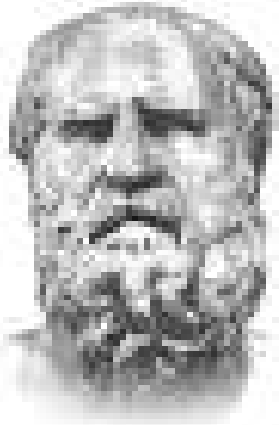


Probing the microstructural origin of complex flow behaviour with in situ Small Angle Neutron and X-ray Scattering

Pavlik Lettinga

What is Rheologie?



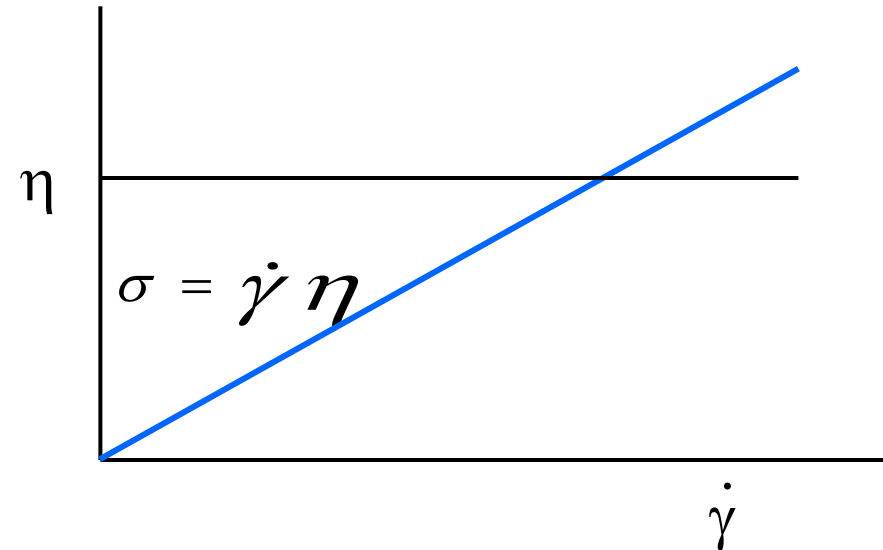
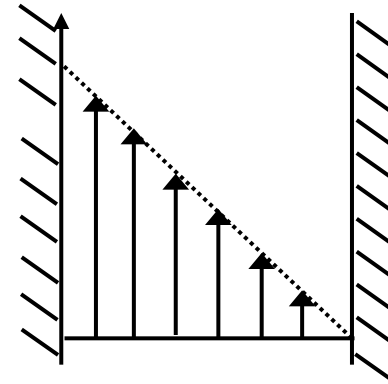
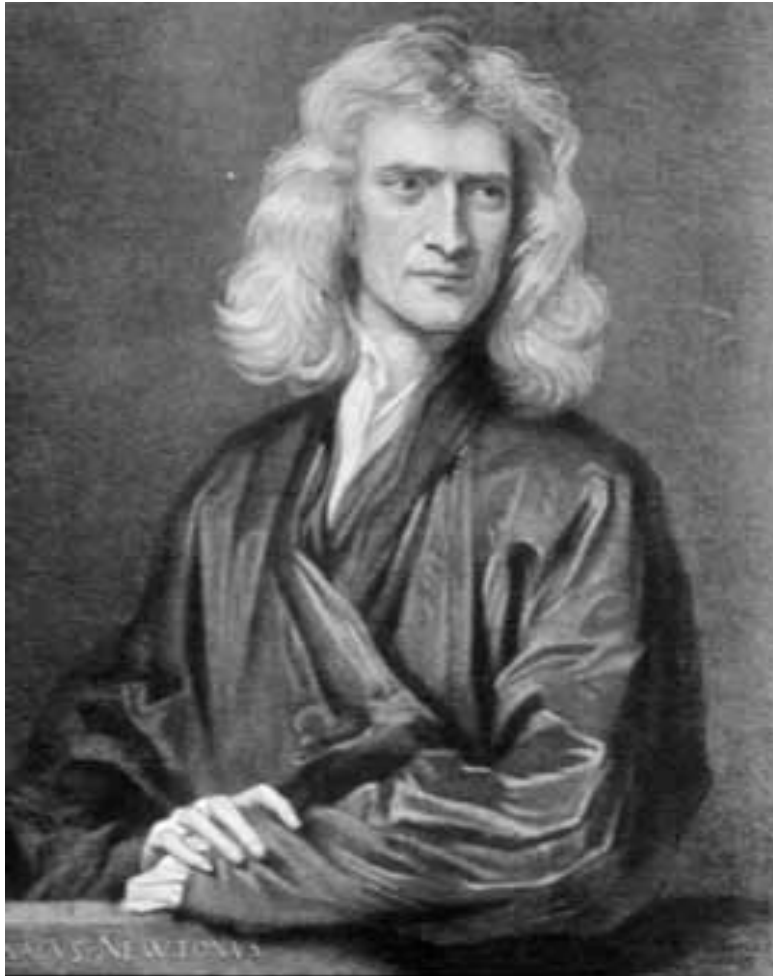
« Science of deformation and flow »

Heraklites : Παντα ρει (= everything flows)

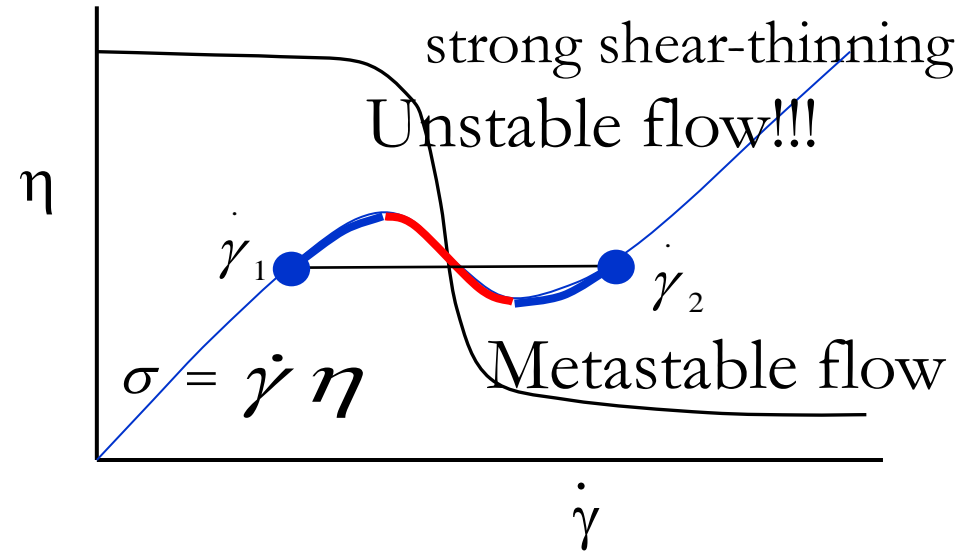
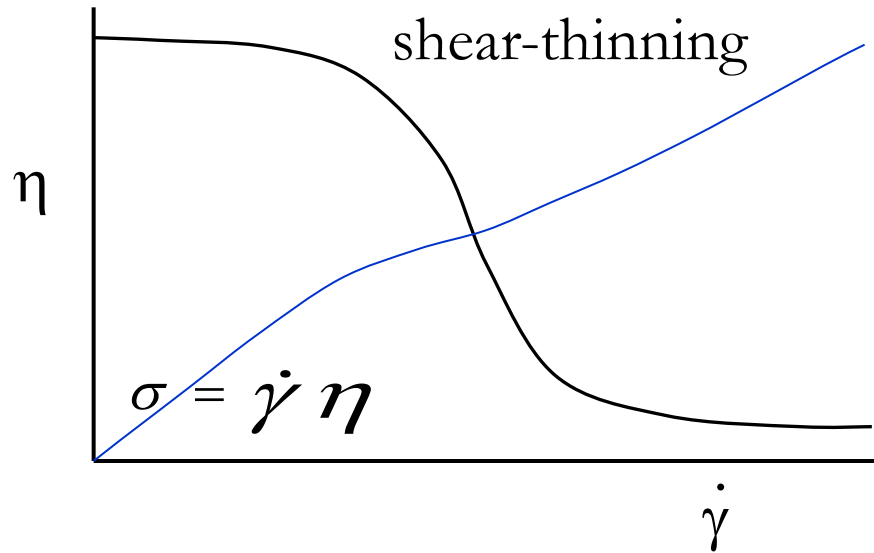
Deborah: "The mountains flowed before the Lord"

in a song by prophetess [Deborah](#) ([Judges](#) 5:5).

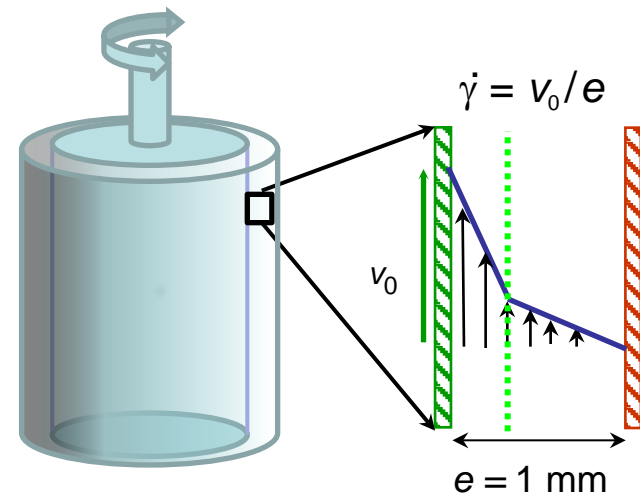
The ideal Newtonian world



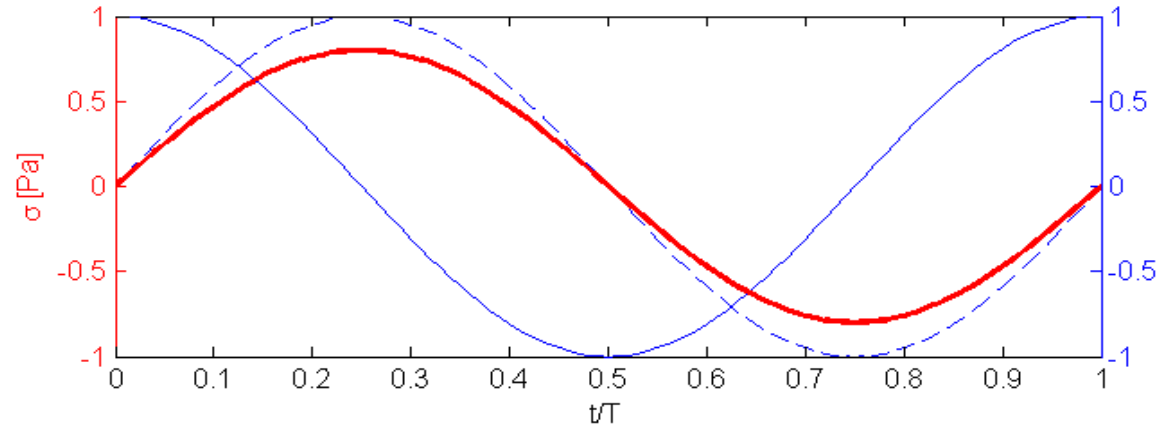
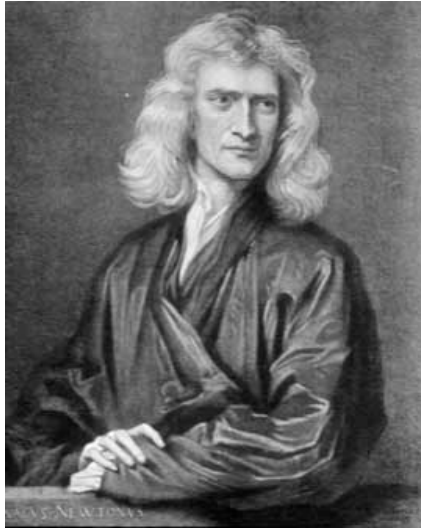
Our real world



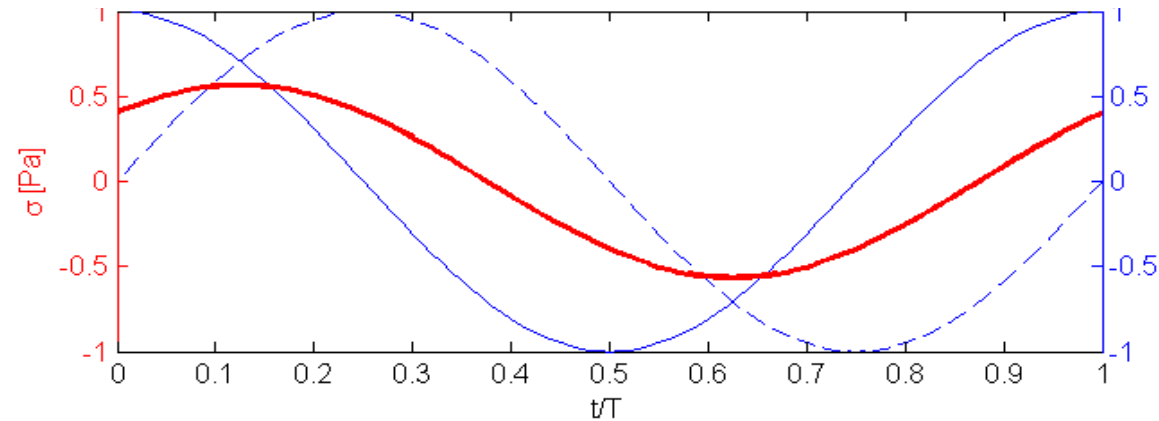
Flow instabilities:
gradient shear banding



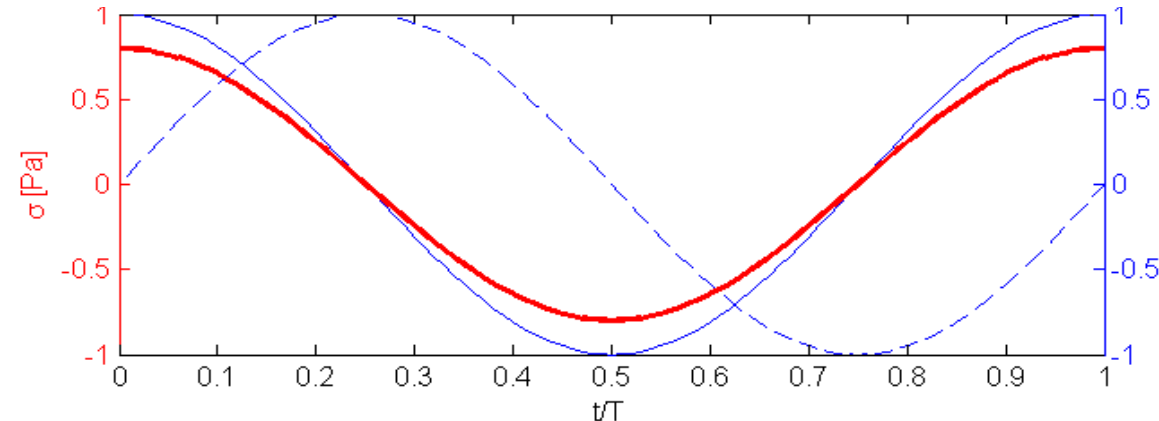
The ideal visco-elastic world: $\sigma = G' \sin(\omega t) + G'' \cos(\omega t)$



$G' \gg G''$



$G' \approx G''$



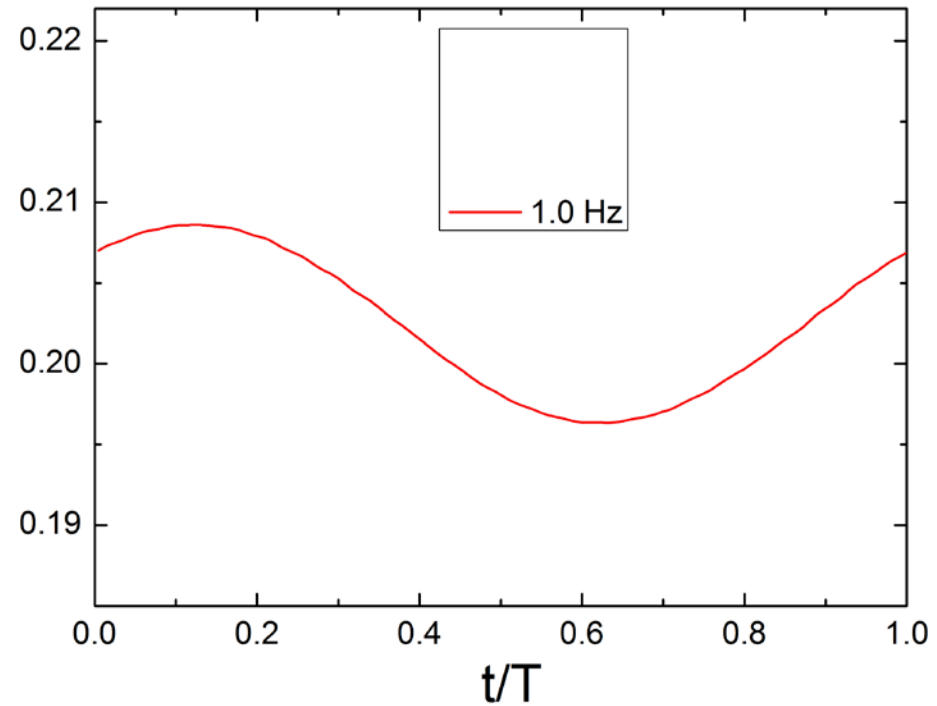
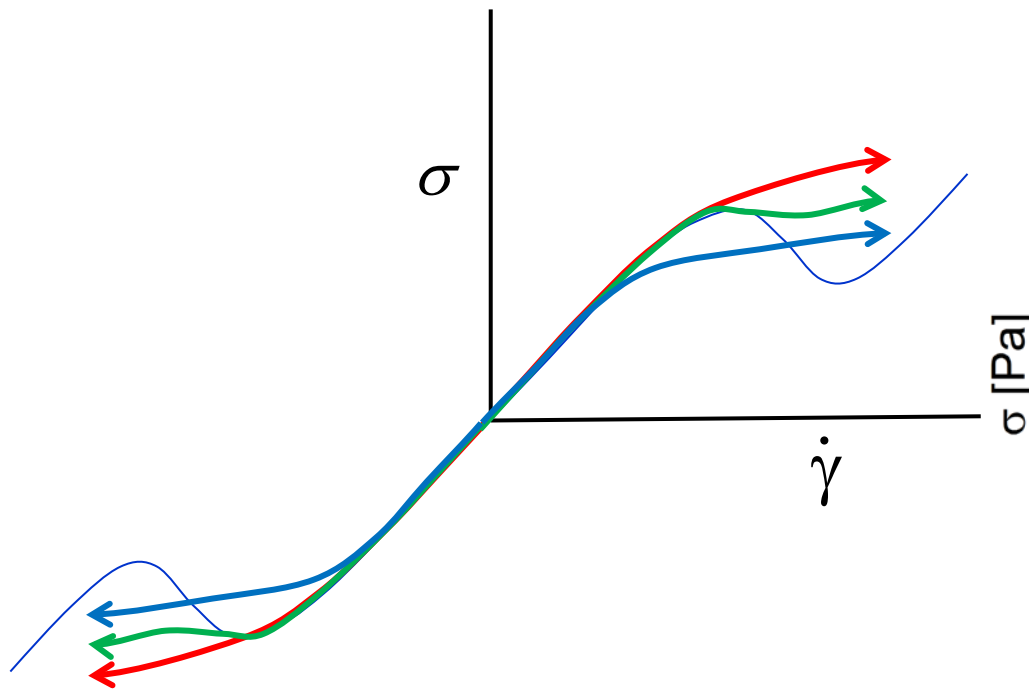
$G' \ll G''$



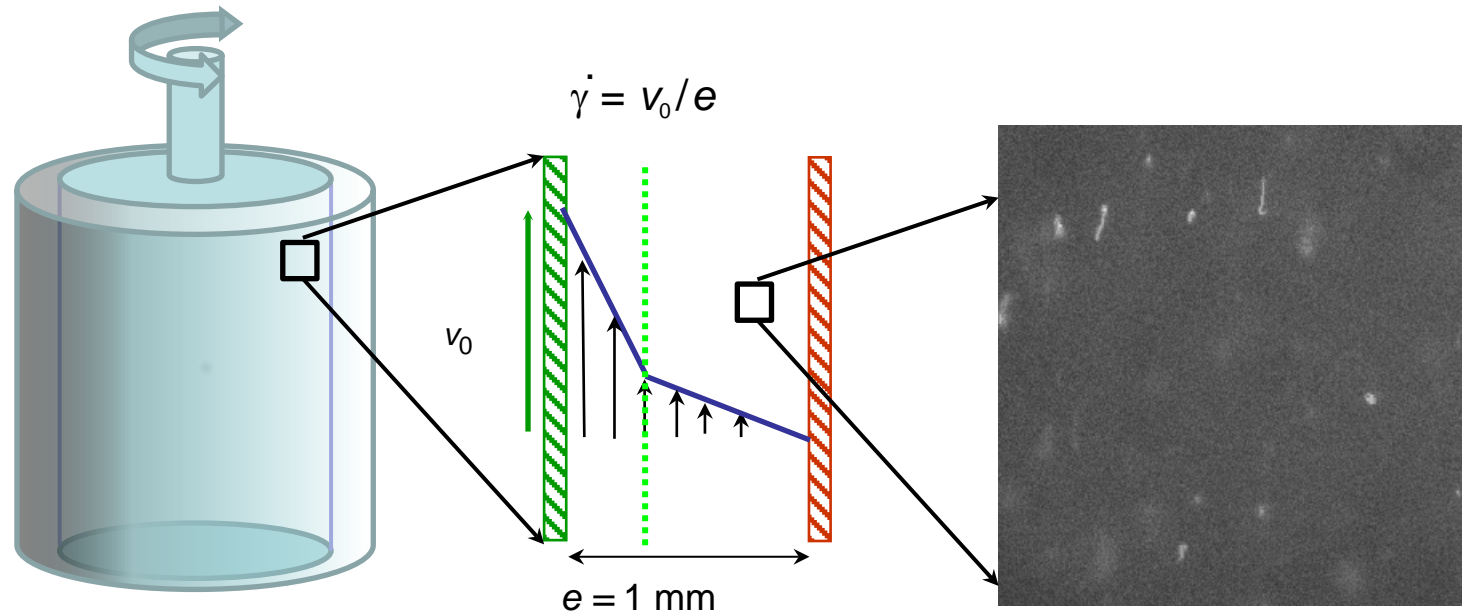
The non-ideal visco-elastic world

- Probe dynamics with Large Amplitude Oscillatory Shear

$$\dot{\gamma}(t) = \frac{A}{l} \omega \cos(\omega t) \text{ ———}$$

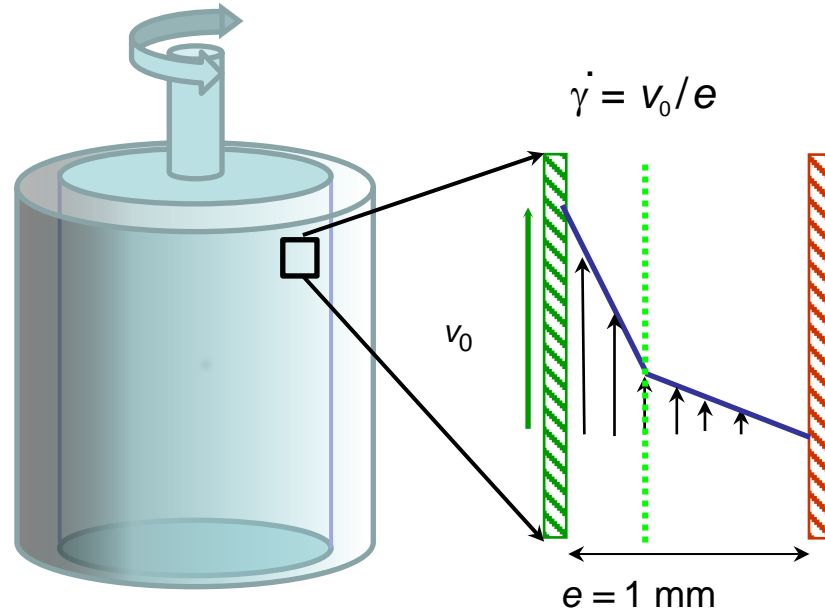


Information needed:



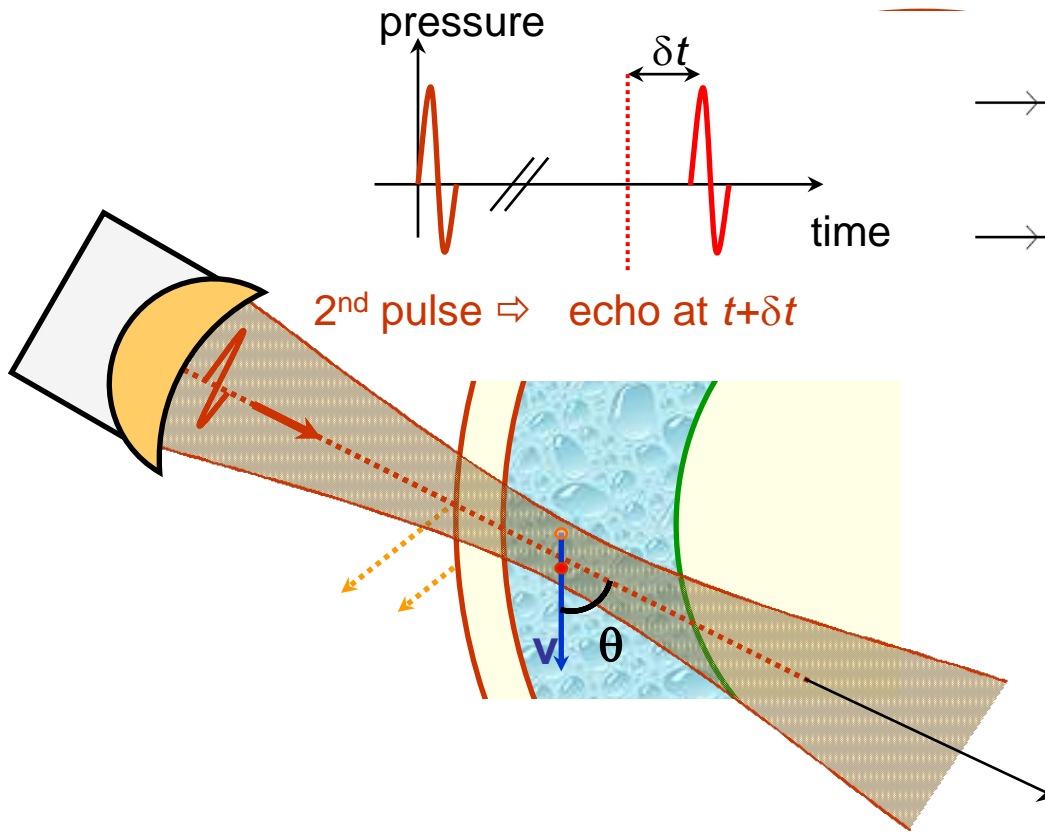
- Probe the mechanical response of the system.
- Probe the stability of the flow.
- Probe structure with *in situ* scattering or imaging methods over broad range of length-scales and time-scales.

Part I: Shear thinning and shear banding systems



Measuring velocity profiles

Ultra-sound velocimetry



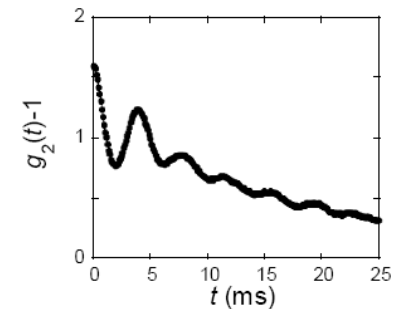
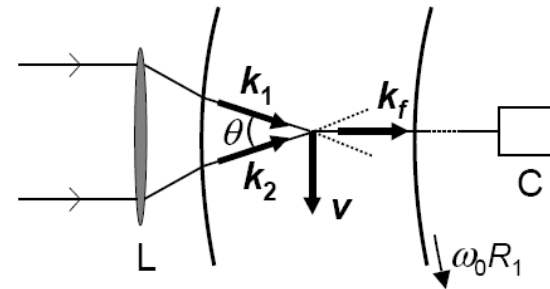
- Echography: arrival time $t \Leftrightarrow$ depth

$$z = \frac{c_0 t}{2}$$

- Velocimetry: time-shift $\delta t \Leftrightarrow$ displacement

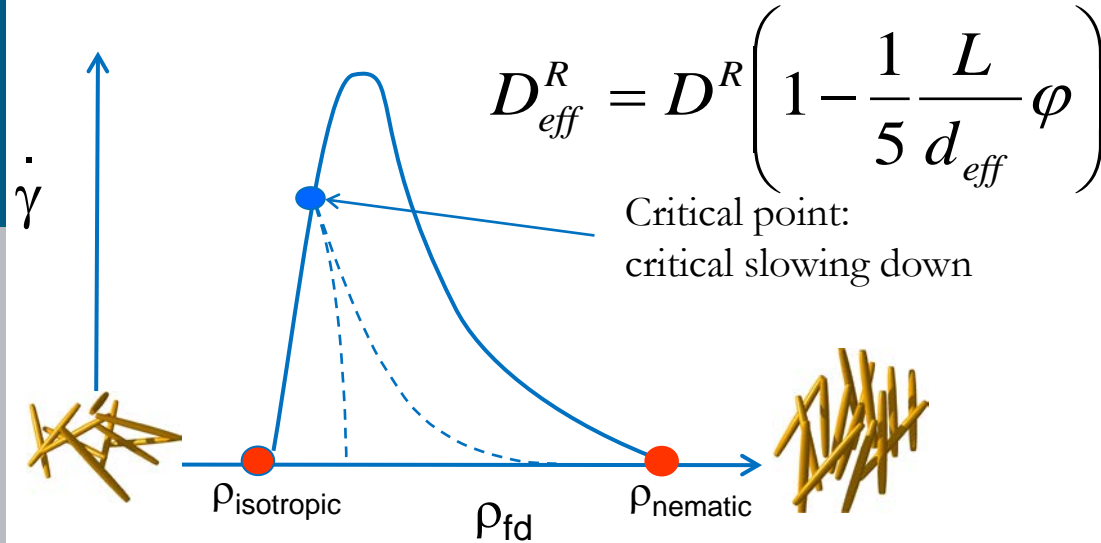
$$\delta z = \frac{c_0 \delta t}{2}$$

Heterodyne dynamic light scattering:

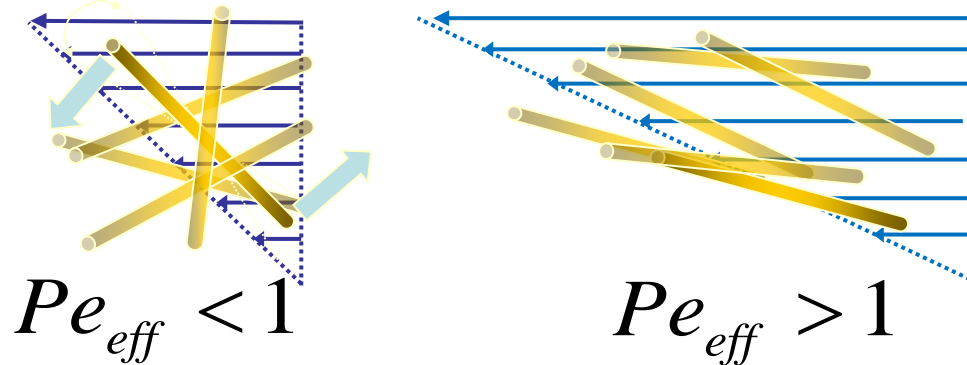


Possible scenarios for shear thinning

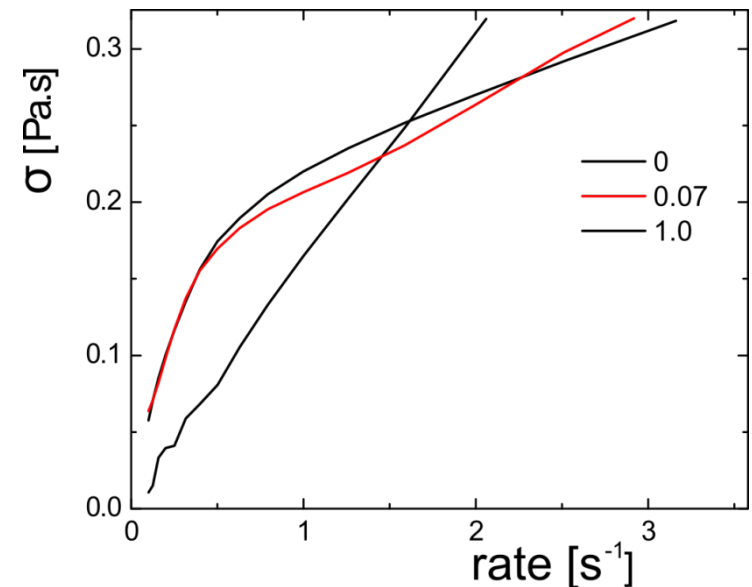
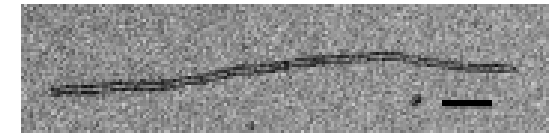
Rods close to I-N coexistence:



$$Pe_{eff} = \dot{\gamma}_0 / D_{eff}^R$$

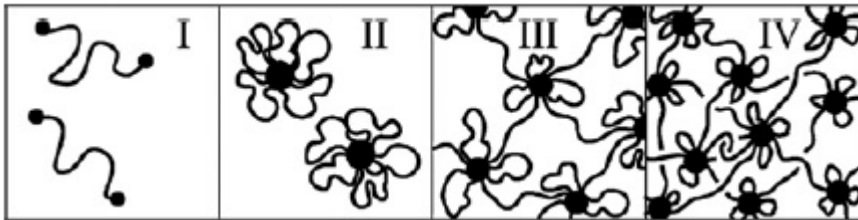


Fd virus: $L=880$ nm
 $D=6.6$ nm
 $P=3.0$ μ m

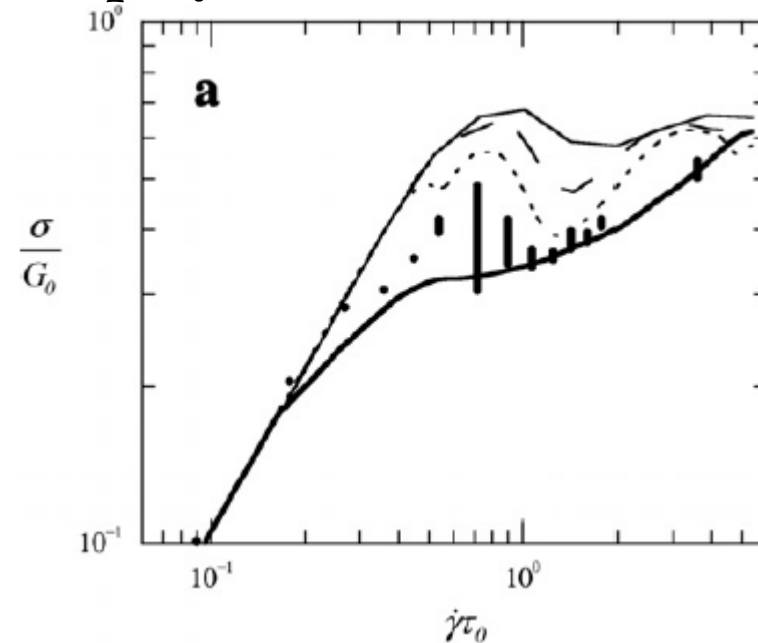
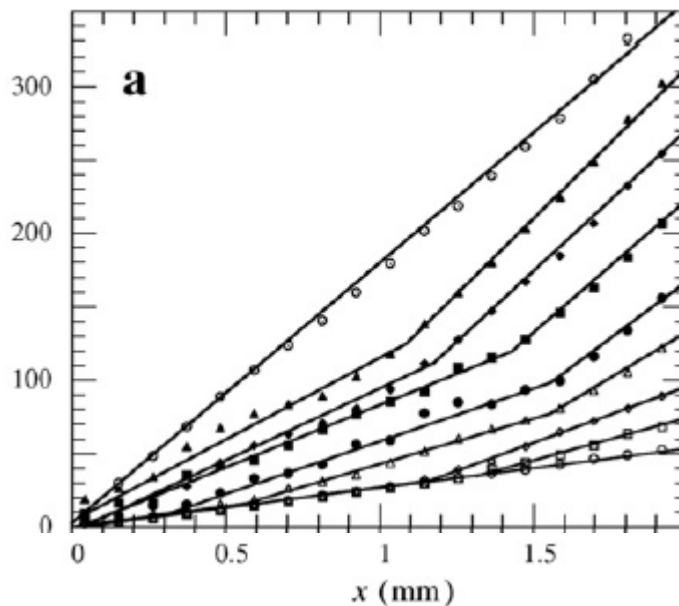


No gradient shear banding!

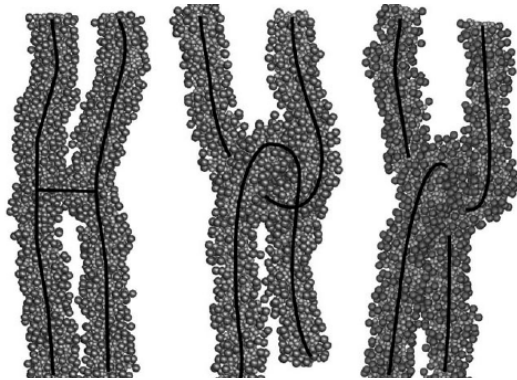
Living networks: tri-block copolymers



$$\tau = \tau_0 \exp\left(\frac{-f\delta}{k_B T}\right)$$



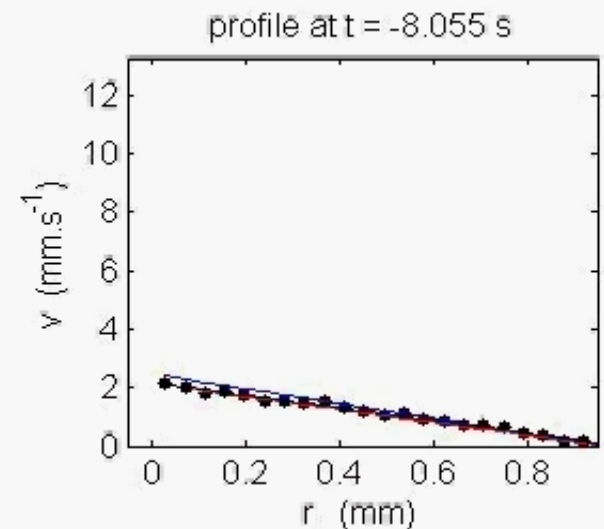
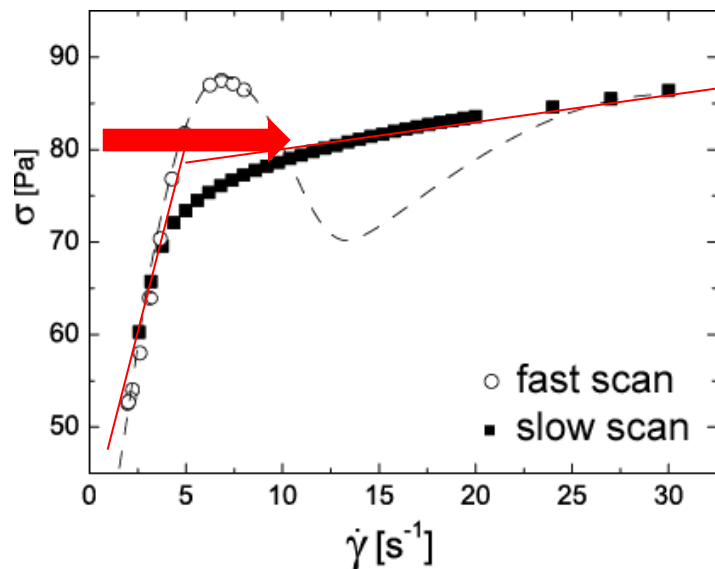
Living polymers: surfactant wormlike micelles



W J Briels, P Mulder and W K den Otter
J. Phys.: Condens. Matter **16** (2004) S3965–S3974

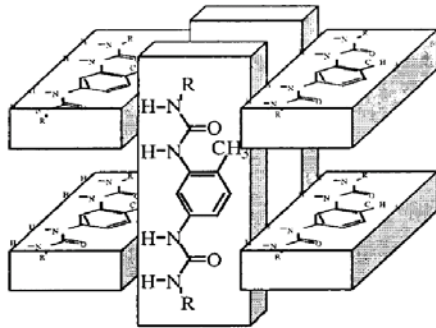
6% ww cetylpyridinium chloride/
sodium salicylate

- Single relaxation time $\lambda = \sqrt{\tau_{rep.} \tau_{break}}$
- (I-N at about 14 % ww)

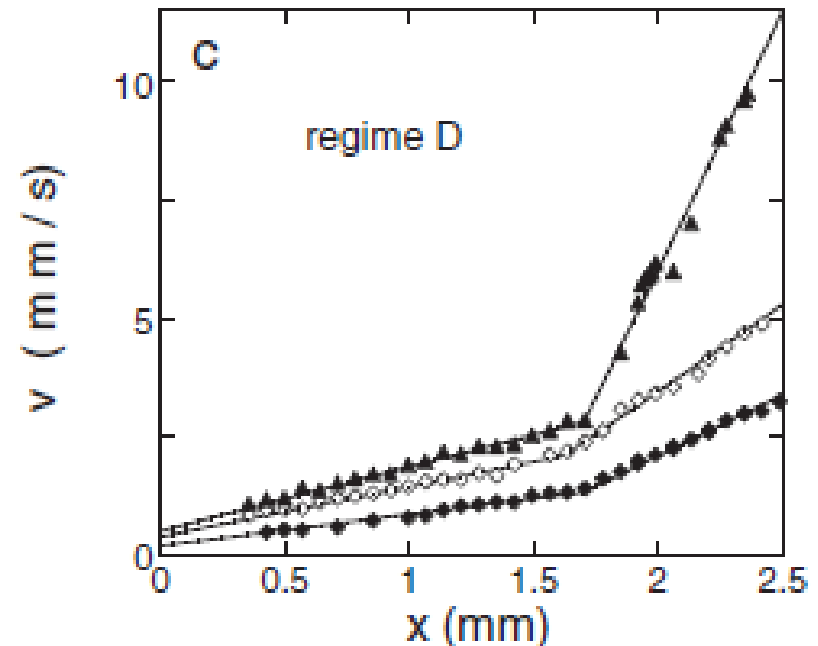
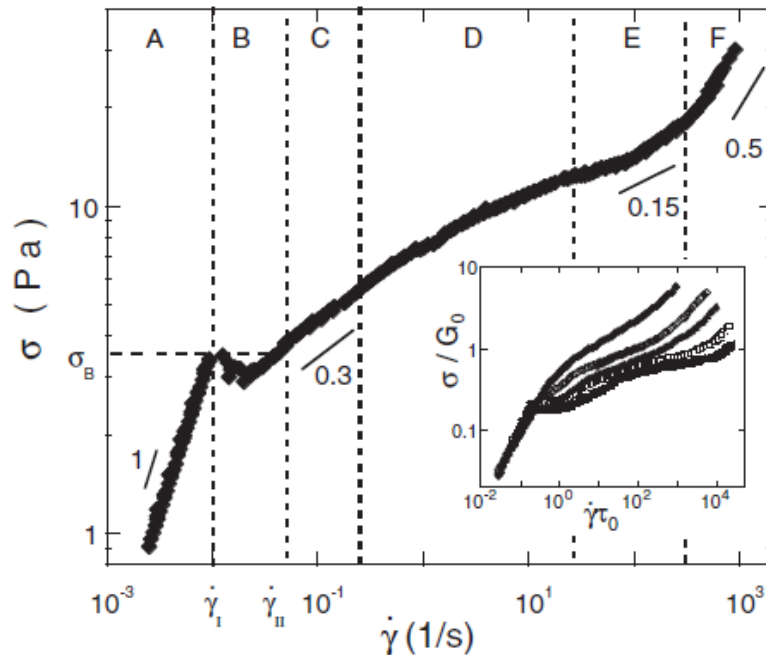


Possible scenarios for shear thinning

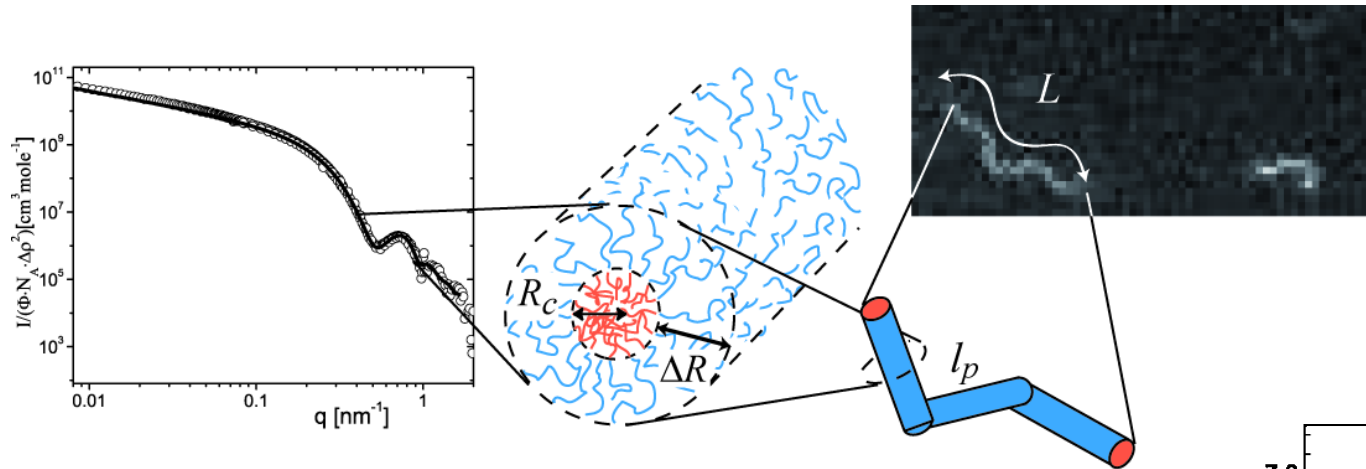
Living polymers: supra-molecular polymers



$$\lambda = \sqrt{\tau_{rep} \cdot \tau_{break}}$$



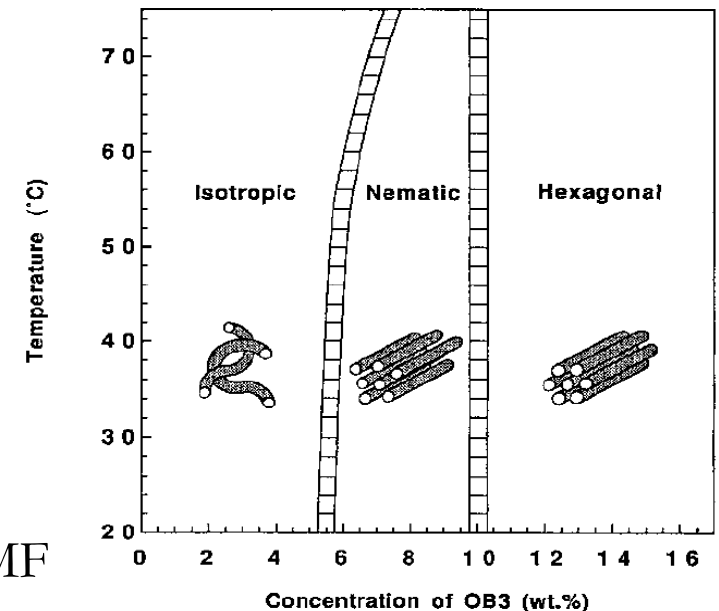
Rods close to I-N coexistence: PB-PEO wormlike micelles



Poly-(butadiene) 2.5 kd – Poly-ethylene glycol 2.5 kd

Advantage Pb-PEO:

- Displays I-N transition
- Wormlike micelles at size ratio of 1:1
- Detectable with fluorescence microscopy
- Tunable system at different length scales using D₂O:DMF

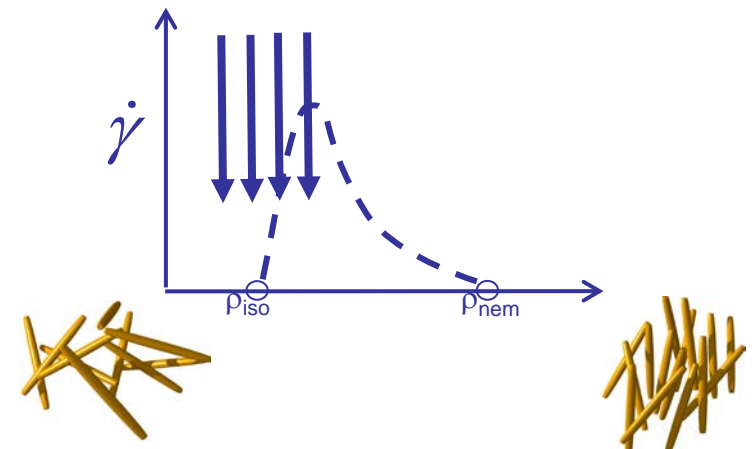
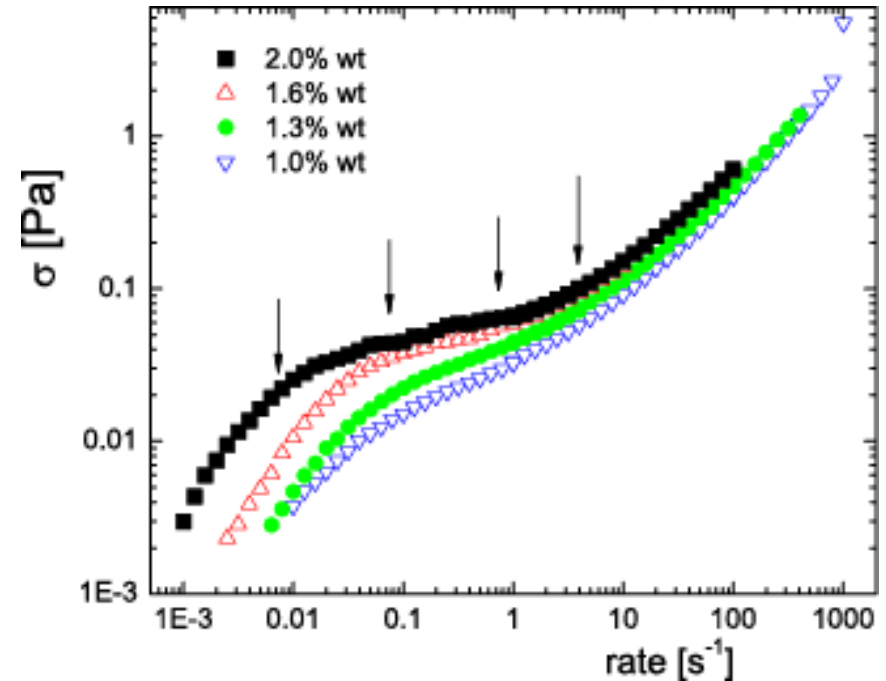
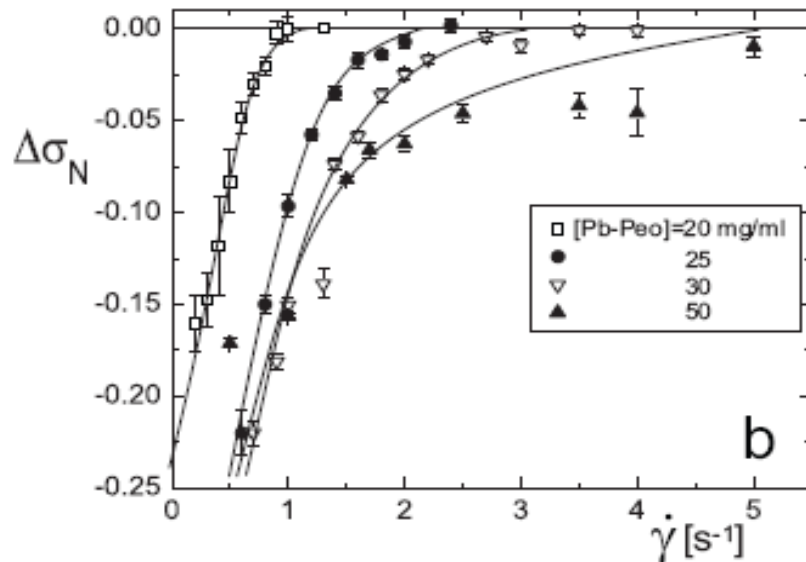
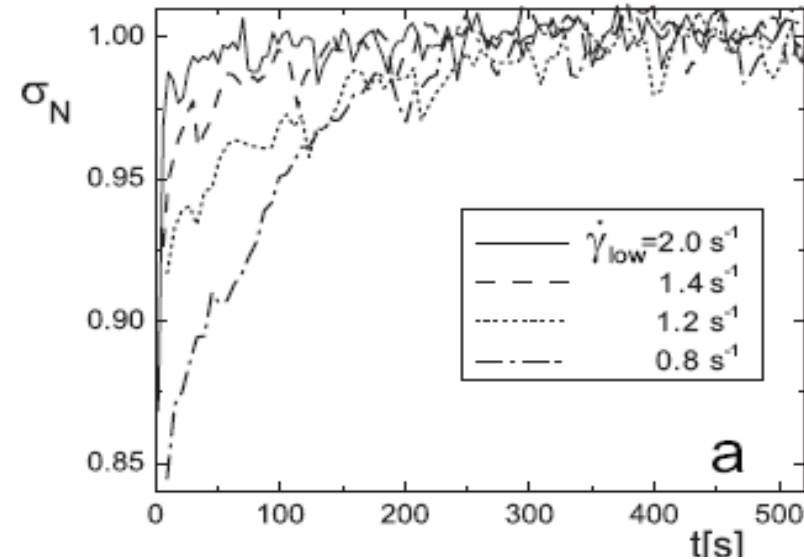


You-Yeon Won, H. Ted Davis, Frank S. Bates*

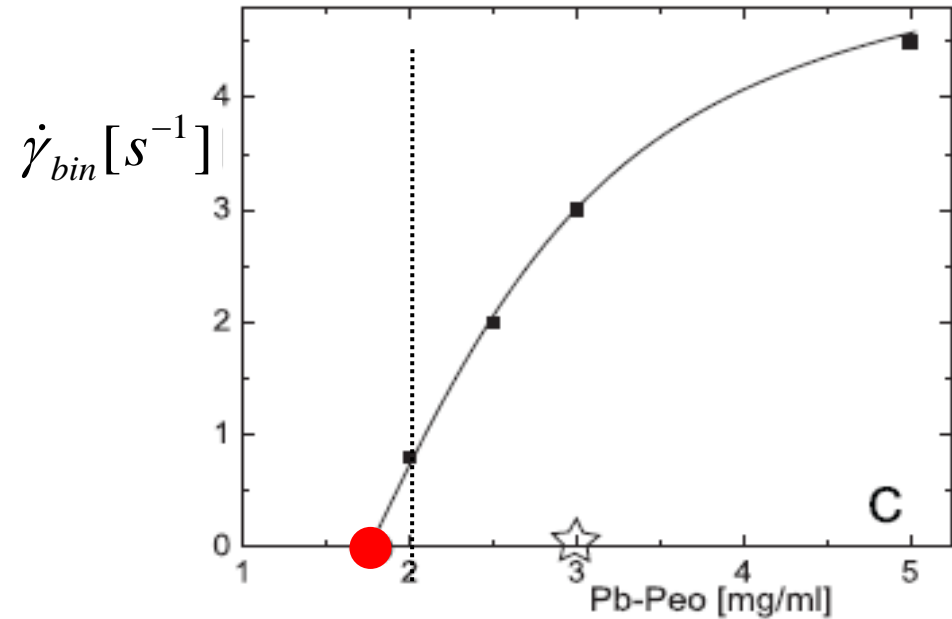
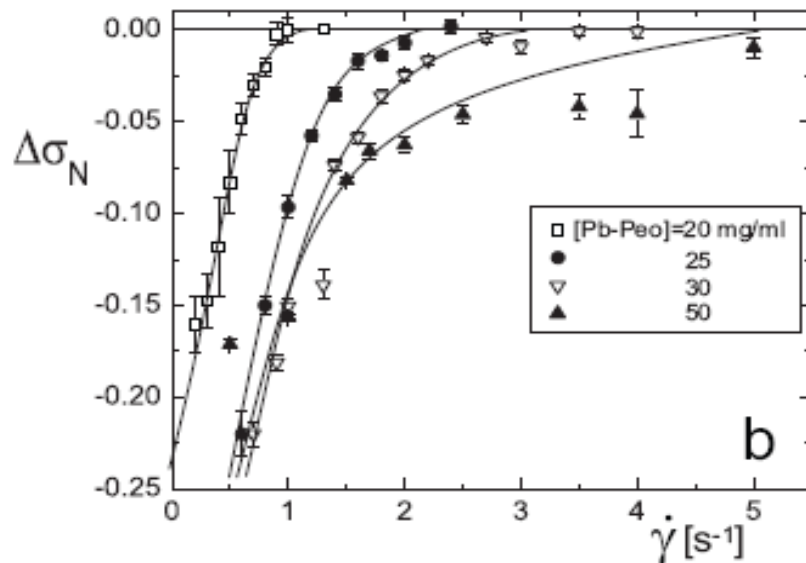
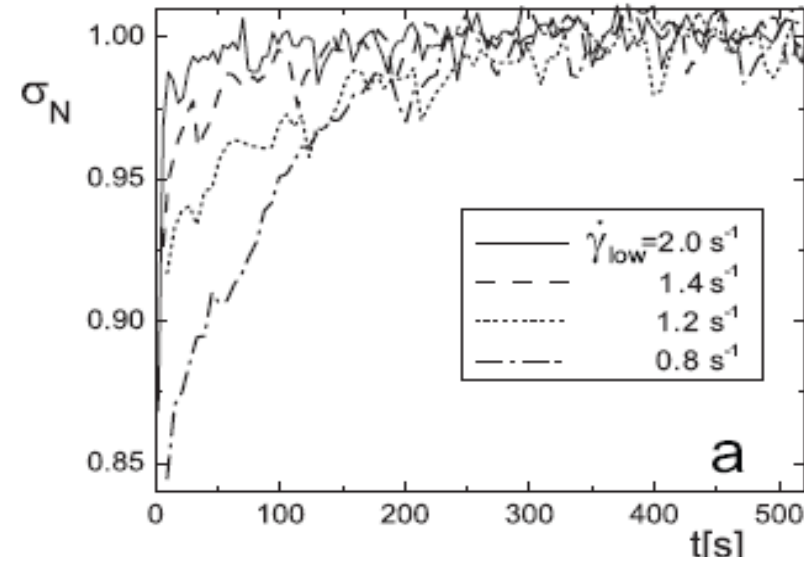
12 FEBRUARY 1999 VOL 283 SCIENCE

Possible scenarios for shear thinning

Test dynamics: quench down experiments

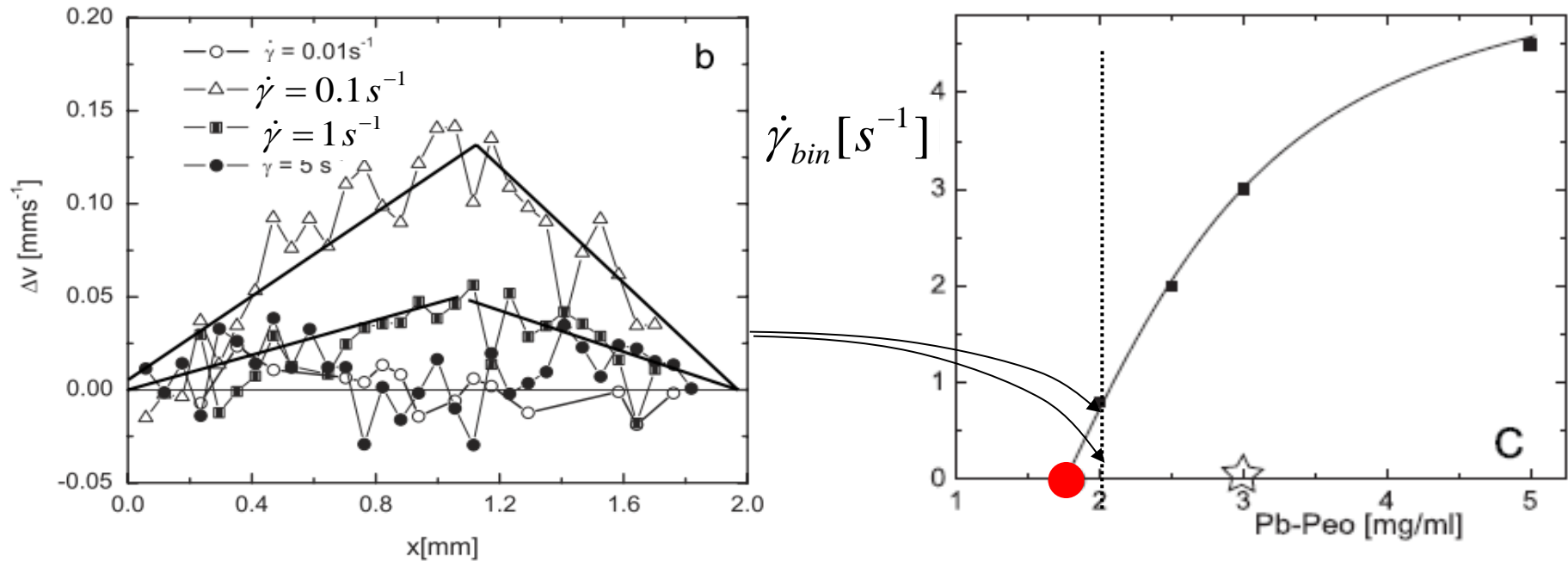


Possible scenarios for shear thinning

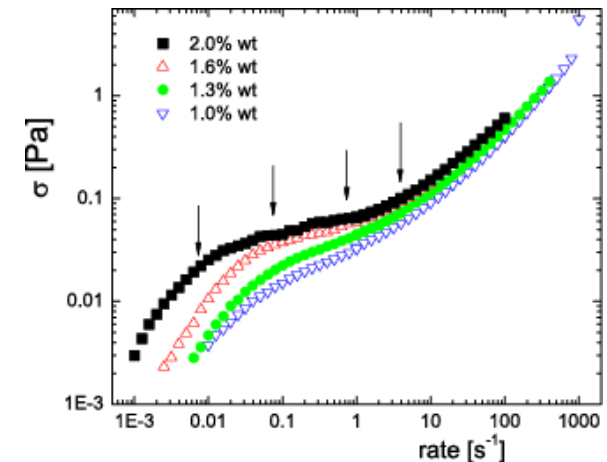


➤ First determination
equilibrium binodal

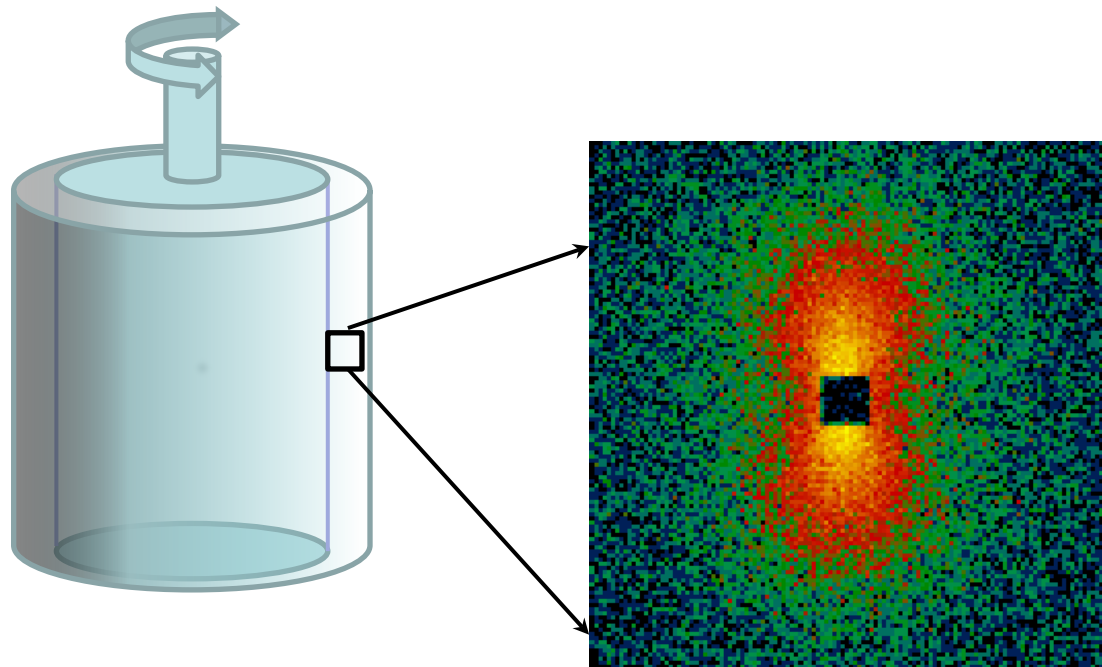
Rods close to I-N coexistence: PB-PEO wormlike micelles



➤ Polydispersity enhances shear banding

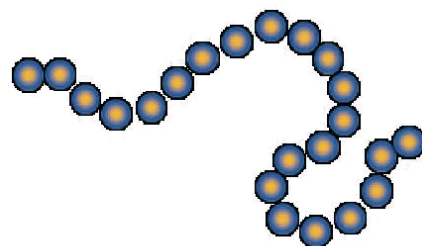


Part II: Molecular origin shear thinning using *in situ* LAOS

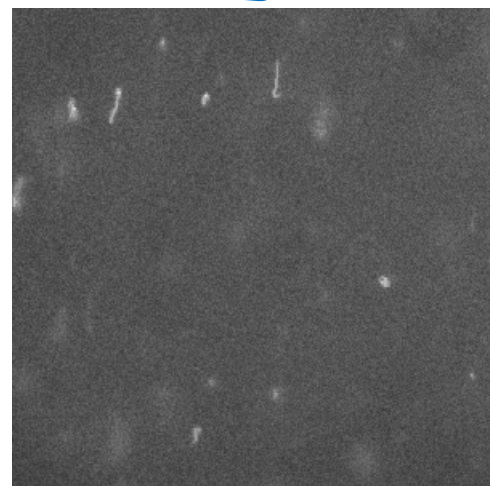
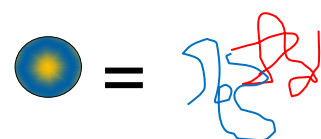


Focus on three systems:

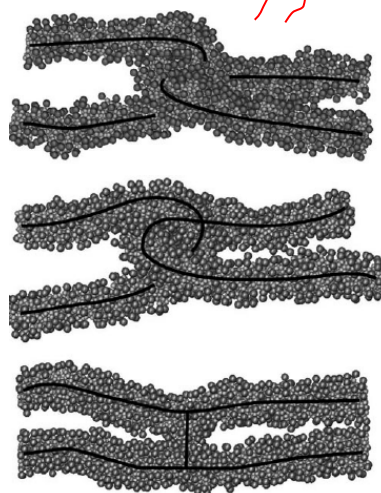
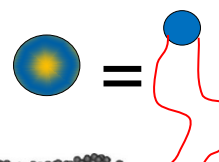
Living Polymers



Rods

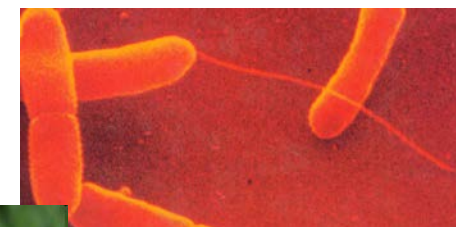


Pb-PEO
= Giant



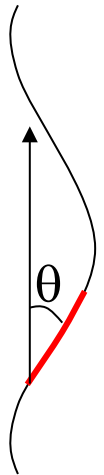
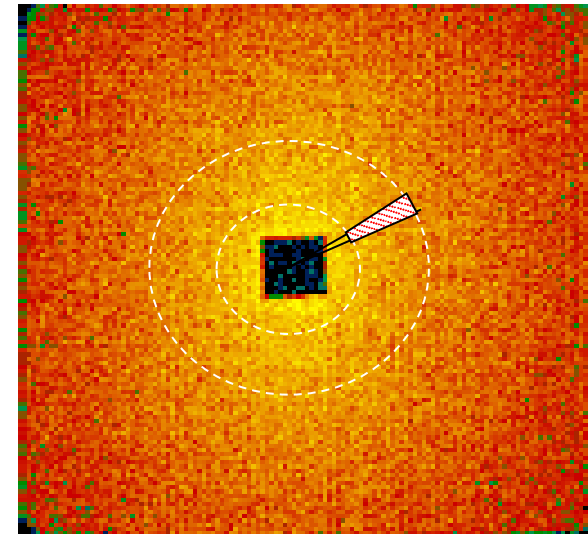
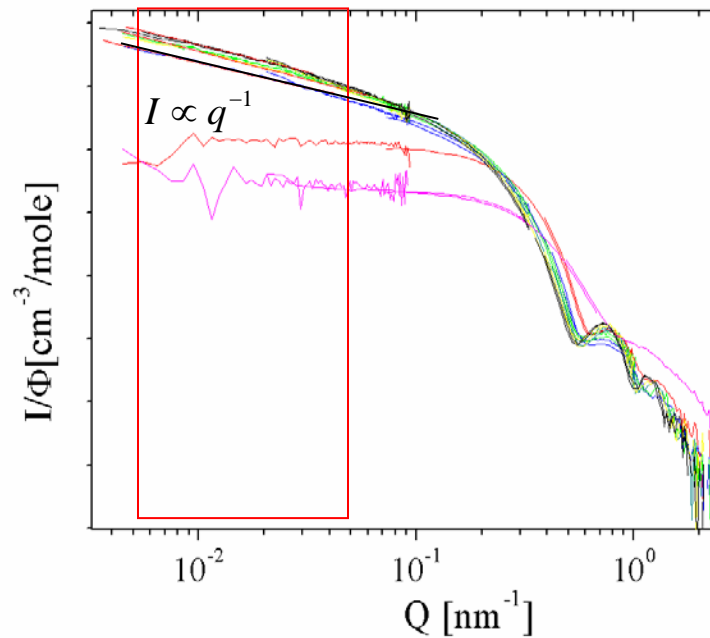
CpCl/NaSyl
= Dwarf

$L=880 \text{ nm}$
 $D=6.6 \text{ nm}$
 $P=3.0 \mu\text{m}$



Fd virus

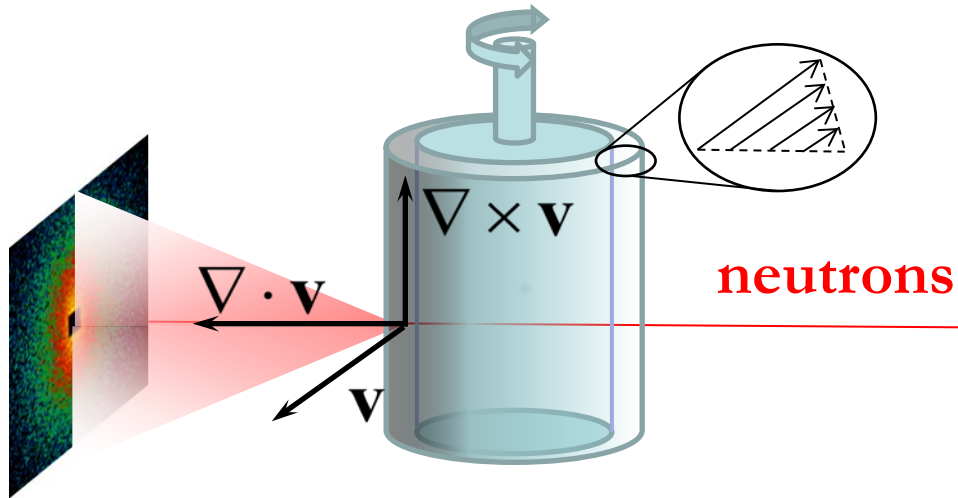
Scattering of a rod: $I \propto q^{-1}$



For the really soft matter (polymers and protein clusters) it is best to use neutrons:

- Better contrast
- No beam damage

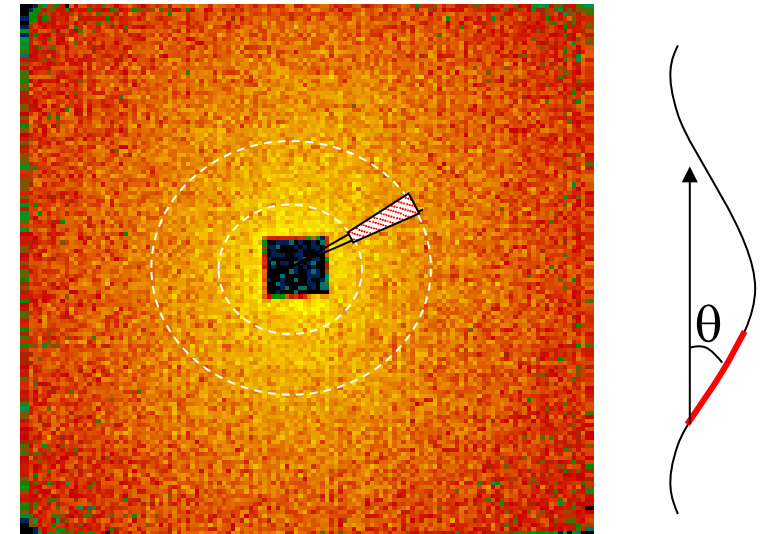
t-SANS to probe segment ordering



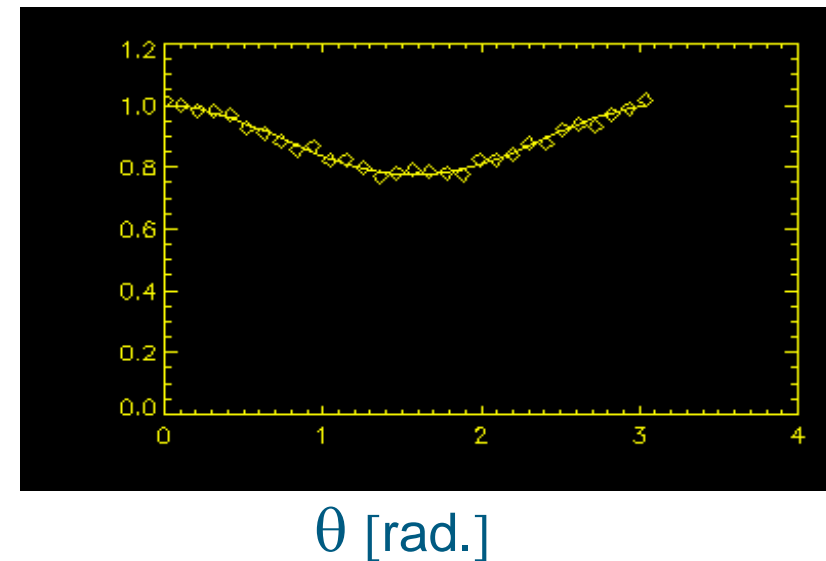
flow-vorticity plane probed

$$\langle P_2(t) \rangle = \frac{\int d\vartheta \sin(\vartheta) f(\vartheta) P_2(\vartheta)}{\int d\vartheta \sin(\vartheta) f(\vartheta)}$$

$f(\theta)$



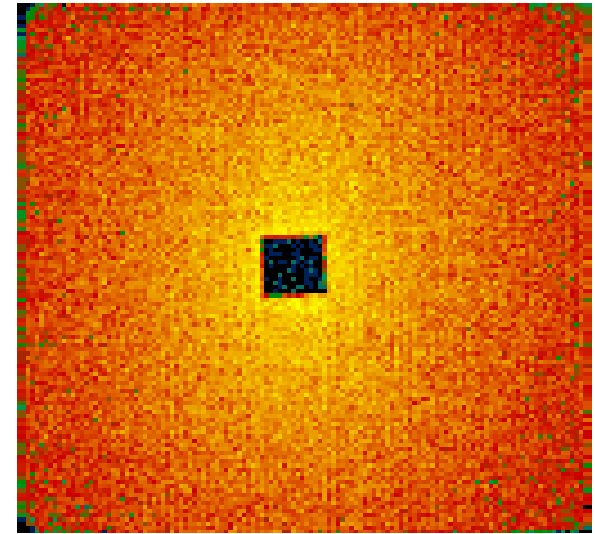
Orientalional disitribution function



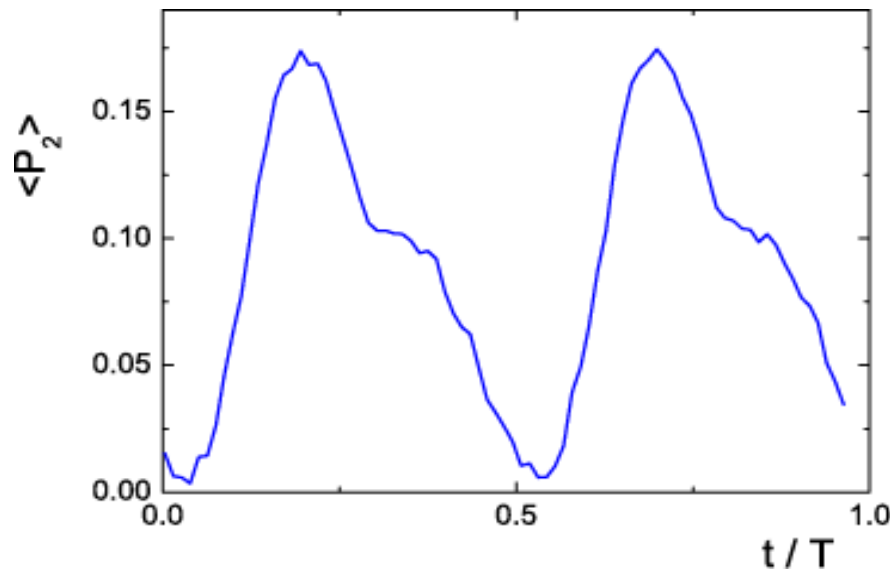
SANS in stroboscopic mode

$$I(t_i, \vec{q}) = \sum_n^{N_{\text{cycle}}} I(t_i + n\Delta t, \vec{q})$$

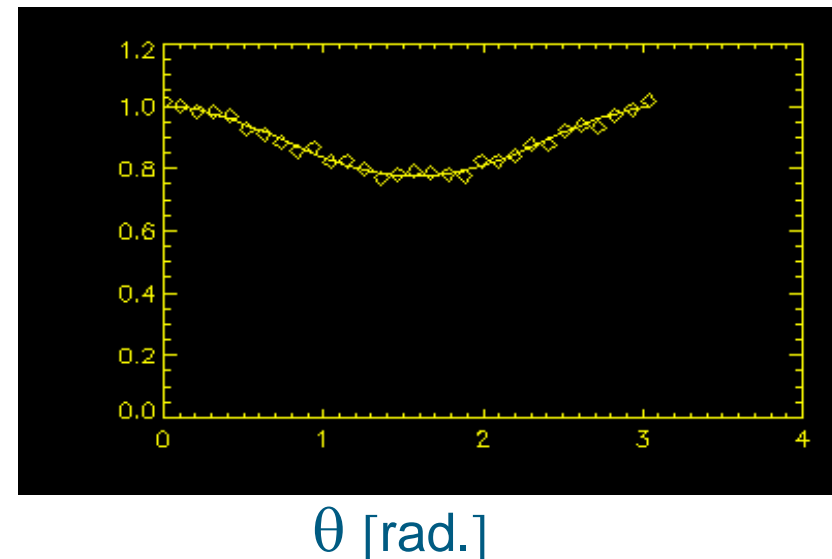
- at least $n = 100$ time channels of widths $\Delta t = (n\omega/2\pi)^{-1}$
- trigger is sent when the maximum shear rate is reached
- Summed over an interval of time ranging from one hour to fifteen minutes



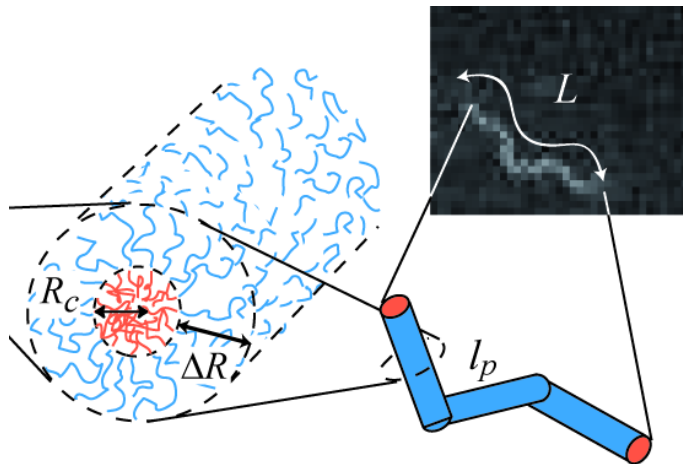
$$f = 0.25 \text{ Hz}; \dot{\gamma}_{\text{max}} = 18 \text{ s}^{-1}$$



$f(\theta)$

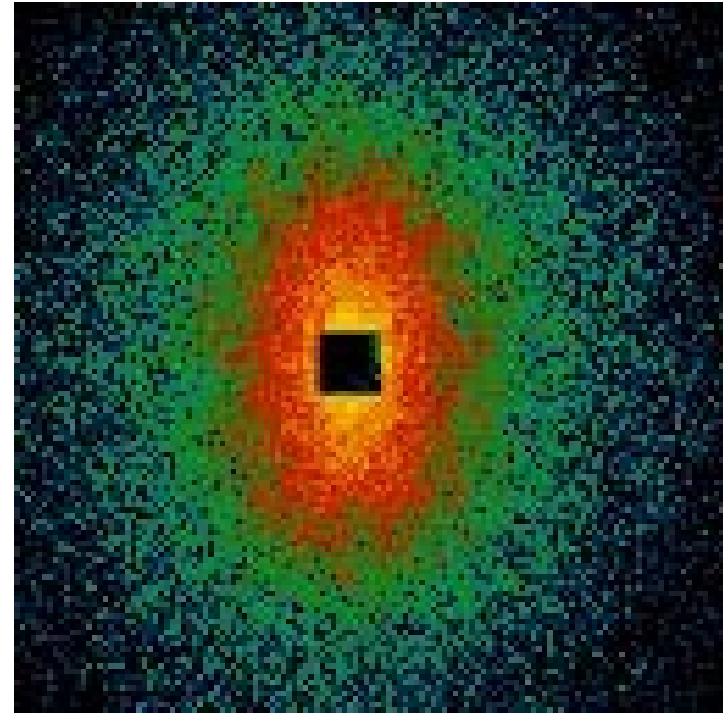


Dynamic response **Pb-Peo** wormlike micelles



$$\nabla \times \mathbf{v}$$

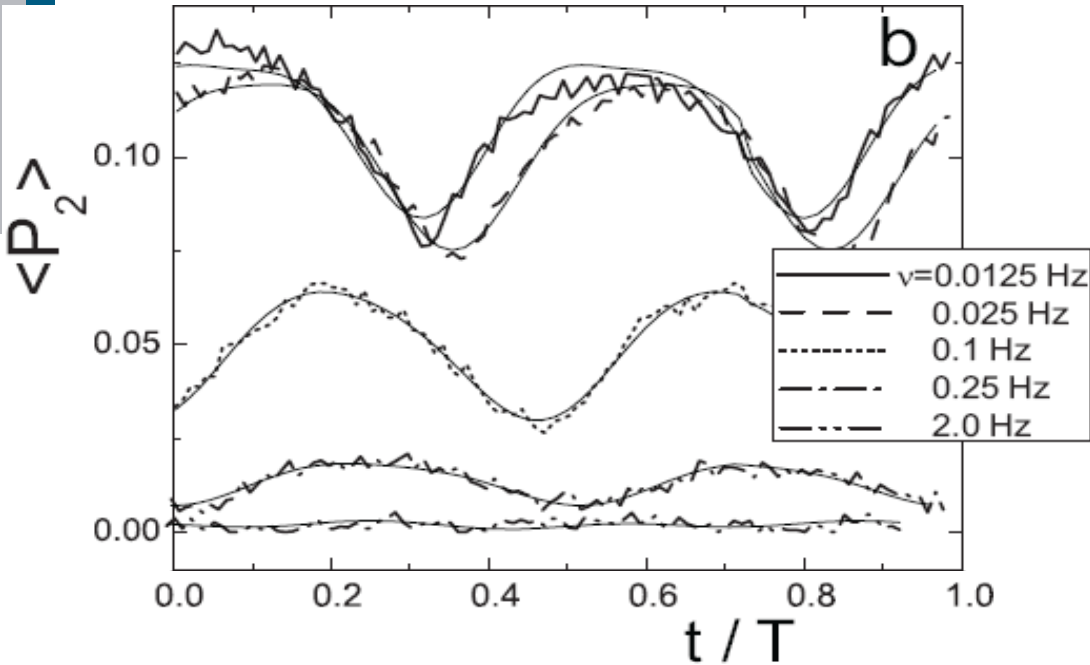
A coordinate system with a vertical arrow pointing upwards labeled $\nabla \times \mathbf{v}$ and a horizontal arrow pointing to the right labeled \mathbf{v} .



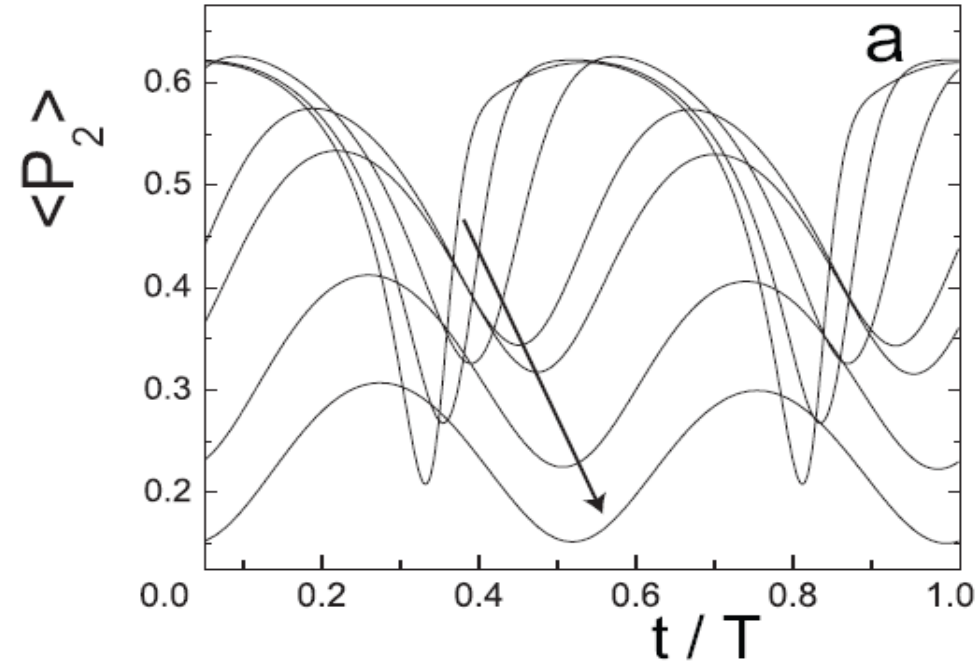
t-SANS @ PSI
 $f=0.025$, $\gamma=7.3$

Oscillatory flow: compare P_2

Fixed shear rate = 1 s^{-1}

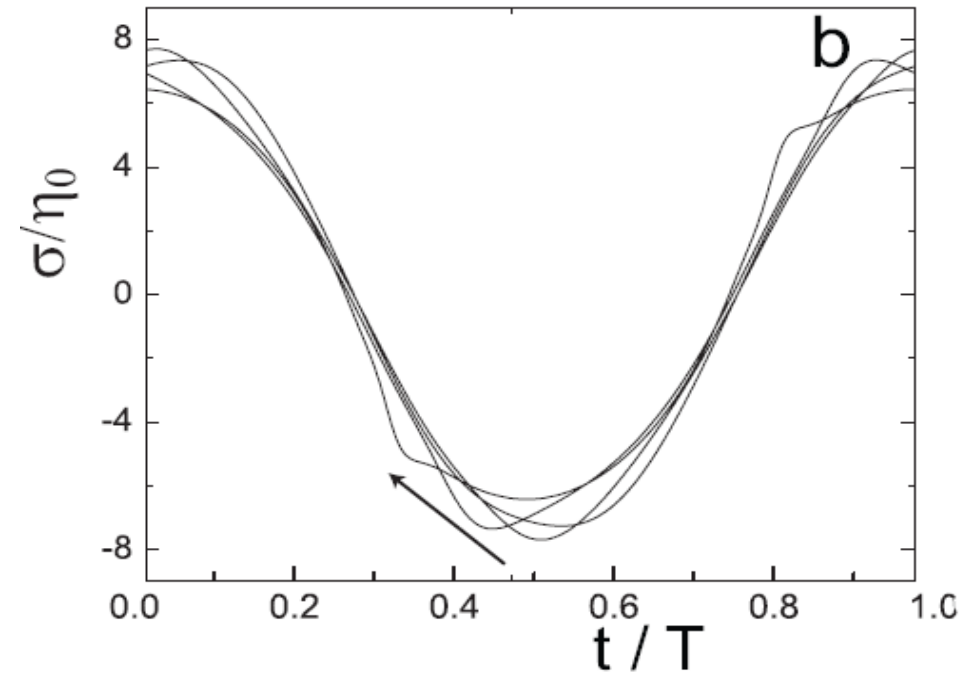
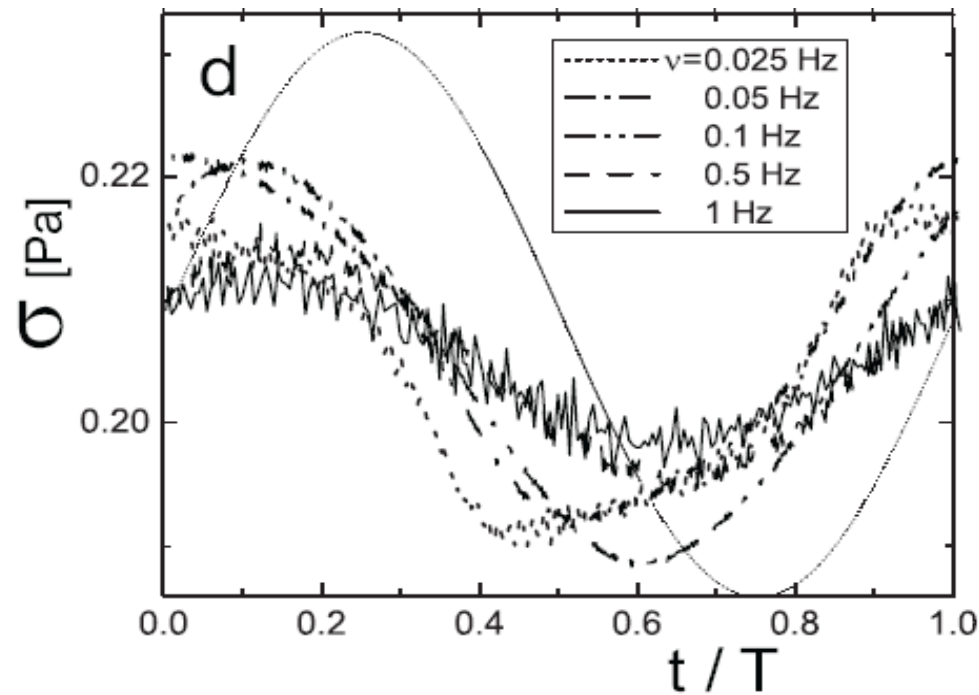


Fixed $Pe_{\text{eff}} = 75$

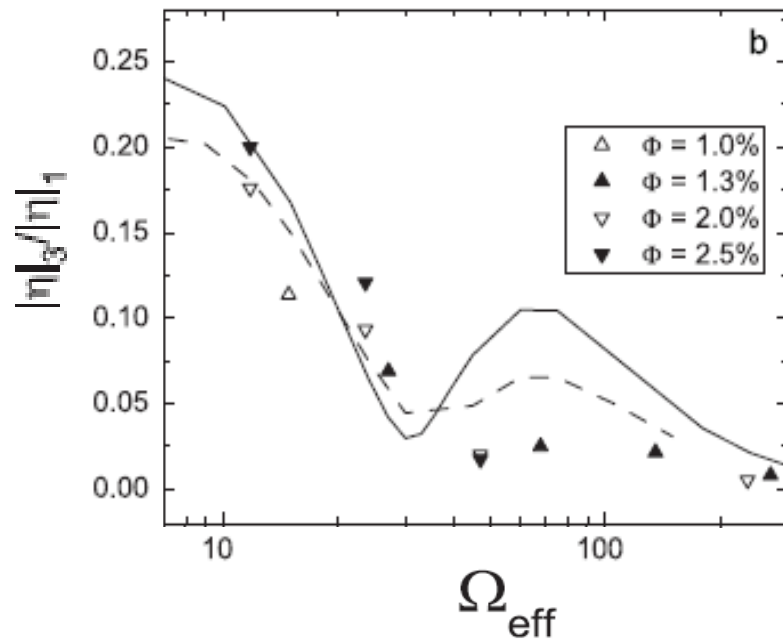


$$\frac{d}{dt}\mathbf{S} = -6D_r \left\{ \mathbf{S} - \frac{1}{3}\hat{\mathbf{I}} + \frac{L}{D}\varphi \left(\mathbf{S}^{(4)} : \mathbf{S} - \mathbf{S} \cdot \mathbf{S} \right) \right\} + \dot{\gamma} \left\{ \hat{\mathbf{r}} \cdot \mathbf{S} + \mathbf{S} \cdot \hat{\mathbf{r}}^T - 2\mathbf{S}^{(4)} : \hat{\mathbf{E}} \right\}$$

Oscillatory flow: compare stress



$$\Sigma_D = 2\eta_0\dot{\gamma} \left[\hat{\mathbf{E}} + \frac{(L/D)^2}{3 \ln\{L/D\}} \varphi \right. \\ \left. \times \left\{ \hat{\mathbf{r}} \cdot \mathbf{S} + \mathbf{S} \cdot \hat{\mathbf{r}}^T - \mathbf{S}^{(4)} : \hat{\mathbf{E}} - \frac{1}{3} \hat{\mathbf{I}} \mathbf{S} : \hat{\mathbf{E}} - \frac{1}{\dot{\gamma}} \frac{d\mathbf{S}}{dt} \right\} \right]$$

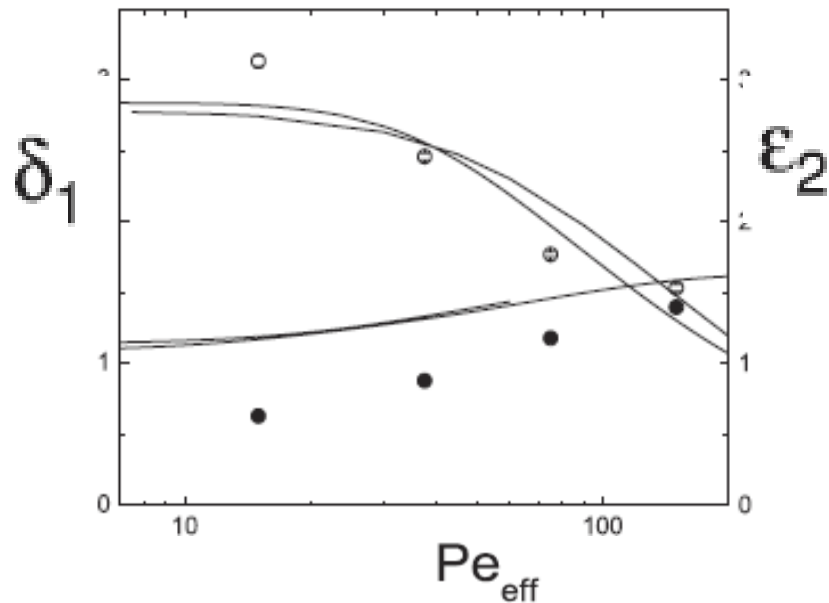


$$\Sigma_D = 2 \dot{\gamma}_0 \hat{\mathbf{E}} \sum_{n=0}^{\infty} |\eta|_n \sin(n\omega t + \delta_n)$$

$$P_2(t) = \sum_n |P_2|_n \cos(\omega t + \epsilon_n)$$

Find scaling for...

$$Pe_{\text{eff}} = \dot{\gamma}_0 / D_r^{\text{eff}} \quad \Omega_{\text{eff}} = \omega / D_r^{\text{eff}}$$



Using...

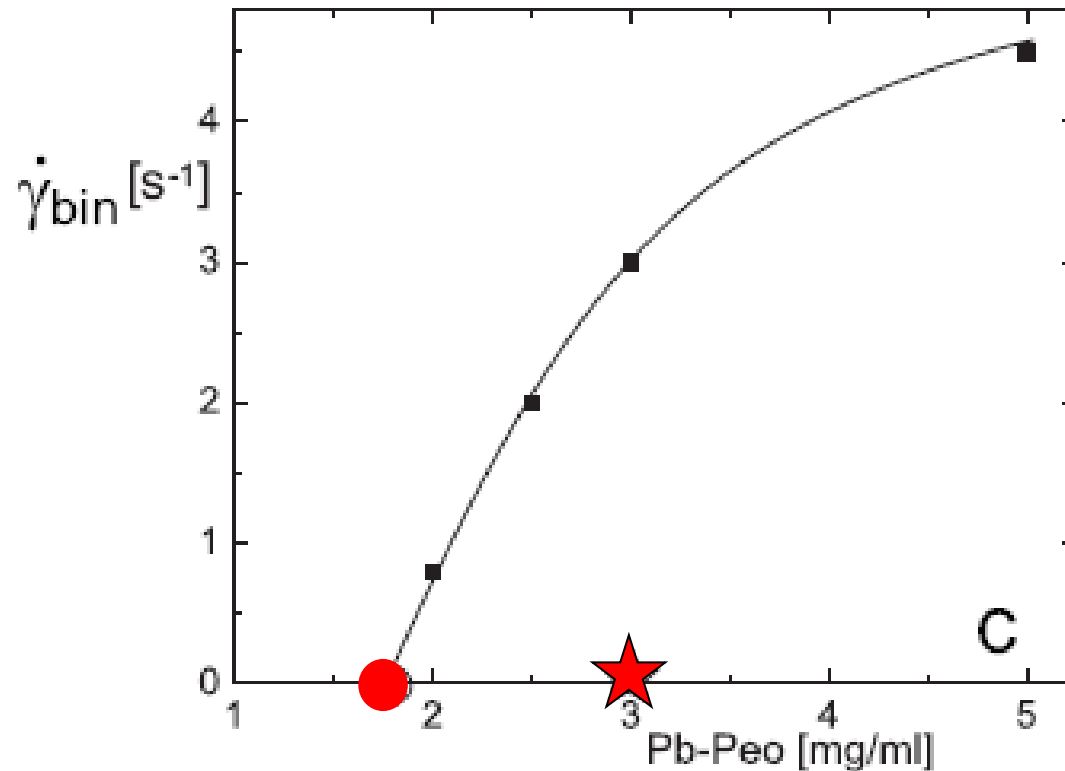
$$D_r^{\text{eff}} = D_r \{1 - [Pb - Peo]/C\}$$

with fit parameters:

$$D_r = 0.04 \text{ s}^{-1} \text{ and } C = 3$$

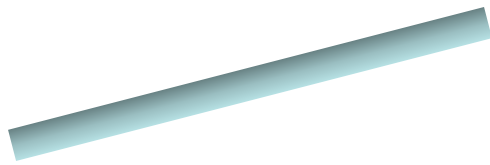
$$C=3 \text{ mg/ml, so } \phi_{I-N} = 4 \frac{L}{d} = 3$$

We obtained the I—N spinodal point!

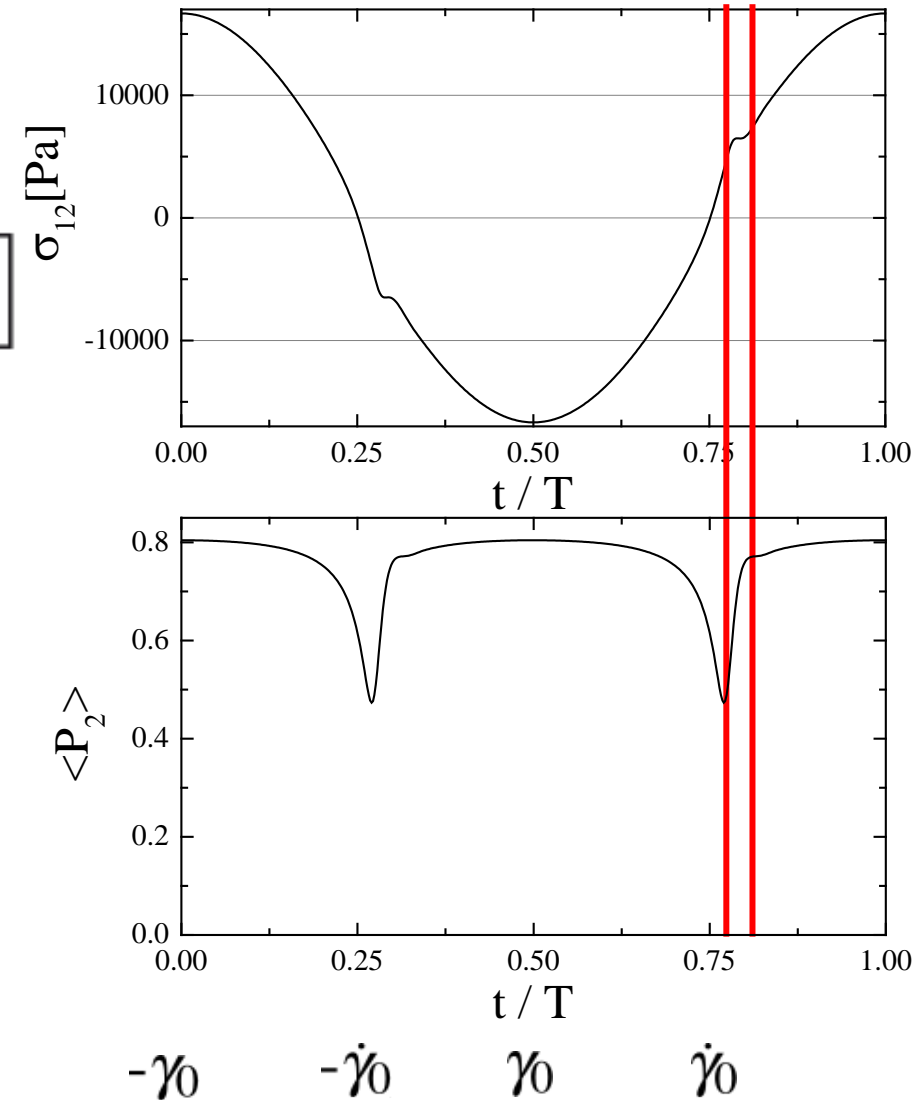


Smoluchowski Theory for hard rods

$$\Sigma_D = 2\eta_0\dot{\gamma} \left[\hat{\mathbf{E}} + \frac{(L/D)^2}{3 \ln\{L/D\}} \varphi \right. \\ \left. \times \left\{ \hat{\mathbf{r}} \cdot \mathbf{S} + \mathbf{S} \cdot \hat{\mathbf{r}}^T - \mathbf{S}^{(4)} : \hat{\mathbf{E}} - \frac{1}{3} \hat{\mathbf{I}} \mathbf{S} : \hat{\mathbf{E}} - \frac{1}{\dot{\gamma}} \frac{d\mathbf{S}}{dt} \right\} \right] \sigma_{12} [\text{Pa}]$$



$$\frac{d}{dt} \mathbf{S} = -6D_r \left\{ \mathbf{S} - \frac{1}{3} \hat{\mathbf{I}} + \frac{L}{D} \varphi \left(\mathbf{S}^{(4)} : \mathbf{S} - \mathbf{S} \cdot \mathbf{S} \right) \right\} \\ + \dot{\gamma} \left\{ \hat{\mathbf{r}} \cdot \mathbf{S} + \mathbf{S} \cdot \hat{\mathbf{r}}^T - 2\mathbf{S}^{(4)} : \hat{\mathbf{E}} \right\}$$

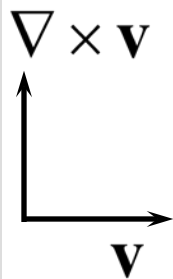
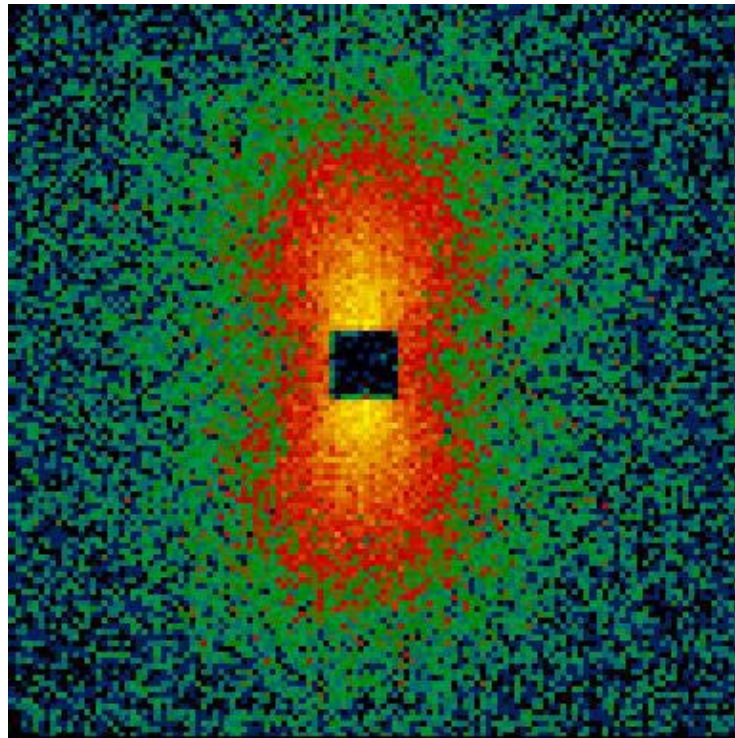
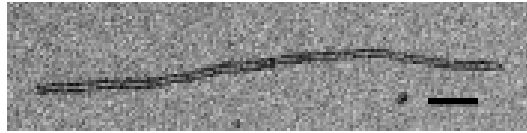


Dynamic response **fd virus** in isotropic phase

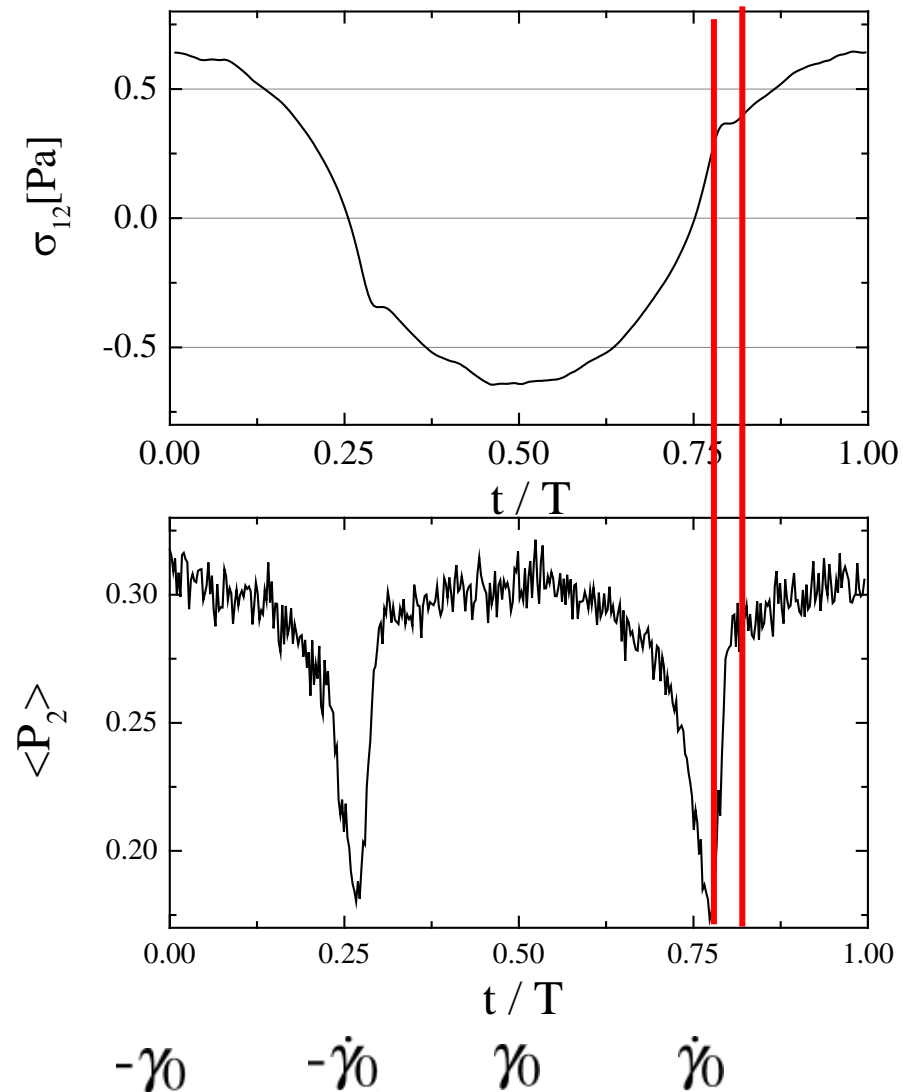
$L=880$ nm

$D=6.6$ nm

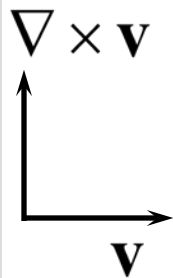
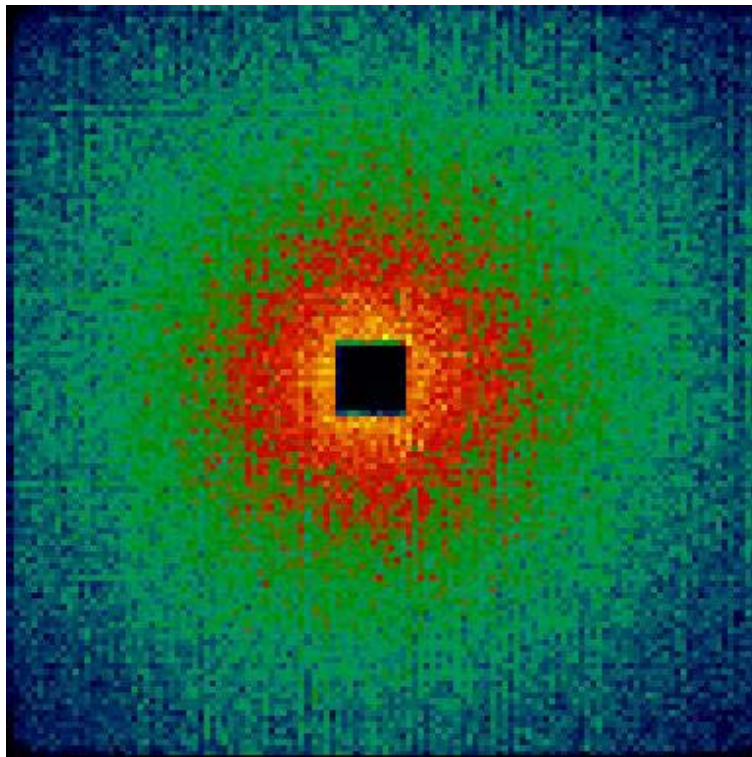
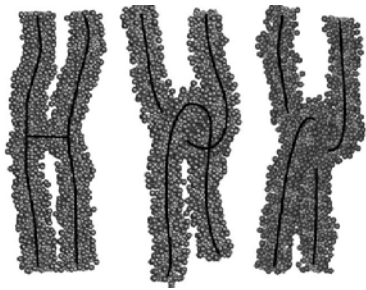
$P=3.0$ μm



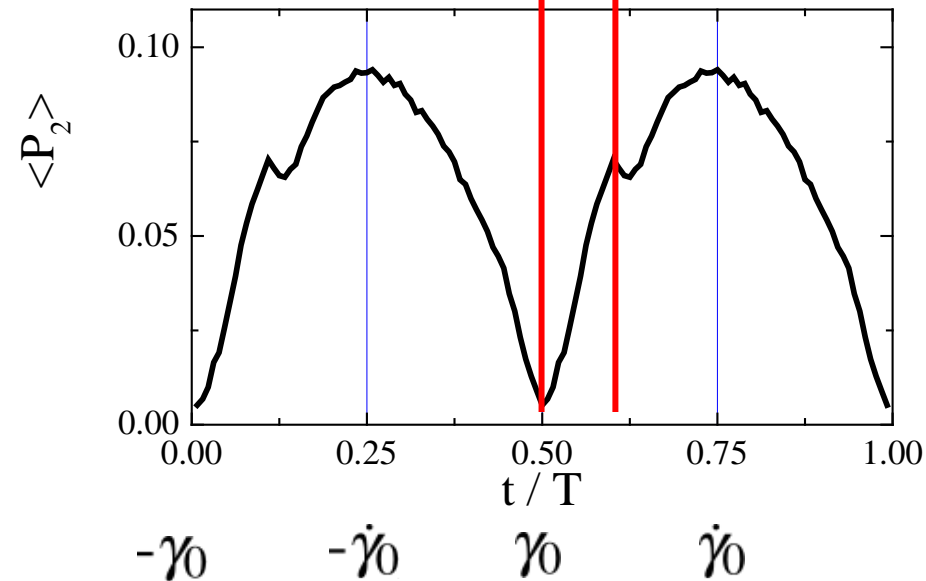
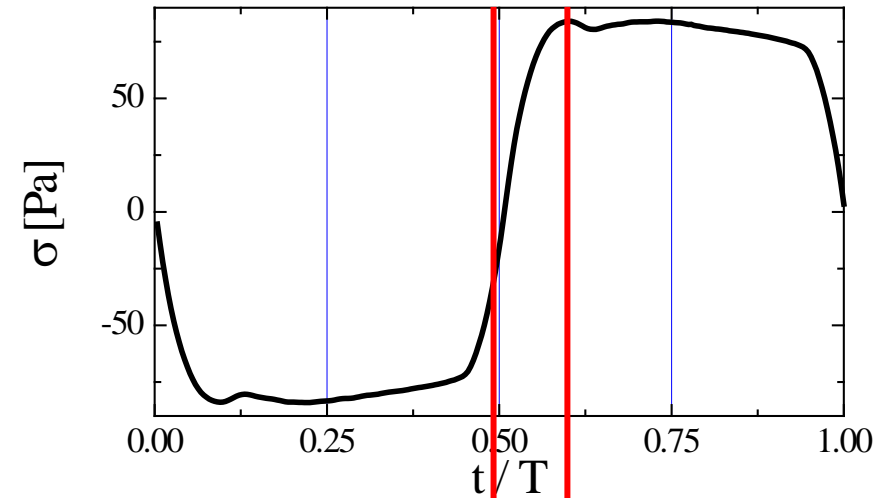
t-SANS @ PSI
 $f=0.01$, $\gamma=25.6$



Dynamic response **CpCl** wormlike micelles

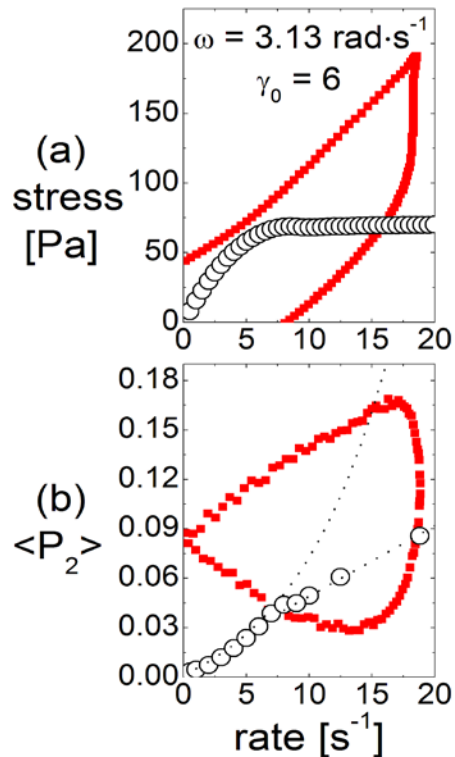


t-SANS @ PSI
 $f=0.0037$, $\gamma=766$



@ constant applied shear rate amplitude $= \gamma_0 \omega$

Rate dependence (positive rate only)



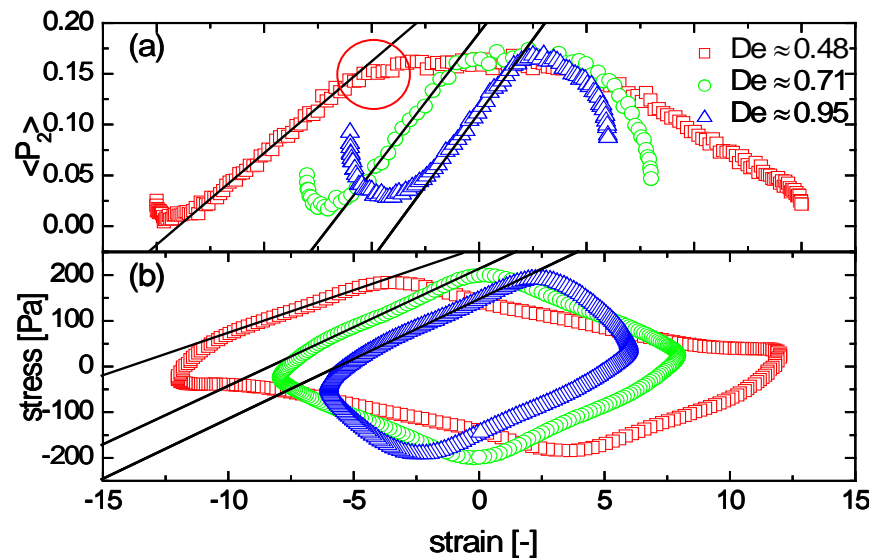
@ constant applied shear rate amplitude = $\gamma_0 \omega$

Strain dependence

$\omega \gg 1/\lambda \rightarrow$ elastic

Stress-optical relation

$$\text{for } \gamma < 2\gamma_0\lambda : \quad \frac{\partial P_2}{\partial \sigma} = \frac{S_K}{G_K}$$



$\omega < 1/\lambda \rightarrow$ yielding

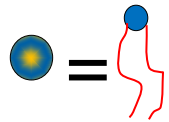
Transform from solid
into liquid

Critical strain:

$$\gamma @ \sigma_{\max} \approx 2\gamma_0\lambda$$

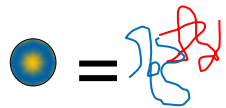
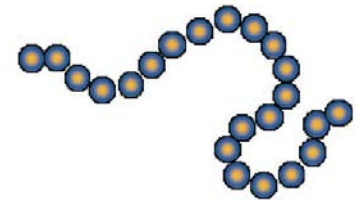
$$\langle P_2 \rangle_{\max} \sim 2\langle P_2 \rangle_{\text{stationary}}$$

Conclusions part I and II



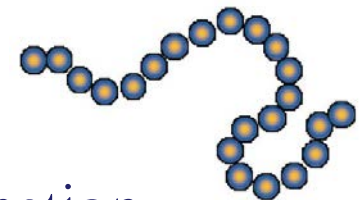
CpCl/NaSyl = Dwarf = Living

- strong shear thinning, fast shear band formation
- stress relaxation probably through break up
- maximum orientation at maximum stress



Pb-PEO = Giant = Dead

- strong shear thinning close to I-N, some shear band formation
- stress relaxation probably through alignment Kuhn segment

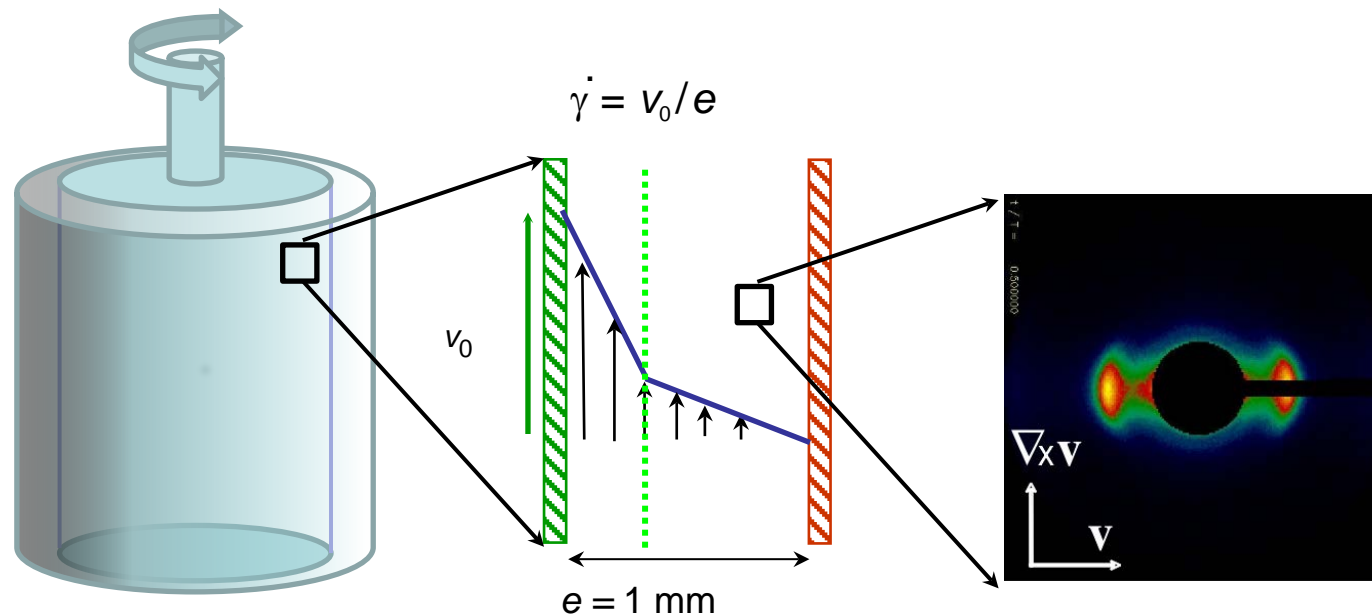


Fd virus = Rods



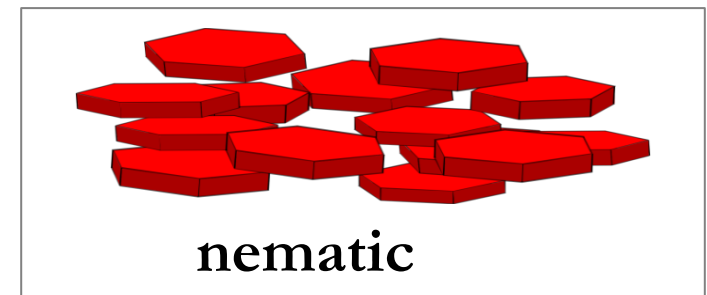
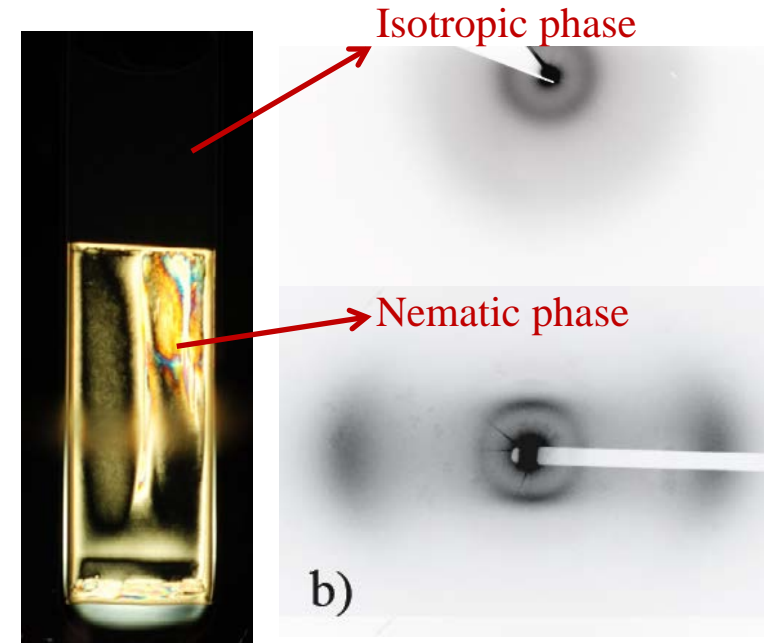
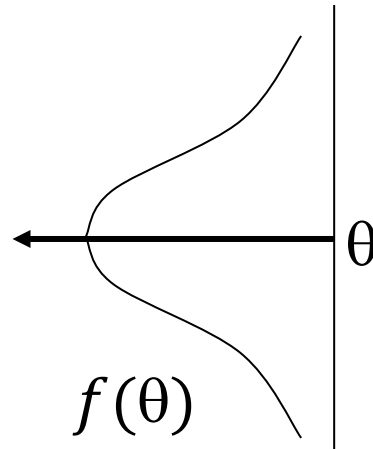
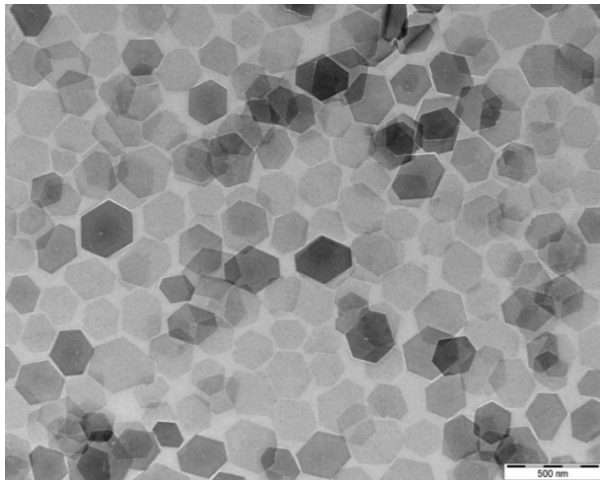
- moderate shear thinning close to I-N, no shear band formation
- stress relaxation probably through alignment rods
- minimum orientation at maximum stress

Part III: 3D reorientational motion in sheared nematic platelet dispersions



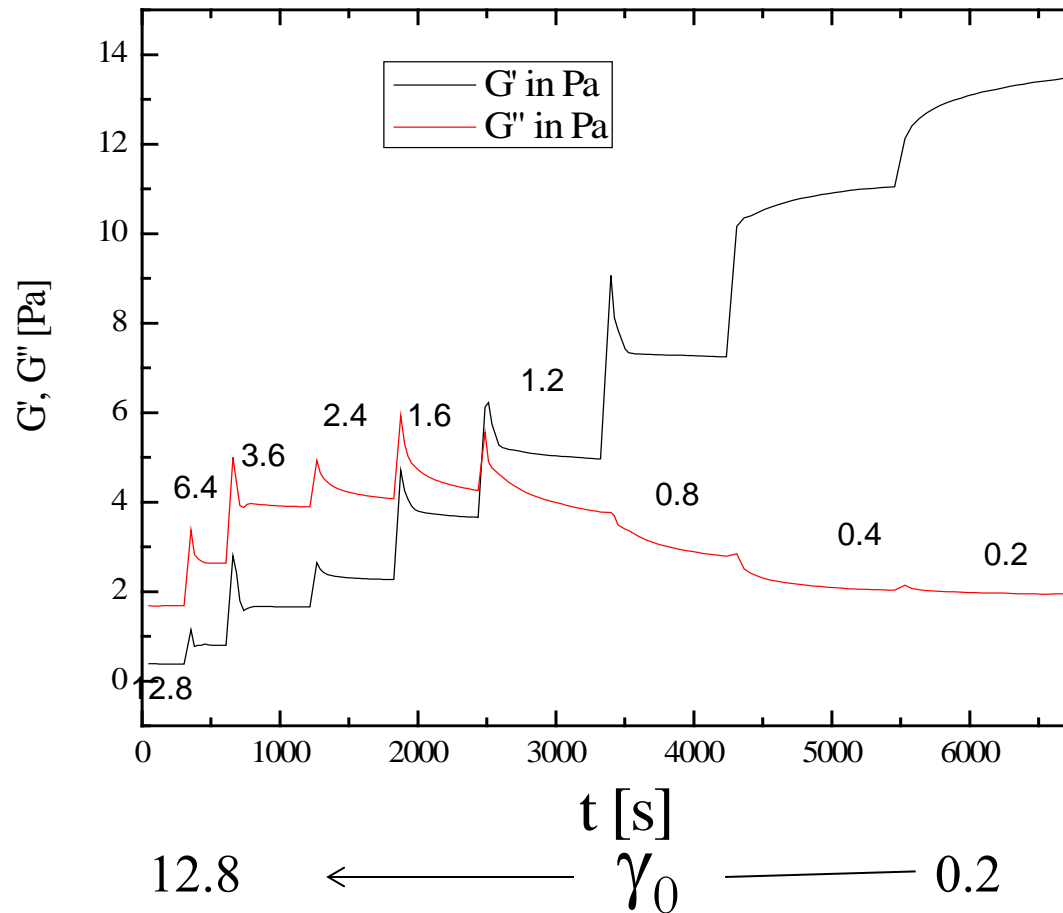
System: Gibbsite (AlOOH)

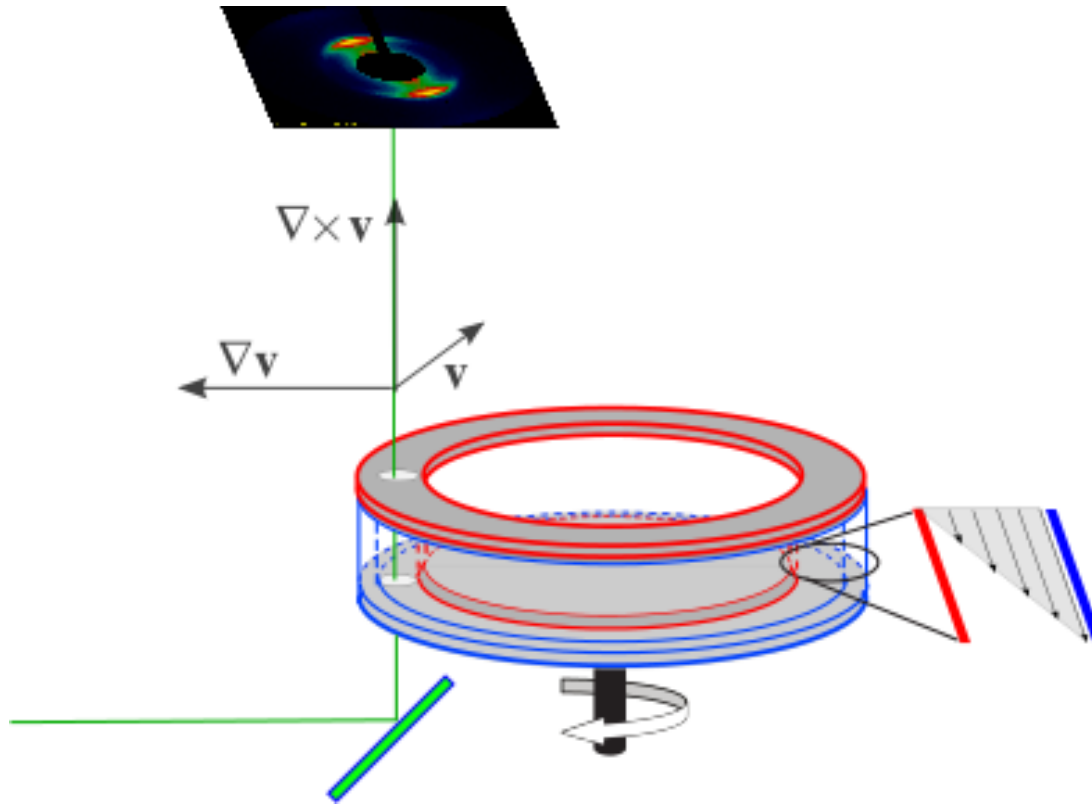
- Charged ,
- Relatively thick ($R=125 \pm 16$ nm, $d=11 \pm 4$ nm)
- Relatively monodisperse ($\sim 13\text{--}20\%$)
- Sides and faces carry the same charges (positive)
- Dispersed in glycerol



How liquid is a liquid crystal?

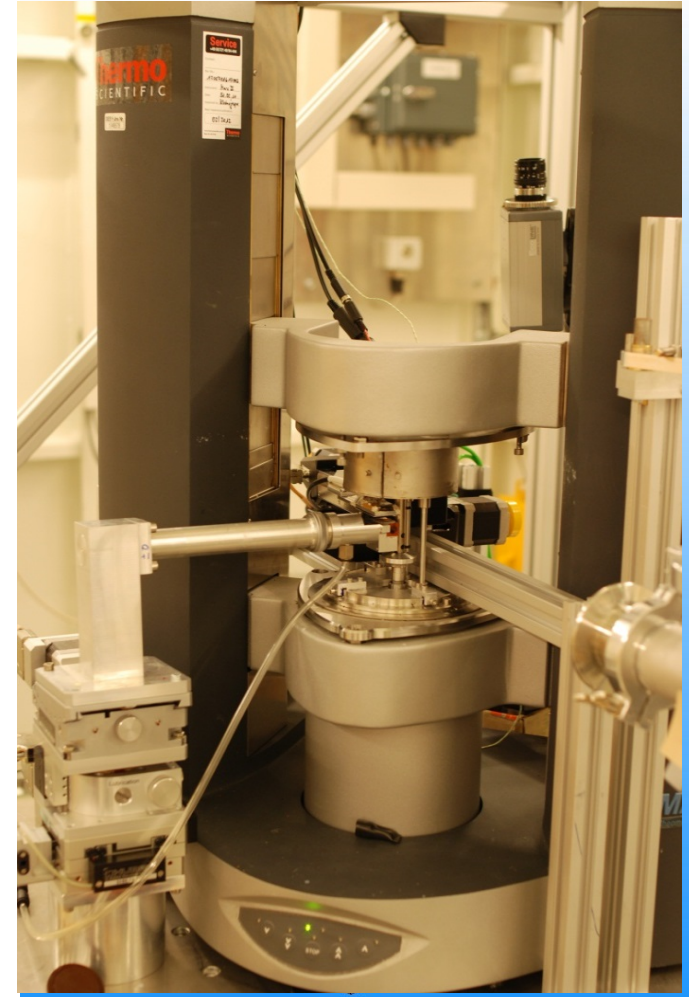
The Experiment: strain sweep at $f=0.04$ Hz



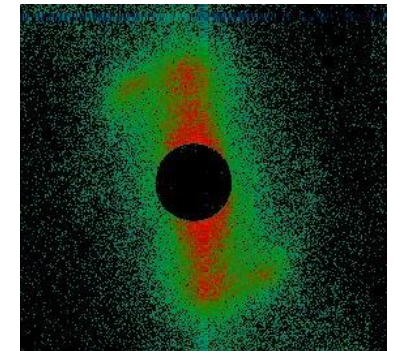
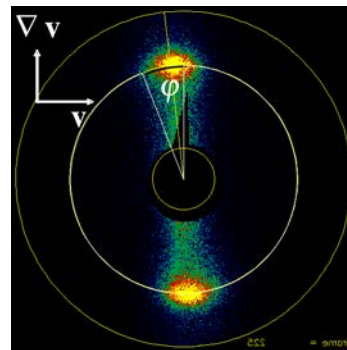
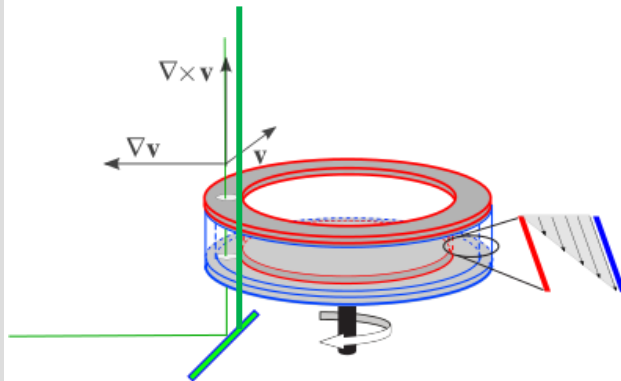
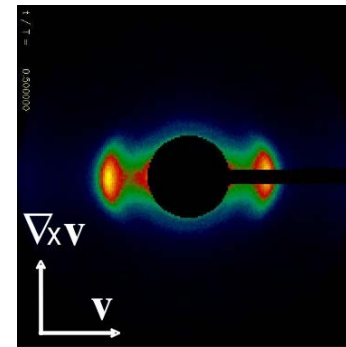
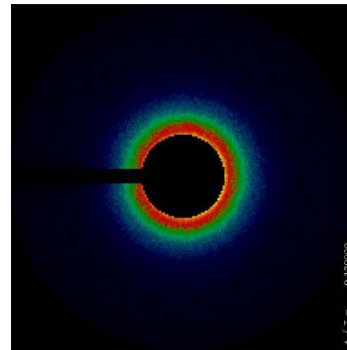
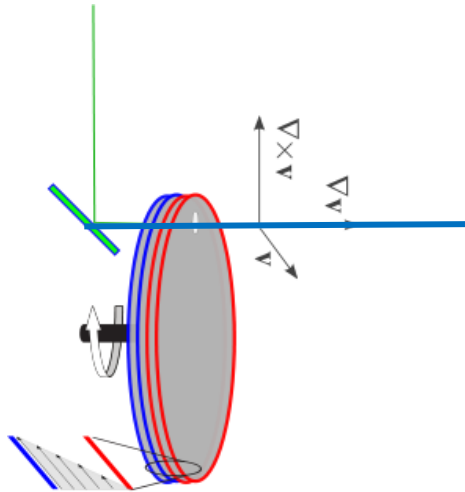
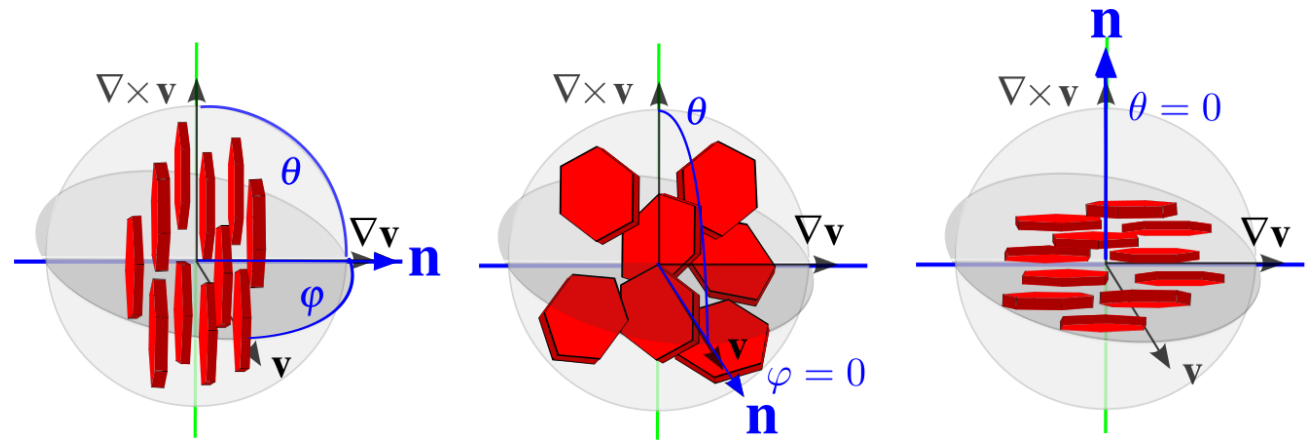


Is flow homogeneous?

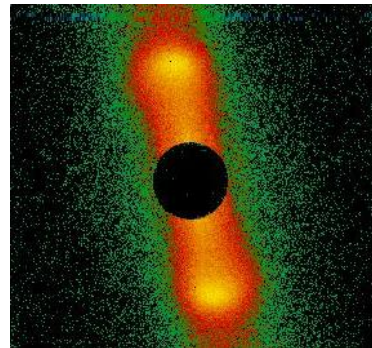
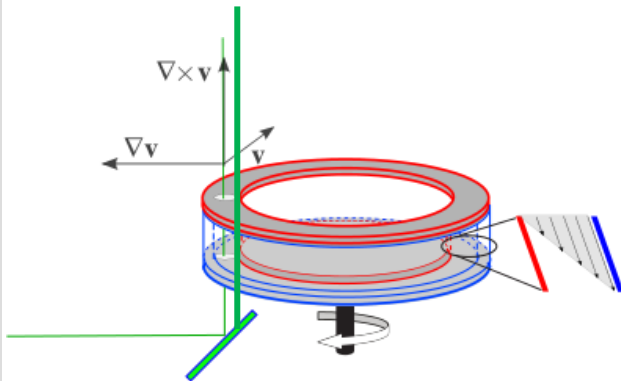
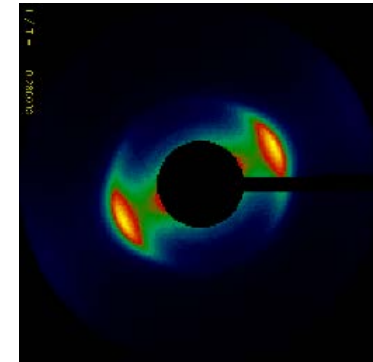
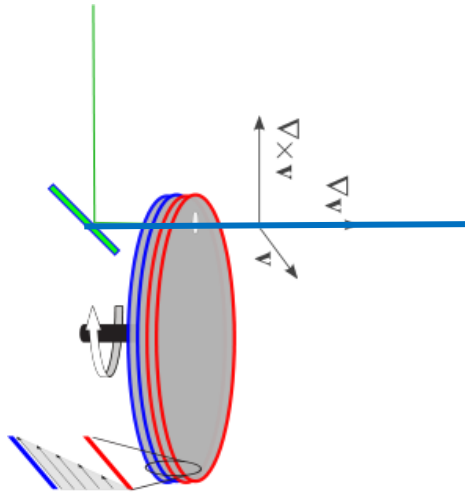
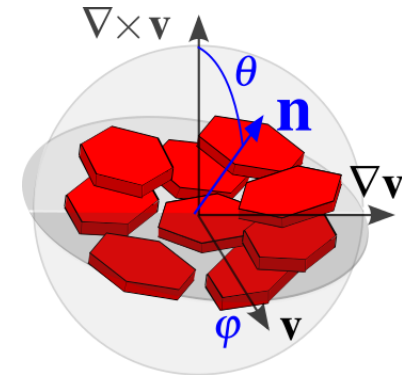
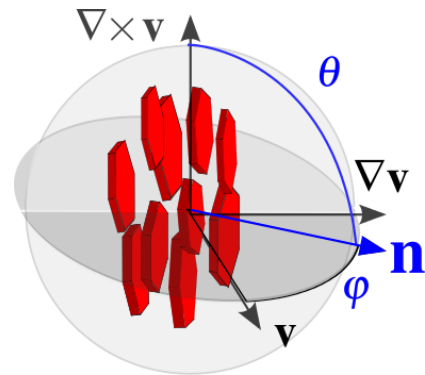
What is the total 3D motion?



Possible configurations



Possible configurations



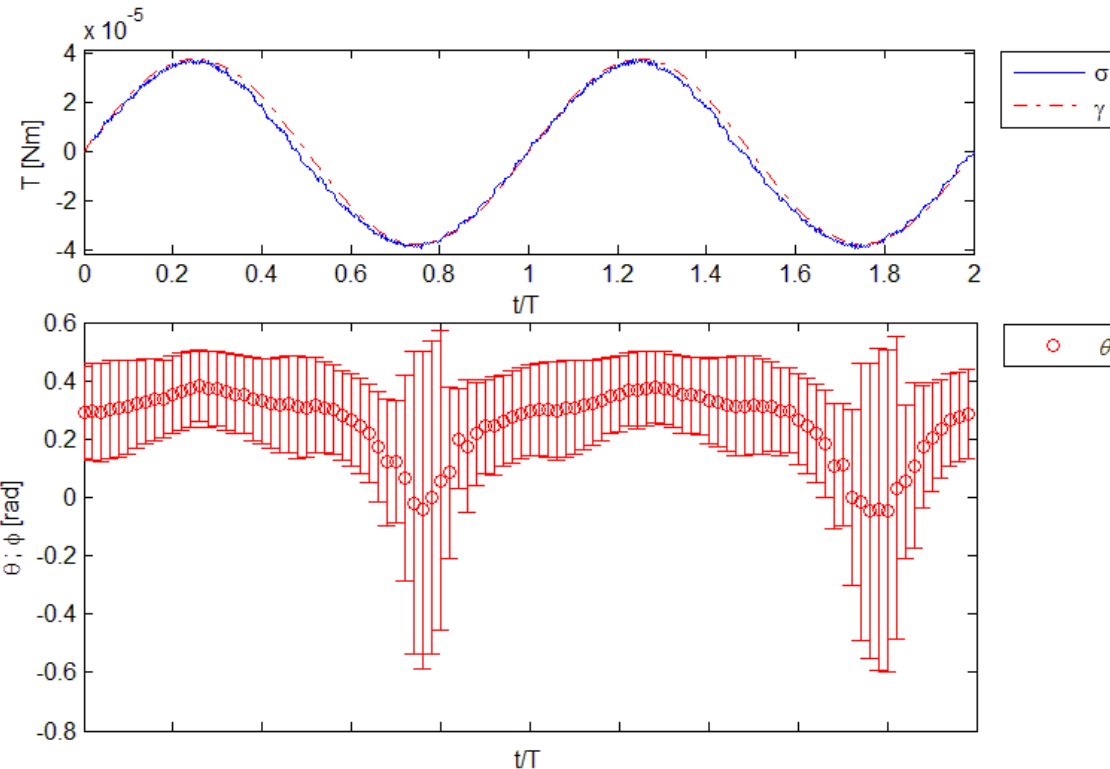
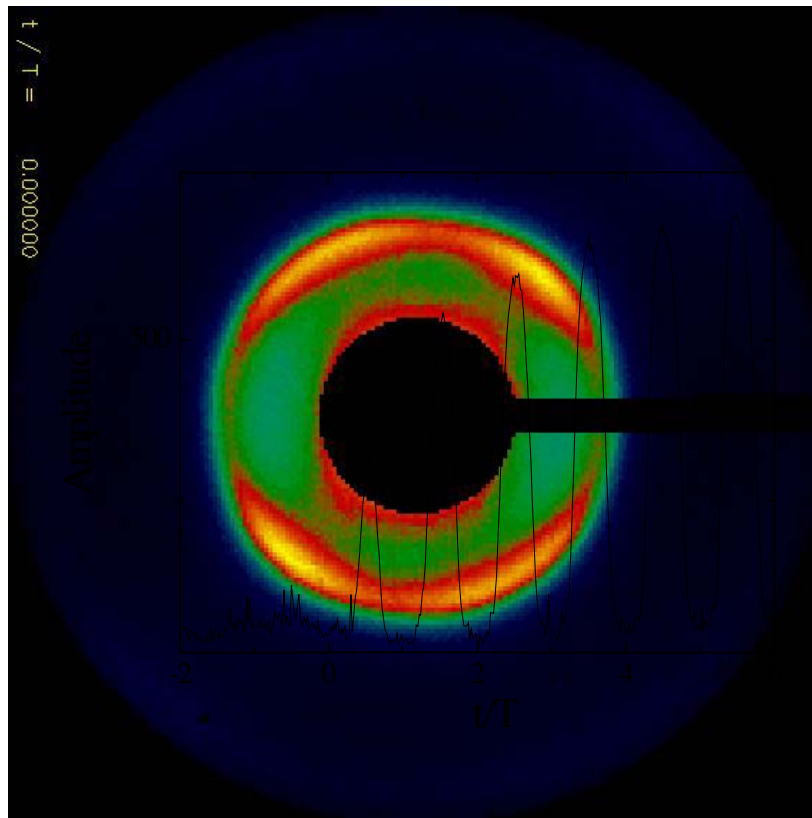
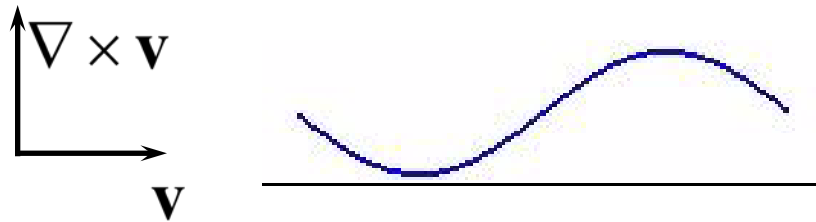
Low strain response: $\gamma=0.4$

1st harmonic response →

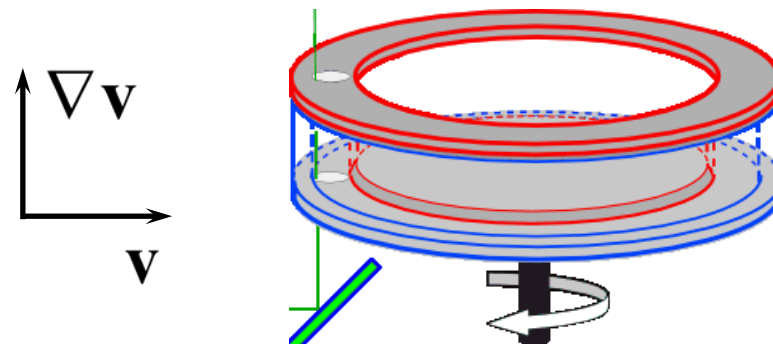
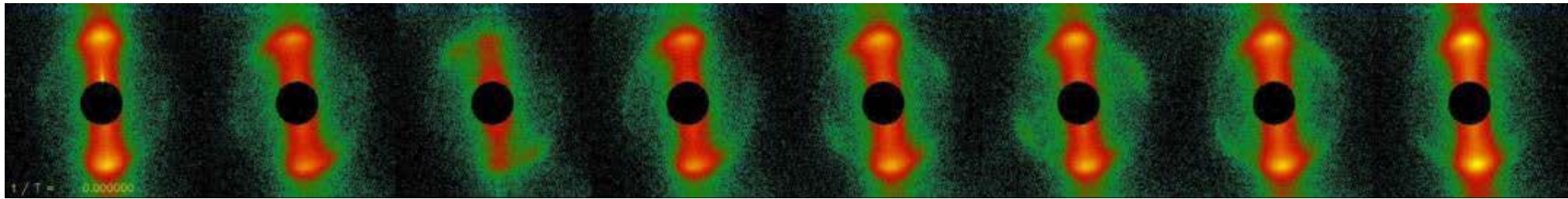
Preferential direction of deformation →

Symmetry breaking →

Dynamic bifurcation

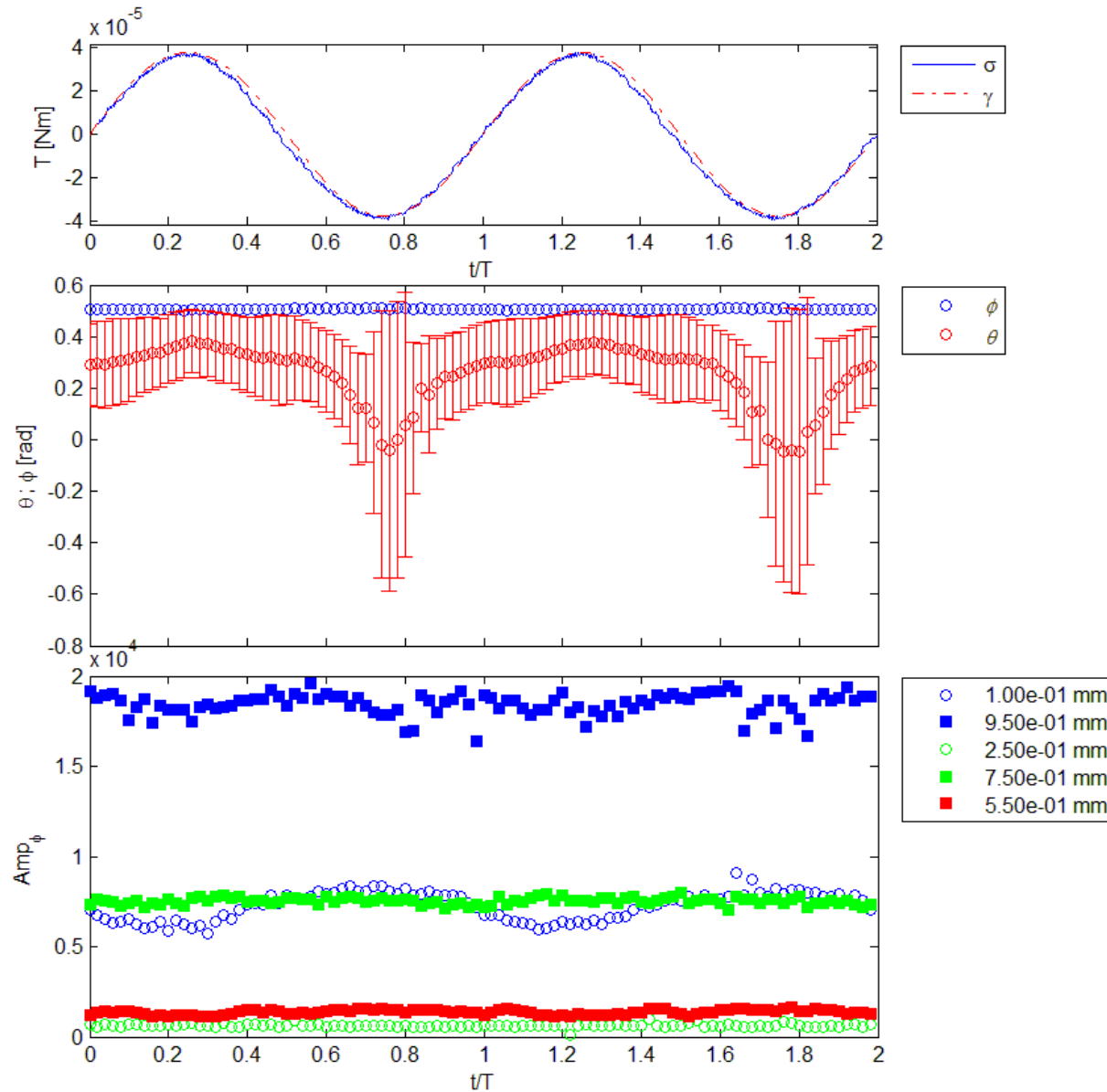
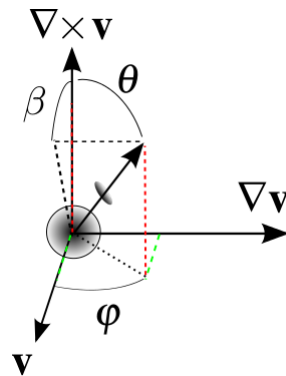


Low strain response: $\gamma=0.4$



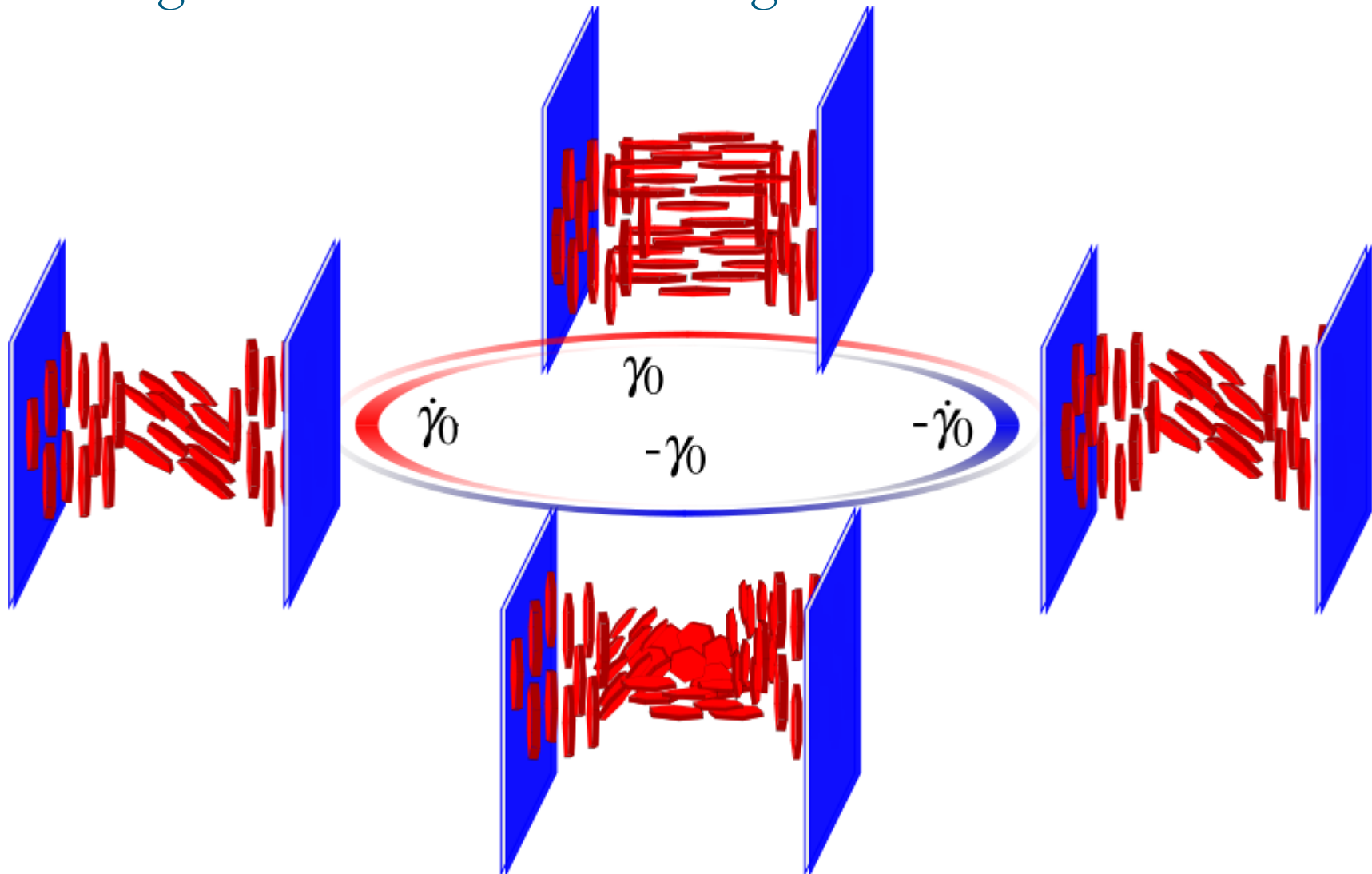
Low strain response: $\gamma=0.4$

No clear structure in
middle of gap

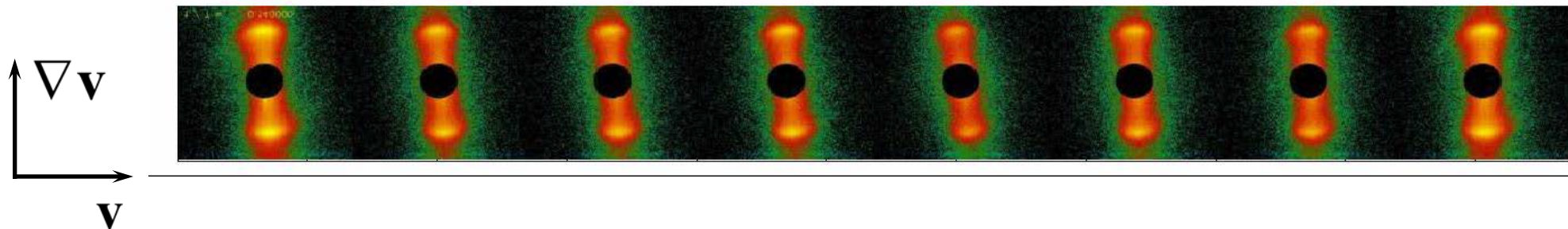
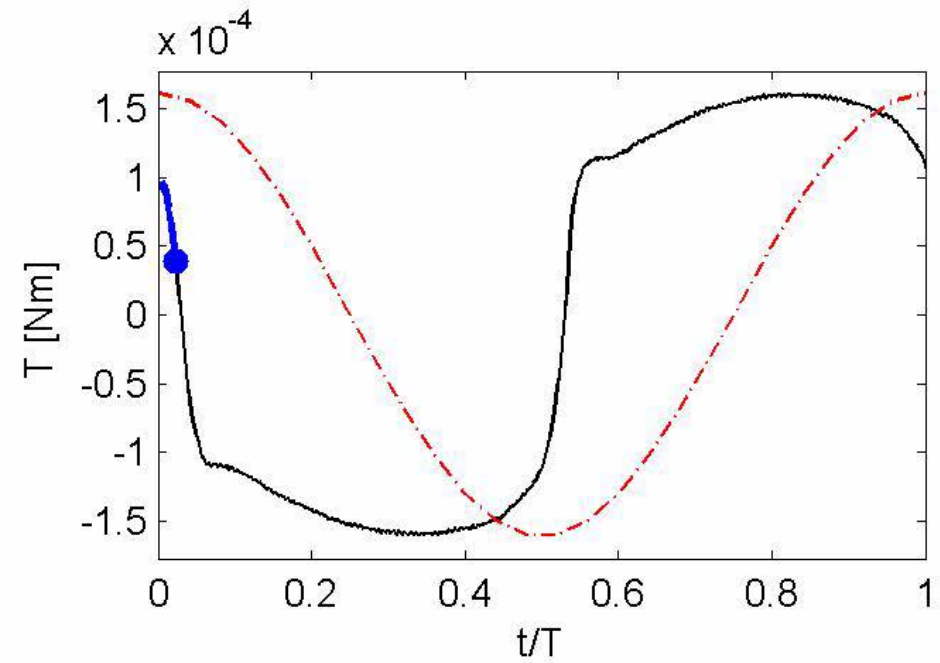
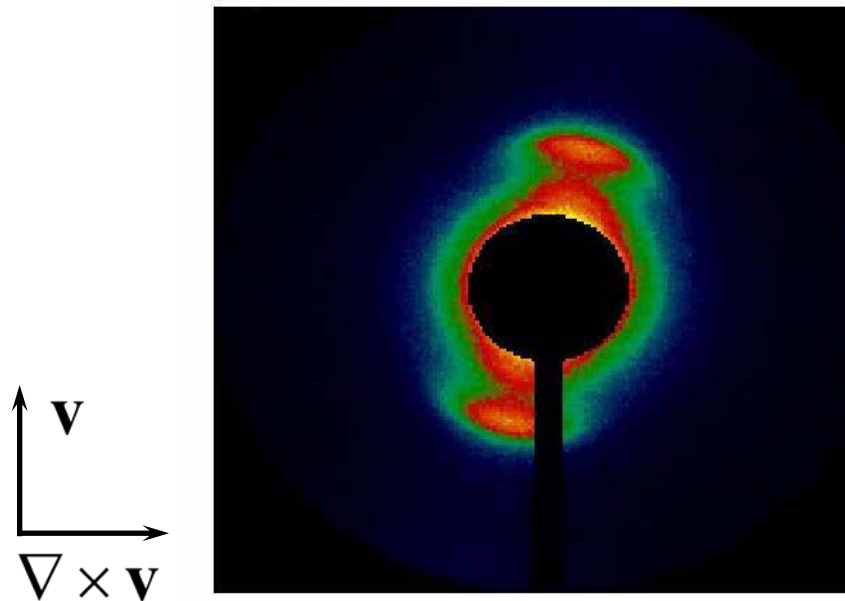


Cartoon low strain response:

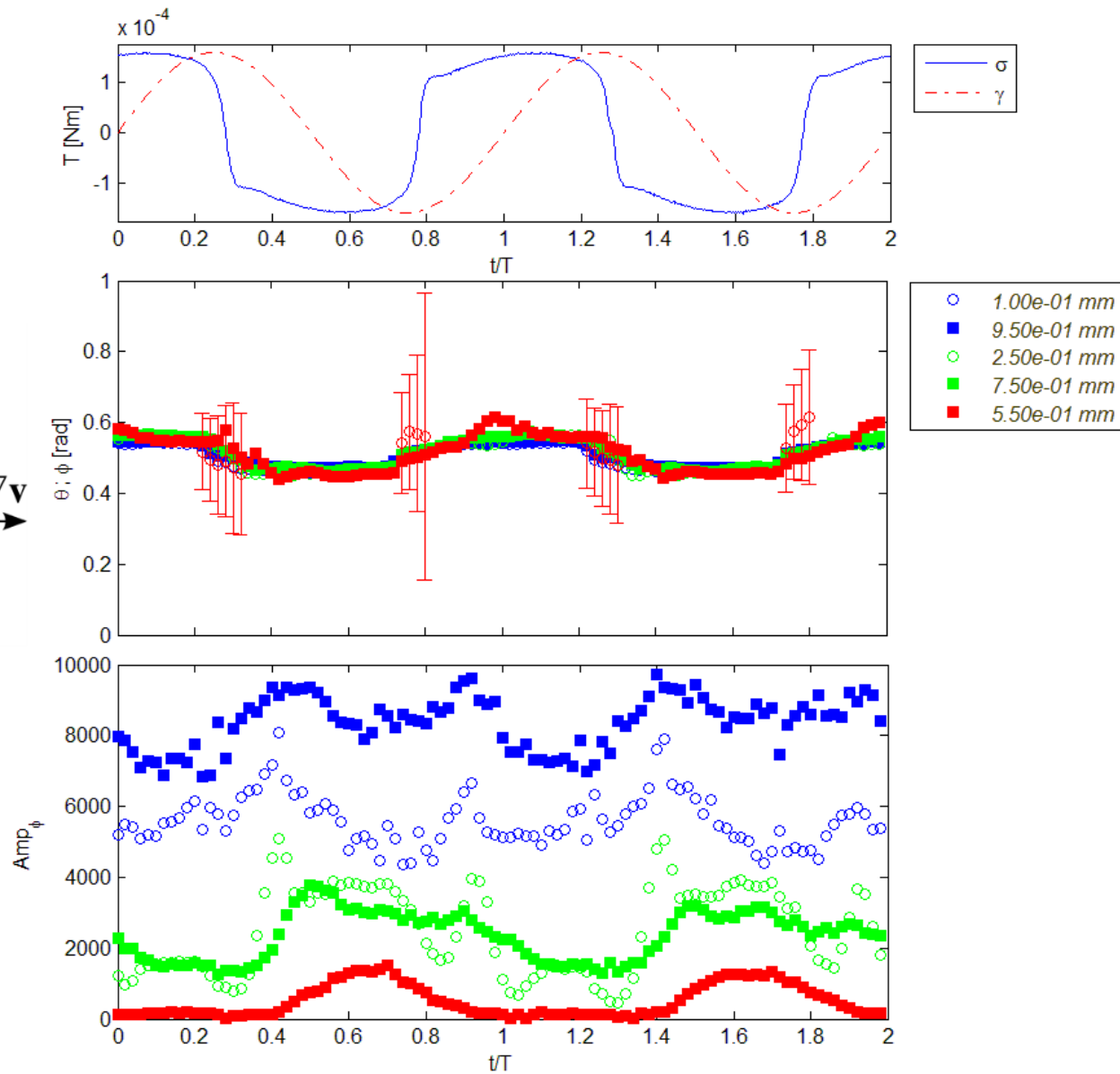
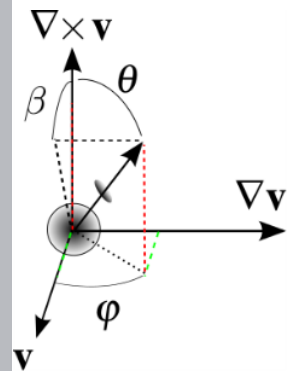
- 1st harmonic response = Dynamic bifurcation
- Huge effect of wall anchoring



High strain response: $\gamma=12.8$

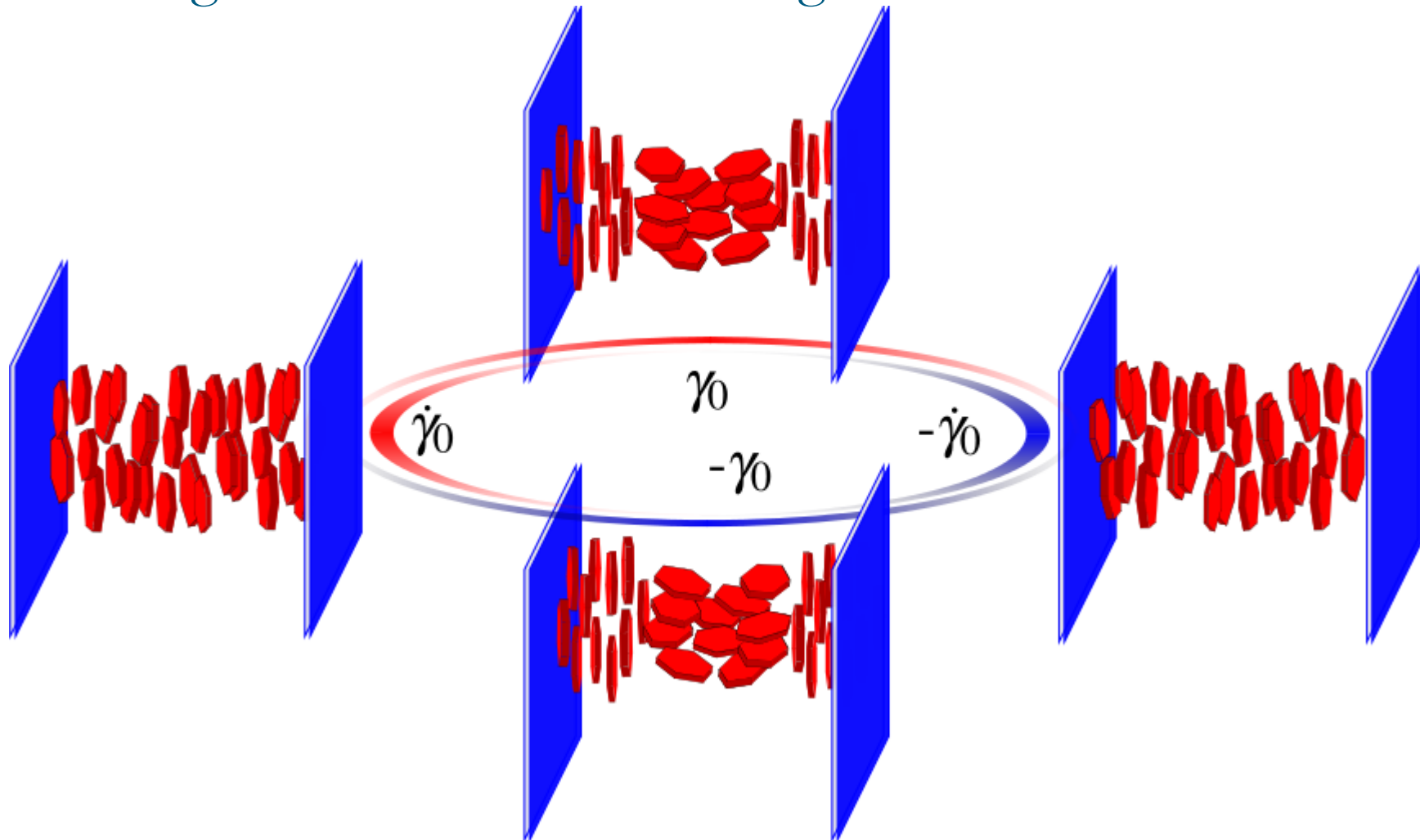


High strain response: $\gamma=12.8$



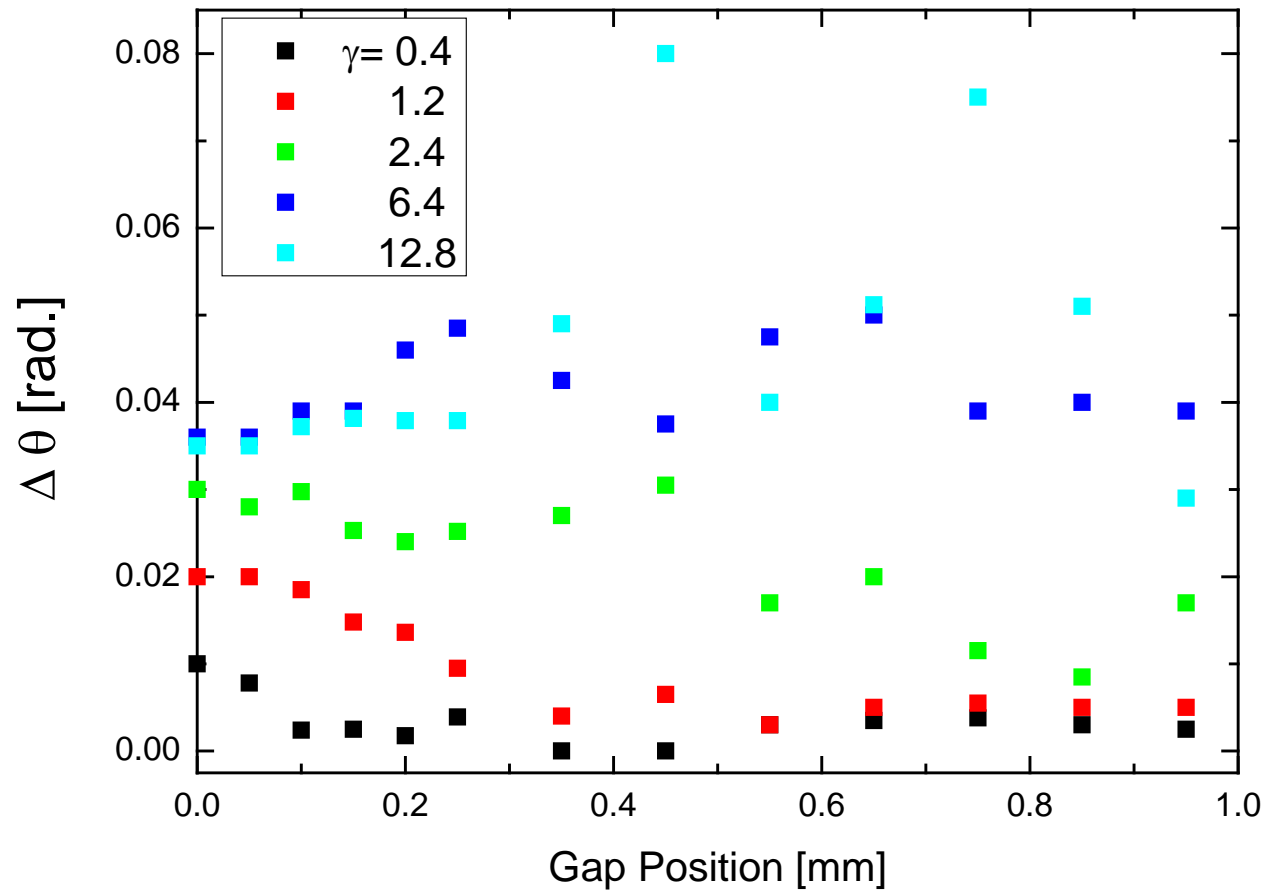
Cartoon high strain response:

- 2nd harmonic response
- Widening followed by flipping
- Huge effect of wall anchoring



How soft is the inside of a shear cell?

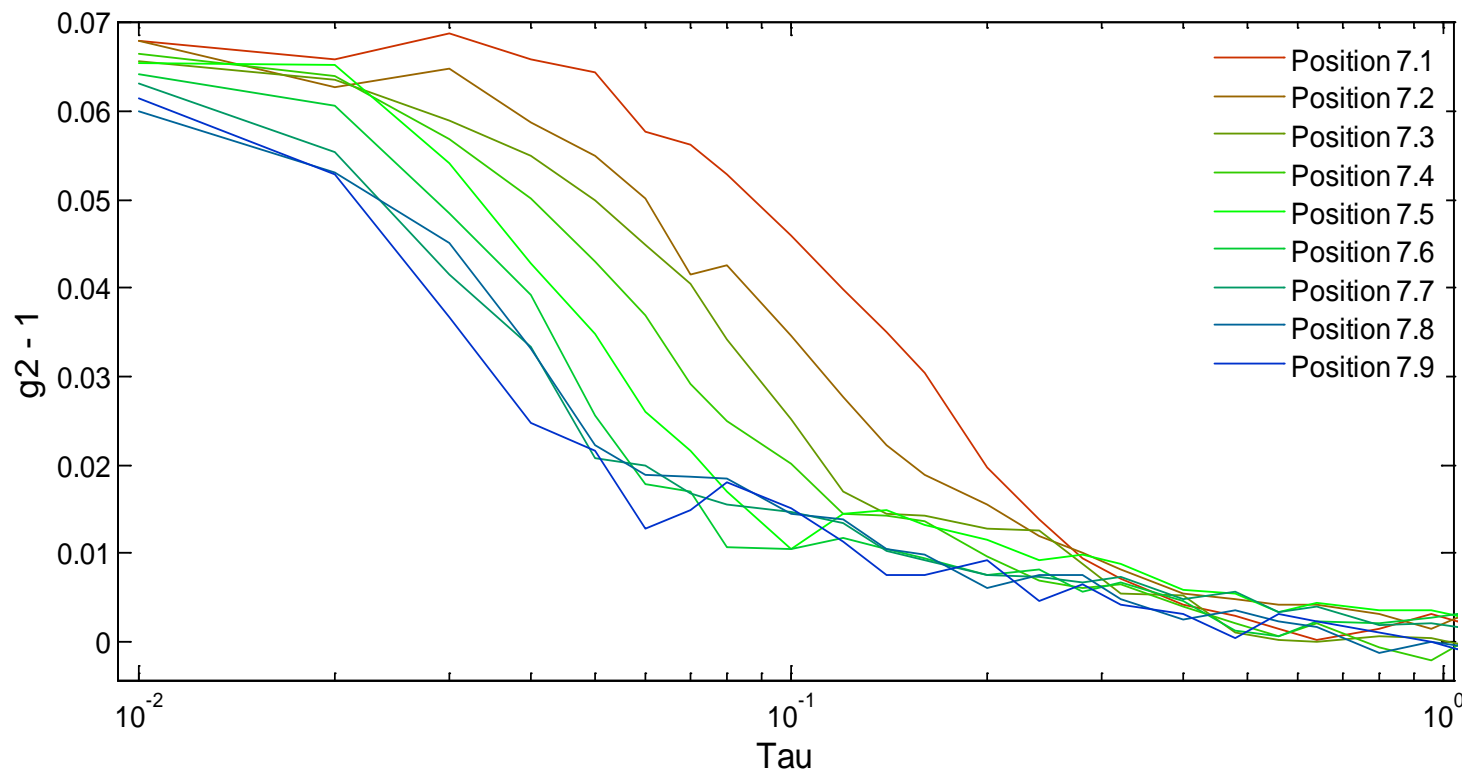
Wall anchoring vs Director motion



- Use the coherence of the squeezed X-rays AND very fast

New Lambda Detector!

Silica particles in highly viscous medium, at shear rate 0.005 s^{-1}



Conclusions Part III

- Low strain: Hookean stress response, **but**

structure response has same frequency as shear field →

Dynamic bifurcation

- High strain: normal double frequency response of structure; on average flow alignment, but system flips at flow reversal.

- BUT: response depends highly on location in the cell:

Strong competition between wall anchoring and director tumbling

Acknowledgements:

Stroboscopic SANS:

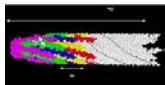
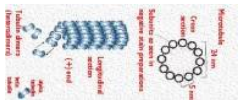
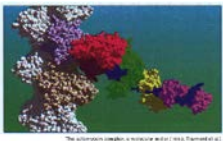
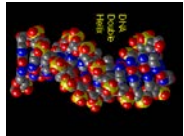
Simon Rogers, Barbara Lonetti, Joachim Kohlbrecher^{PSI}

Platelets:

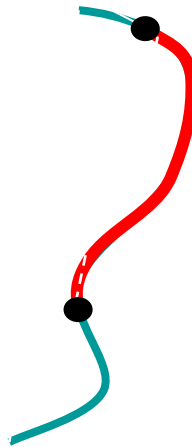
**Peter Holmqvist, Pierre Ballesta, Simon Rogers,
Bernd Struth, Fabian Westermeier, Heinz Graafsma^{DESY}**

Extra: Molecular origin shear thinning using imaging

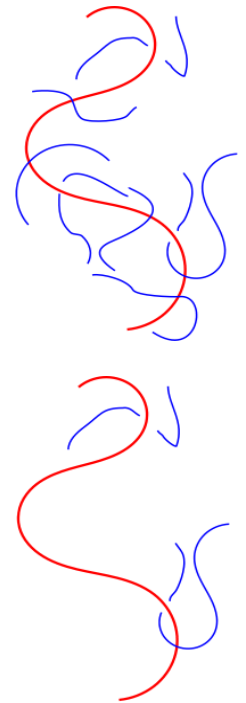
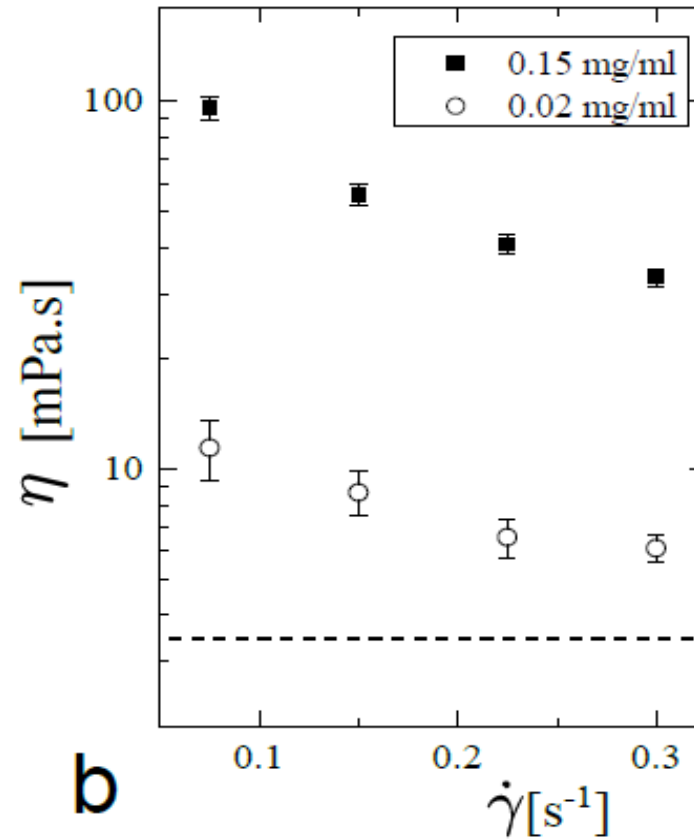
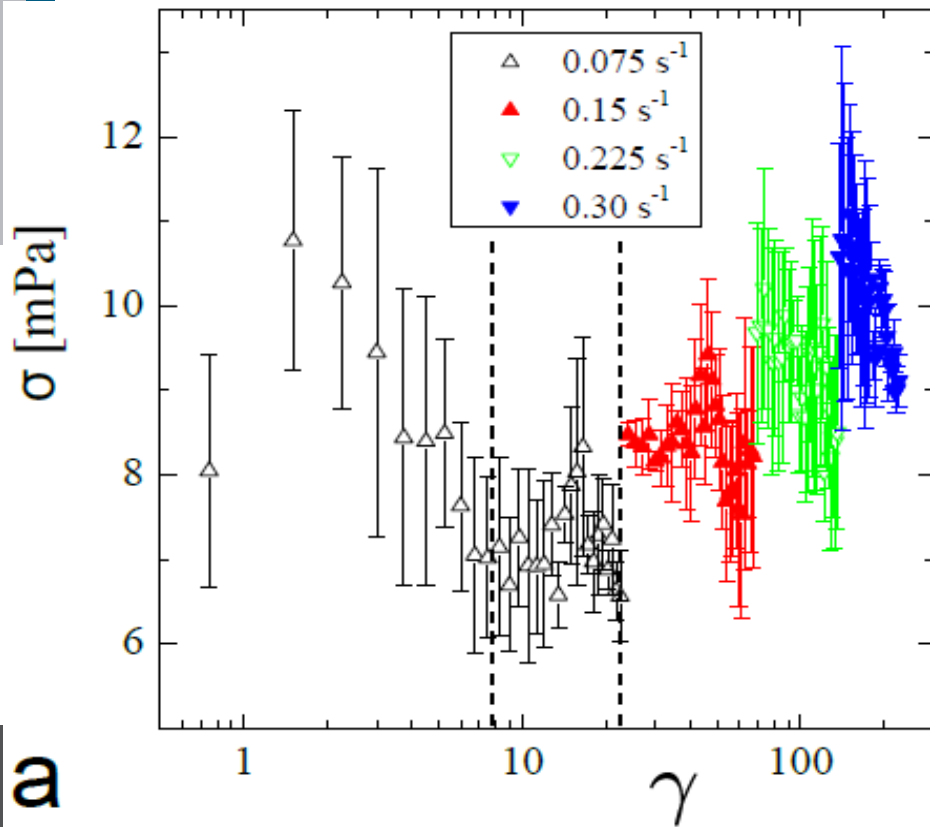
What about semi-flexible systems?



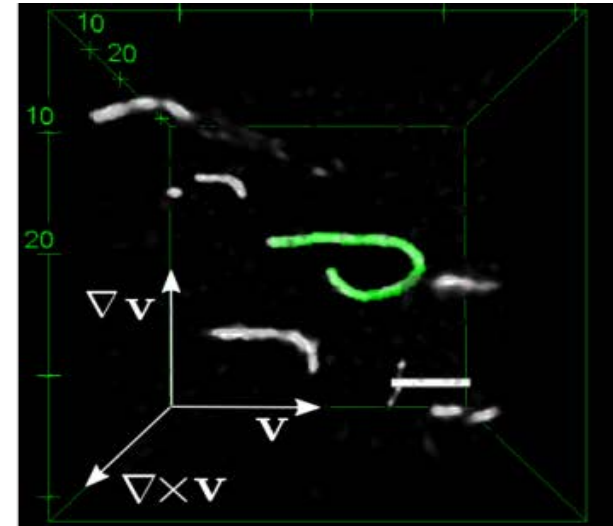
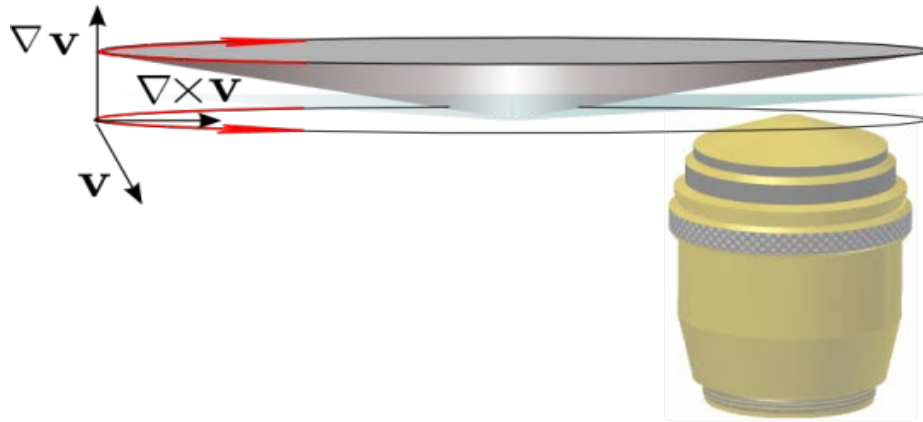
	L	D	P
DNA	variable	3 nm	50 nm
F-actin	variable	7 nm	17 μm
Microtubuli	variable		>>
Fd virus	880 nm	6.6 nm	1.8 μm



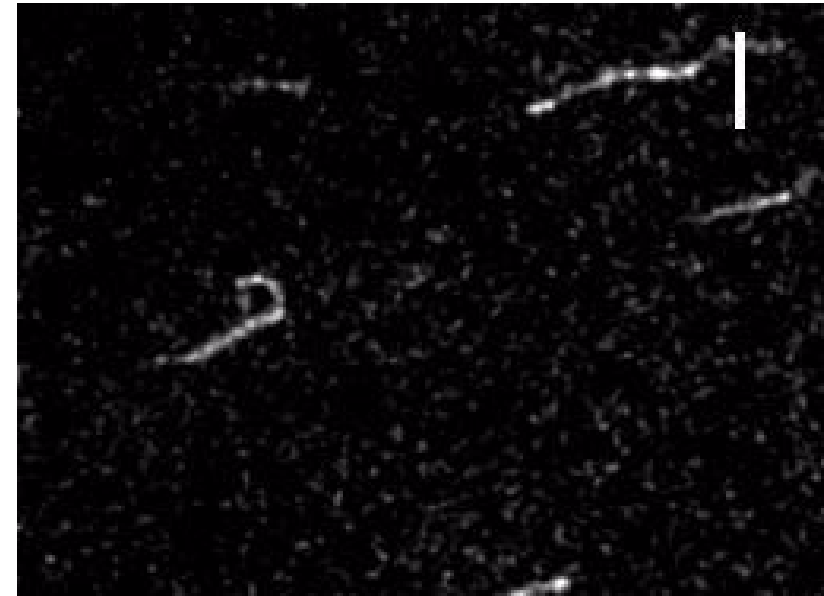
Rheological response of F-actin dispersions



Fast confocal microscopy on F-Actin

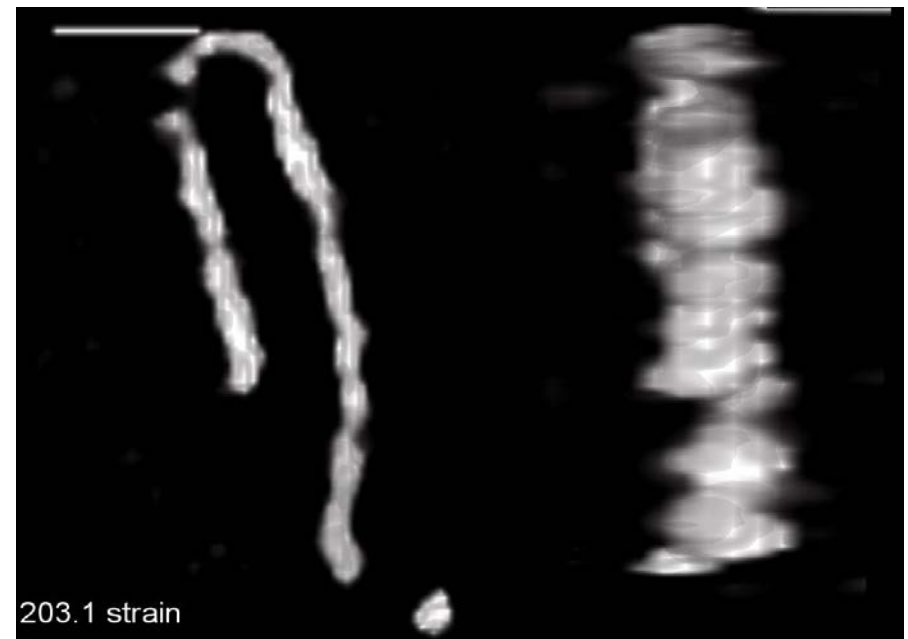
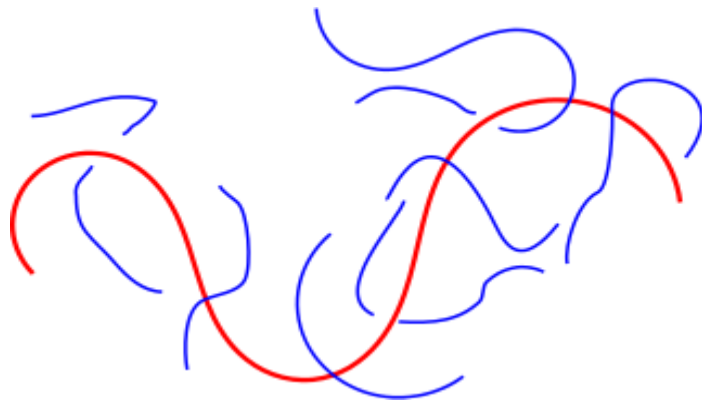
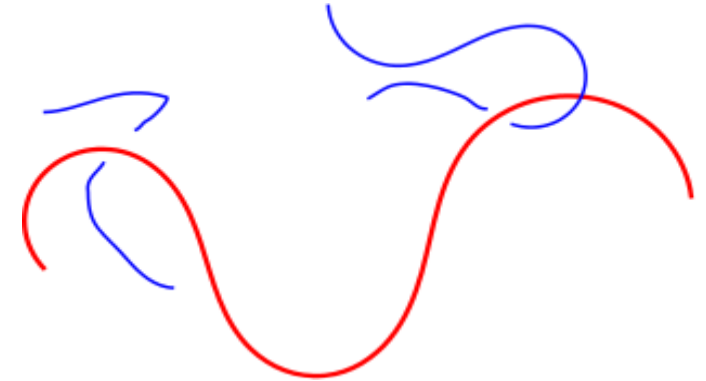


- Image at increasing strain
- Use three concentrations, label 1 per 100 filaments
 $L_c > 21 \mu\text{m}$
- About 100 analyzed filaments per combination

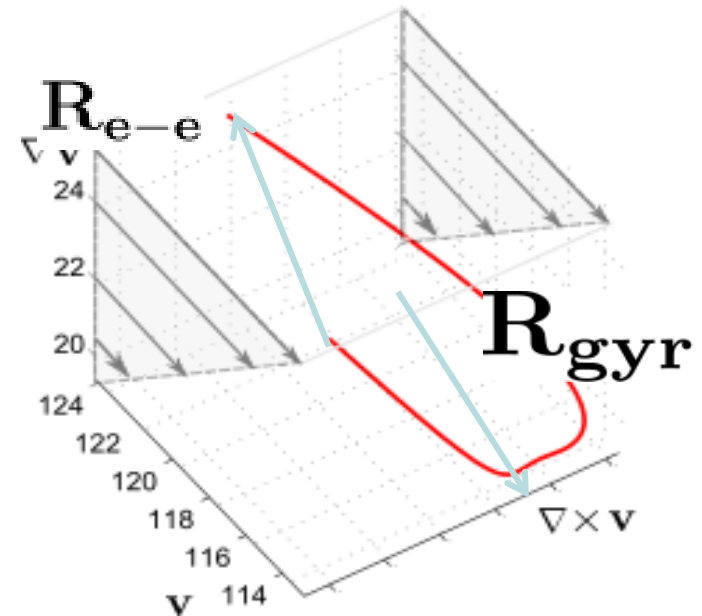
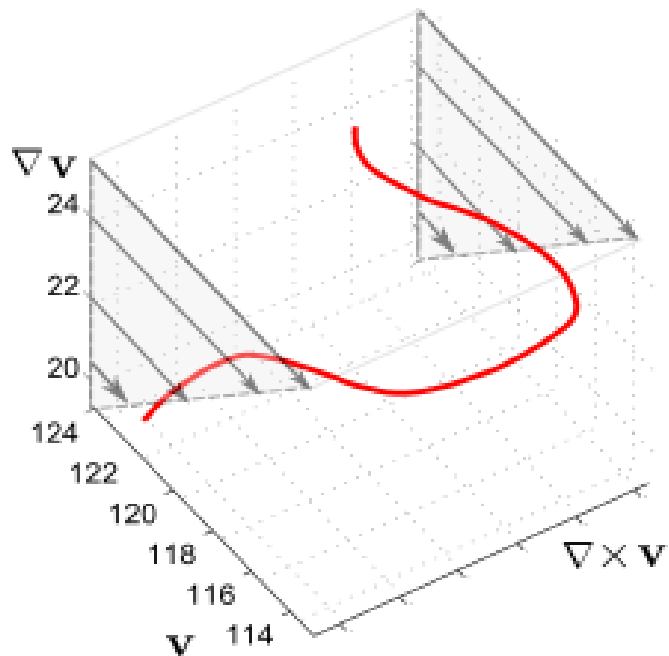


Space bar 10 μm

Sheared F-Actin in 3-D



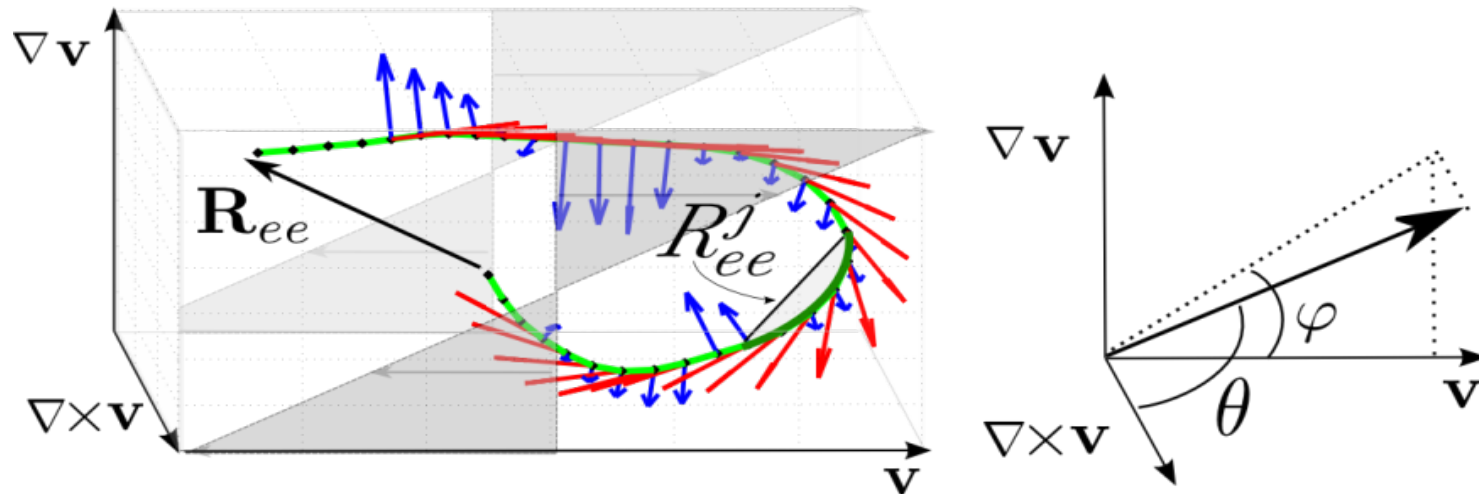
How to analyze?



End-to-end vector \mathbf{R}_{e-e} ; NOT the relevant parameter

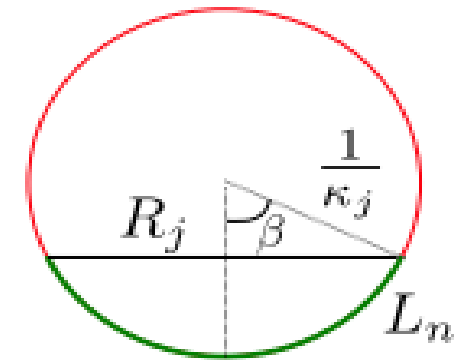
Inertia tensor \mathbf{R}_{gyr} might NOT be the relevant parameter

Analyze local bending and stretching:

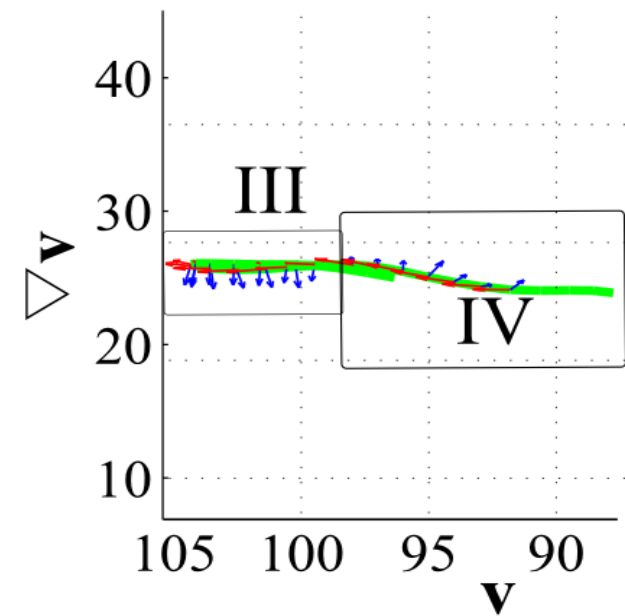
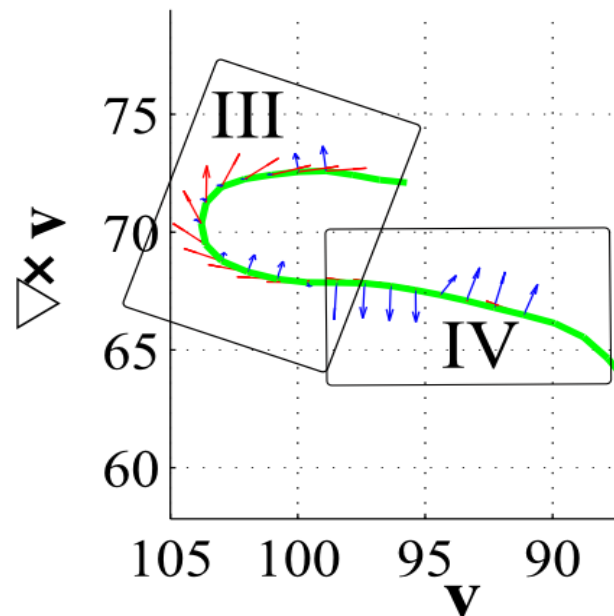
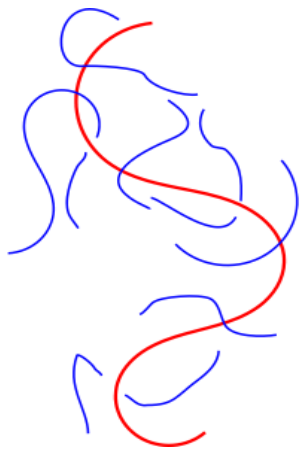
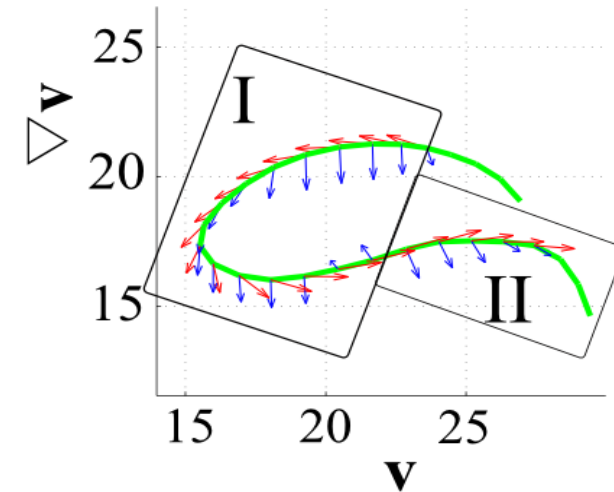
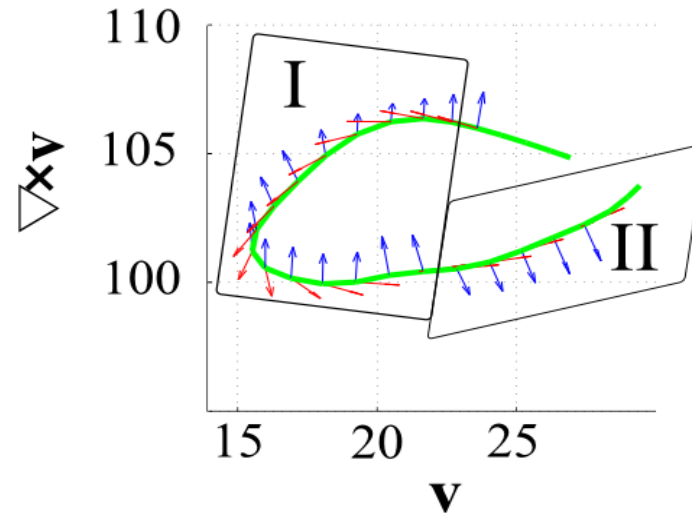
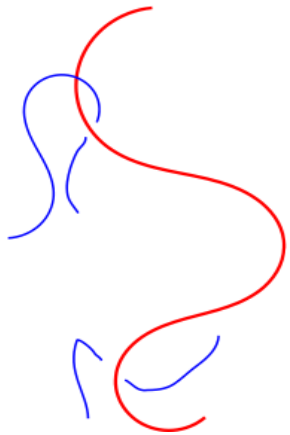


$$x_t(\kappa_t) = \frac{2 \sin\left(\frac{\pi \kappa_t \langle L_n \rangle}{2}\right)}{\pi \kappa_t \langle L_n \rangle}.$$

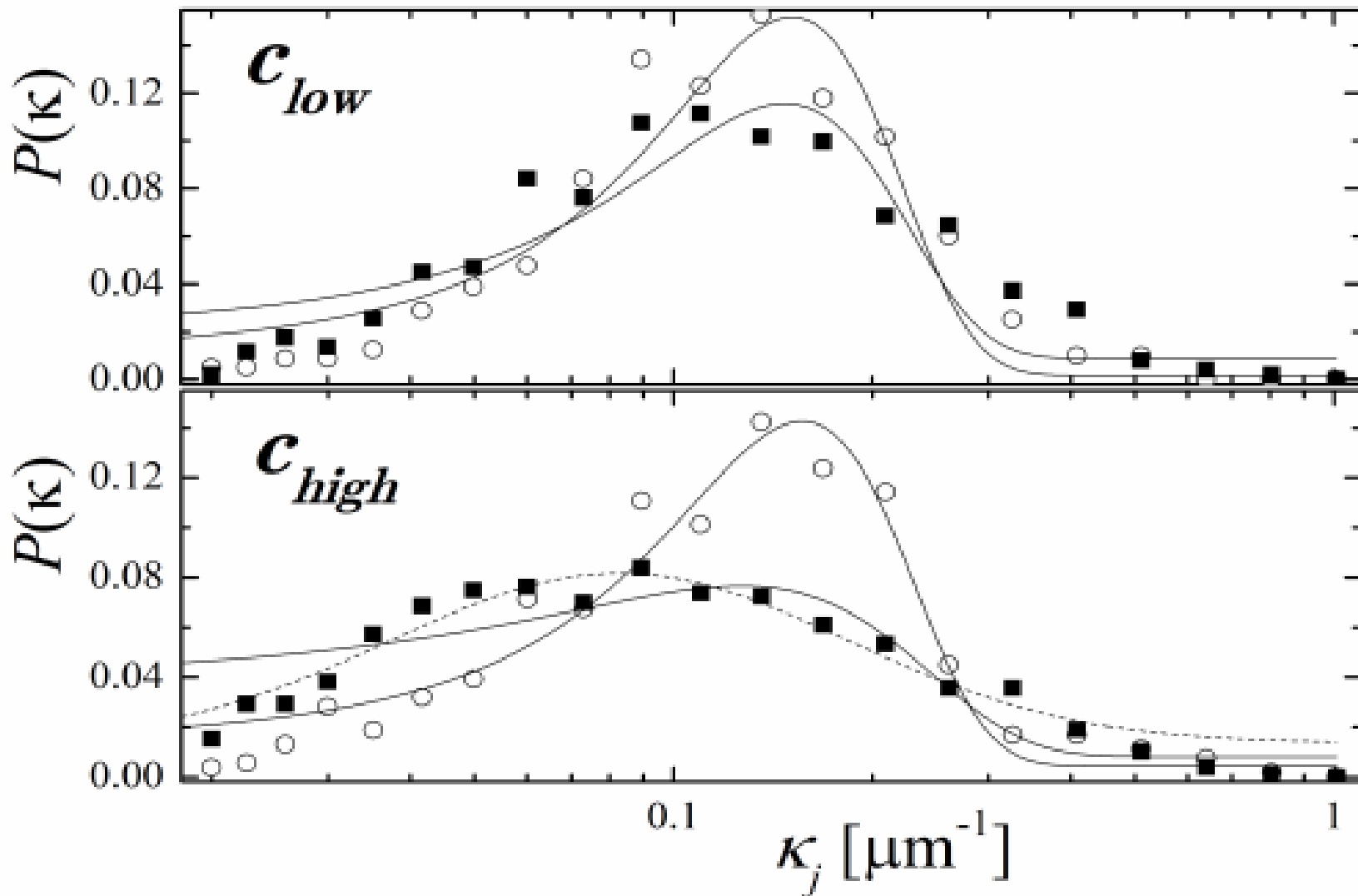
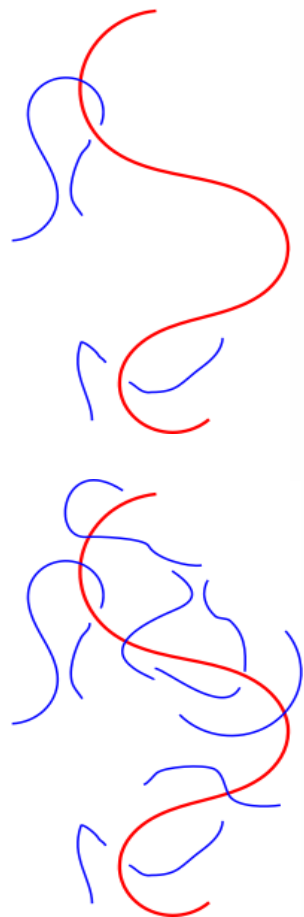
$$\hat{T}_j \equiv \frac{\dot{\mathbf{r}}_j}{|\dot{\mathbf{r}}_j|}; \hat{B}_j \equiv \frac{\dot{\mathbf{r}}_j \times \ddot{\mathbf{r}}_j}{|\dot{\mathbf{r}}_j \times \ddot{\mathbf{r}}_j|}; \kappa_j = \frac{|\dot{\mathbf{r}}_j \times \ddot{\mathbf{r}}_j|}{|\dot{\mathbf{r}}_j|^3}$$



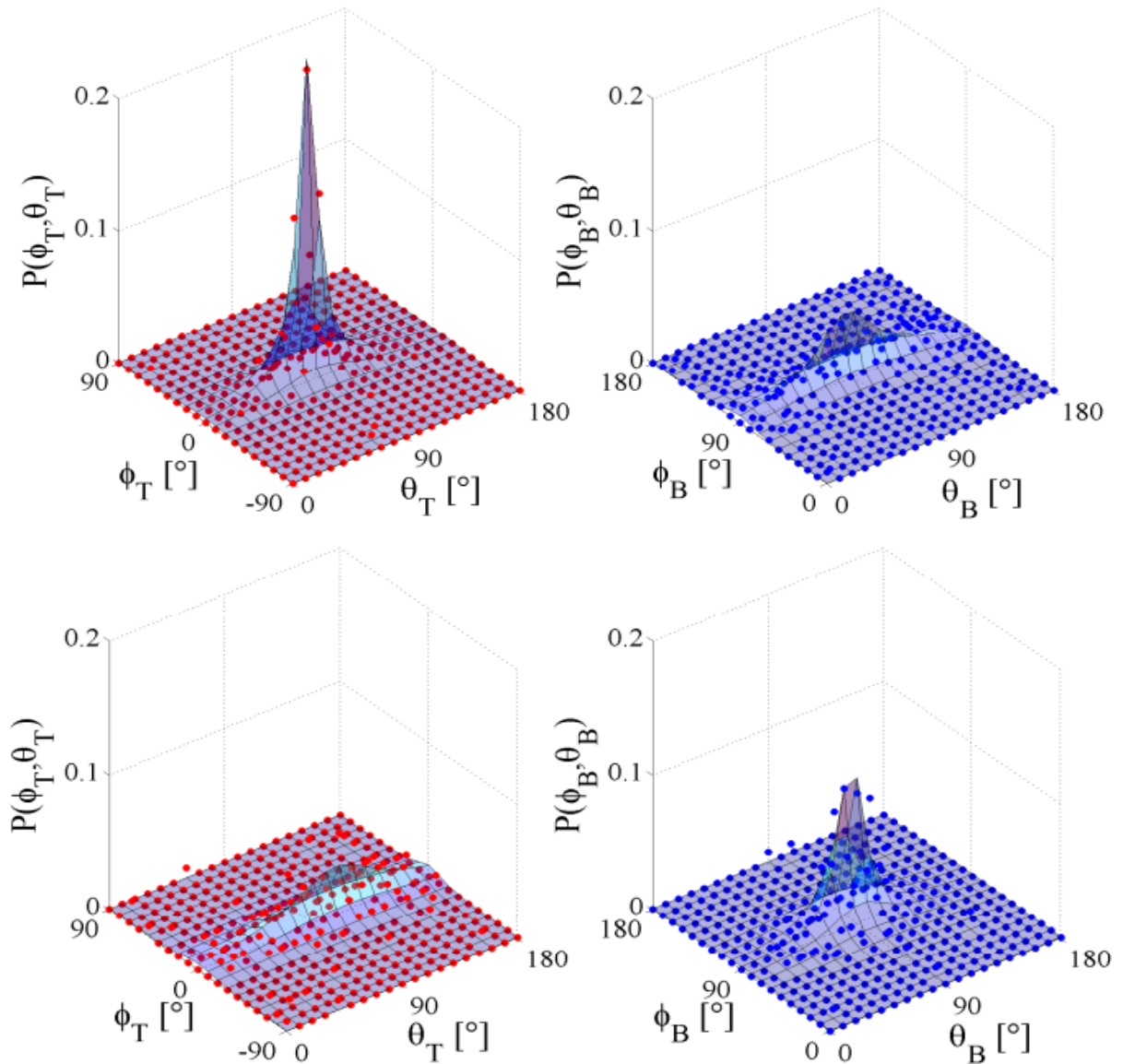
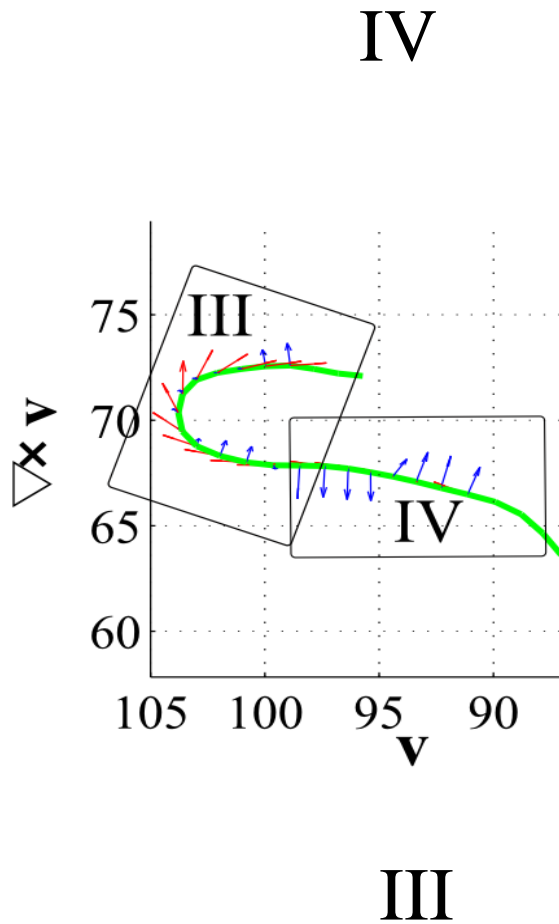
Typical examples:



Distribution of curvatures:



Distribution of angles



Characterizing parameters

$$f(\theta, \phi) = a / \left(\left(\frac{\theta - \Delta\theta}{w_\theta} \right)^2 + \left(\frac{\phi - \Delta\phi}{w_\phi} \right)^2 + 1 \right)$$

$$\bar{S}_T = \int_0^\pi \int_0^{2\pi} d\theta d\phi \sin(\phi) f(\theta_T, \phi_T) \hat{T} \hat{T}$$

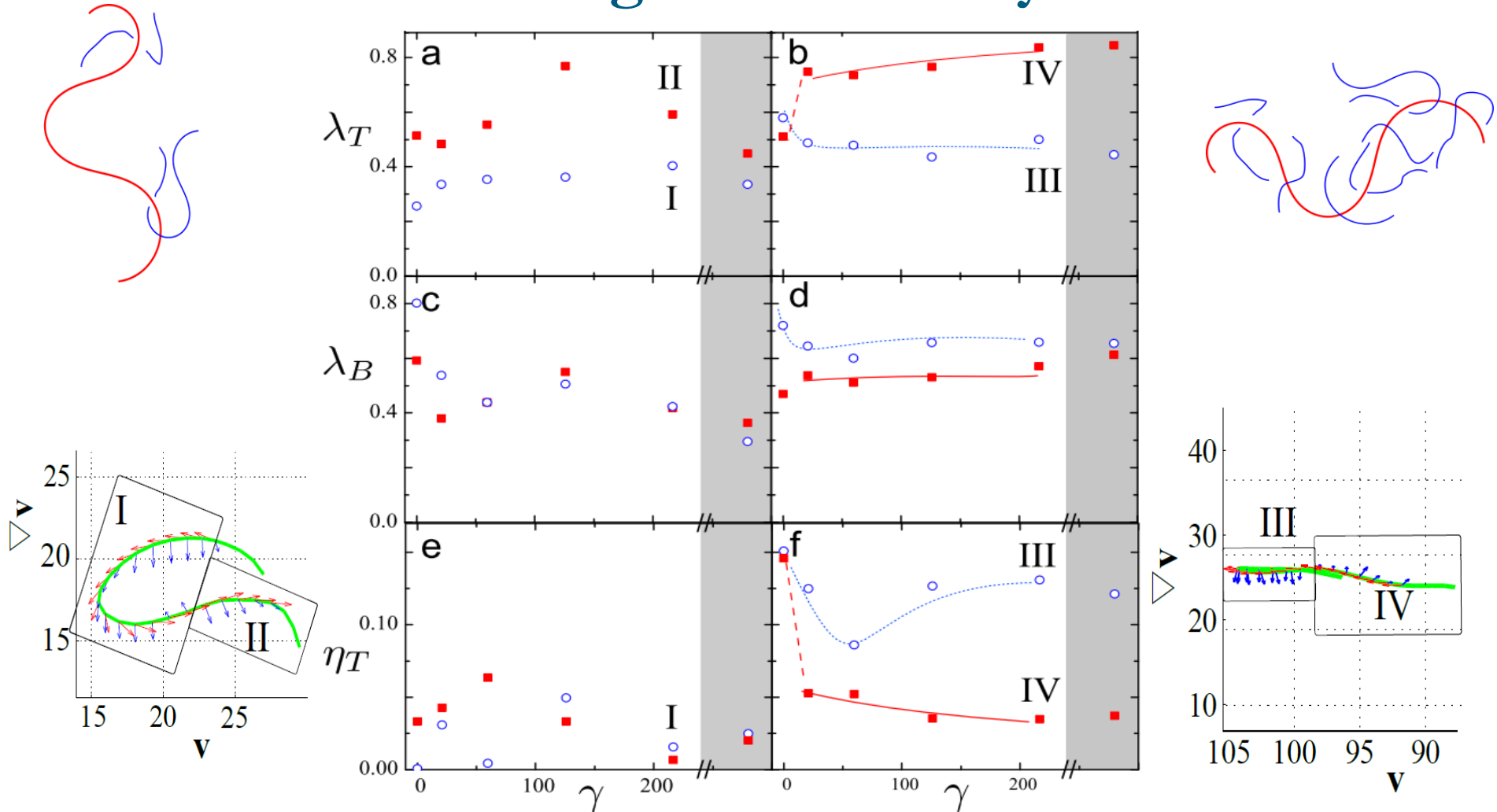
$$\bar{Q} = \frac{1}{2} (3\bar{S} - \mathbf{I})$$

biaxiality

$$\bar{Q}_{T,B} = \begin{pmatrix} -\frac{1}{2}\lambda_{T,B} - \boxed{\eta_{T,B}} & 0 & 0 \\ 0 & -\frac{1}{2}\lambda_{T,B} + \eta_{T,B} & 0 \\ 0 & 0 & \boxed{\lambda_{T,B}} \end{pmatrix}$$

Highest order parameter

Ordering and biaxiality



- We obtain stored energy from microscopic data by analysis on smallest lengthscale
- Alignment higher at high concentration and more along the flow vector
- Orientation does not show strong dependence due to slow relaxation after loading
- Effect of strain mainly in curvature, less in orientation

