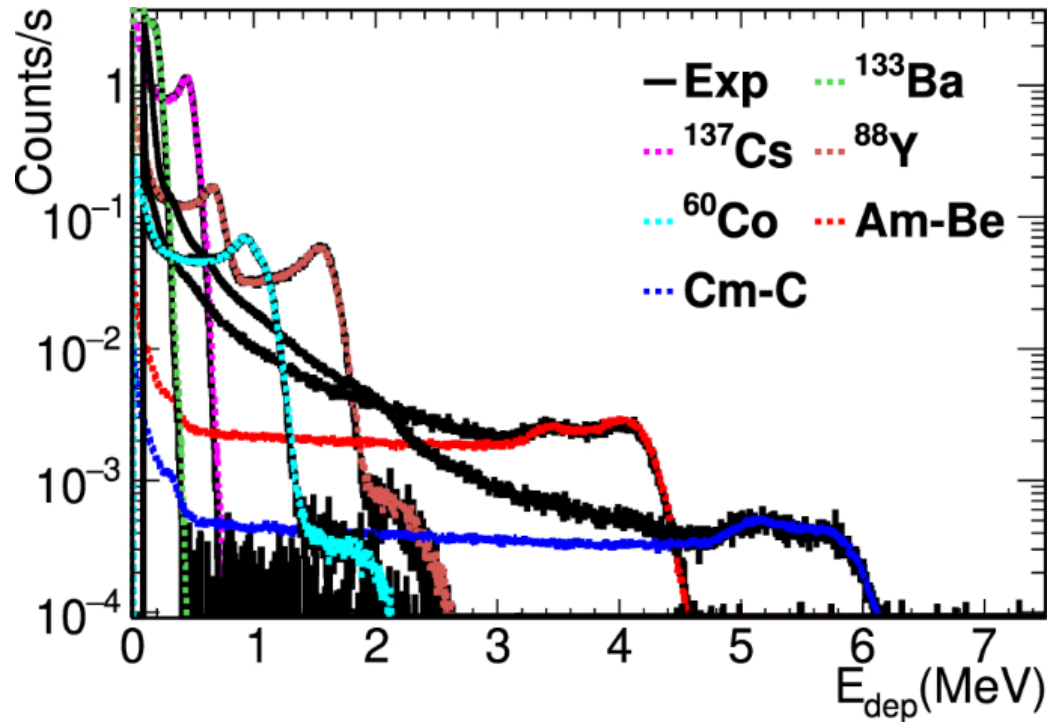
Uncertainties in a $^{241}Am(n,\gamma)$ measurement performed with the TAC. Figure from E. Mendoza et al., PRC 97, 054616 (2018).



Calibrations

Energy calibrations are performed with calibration sources.

The gain of the detectors may vary along time → gain monitoring.

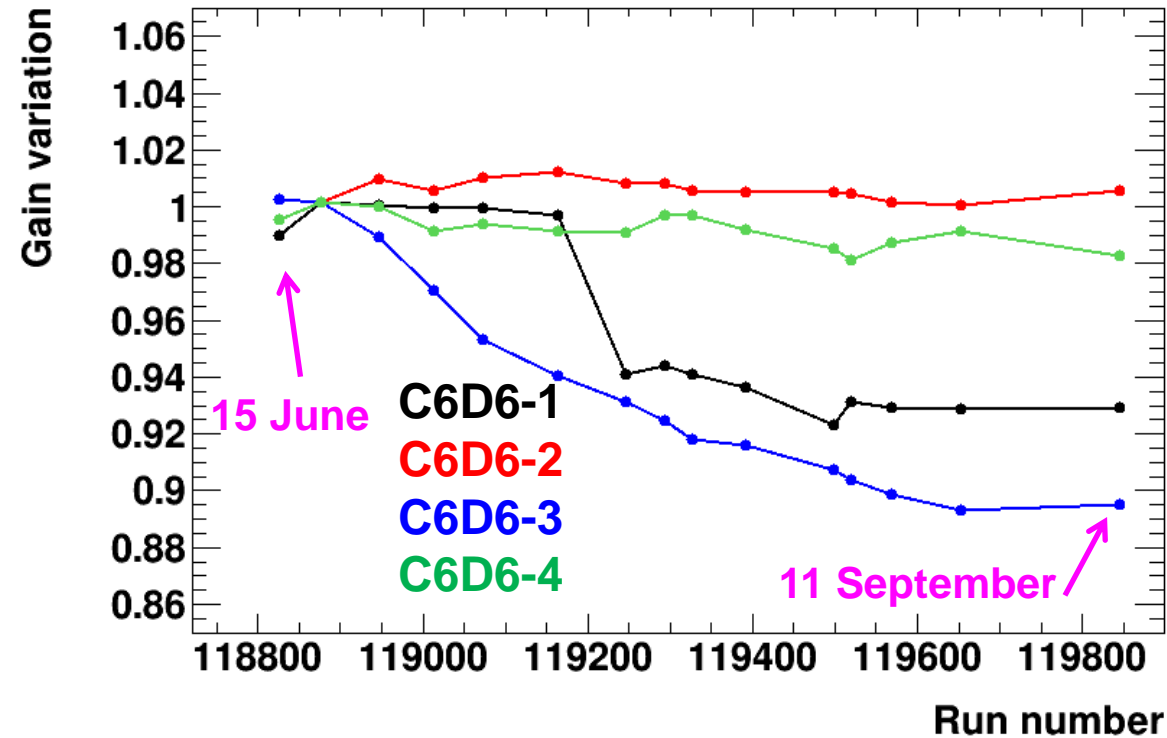


Experimental setup (right) and amplitude spectra when measuring with different calibration γ -ray sources (left). Figures from V. Alcayne et al., Eur. Phys. Jour. 60, 246 (2024) → $^{246,248}\text{Cm}(n,\gamma)$ measurement at n_TOF-EAR2.

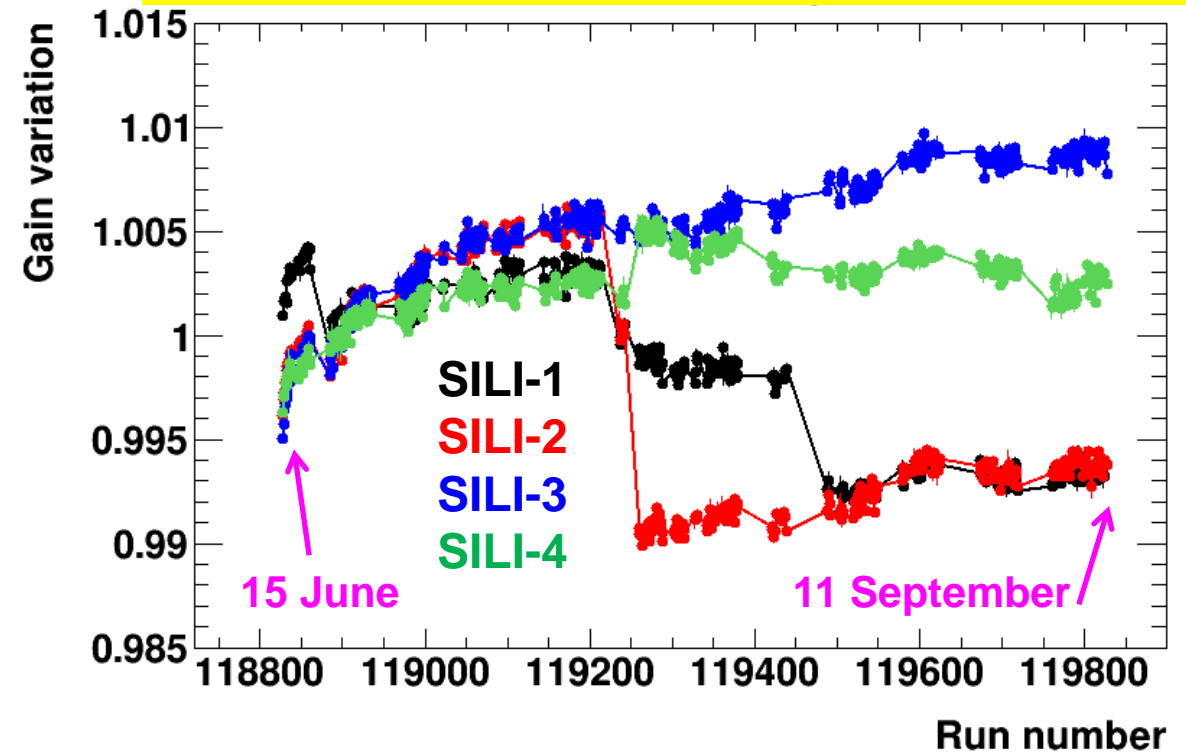


Calibrations

(n,γ) measurements of ^{167}Er , ^{166}Er , ^{63}Cu , ^{238}U - 2024



(n,γ) measurements of ^{167}Er , ^{166}Er , ^{63}Cu , ^{238}U - 2024

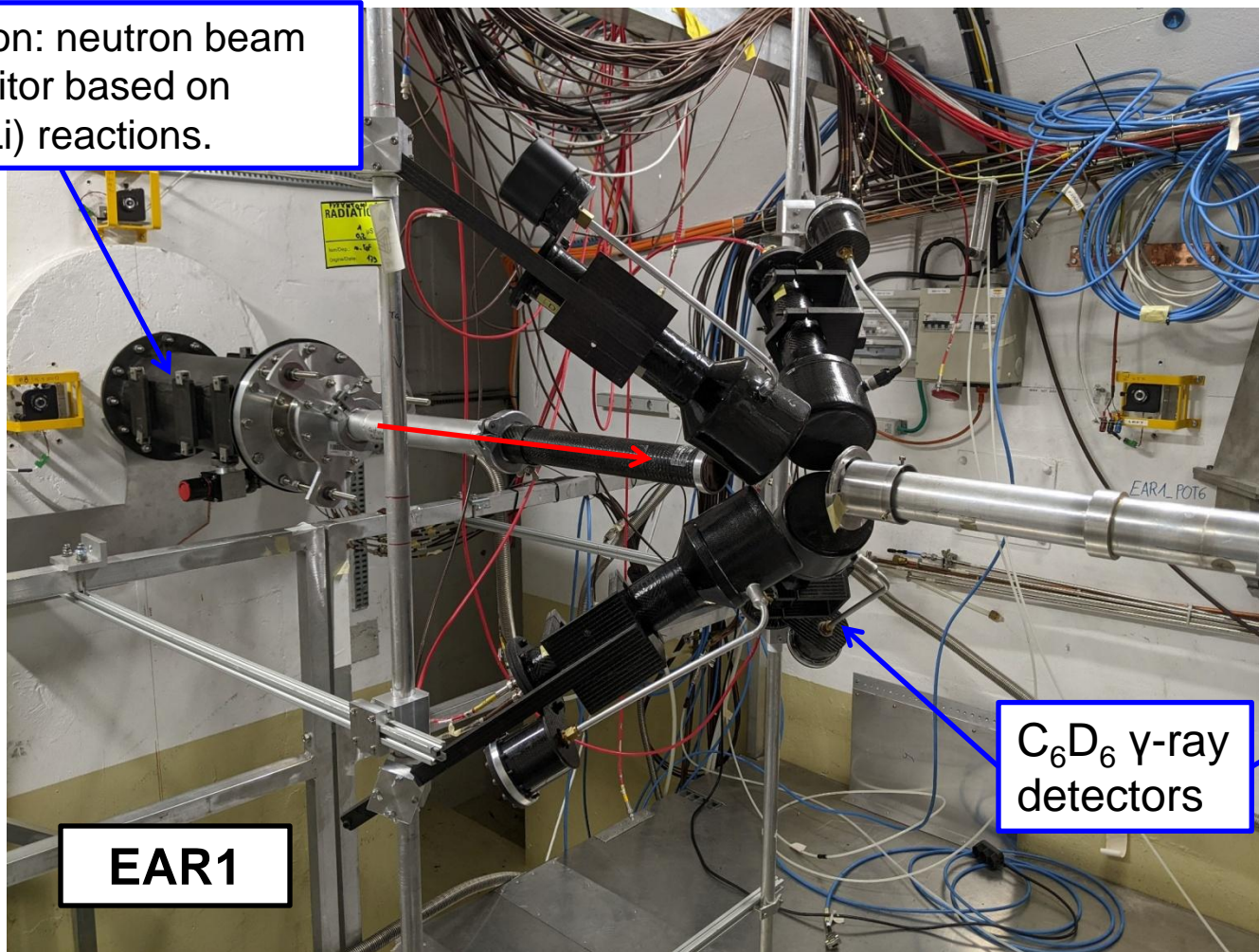


Variations in the gain of four C_6D_6 detectors (left) and four SiMon detectors (right) during 3 months of experimental campaign in 2024.



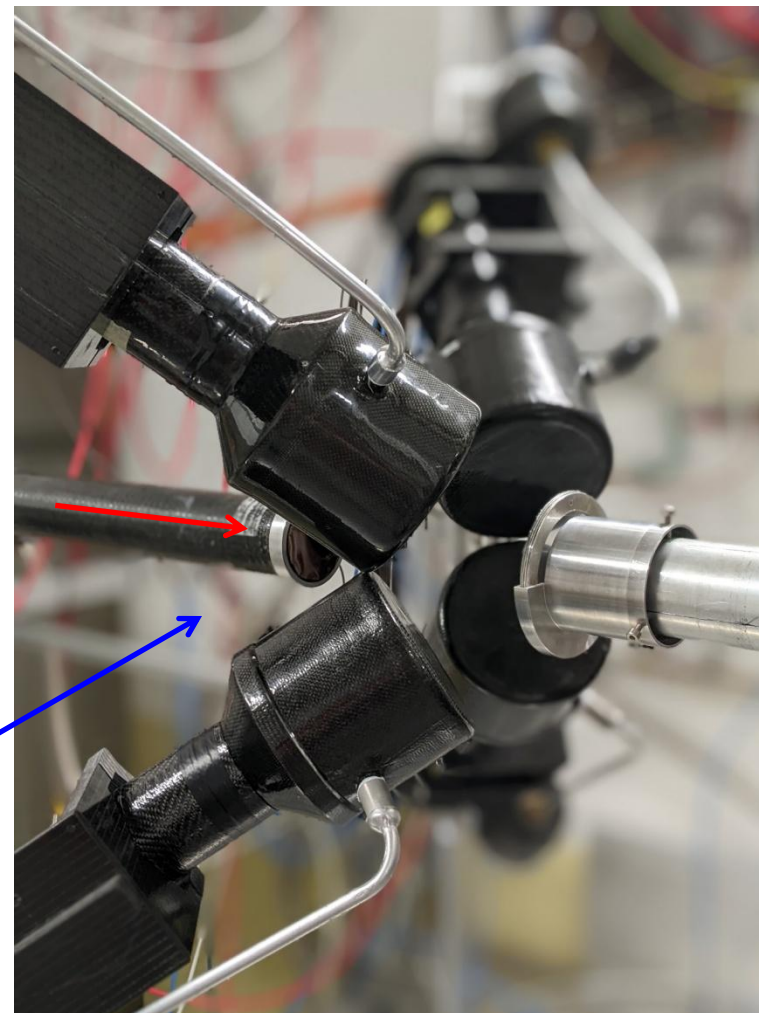
Detection systems

SiMon: neutron beam monitor based on $(n, {}^6\text{Li})$ reactions.



C_6D_6 γ -ray detectors

EAR1



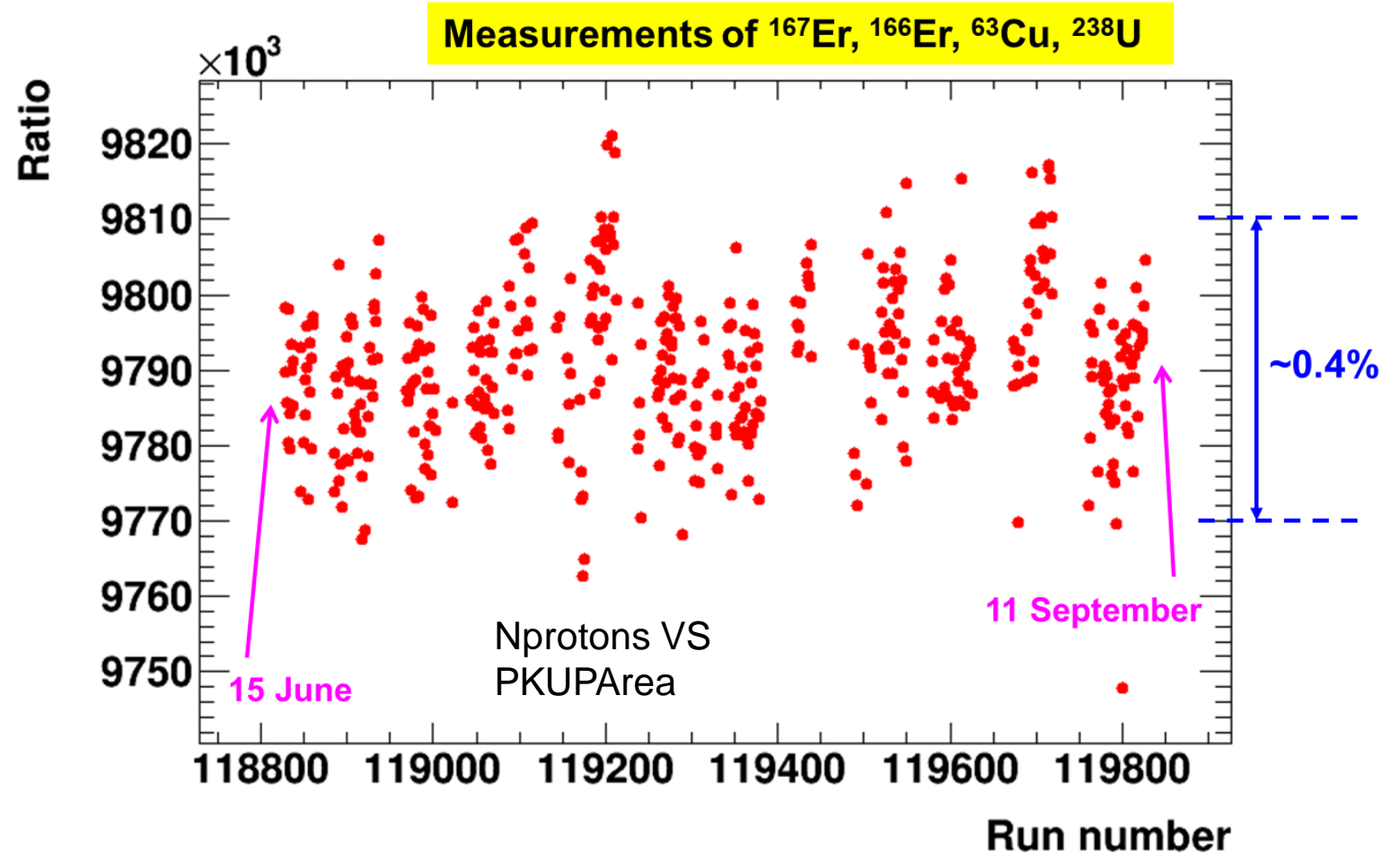
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Consistency checks

After performing calibrations, several consistency checks are performed.
Example: ratios between beam monitors should be constant along time.

$$Y_{(n,\gamma)}(E_n) = \frac{C(E_n) - B(E_n)}{\varepsilon(E_n) \cdot \phi(E_n)}$$



Background

The calculation of the background is one of the most difficult tasks in neutron capture measurements.

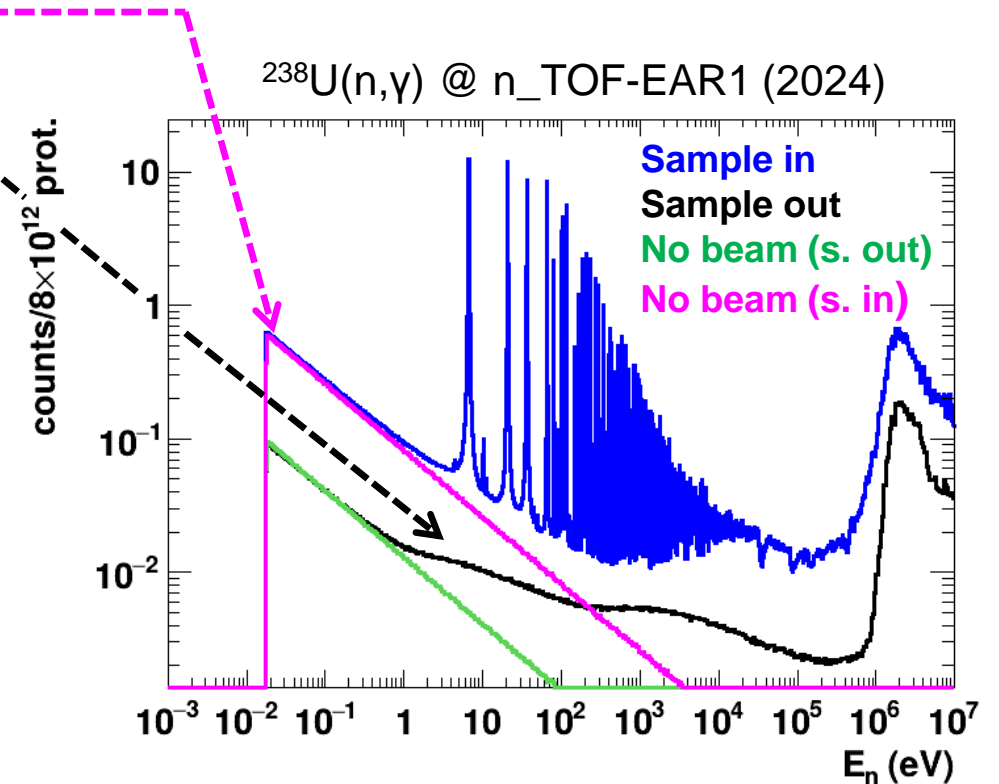
We can classify the background in different components:

1. Background not related with the neutron beam (ambient background + sample activity).
2. Background related with the neutron beam, excluding the interaction of the neutron beam with the sample.
3. Background related with the interaction of the neutron beam with the sample.

Backgrounds (1) and (2) are obtained from dedicated measurements.

Background (3) is more difficult to obtain.

$$Y_{(n,\gamma)}(E_n) = \frac{C(E_n) - B(E_n)}{\varepsilon(E_n) \cdot \phi(E_n)}$$



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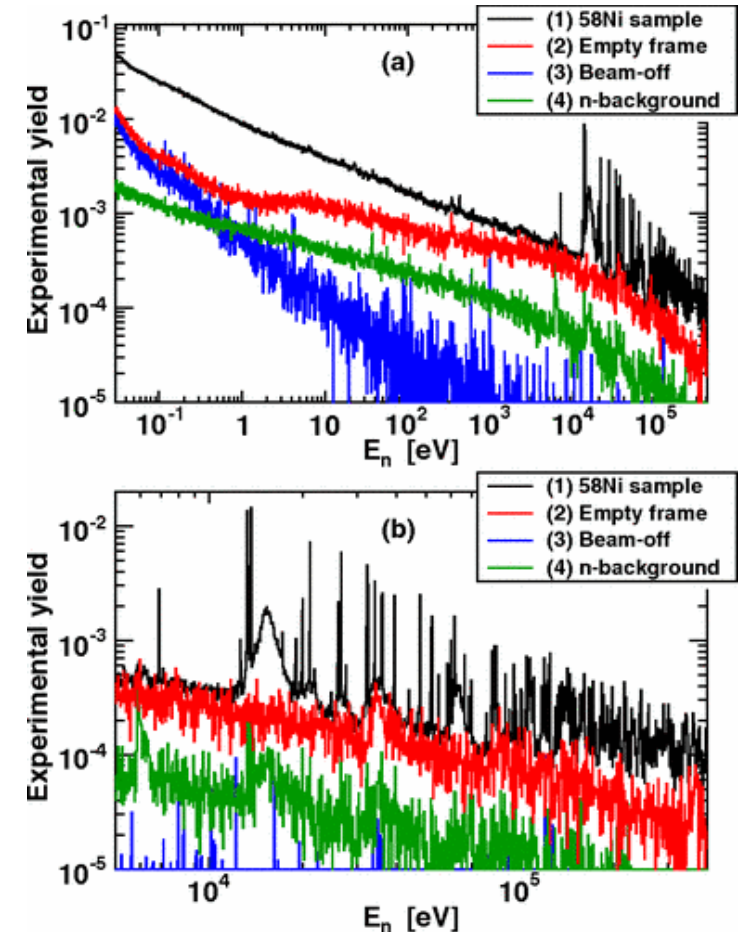
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Background (3) example: in the analysis of $^{58}\text{Ni}(n,\gamma)$ shown on the right this component has been obtained from MC calculations.

Sometimes background (3) is not taken much into account in the analysis ...



$^{58}\text{Ni}(n,\gamma)$: P. Žugec et al. Phys. Rev. C 89, 014605 (2014).



Background (3)

The relevance of background (3) is very case-dependent, and its origin may come from:

- In-beam γ -rays: relevant for heavy nuclei.
- Elastic scattered neutrons in the sample: more relevant for nuclei with large elastic to capture cross section ratios.
- Fission
- Inelastic channels (\rightarrow increasing threshold)
- Activation (\rightarrow MC calculations)
- Impurities (\rightarrow SAMMY analysis)



Background (3) – In beam γ -rays and (n,n)

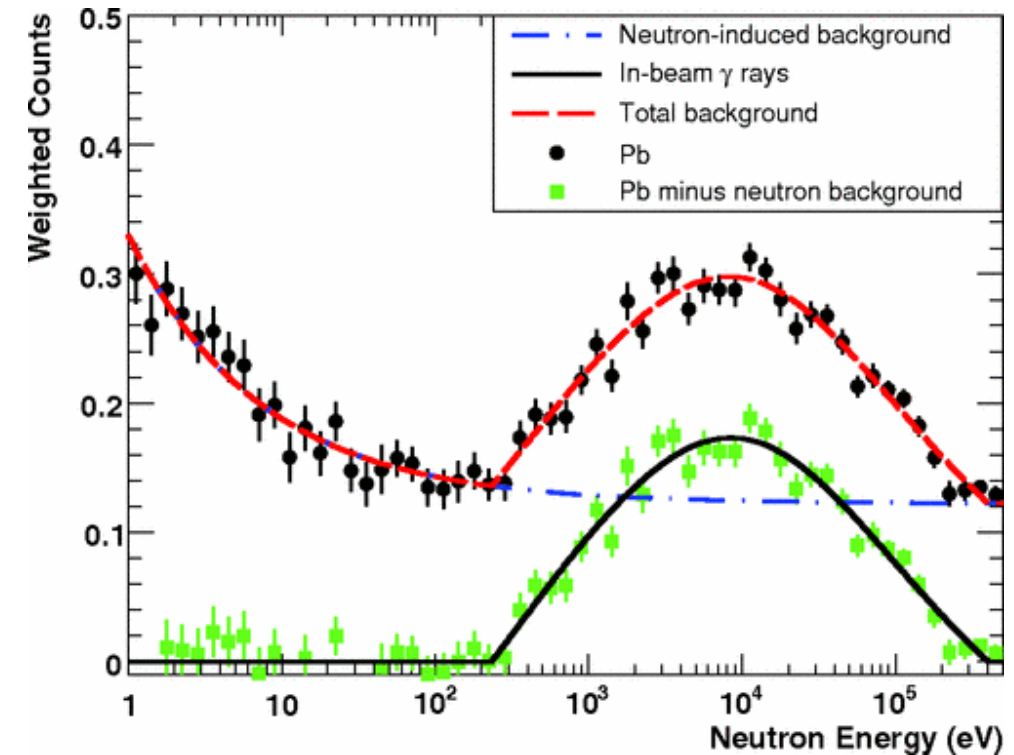
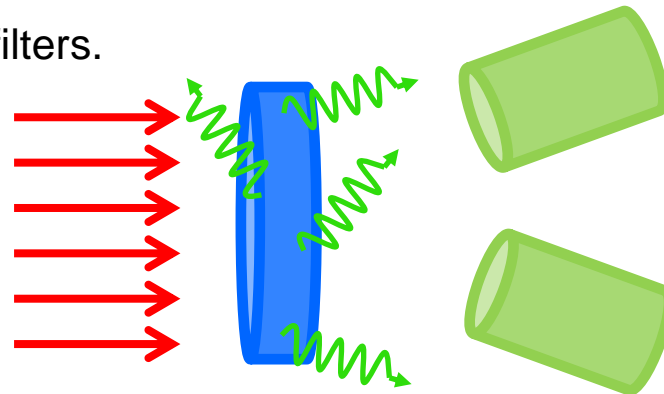
There are γ -rays coming along the beam together with the neutrons from the spallation target and surroundings that may be scattered in the sample \rightarrow in beam γ -ray background.

Elastic scattered neutrons: there is a *prompt* and a *delayed* contribution.

- Prompt: resonant structure.
- Delayed: energy-TOF relation is lost.

To determine these components:

- Measurements with a graphite and a Pb samples.
- Monte Carlo calculations.
- Measurements with filters.



$^{197}\text{Au}(n,\gamma)$ with C_6D_6 : C. Lederer et al.
Phys. Rev. C 83, 034608 (2011).



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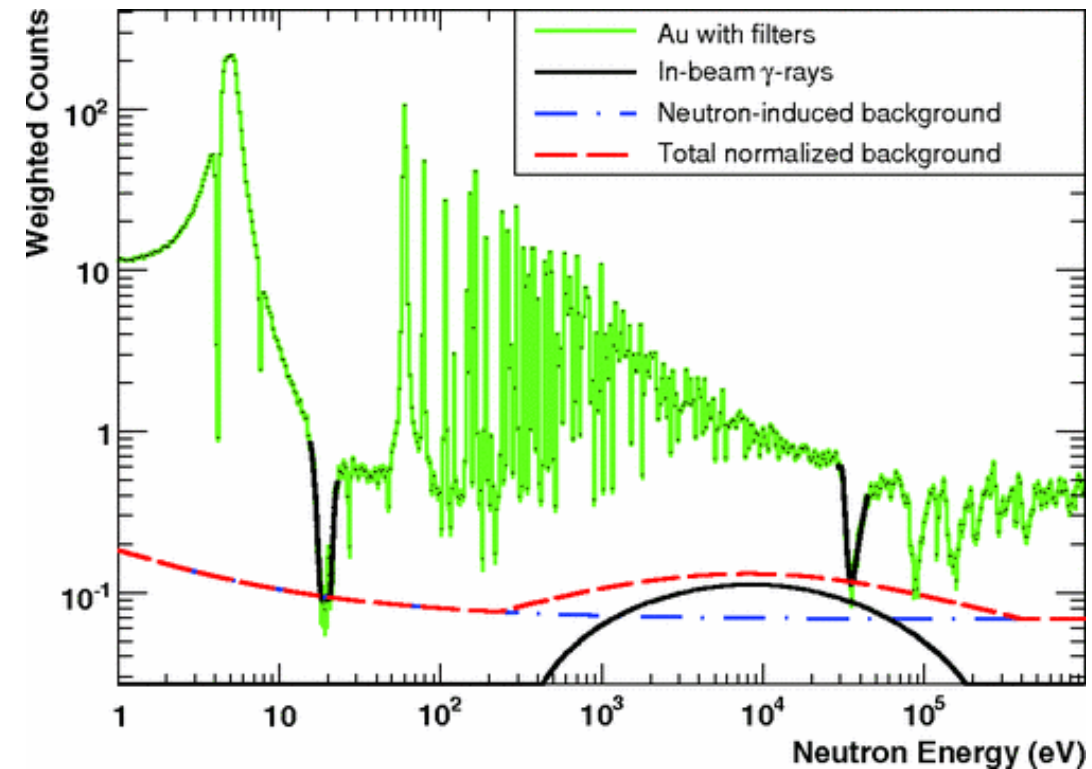
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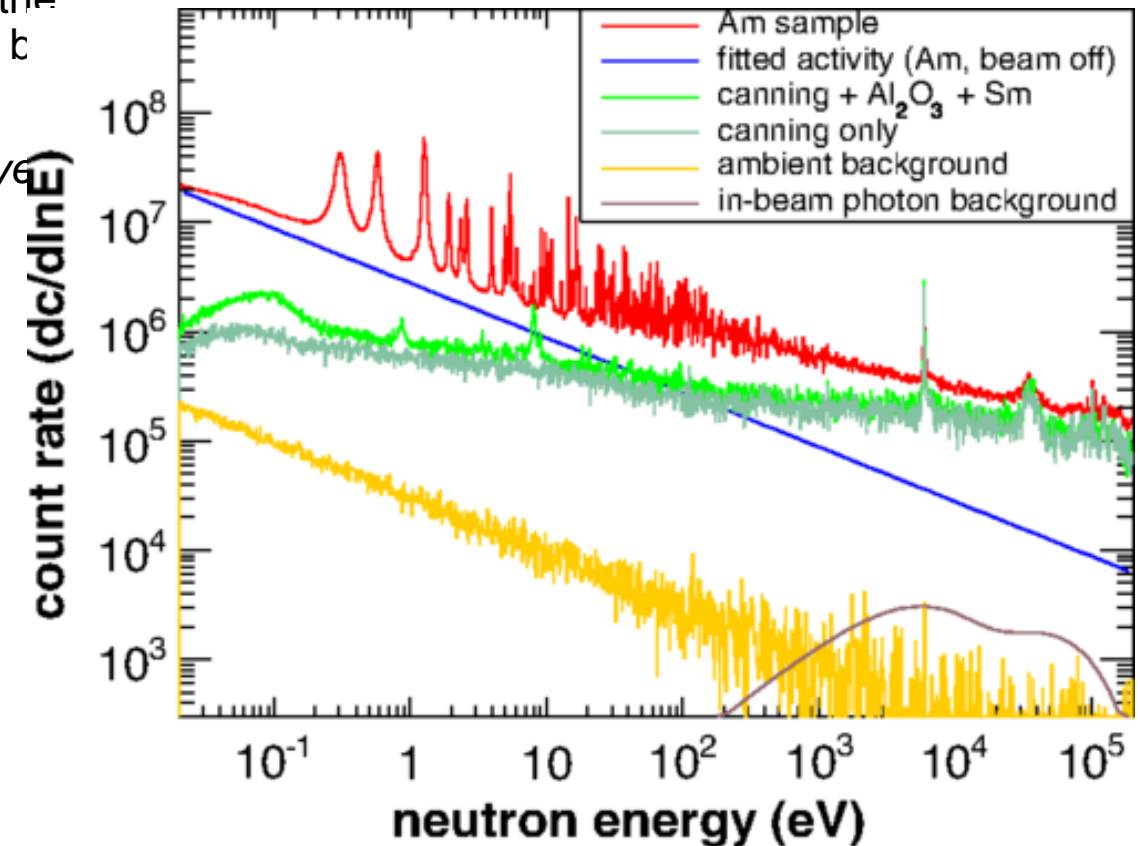
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$^{241}\text{Am}(n,\gamma)$ with C_6D_6 : K. Fraval et al.
Phys. Rev. C 89, 044609 (2014).



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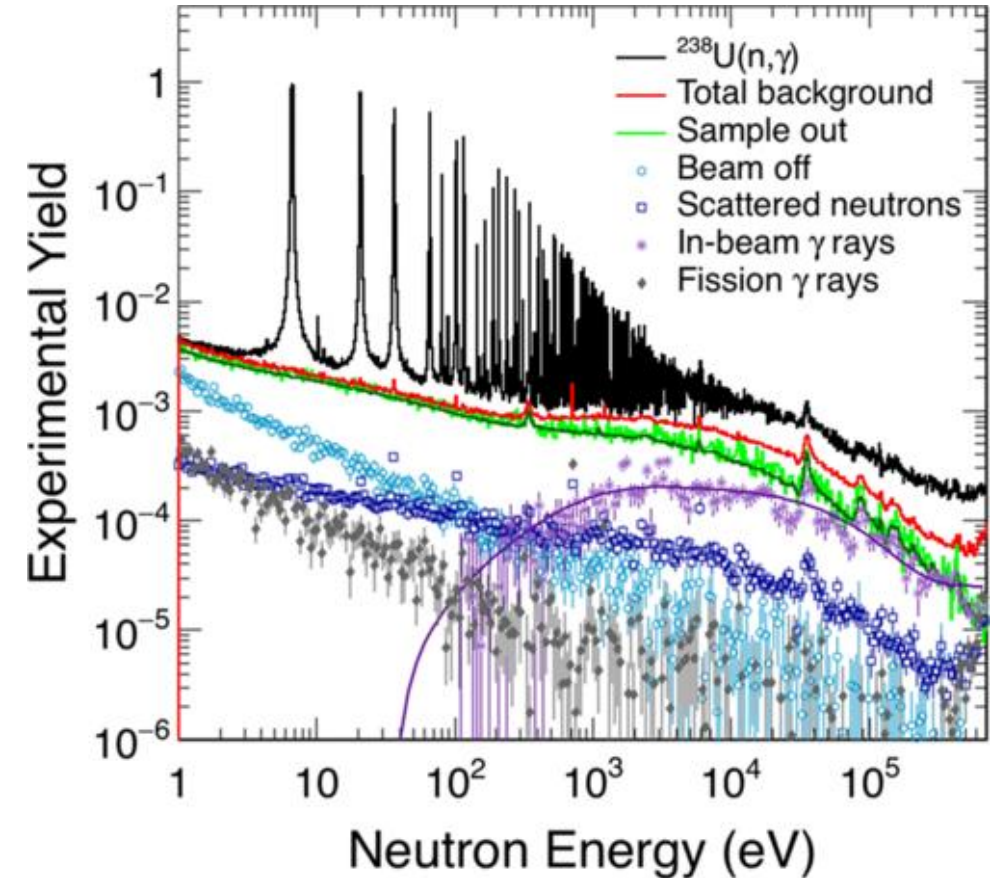
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$^{238}\text{U}(n,\gamma)$ with C_6D_6 : F. Mingrone et al.
Phys. Rev. C 95, 034604 (2017).



Background (3) – In beam γ -rays and (n,n)

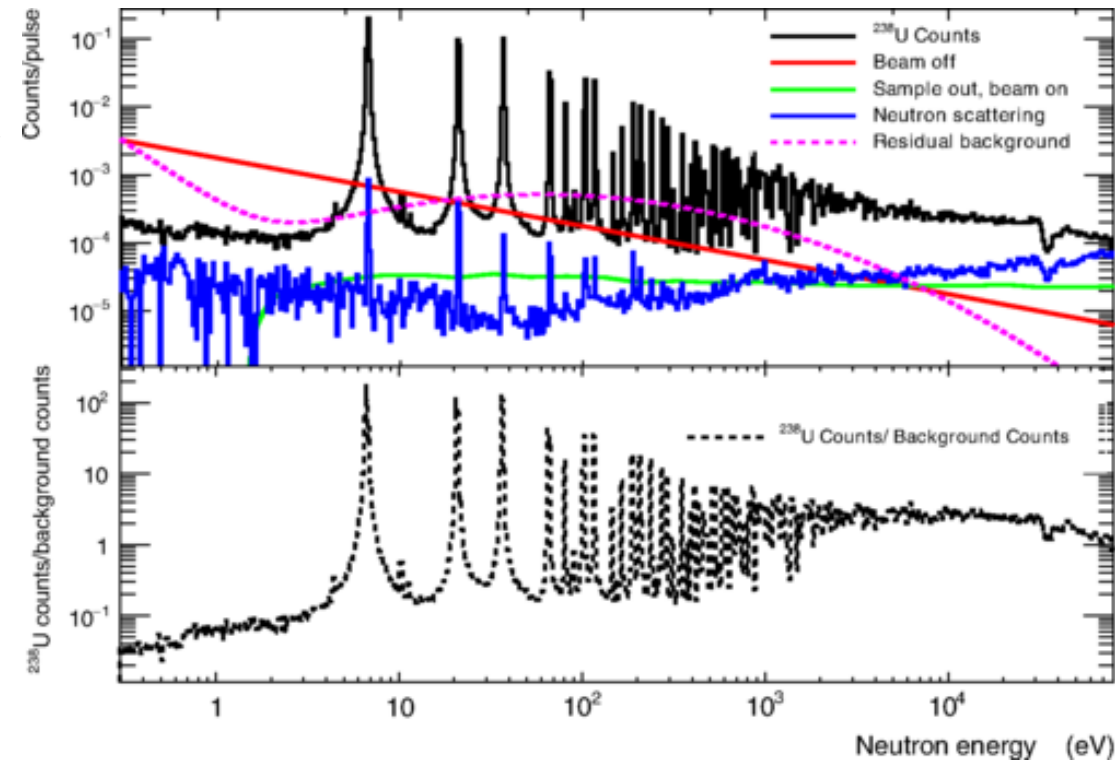
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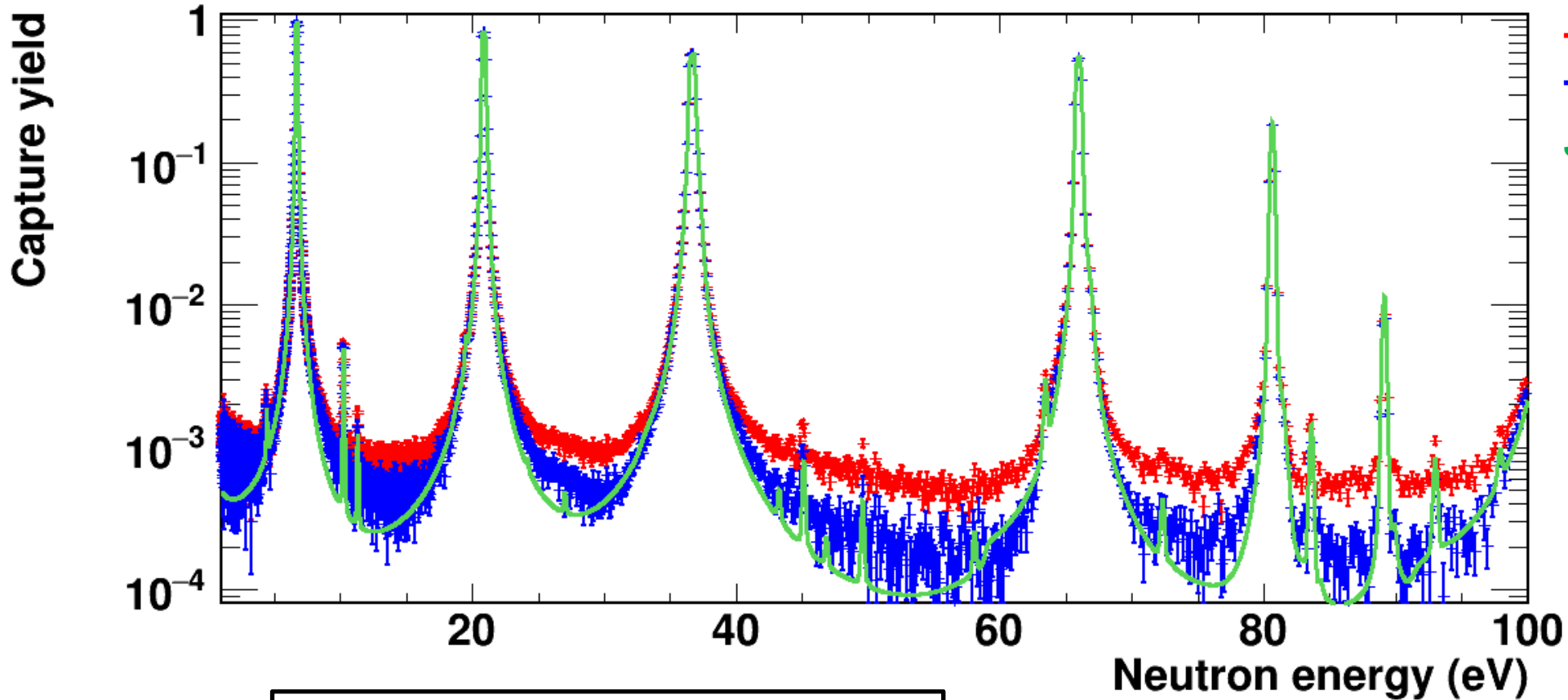
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²³⁸U(n, γ) with the TAC: T. Wright et al.
Phys. Rev. C 96, 064601(2017).



Background (3) – In beam γ -rays and (n,n)

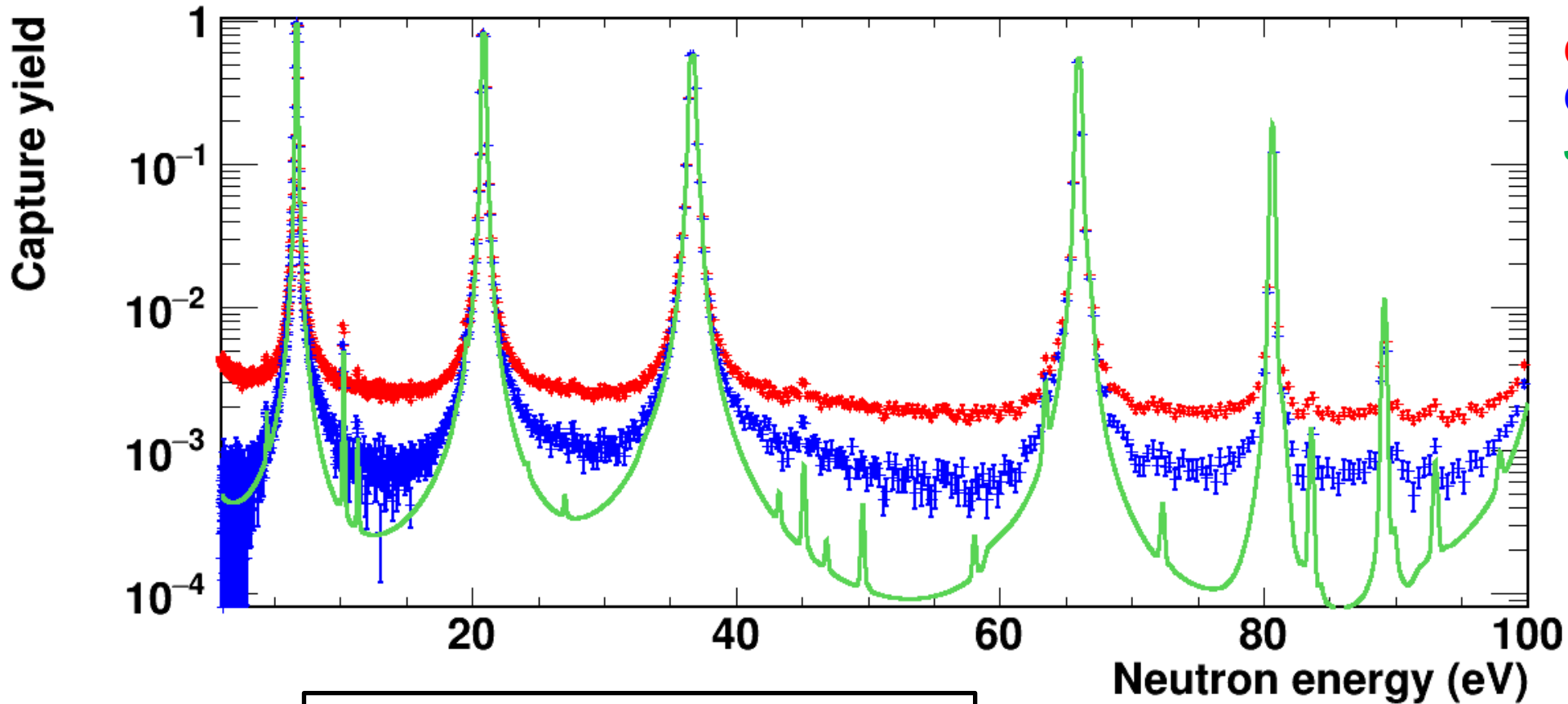


TAC – no *residual bkg.*
TAC – with *residual bkg.*
JEFF-3.3

$^{238}\text{U}(n,\gamma)$ with the TAC: T. Wright et al.
Phys. Rev. C 96, 064601(2017).



Background (3) – In beam γ -rays and (n,n)



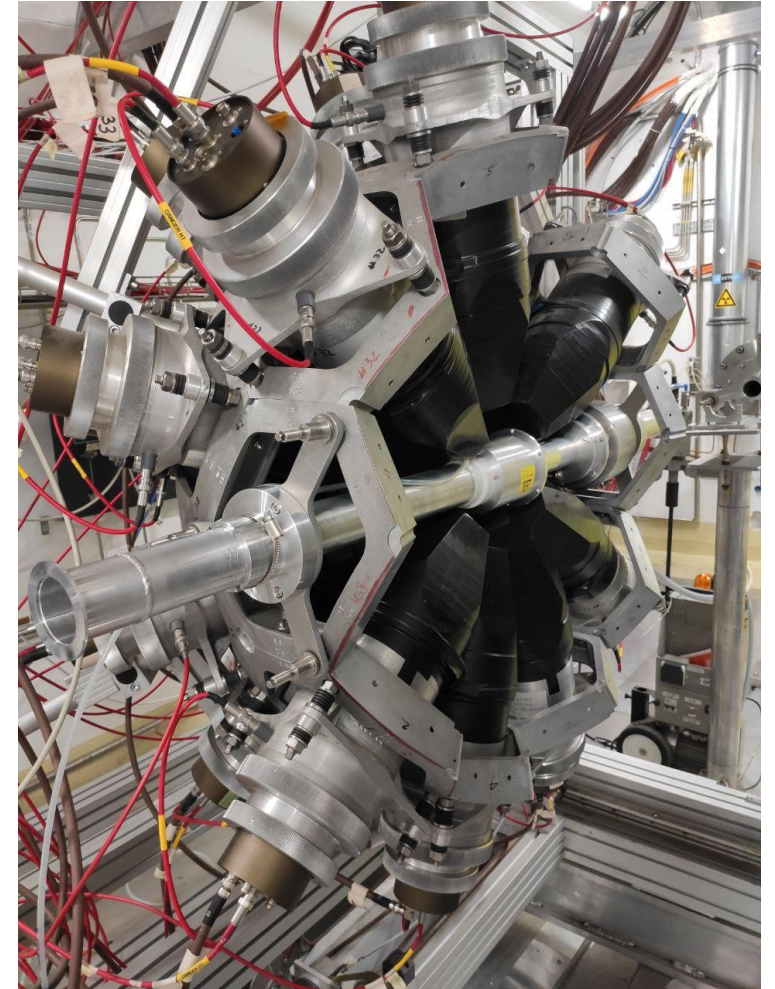
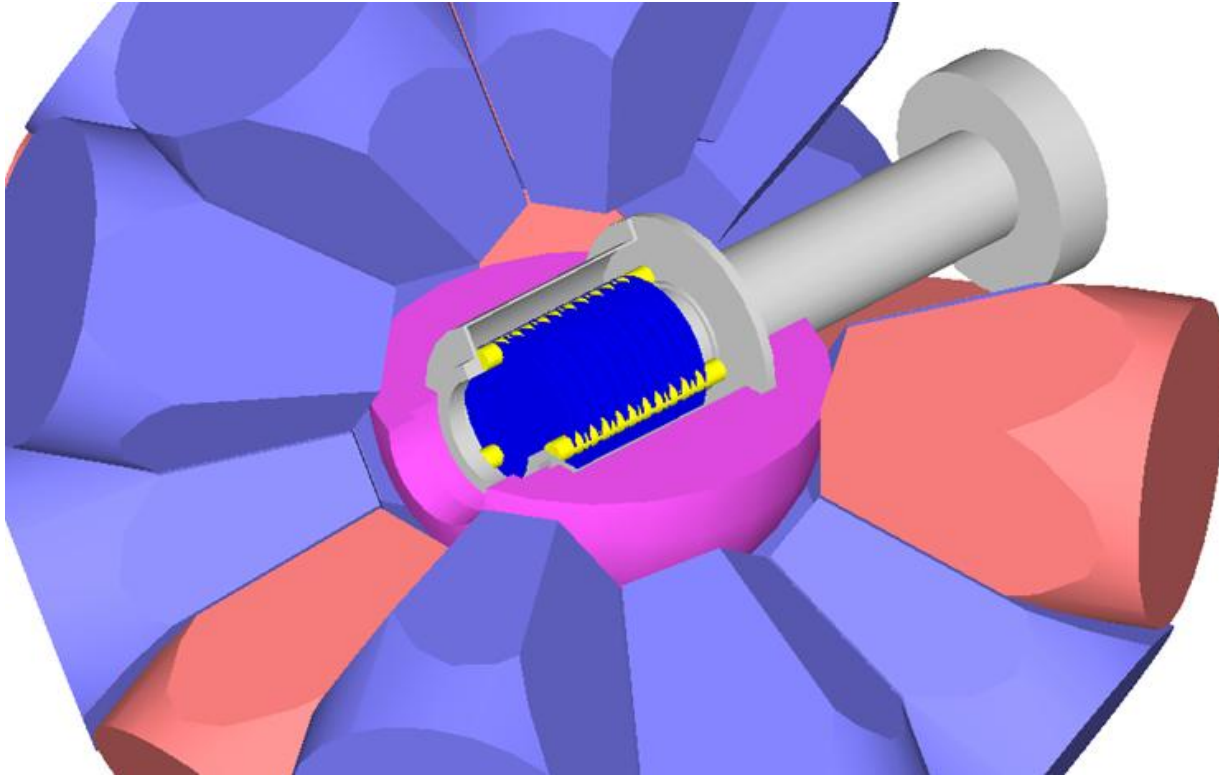
C_6D_6 – no residual bkg.
 C_6D_6 – with residual bkg.
JEFF-3.3

$^{238}U(n,\gamma)$ with C_6D_6 : F. Mingrone et al.
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Background (3) - Fission

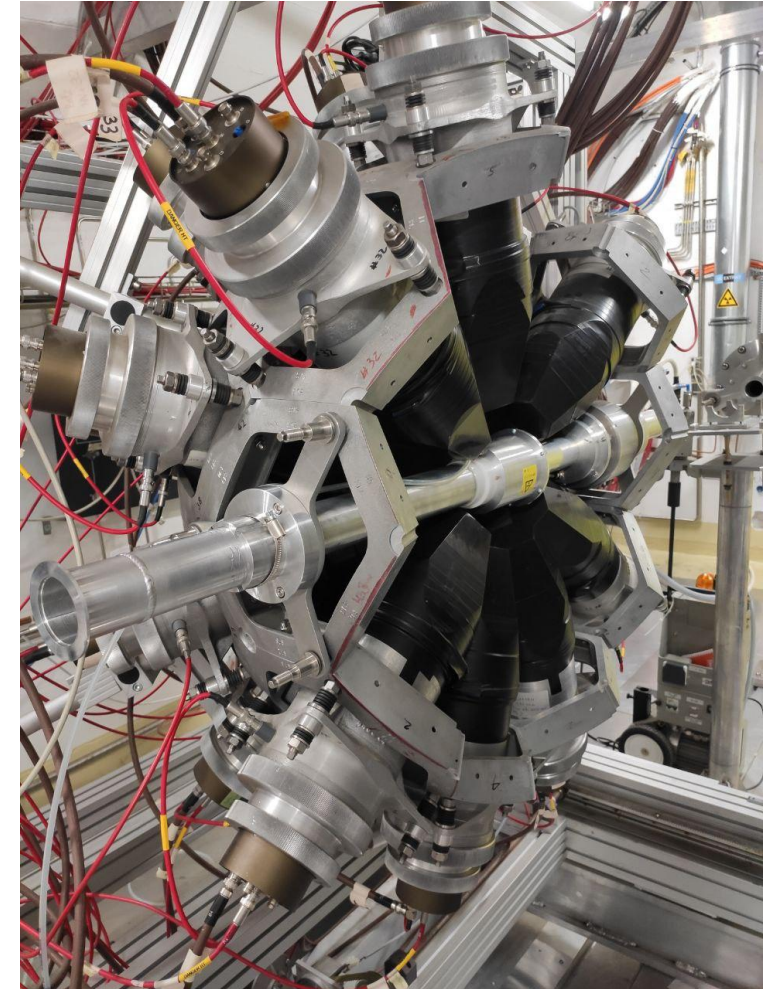
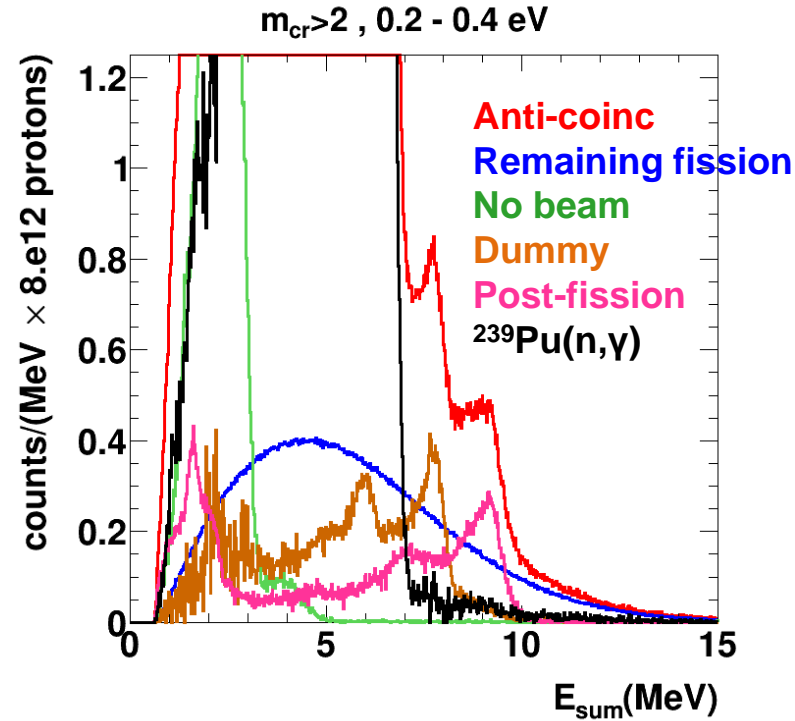
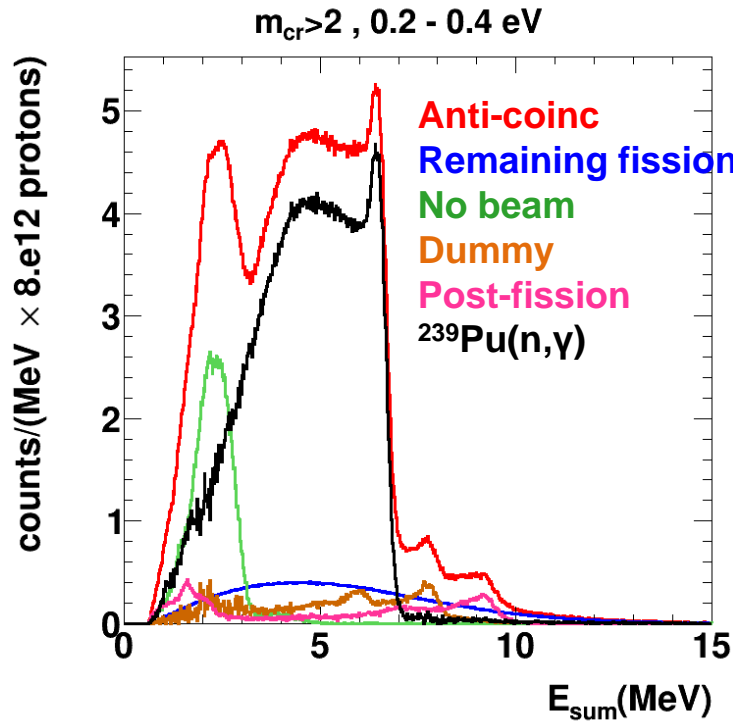
Neutron capture measurements of fissile isotopes ($^{233,235}\text{U}$, $^{239,241}\text{Pu}$) have been performed at n_TOF with the TAC and using a fission fragment detector in anticoincidence.



Background (3) - Fission

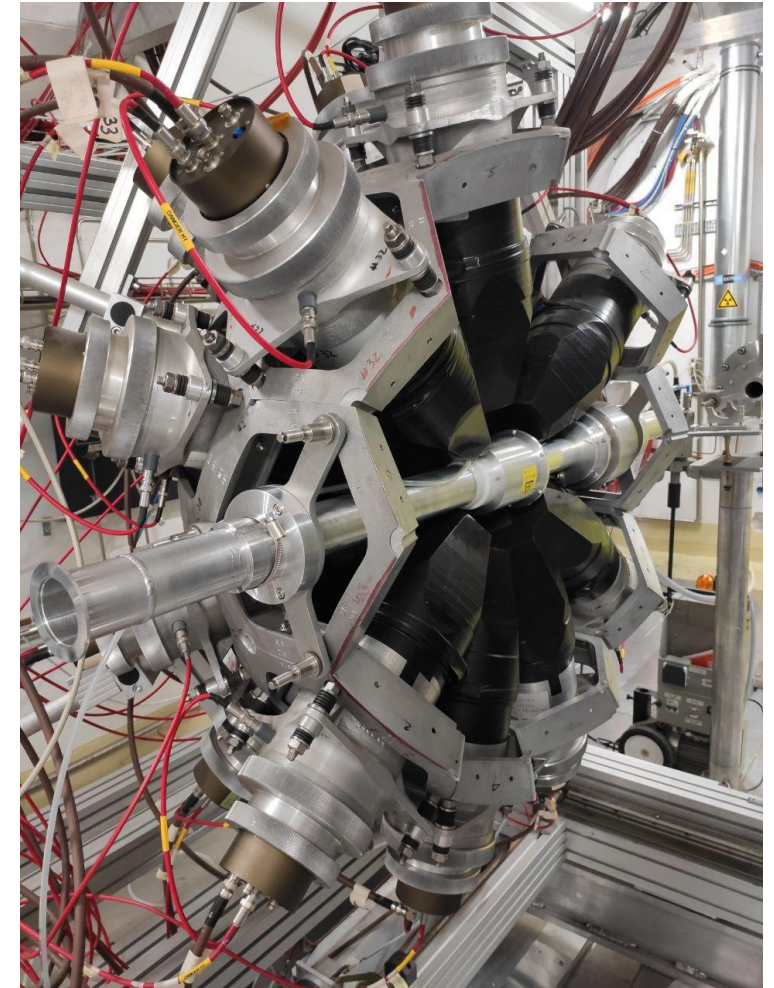
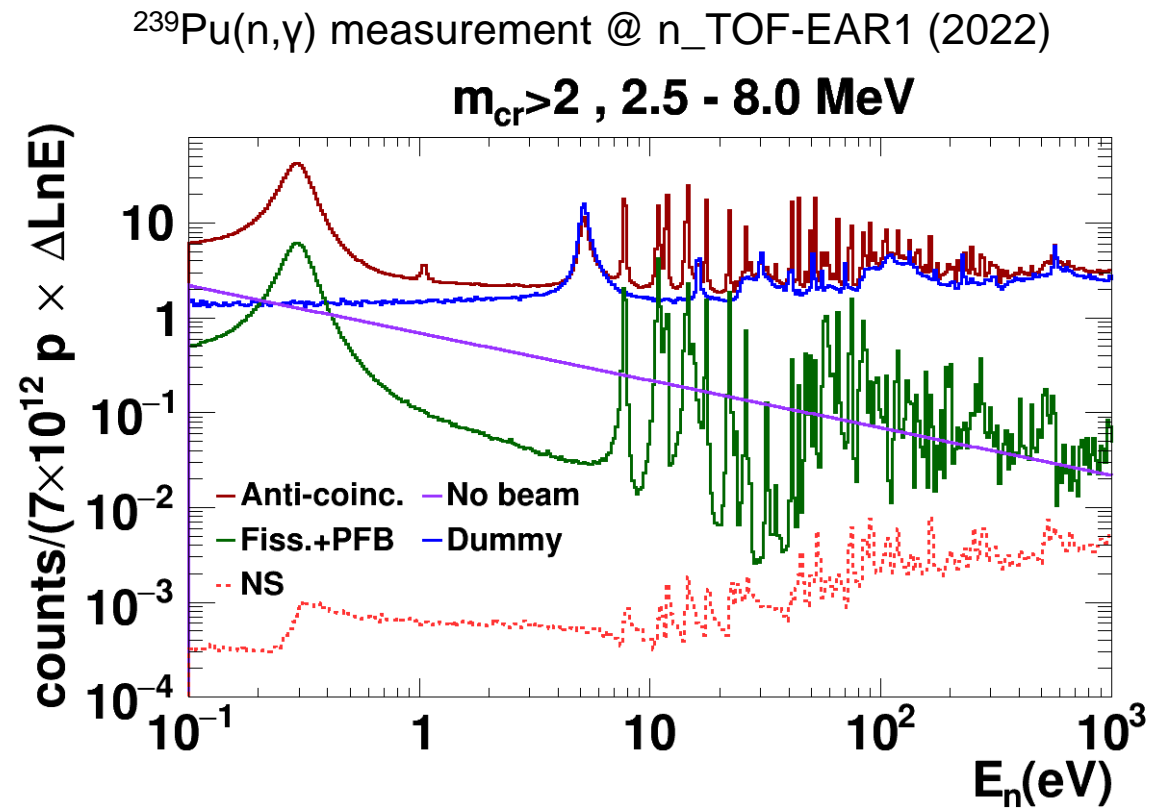
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$^{239}\text{Pu}(n,\gamma)$ measurement @ n_TOF-EAR1 (2022)



Background (3) - Fission

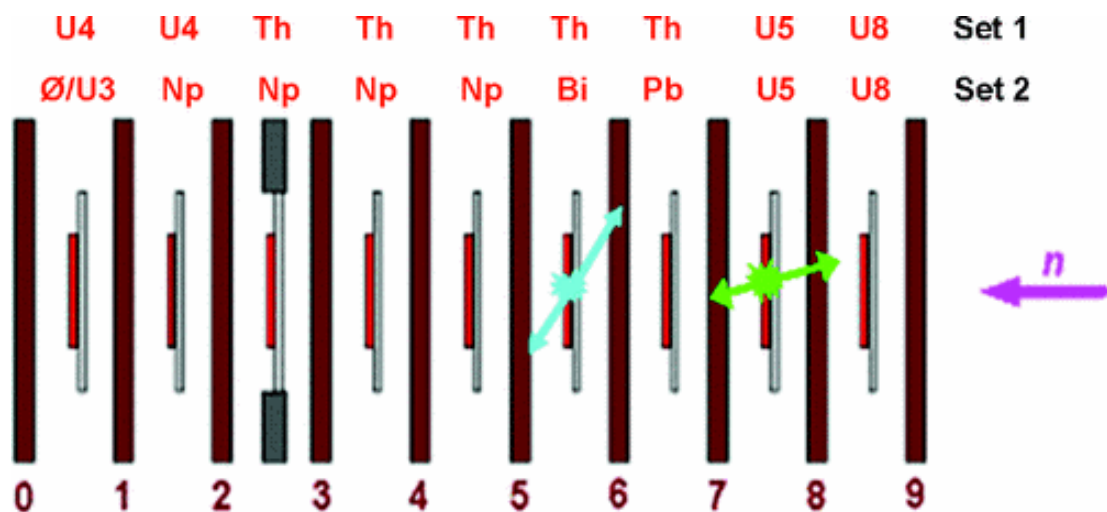
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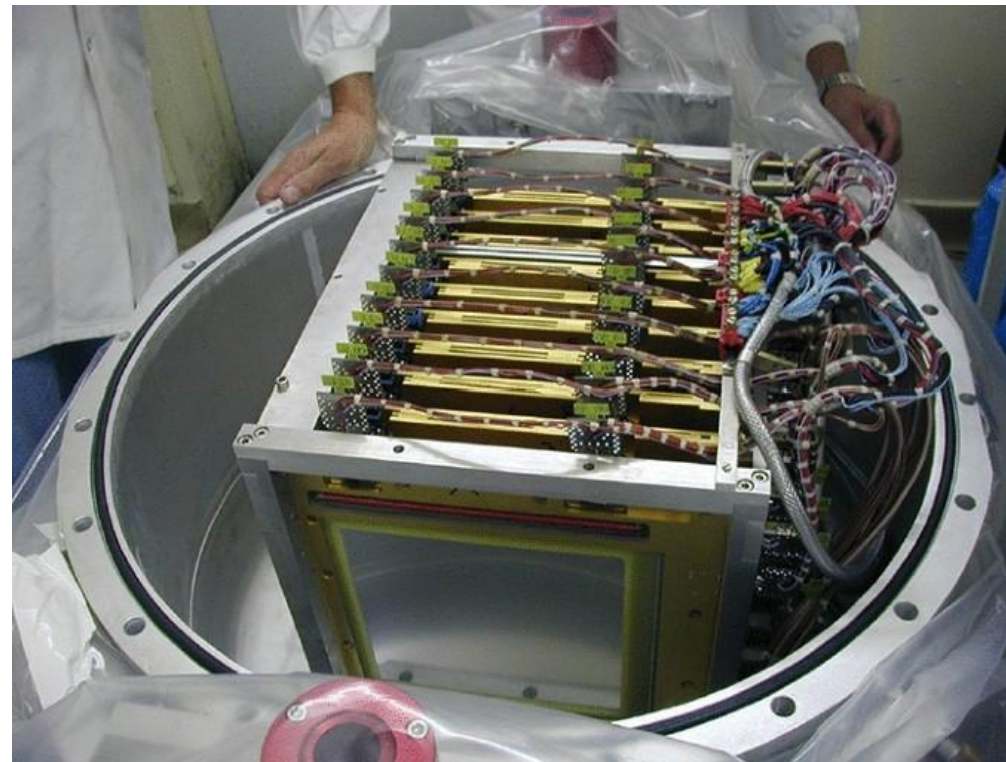
Incident neutrons

(n,f) measurements at n_TOF → cross section ratios.

$$Y_{(n,\gamma)}(E_n) = \frac{C(E_n) - B(E_n)}{\varepsilon(E_n) \cdot \phi(E_n)}$$



$^{234}\text{U}(n,f)$ and $^{237}\text{Np}(n,f)$ with PPAC:
C. Paradela et al. Phys. Rev. C 82,
034601(2010).

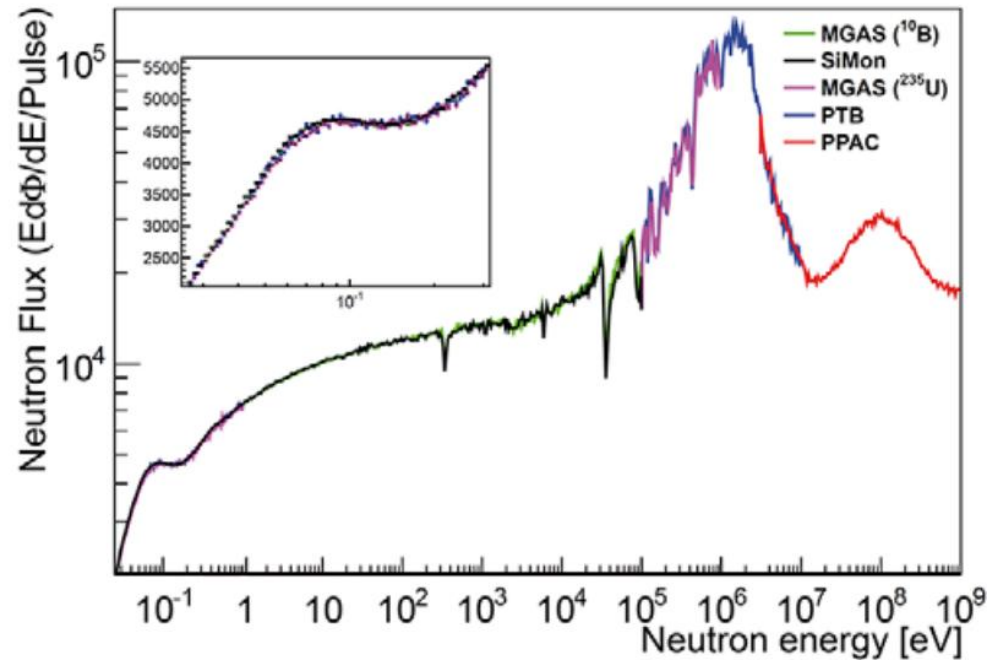


Incident neutrons

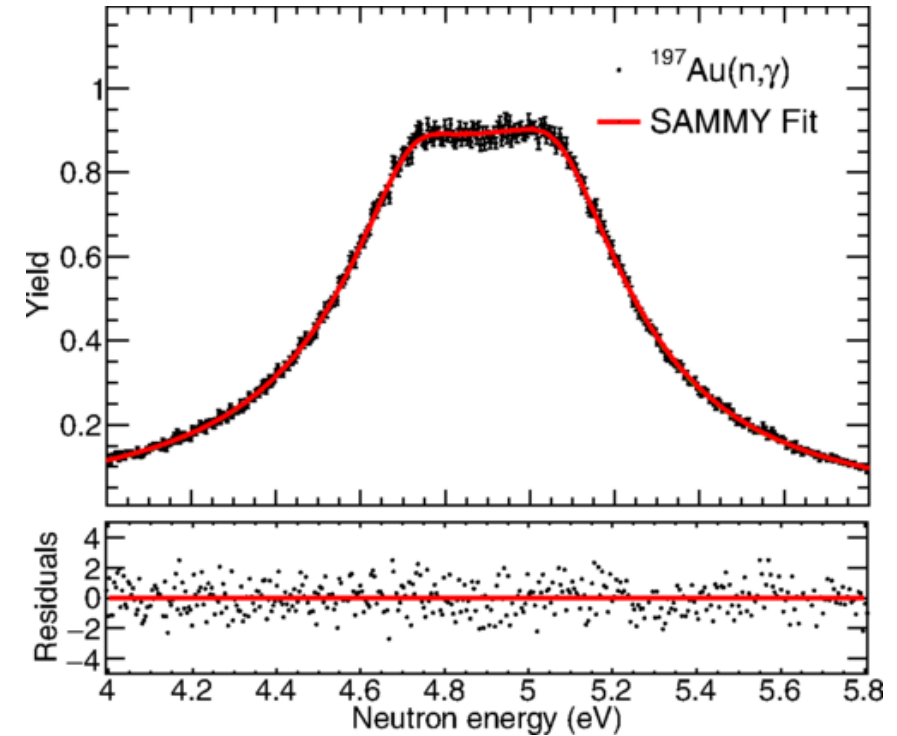
(n,γ) measurements at n_TOF → evaluated neutron fluence/flux.

Detectors based on ${}^6\text{Li}(n,t)$, ${}^{10}\text{B}(n,\alpha)$ and ${}^{235}\text{U}(n,f)$ reactions.

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+



M. Barbagallo et al., Eur. Phys. Jour. A
49, 156 (2013).

${}^{242}\text{Pu}(n,\gamma)$ with C_6D_6 : J. Lerendegui et al.
Phys. Rev. C 97, 024605 (2018).

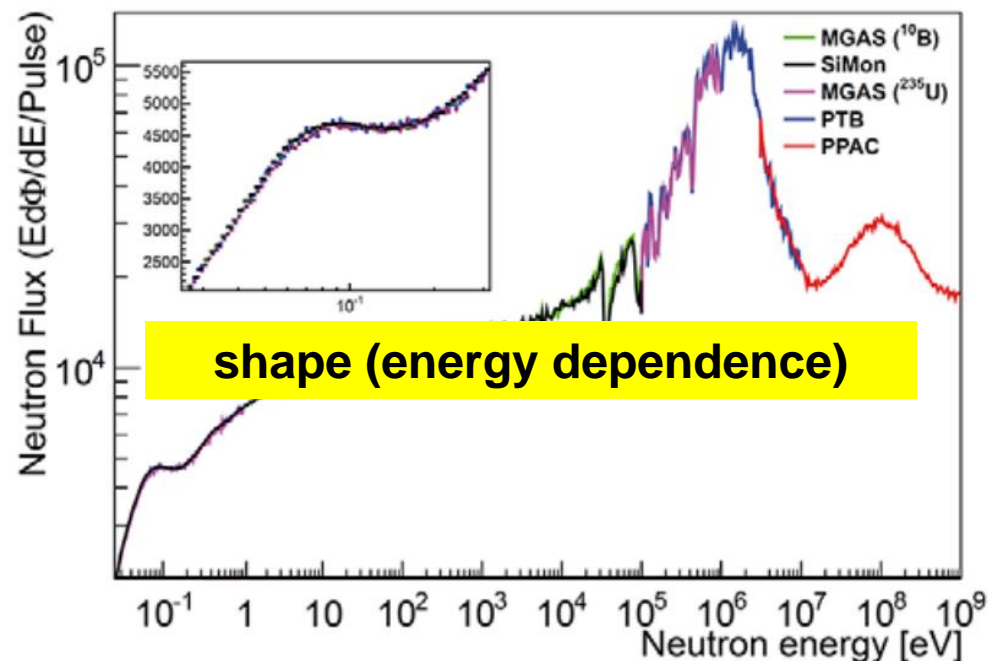


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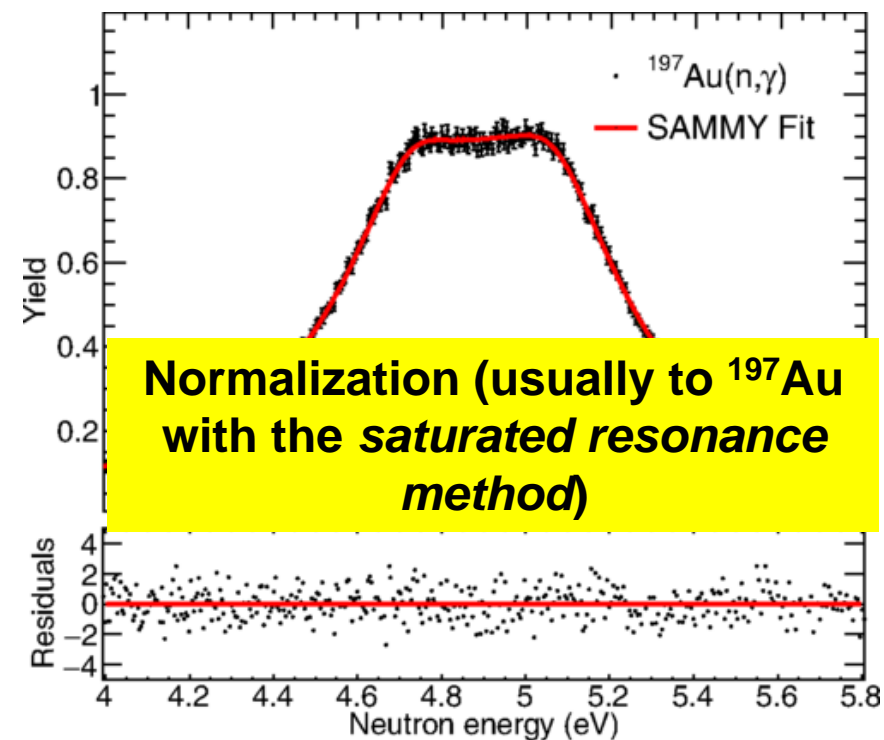
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Funded by
the European Union

Views and opinions expressed are however those of the author(s) only and do not necessarily reflect those of the European union or the European Climate, Infrastructure and Environment Executive Agency (CINEA). Neither the European Union nor the granting authority can be held responsible for them.

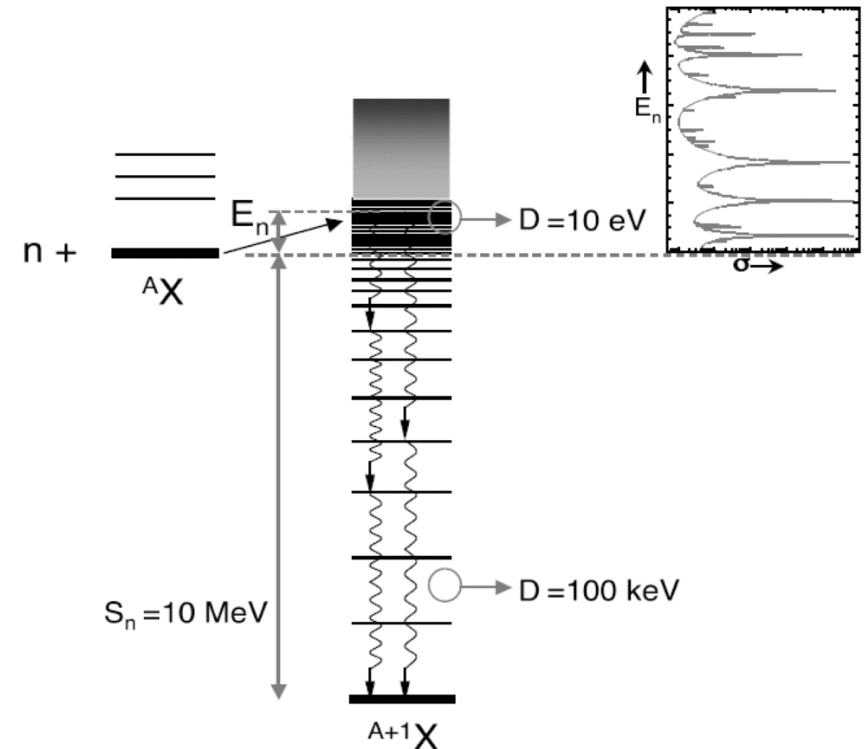
Detection efficiency

The (n,γ) detection efficiency may vary from one resonance to other.

To avoid this:

- **C₆D₆**: Total Energy Detection technique + Pulse Height Weighting Technique.
 - M. Moxon et al., Nucl. Instrum. Methods 24, 445 (1963).
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 - E. Mendoza et al., Nucl. Instrum. Methods A 1047, 167894 (2023).
- Idea: replace $\varepsilon(E_n)$ by other variable (estimator) which does not depend (i.e. has a small dependency) on the (n,γ) cascade.
- **TAC**: measurements of heavy nuclei, with large level densities
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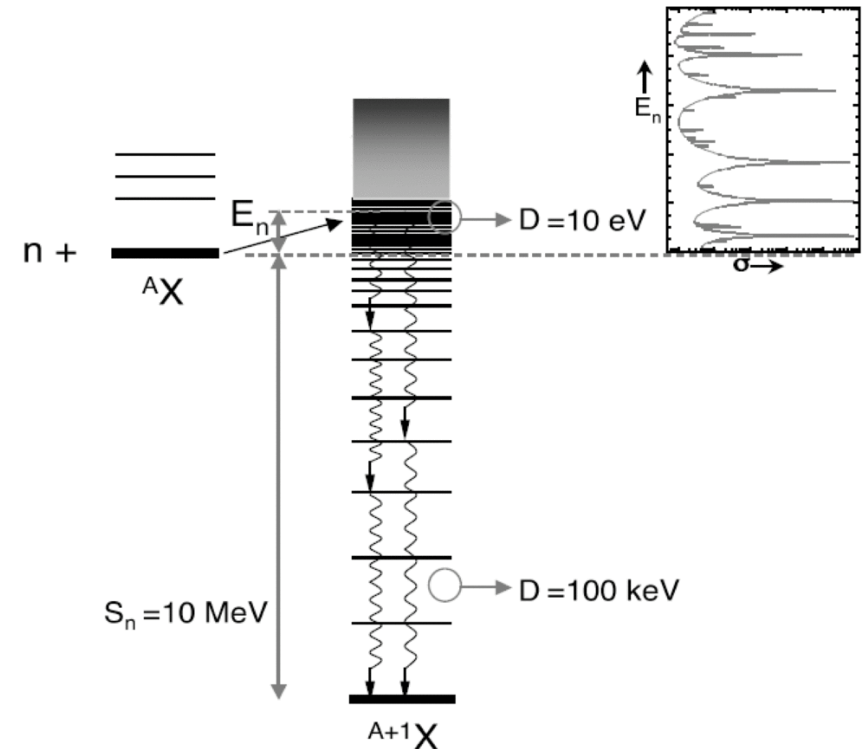
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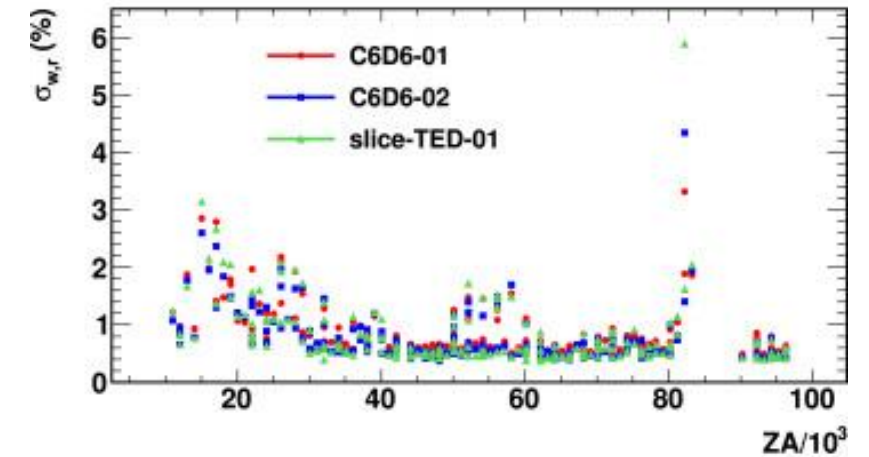
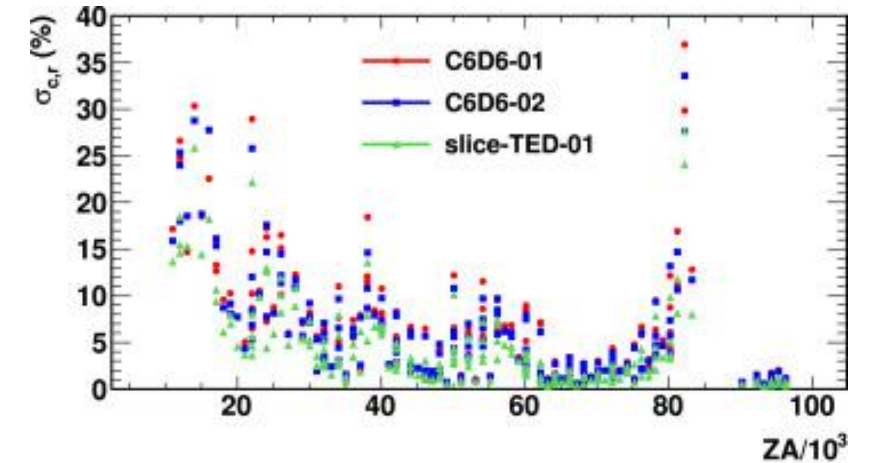


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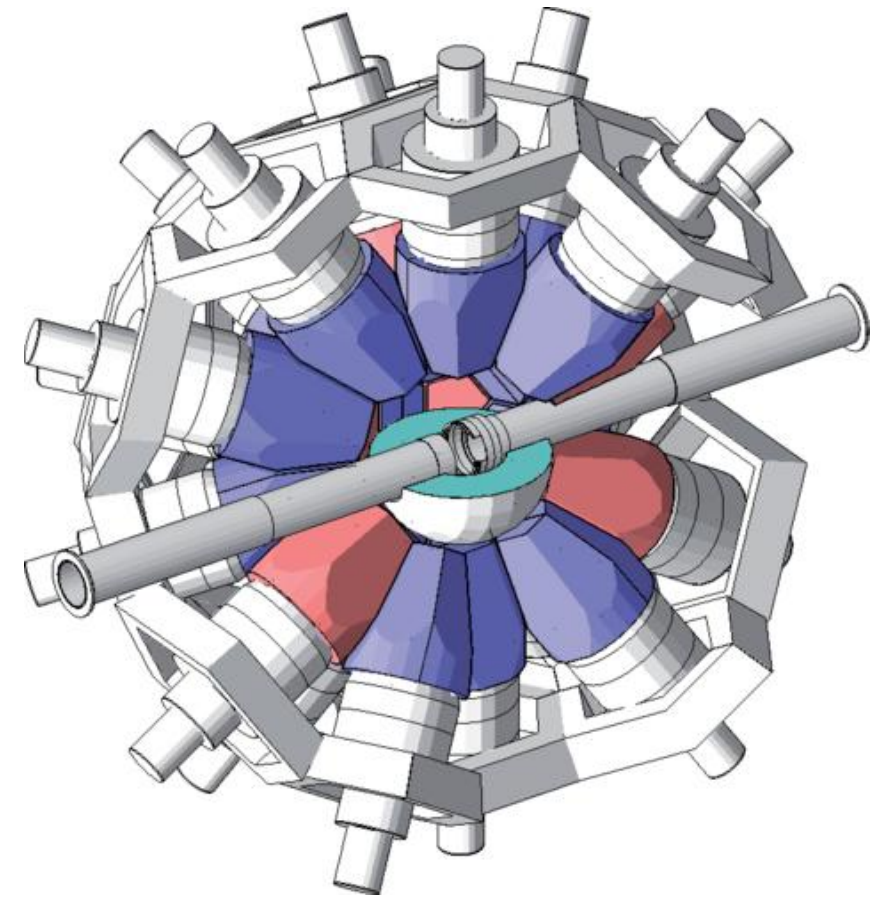


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²⁴¹Am(n, γ) with the TAC: E. Mendoza et al. Phys. Rev. C 97, 054616 (2018).

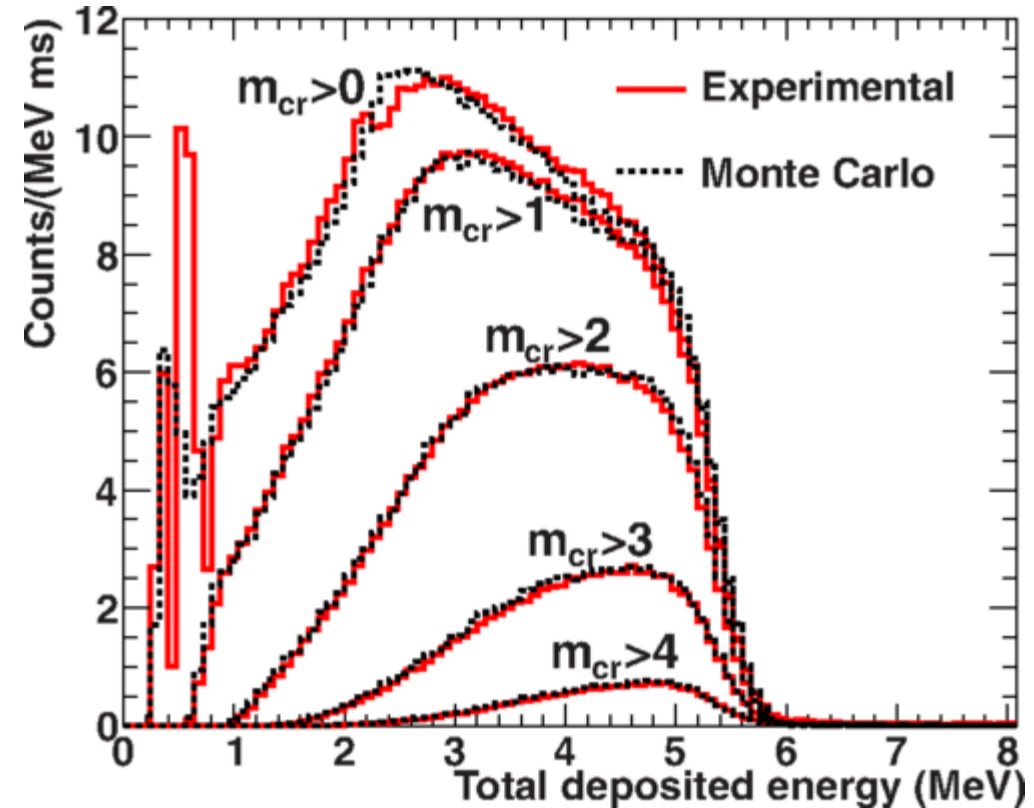


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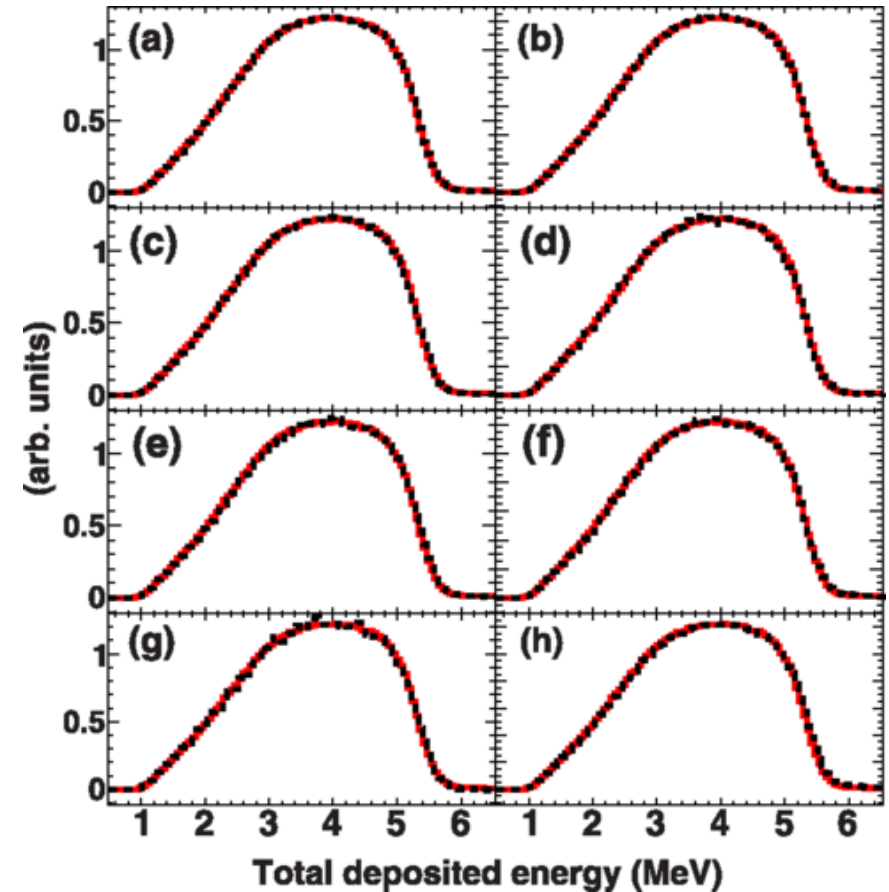


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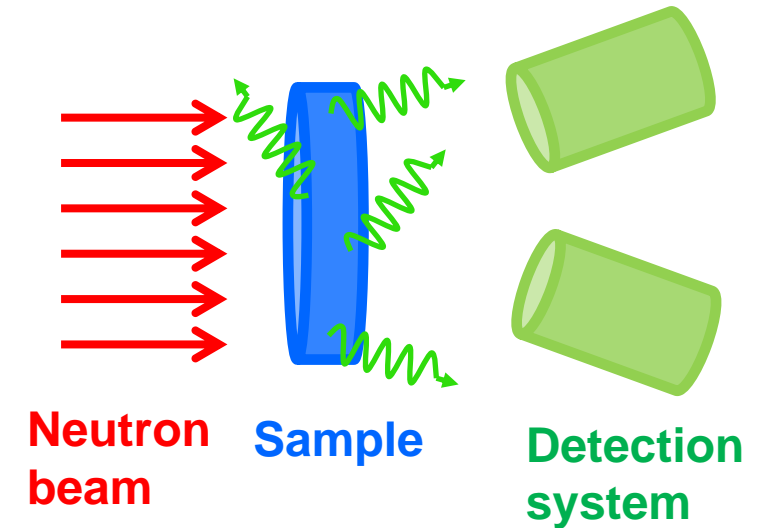
Neutron capture measurements

The experimental capture yield is usually calculated from an expression similar to this:

$$Y_{(n,\gamma)}(E_n) = \frac{C(E_n) - B(E_n)}{\varepsilon \cdot \phi(E_n)}$$

where:

- $Y_{(n,\gamma)}$ is the capture yield.
- E_n is the neutron energy.
- $C(E_n)$ are the number of counts in the detector [per unit something].
- $B(E_n)$ is the background.
- $\varepsilon(E_n)$ is the detection efficiency.
- $\phi(E_n)$ are the number of neutrons impinging in the sample [per unit something].



Finally, I will talk about two issues that do not appear explicitly in the previous equation:

- Resolution function
- Inhomogeneities

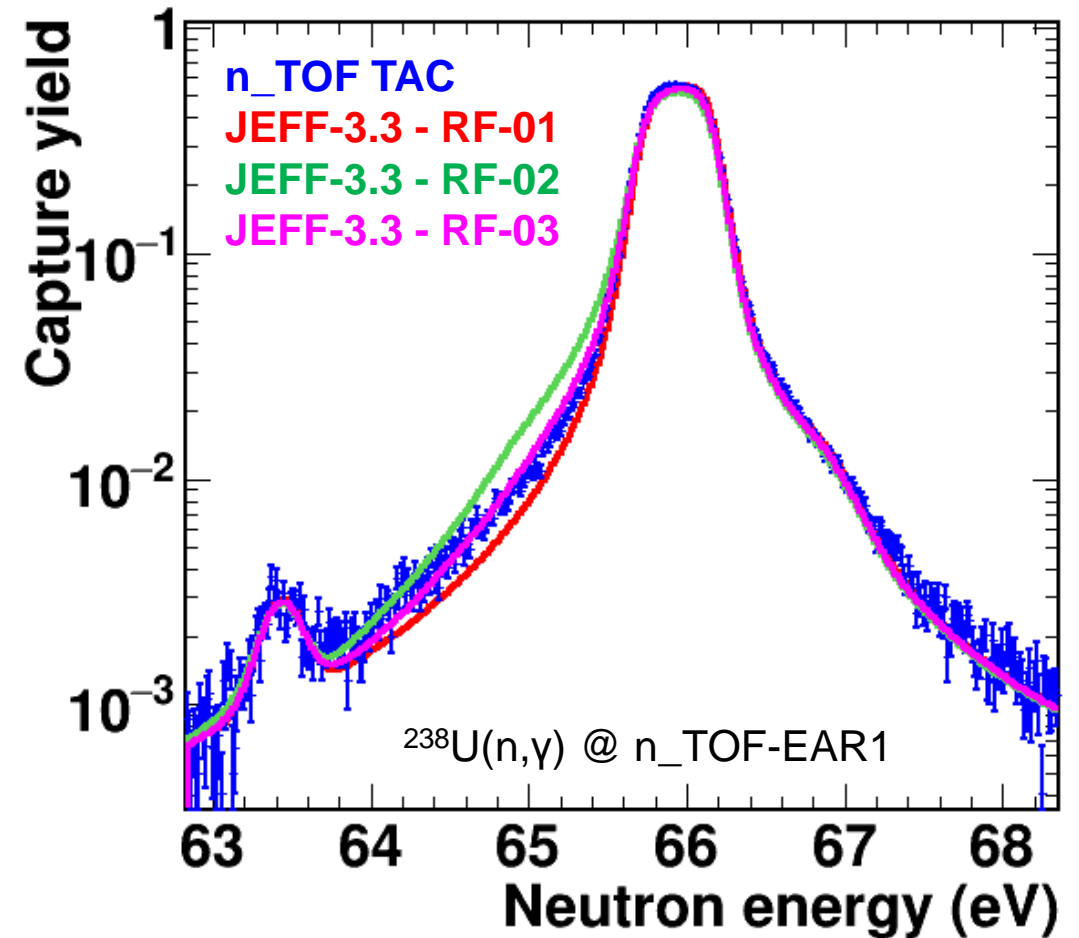
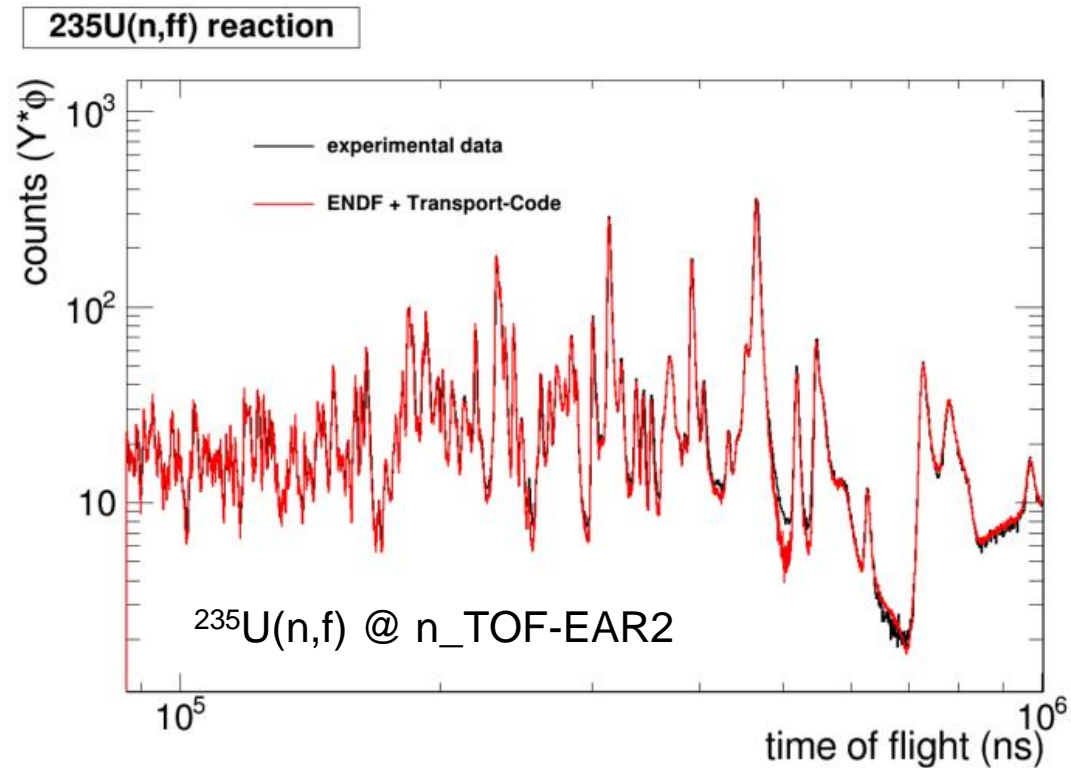


Resolution function

The n_TOF resolution function (RF) is obtained from MC calculations

The RF is validated with experimental data.

But there are some limitations ...



Inhomogeneities

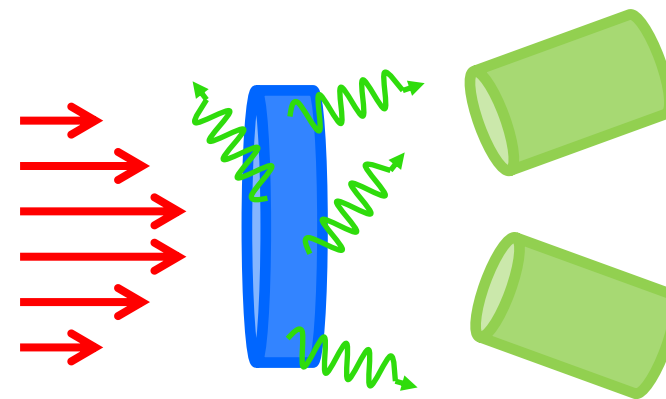
Sample inhomogeneities may appear when measuring oxides.

Sample inhomogeneities may have a sizeable impact in the measured yield.

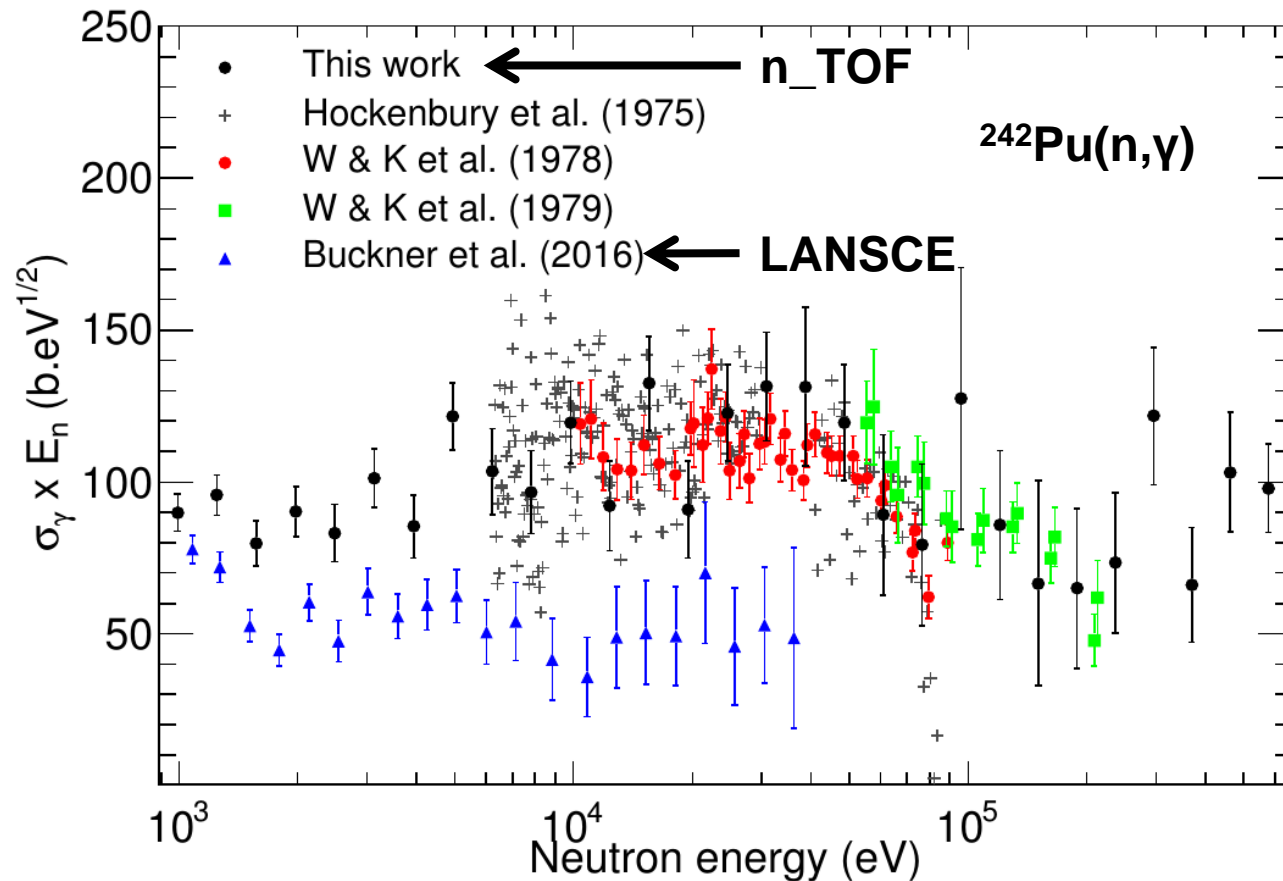
We may have two types of inhomogeneities:

- Local inhomogeneities: impact in the self shielding → affects the shape of the strongest resonances.
- Global inhomogeneities: affects the normalization.

Inhomogeneities do not affect the shape of the yield (or transmission) if the sample is thin (i.e. no self shielding).



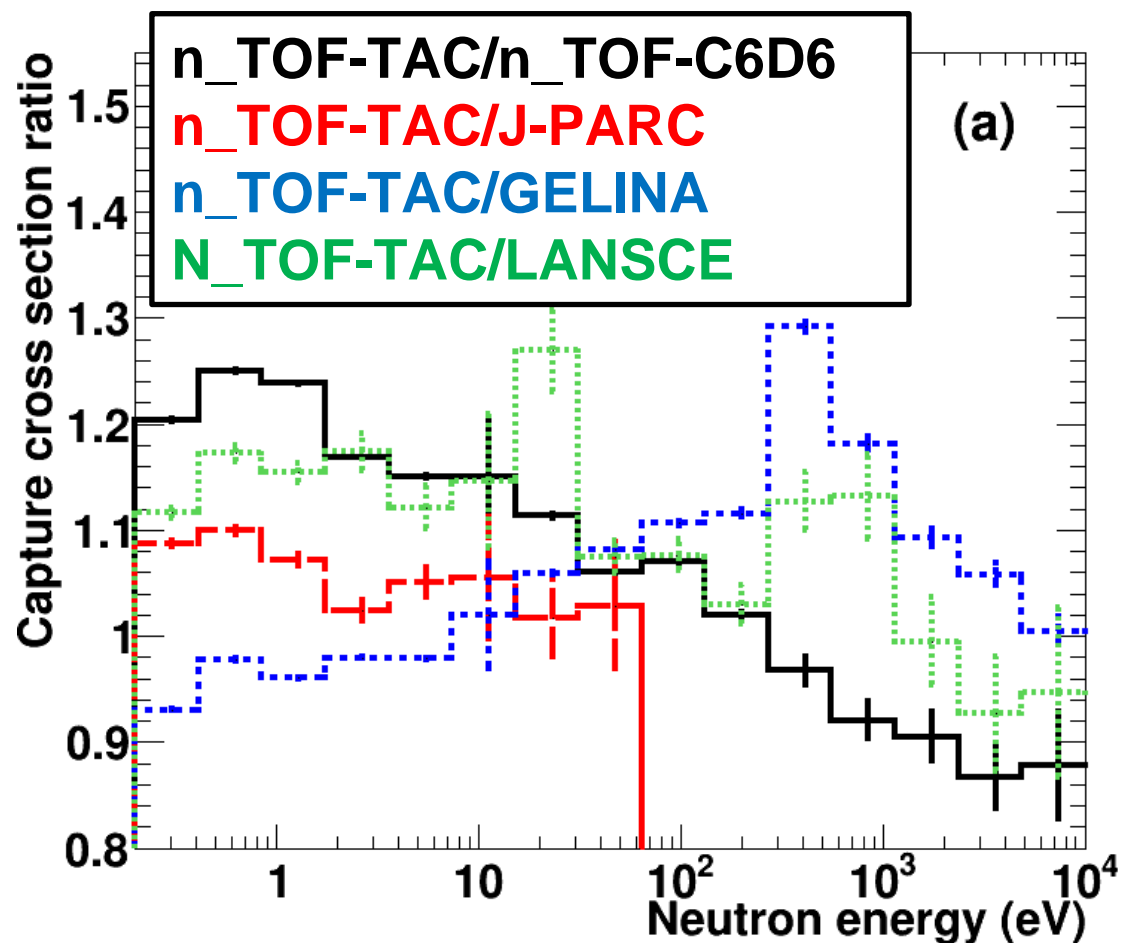
Discrepancies between measurements



J. Lerendegui-Marco, PhD Thesis
<https://cds.cern.ch/record/2661485>



Discrepancies between measurements



Conclusions

I have presented an overview of an analysis of a (neutron capture) measurement in n_TOF.

Understanding the entire analysis process may help to understand different biases, what data may be more questionable, etc.

In case something is not clear in the EXFOR dataset (+ paper), or you think there is some problem with the data ... please do not hesitate in contacting us for clarification. This will also help us to detect possible issues in our data.

