

Cyclotrons for Ions

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The mass unit of a nucleon is defined with: $m_0 = m(C_{12}^{6+}) / 12$

mass m and charge q with dimensionless quantities:

$$m \equiv A m_0, \quad q \equiv Z e \quad (\text{for protons: } A=1.007)$$

Rest energy E_0 of a nucleon : $E_0 = m_0 c^2 = 931.5 \text{ MeV}$ (938.27 MeV for protons)

kinetic energy of an ion : $(E/A) = \text{energy/nucleon}$

For numerical calculations in cyclotrons only **2 basic numbers** are needed:

$$c \equiv 299.792458 \text{ m}/\mu\text{s}$$

$$U_0 \equiv (E_0/e) = 931.5 \text{ MV} \quad (=> \text{the unit charge } e \text{ is not needed!})$$

It is helpful to use the following derived constants:

$$(1) \quad \left(\frac{c}{2\pi}\right) = 47.71 \text{ MHz} \cdot \text{m}$$

$$(2) \quad \left(\frac{e}{2\pi m_0}\right) = \frac{c^2}{2\pi U_0} = 15.356 \text{ MHz} / \text{T} \quad (\text{reference for cyclotron frequency})$$

$$(3) \quad (B\rho)_0 \equiv \left(\frac{m_0 c}{e}\right) \equiv \frac{U_0}{c} = 3.107 \text{ Tm} \quad (\text{reference for magnetic rigidity } B\rho)$$

$$(4) \quad \xi \equiv \frac{e^2}{2m_0} \equiv e \frac{c^2}{2U_0} \equiv \frac{E_0}{2(B\rho)_0^2} = 48.24 \frac{\text{MeV}}{(\text{Tm})^2} \quad (\text{reference for kinetic energy})$$

$$\text{note: } (1)/(2) = (3) \quad (47.71/15.356 = 3.107)$$

$$(2) \times \pi = (4) \quad (15.356 \times \pi = 48.24)$$

For protons the corresponding constants are:

$$15.245 \text{ MHz/T}, \quad 3.130 \text{ Tm}, \quad 47.89 \text{ MeV}/(\text{Tm})^2$$

use dimensionless units for

velocity v , total Energy $E_{\text{tot}} = \gamma E_0$ and momentum p :

$$(5) \quad \beta \equiv \frac{v}{c} \quad (\text{can be parametrized as } \beta = \sin \varphi)$$

$$(6) \quad \gamma \equiv \frac{1}{\sqrt{1 - \beta^2}} \quad \left(\gamma = \frac{1}{\cos \varphi} \right)$$

$$(7) \quad \tilde{p} \equiv \beta\gamma = \frac{p}{mc}, \quad (\tilde{p}^2 = \gamma^2 - 1) \quad (\tilde{p} = \tan \varphi)$$

The two basic equations for Cyclotrons are :

$$(8) \quad p = q(B\rho), \quad (B\rho) = \text{magnetic rigidity}$$

$$(9) \quad \omega_0 = \left(\frac{q}{m\gamma} \right) \langle B \rangle, \quad \text{Larmor frequency}$$

For each energy exists an equilibrium orbit with circumference $L \equiv 2\pi R$ and revolution frequency ω_0 . This defines the average radius R , the average magnetic field $\langle B \rangle \equiv B_0(R)$ and the magnetic rigidity ($B\rho$) :

$$(10) \quad B_0(R) \equiv \frac{1}{2\pi R} \int B_z(s) ds$$

$$(11) \quad p = q(B\rho)$$

$$(12) \quad (B\rho) = B_0(R) \cdot R$$

since $\omega_0 = v/R$, ($v \equiv \beta c$)

=> scaling laws for cyclotrons :

$$13) \quad R = \beta R_\infty, \quad R_\infty \equiv \frac{c}{\omega_0}$$

$$14) \quad B_0(R) = B_{center} \gamma(R)$$

using reference value $(B\rho)_0$ gives the formula
connecting $(B\rho)$ and (E/A) :

$$(15) \quad \tilde{p} = \left(\frac{Z}{A}\right) \frac{(B\rho)}{(B\rho)_0}$$

$$(16) \quad \gamma = \sqrt{1 + \tilde{p}^2}$$

$$(17) \quad (E/A) = (\gamma - 1) E_0$$

The K_B - value of a cyclotron is given by the maximum
bending power $(B\rho)$ of an ion before extraction :

$$(18) \quad (B\rho)_{\max} = B_0(R_{\max}) \cdot R_{\max}$$

$$(19) \quad K_B \equiv \xi (B\rho)_{\max}^2, \quad \xi = 48.24 \frac{\text{MeV}}{(\text{Tm})^2}$$

for nonrelativistic energies :

$$(20) \quad E = \frac{p^2}{2m}$$

$$(21) \quad (E/A)_{nonrel} = \frac{1}{2} \left(\frac{q}{m} \right)^2 (B\rho)^2$$

and for the maximum energy :

$$(22) \quad (E/A)_{\max} = \varepsilon \equiv \left(\frac{Z}{A} \right)^2 K_B$$

the relativistically correct maximum energy is :

$$(23) \quad (E/A) = (\gamma - 1) E_0, \quad \gamma = \sqrt{1 + \frac{2\varepsilon}{E_0}}$$

series expansion for low energies :

$$(24) \quad (E/A) = \varepsilon \left[1 - \frac{1}{2} \left(\frac{\varepsilon}{E_0} \right) + \frac{1}{2} \left(\frac{\varepsilon}{E_0} \right)^2 - \dots \right]$$

inverse formula for a given energy E/A :

$$(25) \quad \gamma = 1 + \frac{(E/A)}{E_0}, \quad \tilde{p} = \sqrt{\gamma^2 - 1}, \quad \beta = \frac{\tilde{p}}{\gamma}$$

$$(26) \quad \varepsilon = \frac{E_0}{2} \tilde{p}^2 = (E/A) \left[1 + \frac{(E/A)}{2E_0} \right]$$

$$(27) \quad K_B = \varepsilon \left(\frac{A}{Z} \right)^2$$

$$(28) \quad (B\rho) = \frac{A}{Z} \tilde{p} (B\rho)_0 = \sqrt{\frac{K_B}{\xi}} = \sqrt{\frac{K_B}{48.24 \text{ MeV}}} [Tm]$$

Focusing Limit K_F :

For light ions with a high (Z/A) ratio, the vertical focusing becomes critical and requires large spiral angles for the magnet sectors.

There exists an energy limit, which depends linearly on (Z/A) and can be lower than the bending limit.

Formula (22) for the non relativistic energy limit ε is then replaced by :

$$(29) \quad (E/A)_{\max} = \varepsilon_F \equiv \left(\frac{Z}{A}\right) K_F$$

and the relativistically correct maximum energy is given as in (23)

$$(30) \quad (E/A) = (\gamma - 1) E_0, \quad \gamma = \sqrt{1 + \frac{2\varepsilon_F}{E_0}}$$

cyclotron frequency ω_0 :

$$(31) \quad \omega_0 \equiv 2\pi\nu_0 = \frac{q}{m} \frac{B_0(R)}{\gamma(R)}$$

$$(32) \quad \nu_0 = \frac{Z}{A} \left(\frac{e}{2\pi m_0} \right) \frac{B_0(R)}{\gamma(R)} = \frac{Z}{A} \frac{B_0(R)}{\gamma(R)} 15.356 \text{ MHz/T}$$

the RF - frequency can be a harmonic h of the revolution frequency ν_0 :

$$(33) \quad \nu_{RF} \equiv h\nu_0$$

the reference value R_∞ for the scaling of the radius is then :

$$(34) \quad R_\infty \equiv \frac{c}{\omega_0} = \frac{(c/2\pi) \cdot h}{\nu_{RF}} = h \frac{47.71 \text{ MHz} \cdot \text{m}}{\nu_{RF}}$$

$$R = \beta R_\infty$$

to have low losses at extraction it is important to have a large turn separation dR/dn

$$(35) \quad \frac{dR}{dn} = \frac{\gamma}{\gamma + 1} R \frac{E_G}{(E/A)} \frac{f^2}{v_r^2}, \quad E_G = \frac{Z}{A} e \hat{V}$$

\hat{V} = peak voltage/turn, v_r = radial tune, $f = \frac{\nu_r(\text{isochr.})}{\gamma}$ (≈ 1.2 in ring cyclotrons)

This shows the advantage of Ring Cyclotrons with a large radius R:

- space for many high voltage cavities \Rightarrow large \hat{V}
- the turn separation scales with R

there are two further effects, which can enhance dR/dn :

1. ν_r drops in the radial fringe field region

\Rightarrow this drop is fast in a warm magnet with a narrow gap (tolerable phase slip)

2. If ν_r is around 1.4-1.6 at extraction, one can double dR/dn

between the last 2 turns with excentric injection

and preserving the induced amplitude with single turn extraction

Example :

RIKEN Superconducting Ring Cyclotron SRC :

Uranium : $A = 238$, $Z = 86$, $E/A = 345 \text{ MeV/N}$

$$(25) \quad \gamma = 1 + \frac{(E/A)}{E_0} = 1.37, \quad \tilde{p} = 0.94, \quad \beta = 0.68$$

$$(26) \quad \varepsilon = \frac{E_0}{2} \tilde{p}^2 = (E/A) \left[1 + \frac{(E/A)}{2E_0} \right] = 1.185 (E/A) = 409 \text{ MeV/N}$$

$$(27) \quad K_B = \varepsilon \left(\frac{A}{Z} \right)^2 = 409 \left(\frac{238}{86} \right)^2 = 3130 \text{ MeV}$$

$$(28) \quad (B\rho) = \sqrt{\frac{K_B}{48 \text{ MeV}}} = \sqrt{\frac{3130}{48}} = 8.06 \text{ Tm}$$

Magnet weight W of a cyclotron: => Plot by W.Joho

("Modern Trends in Cyclotrons", CERN Accelerator School, Aarhus Denmark 1986):

$$(36) \quad W = W_1 (K_B / 1000 \text{ MeV})^{3/2} = W_2 (B\rho)^3$$

warm magnets: $W_1 \approx 5'000 \text{ t}$, $W_2 \approx 50 \text{ t}$

But PSI Ring Cyclotron with 590 MeV protons, $K_B=775 \text{ MeV}$, $B\rho=4\text{Tm}$,

$B_{\text{max}}=2.1\text{T}$ weighs only 2'000 t; => $W_2 = 31 \text{ t}$.

sc-magnets: only a few operating cyclotrons in 1986. These magnets were about a factor 12 lighter than warm magnets => $W_1=380 \text{ t}$, $W_2=4 \text{ t}$

Today more emphasis on low stray fields => more iron in the yoke.

Examples:

1) COMET Cyclotron PSI, 250 MeV protons, $B\rho=2.4\text{Tm}$, $K_B=283 \text{ MeV}$, $B_{\text{max}}=5\text{T}$

$W=90 \text{ t}$, => $W_2 = 6 \text{ t}$

2) SRC RIKEN, 345 MeV/N , $B\rho=8.06 \text{ Tm}$, $B_{\text{max}}=4\text{T}$, $W=8'300 \text{ t}$, => $W_2 = 16 \text{ t}$

Thus here the weight advantage over the warm PSI Ringcyclotron is only a factor of 2.