# Parity-odd Neutron Spin Rotation 

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Physics of Fundamental Symmetries and Interactions
PSI, September 10, 2013 National Institute of Standards and Technology $\square$

| $\overline{\overline{\text { THE GEORGE }}}$ |
| :--- |
| WASHINGTON |
| UNIVERSITY |



## Why Hadronic Weak Interactions?

We know there is a quark-quark Weak interaction:

$$
\begin{aligned}
& A_{\gamma}\left({ }^{180} \mathrm{Hf}^{*} \rightarrow{ }^{180} \mathrm{Hf}+\gamma\right)=-(1.66 \pm 0.18) \times 10^{-2} \\
& A_{h}\left(\vec{n}+{ }^{139} \mathrm{La}\right)=(9.55 \pm 0.35) \times 10^{-2}
\end{aligned}
$$

Enhanced PV signal from compound nuclear resonances
...but we don't fully understand observations because the low-energy (nonperturbative) limit of QCD controls how the HWI manifests itself. Need to study simpler systems:

Non-leptonic Flavor-changing
meson decays: EFT is effective, but, e.g., $\Delta \mathrm{I}=1 / 2$ rule not fully understood

$$
\mathrm{K} \rightarrow \pi \pi
$$

baryon decays: $\Delta S=1$ hyperon decay properties don't follow from QCD symmetries
Evidence of non-trivial QCD ground state dynamical phenomena? $\quad \Lambda \rightarrow N \pi$
Flavor-conserving NN Interactions $\frac{\mathcal{M}_{\text {weak }}}{\mathcal{M}_{\text {strong }}} \sim \frac{e^{2} / M_{\mathrm{W}}^{2}}{g^{2} / M_{\pi}^{2}} \sim 10^{-7}$ Hallmark is use of PV to single out weak part of interaction

- $\Delta S=0$ : is strange quark or light quark dynamics in general responsible for hyperon decay?
- Sensitive to NC effects at low energy, probes the QCD ground state
- NN and few-body systems


## SM Framework for HWI

Current-current form at low energy ( $\Lambda_{\mathrm{QCD}}<\mathrm{q}<v$ )

$$
\begin{aligned}
& \mathcal{H}_{\text {weak }}=\frac{G_{\mathrm{F}}}{\sqrt{2}}\left(J_{\mu}^{(\mathrm{C}) \dagger} J^{(\mathrm{C}) \mu}+\frac{1}{2} J_{\mu}^{(\mathrm{N}) \dagger} J^{(\mathrm{N}) \mu}\right)+\text { h.c. } \\
& J_{\mu}^{(\mathrm{C})}=\mathcal{U}_{u d} \bar{u} \gamma_{\mu}\left(1+\gamma_{5}\right) d+\mathcal{U}_{u s} \bar{u} \gamma_{\mu}\left(1+\gamma_{5}\right) s \\
& J_{\mu}^{(\mathrm{N})}=\bar{u} \gamma_{\mu}\left(1+\gamma_{5}\right) u-\bar{d} \gamma_{\mu}\left(1+\gamma_{5}\right) d-\bar{s} \gamma_{\mu}\left(1+\gamma_{5}\right) s-4 \sin ^{2} \theta_{\mathrm{W}} J_{\mu}^{\mathrm{em}} \\
& v \sim 246 \mathrm{GeV} \\
& \Lambda_{\mathrm{HAD}} \sim 1 \mathrm{GeV} \\
& \Lambda_{\mathrm{QCD}} \sim 200 \mathrm{MeV} \\
& m_{\pi} \sim 140 \mathrm{MeV} \\
& \Delta I=2 \longrightarrow J_{I=1}^{(\mathrm{C})} J_{I=1}^{(\mathrm{C})} \\
& \Delta I=1 \longrightarrow J_{I=1 / 2}^{(\mathrm{C})} J_{I=1 / 2}^{(\mathrm{C})}+J_{I=1}^{(\mathrm{N})} J_{I=1}^{(\mathrm{N})} \\
& \Delta I=0 \longrightarrow J_{I=0}^{(\mathrm{C})} J_{I=0}^{(\mathrm{C})}+J_{I=0}^{(\mathrm{N})} J_{I=0}^{(\mathrm{N})}+J_{I=1}^{(\mathrm{N})} J_{I=1}^{(\mathrm{N})}
\end{aligned}
$$

Calculation of the NN Weak Interaction ( $q<\Lambda_{\mathrm{QCD}}$ )
DDH: Model dependent, seven (non-elementary) meson-nucleon couplings: $\quad h_{\pi}^{1}, h_{\rho}^{0,1,2}, h_{\rho}^{1^{\prime}}, h_{\omega}^{0,1}$
EFT: Model-independent, QCD symmetries built in, perturbative expansion in $\mathrm{q} / \Lambda$, pionful ( $\mathrm{q} \sim 140 \mathrm{MeV}$ ) theory has six parameters

$$
\lambda_{s}^{0,1,2}, \lambda_{t}, \rho_{t}, \widetilde{C}_{6}^{\pi} \propto h_{\pi}^{1}
$$

Lattice QCD: Gold standard

$$
\begin{aligned}
& (I=1)^{+} \otimes L^{+} \otimes(S=0)^{-} \rightarrow{ }^{1} S_{0}, \ldots \\
& (I=1)^{+} \otimes L^{-} \otimes(S=1)^{+} \rightarrow{ }^{3} P_{0,1,2}, \ldots \\
& (I=0)^{-} \otimes L^{-} \otimes(S=0)^{-} \rightarrow{ }^{1} P_{1}, \ldots \\
& (I=0)^{-} \otimes L^{+} \otimes(S=1)^{+} \rightarrow{ }^{3} S_{1}, \ldots
\end{aligned}
$$

## PV Observables in terms of DDH Parameters

|  | $h_{\pi}^{1}$ | $h_{\rho}^{0}$ | $h_{\rho}^{1}$ | $h_{\rho}^{2}$ | $h_{\omega}^{0}$ | $h_{\omega}^{1}$ | Status |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\begin{gathered} \phi_{\mathrm{PV}}, n^{4} \mathrm{He} \\ \sigma_{n} \cdot k_{n} \end{gathered}$ | -0.97 | -0.32 | 0.11 | 0 | -0.22 | 0.22 | will run at NIST |
| $\begin{gathered} A_{\gamma}, n p d \gamma \\ \sigma_{n} \cdot k_{\gamma} \end{gathered}$ | -0.107 | 0 | -0.001 | 0 | 0 | 0.004 | running at SNS |
| $\begin{gathered} A_{p}, n^{3} H e \\ \sigma_{n} \cdot k_{p} \end{gathered}$ | -0.189 | -0.036 | 0.019 | -0.001 | -0.033 | 0.041 | proposed at SNS |
| $A_{L}, p \alpha$ $\sigma_{p} \cdot k_{p}$ | -0.333 | 0.140 | 0.047 | 0 | 0.059 | 0.059 | $\begin{gathered} 1982- \\ 1985 \end{gathered}$ |
| $\begin{gathered} A_{L}, p p \\ \sigma_{p} \cdot k_{p} \\ (45 \mathrm{MeV}) \end{gathered}$ | 0 | 0.074 | 0.074 | 0.030 | 0.067 | 0.067 | $\begin{aligned} & 1991- \\ & 2003 \end{aligned}$ |
| $P_{\gamma},{ }^{18} F$ | 4385 | 0 | 0 | 0 | 0 | 0 | $\begin{gathered} 1978- \\ 1987 \end{gathered}$ |

- Recent experiments focus on simple NN and few-nucleon systems
- Many-body systems are also beginning to contribute via measurements of the nuclear anapole moment


## The Current Experímental Landscape, DDH-Style

Weak NN iso-scalar, iso-vector coupling subspace


## PV Spin Rotation in ${ }^{4} \mathrm{He}$


4K LHe
$f(0)=f_{\mathrm{PC}}+f_{\mathrm{PV}}\left(\vec{\sigma}_{n} \cdot \vec{k}_{n}\right)$
$\left\langle\vec{S}_{4}{ }_{\mathrm{He}}\right\rangle=0$


$$
\psi_{n} \sim \exp \left\{i\left[1+\frac{2 \pi}{k^{2}} \rho f(0)\right] \vec{k}_{0} \cdot \vec{z}\right\}
$$

Passage through the ${ }^{4} \mathrm{He}$ optical potential results in a helicity-dependent total phase
which causes rotation of the neutron polarization

$$
\phi_{ \pm}=\underbrace{k z\left(1+\frac{2 \pi \rho}{k^{2}} f_{\mathrm{PC}}\right)}_{\varphi_{\mathrm{PC}}} \pm \underbrace{2 \pi \rho z f_{\mathrm{PV}}}_{\varphi_{\mathrm{PV}}}
$$

$$
|\uparrow\rangle=\frac{1}{\sqrt{2}}[|+\rangle+|-\rangle] \longrightarrow \frac{1}{\sqrt{2}}\left[e^{i\left(\varphi_{\mathrm{PC}}+\varphi_{\mathrm{PV}}\right)}|+\rangle+e^{i\left(\varphi_{\mathrm{PC}}-\varphi_{\mathrm{PV}}\right)}|-\rangle\right]
$$

$$
\phi_{\mathrm{PV}}=\varphi_{+}-\varphi_{-}=2 \varphi_{\mathrm{PV}}=4 \pi \rho z f_{\mathrm{PV}}
$$

DDH prediction is $[-16,12] \times 10^{-7} \mathrm{rad} / \mathrm{m}$

## PV Spin Rotation Measurement Technique



Two measurements produce a spin rotation angle $P A \sin \phi=\frac{I_{+}-I_{-}}{I_{+}+I_{-}}$ PV rotation $\sim 10^{-7} \mathrm{rad} / \mathrm{m}$, but earth's field would produce a rotation $\sim 10 \mathrm{rad} / \mathrm{m}$ !

## PV Spin Rotation Measurement Technique



## The Neutron Spin Rotation Experiment




## NSR-III: NIST NG~C Beam Líne

- NG-C beam line completed
- NG-C end station ready early 2014
- SM polarizer and analyzer arriving Oct.



## The Neutron Spin Rotation Experiment




## NSR-III Neutron Guides

- $10 \mathrm{~cm} \times 10 \mathrm{~cm}, 1.25$ and 2.0 nonmagnetic supermirror ( NiMo / T) neutron guides
- 1 cm thick borofloat glass
- $\mathrm{M}=2.0, \mathrm{R}>90 \%$
- depolarization probability per bounce $<1 \%$



## The Neutron Spin Rotation Experiment




## NSR-III Ionization Chamber

- New ionization chamber constructed and will be tested at Indiana University in October
- Based on the radially and longitudinally segmented detector used for NSR-II



## The Neutron Spin Rotation Experiment




## NSR-II

$$
\frac{\mathrm{d} \phi_{\mathrm{PV}}}{\mathrm{~d} z}=[1.7 \pm 9.1(\text { stat }) \pm 1.4(\text { sys })] \times 10^{-7} \mathrm{rad} / \mathrm{m}
$$

2008 NIST Fundamental
Physics Beam Line NG-6

C. D. Bass et al., Nucl. Instrum. Meth. A612, 69-82 (2009).
A. M. Micherdzinska et al., Nucl. Instrum. Meth. A631, 80 (2011). W. M. Snow et al., Phys. Rev. C83, 022501(R) (2011).

## NSR-II

$\frac{\mathrm{d} \phi_{\mathrm{PV}}}{\mathrm{d} z}=[1.7 \pm 9.1($ stat $) \pm 1.4($ sys $)] \times 10^{-7} \mathrm{rad} / \mathrm{m}$

NSR-II

| Source | Uncertainty <br> $(\mathrm{rad} / \mathrm{m})$ | Method | Uncertainty <br> $(\mathrm{rad} / \mathrm{m})$ |
| :--- | :---: | :--- | :--- |
| L 4He diamagnetism | $2 \times 10^{-9}$ | calc | $2 \times 10^{-10}$ |
| L 4He optical potential | $3 \times 10^{-9}$ | calc | $3 \times 10^{-10}$ |
| Neutron E spectrum shift | $8 \times 10^{-9}$ | calc | $8 \times 10^{-10}$ |
| Refraction/reflection | $3 \times 10^{-10}$ | calc | $3 \times 10^{-11}$ |
| Non-forward scattering | $2 \times 10^{-8}$ | calc | $2 \times 10^{-9}$ |
| Polarimeter non-uniformity | $1 \times 10^{-8}$ | meas | $<1 \times 10^{-8}$ |
|  |  |  |  |
| B amplification | $<4 \times 10^{-8}$ | meas | $<4 \times 10^{-9}$ |
| B gradient amplification | $<3 \times 10^{-8}$ | meas | $<3 \times 10^{-9}$ |
| PA/target nonuniformity | $<6 \times 10^{-8}$ | meas | $<6 \times 10^{-8}$ |
| Total | $1.4 \times 10^{-7}$ |  | $\leq 1 \times 10^{-7}$ |

$x 10$ in neutron flux
x 40 effective increase in polarized cold neutron flux from enhanced neutron transport

## NSR-III Magnetic Shieldíng

- Three nested layers of $\mu$ metal shielding (two external and one internal); demonstrated to reduce B-field in target region to $<10 \mu \mathrm{G}$, a factor of ten improvement from NSR-II.
- In-situ de-gaussing, internal fluxgate magnetometers, trim coils, and active field cancelation supplement the passive shields.





## NSR-II

$\frac{\mathrm{d} \phi_{\mathrm{PV}}}{\mathrm{d} z}=[1.7 \pm 9.1($ stat $) \pm 1.4($ sys $)] \times 10^{-7} \mathrm{rad} / \mathrm{m}$

NSR-II

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## NSR-III Cryogenics

- Improved cryogenic design for reduced heat load, simpler assembly / disassembly, and more robust operation
- He re-liquefier removes necessity of LHe fills
- R\&D on new LHe pump to reduce target change time



## NSR-III Cryogenics

## Re-Liquefier

## R\&D on New LHe Pump

- Cryomech pulse tube re-liquefier tested for three months of continuous operation
- Observed a liquefaction rate from warm gas of 12L / day
- Demonstrated ability to continuously maintain LHe volume comparable to spin rotation cryostat with little intervention under a heat load of $\sim 600 \mathrm{~mW}$

- Bellows pump demonstrated equivalent of 100 days of data-taking in water
- Pumping rate in water is ten times larger than previous centrifugal pump
- Cryogenics tests are underway



## Additional Physics with our Apparatus

There is vigorous interest in placing constraints on spin- and momentum- dependent mesoscopic exotic forces:

A New Limit on Possible Long-Range Parity-odd Interactions of the Neutron from Neutron Spin Rotation in Liquid ${ }^{4} \mathrm{He}$
H.Yan and W. M. Snow*

Indiana University, Bloomington, Indiana 47408, USA and
Center for Exploration of Energy and Matter, Indiana University, Bloomington, IN 47408
(Dated: November 23, 2012)
Phys. Rev. Lett. 110, 082003 (2013)


Improved Limits on Long-Range Parity-Odd Interactions of the Neutron
E. G. Adelberger and T. A. Wagner

Center for Experimental Nuclear Physics and Astrophysics, Box 354290, University of Washington, Seattle, WA 98195-4290 (Dated: July 26, 2013)
arXiv:1307.6602 [hep-ex]


Such interactions can be generated by a new light vector boson $X_{\mu}$ coupling to a fermion

$$
\mathcal{L}=\bar{\psi}\left(g_{V} \gamma^{\mu}+g_{A} \gamma^{\mu} \gamma_{5}\right) \psi X_{\mu}
$$

Examples of the interaction potentials generated by such a coupling include

$$
V \sim g_{V} g_{A} \vec{\sigma} \cdot \vec{v} \quad V \sim g_{A}^{2} \vec{\sigma} \cdot(\vec{v} \times \hat{r})
$$

## NSR Collaboration

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## NSR-III Input, Output, and Pí Coils

Pi-Coil design by Chris Crawford (University of Kentucky)


Input and output by Libertad
Barrón Palos (Universidad
Nacional Autónoma de México Instituto de Física)


