

NNLO corrections to jet production in deep inelastic scattering

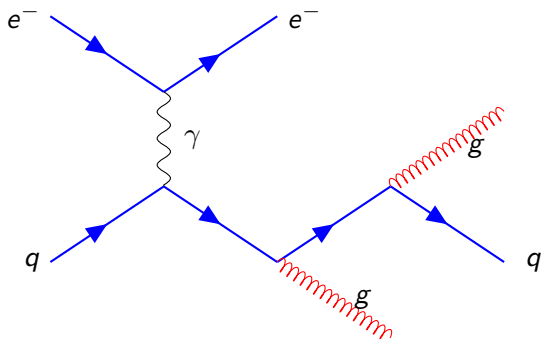
Jan Niehues

August 27, 2015

Overview

- Motivation
 - What is DIS?
 - Why QCD to higher orders?
- Theoretical Aspects
 - Infrared divergences
 - Antenna subtraction
- Outlook

Deep Inelastic Scattering



HERA



Why QCD to higher orders?

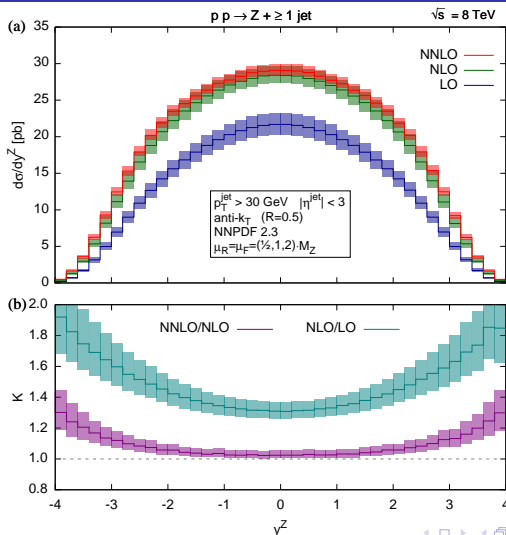
Last HERA analysis (2010):

$$\sigma_{meas}^{dijet}(ep \rightarrow 2\text{jets}) = 73 \pm 2(\text{stat.}) \pm 7(\text{syst.})\text{pb},$$

$$\sigma_{theo}^{dijet}(ep \rightarrow 2\text{jets}) = 77_{-20}^{+25}(\text{scale})_{-14}^{+4}(\text{PDF}) \pm 3(\text{had})\text{pb}.$$

- Differential cross sections are important for determination of PDFs.
- Need to reduce renormalisation and factorisation uncertainties on theoretical predictions.

Visualisation



Jet Cross Sections

Cross sections at different orders:

$$d\sigma_{\text{LO}} = \int_{d\Phi_m} d\sigma_{\text{B}}$$

$$d\sigma_{\text{NLO}} = \int_{d\Phi_{m+1}} d\sigma_{\text{NLO}}^{\text{R}} + \int_{d\Phi_m} d\sigma_{\text{NLO}}^{\text{V}}$$

$$d\sigma_{\text{NNLO}} = \int_{d\Phi_{m+2}} d\sigma_{\text{NNLO}}^{\text{RR}} + \int_{d\Phi_{m+1}} d\sigma_{\text{NNLO}}^{\text{RV}} + \int_{d\Phi_m} d\sigma_{\text{NNLO}}^{\text{VV}}$$

IR divergences and KLN theorem

- IR divergences appear in dimensional regularisation as poles in $\frac{1}{\epsilon^n}$.
- They come from:
 - Virtual loops in Feynman diagrams,
 - Phase space integration: Regions where particle momenta are collinear or go to zero.
- Kinoshita-Lee-Nauenberg (KLN) theorem:
 - Sufficient inclusive quantities are IR finite, i.e. IR divergences cancel overall when all corrections are added together.

Problem: Direct numerical integration of phase space integrals is not possible due to divergences.

Jet Cross Sections

Cross sections at different orders:

$$d\sigma_{\text{LO}} = \int_{d\Phi_m} d\sigma_{\text{B}}$$

$$d\sigma_{\text{NLO}} = \int_{d\Phi_{m+1}} d\sigma_{\text{NLO}}^{\text{R}} + \int_{d\Phi_m} d\sigma_{\text{NLO}}^{\text{V}}$$

$$d\sigma_{\text{NNLO}} = \int_{d\Phi_{m+2}} d\sigma_{\text{NNLO}}^{\text{RR}} + \int_{d\Phi_{m+1}} d\sigma_{\text{NNLO}}^{\text{RV}} + \int_{d\Phi_m} d\sigma_{\text{NNLO}}^{\text{VV}}$$

The idea: Adding a clever zero

$$\begin{aligned}
 \overbrace{d\sigma_{\text{NNLO}}}^{\text{IR finite}} &= \int_{d\Phi_{m+2}} \overbrace{\left[d\sigma_{\text{NNLO}}^{\text{RR}} + d\sigma_{\text{NNLO}}^{\text{RR,S}} \right]}^{\text{IR finite}} \\
 &\quad + \int_{d\Phi_{m+1}} \underbrace{\left[d\sigma_{\text{NNLO}}^{\text{RV}} + d\sigma_{\text{NNLO}}^{\text{RV,T}} \right]}_{\text{IR finite}} \\
 &\quad + \int_{d\Phi_m} d\sigma_{\text{NNLO}}^{\text{VV}} + \int_{d\Phi_m} d\sigma_{\text{NNLO}}^{\text{VV,U}} + d\sigma_{\text{MF}}.
 \end{aligned}$$

Subtraction terms have the same IR divergences as MEs

→ integrand is IR finite

→ numerical integration is possible.

Antenna formalism

How does this work in detail?

Factorisation of IR divergences

- MEs in colour-order form follow universal IR divergence factorisation:

- 1 Soft radiation, $p_j \rightarrow 0$:

$$|\mathcal{M}_{m+1}^0(\cdots, p_i, p_j, p_k, \cdots)|^2 \rightarrow S_{ijk} |\mathcal{M}_m^0(\cdots, p_i, p_k, \cdots)|^2$$

- 2 Collinear splitting $p_i || p_j$:

$$|\mathcal{M}_{m+1}^0(\cdots, p_i, p_j, p_k, \cdots)|^2 \rightarrow \frac{P_{ij \rightarrow l}}{s_{ij}} |\mathcal{M}_m^0(\cdots, p_{i+j}, p_k, \cdots)|^2$$

Same idea at NNLO where up to two particles can be unresolved
 \rightarrow more soft kernels and collinear splitting functions, e.g. $P_{ijk \rightarrow K}$.

Factorisation of phase space

Phase space factorises for appropriate linear mappings

$p_i, p_j, p_k \rightarrow p_I, p_K$:

$$d\Phi_{m+1}(p_1, \dots, p_{m+1}; q) \\ = \underbrace{d\Phi_m(p_1, \dots, p_I, p_K, \dots, p_{m+1}; q)}_{\text{reduced phase space}} \times \underbrace{d\Phi_3(p_i, p_j, p_k; p_I + p_K)}_{\text{antenna phase space}}$$

- Mapping conserves 4 momentum.
- p_I and p_K "on mass-shell".
- Analogous factorisation of 4 particle phase space.

Subtraction method

Using momentum map from $\{p_{m+1}\} \rightarrow \{\widetilde{p}_m\}$ with $\{p_X\} \subset \{p_{m+1}\}$, construct subtraction terms according to factorisation:

$$d\sigma_{NNLO}^{RR,S} \approx \underbrace{X(\{p_X\})}_{\text{antenna}} d\Phi_3(\{p_X\}) \times \overbrace{|\mathcal{M}(\{\widetilde{p}_m\})|^2}^{\text{reduced ME}} d\Phi_m(\{\widetilde{p}_m\}) \times \underbrace{\mathcal{J}(\{\widetilde{p}_m\})}_{\text{jet function}}$$

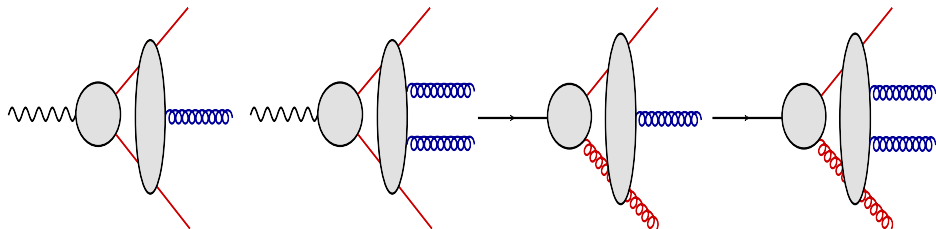
Analytic Integration of antenna phase space

→ IR divergences as explicit ϵ poles

→ analytic cancellation with poles of virtual corrections.

Antenna functions

- Antennae have two hard radiators
- Antennae reproduce all IR divergences of QCD.
- Have all been analytically integrated over antenna phase space!
- Have some none QCD divergences.



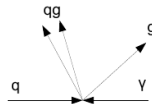
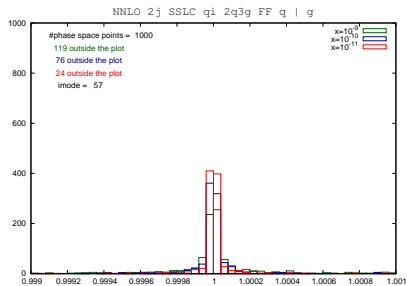
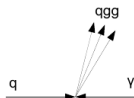
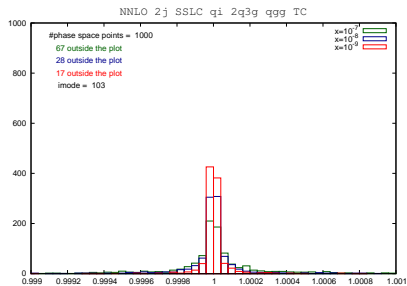
Recipe

- construct RR subtraction terms. ✓
- construct RV subtraction terms. ✓
- construct VV subtraction terms. ✓
- do numerical integration; in progress.

Correctness of RR,RV,VV are checked by pole cancellation and spike plots:

- spikeplots examine ratio of e.g. $\sigma_{\text{NNLO}}^{\text{RR,S}}/\sigma_{\text{NNLO}}^{\text{RR}}$.
- pole cancellation is analytic.

“Spike Plots”



Outlook

- Results for NNLO DIS are expected to appear very soon!
- Other processes that are calculated using antennae are:
 - V+jet,
 - H+jet,
 - jet+jet,