



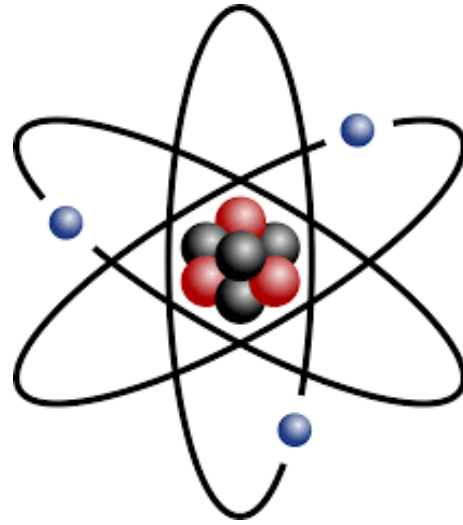
# Nuclear structure corrections in light muonic atoms

Sonia Bacca

Johannes Gutenberg University, Mainz

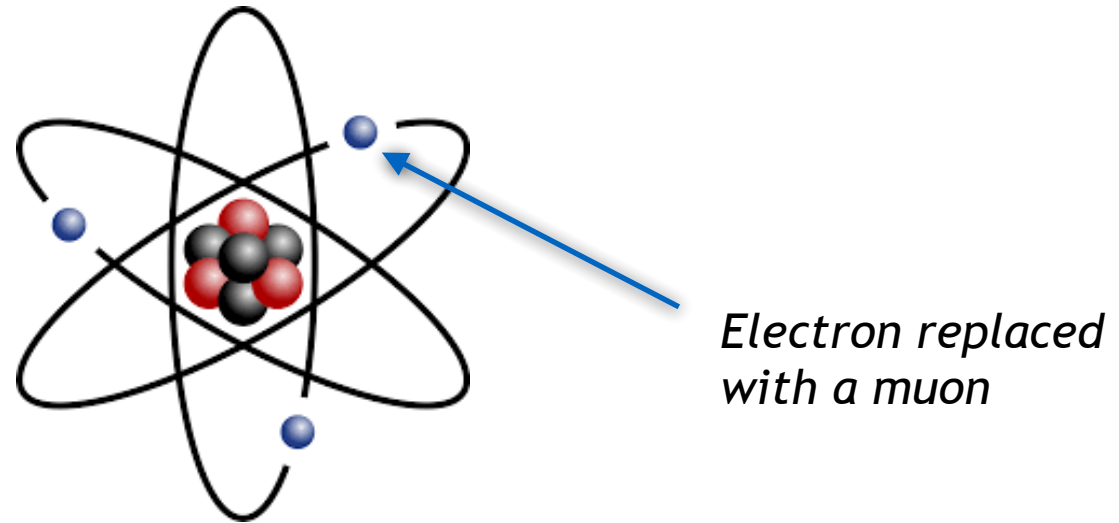
# What are muonic atoms?

*Exotic atoms*



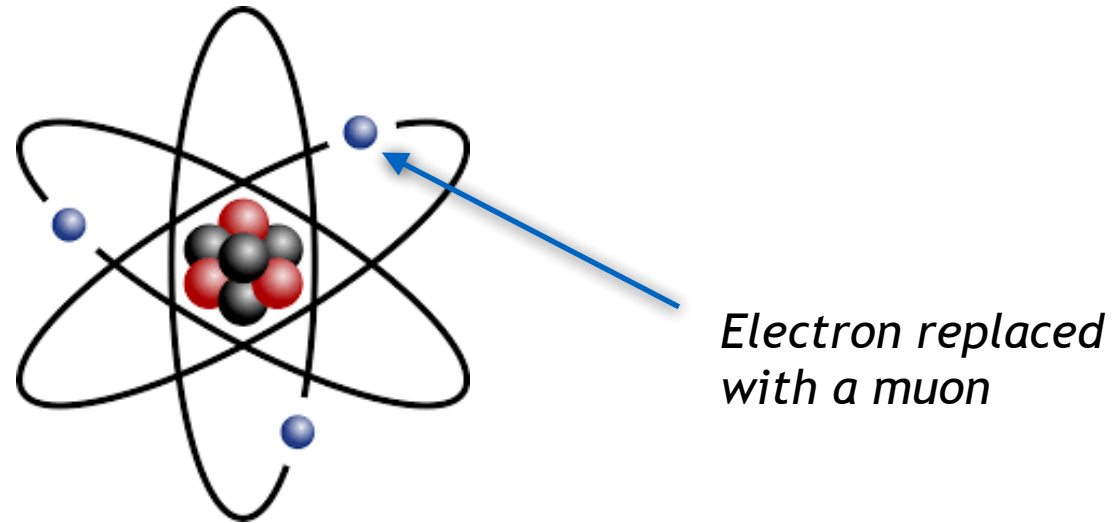
# What are muonic atoms?

*Exotic atoms*



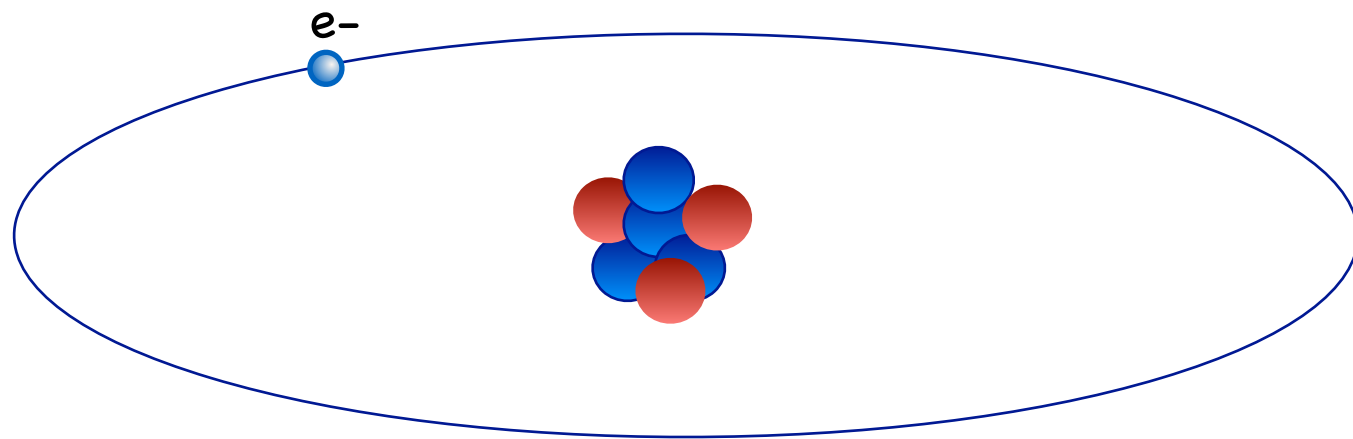
# What are muonic atoms?

*Exotic atoms*

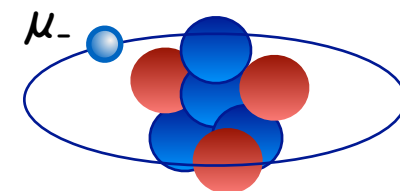


## Hydrogen-like systems

*Ordinary atoms*



*Muonic atoms*

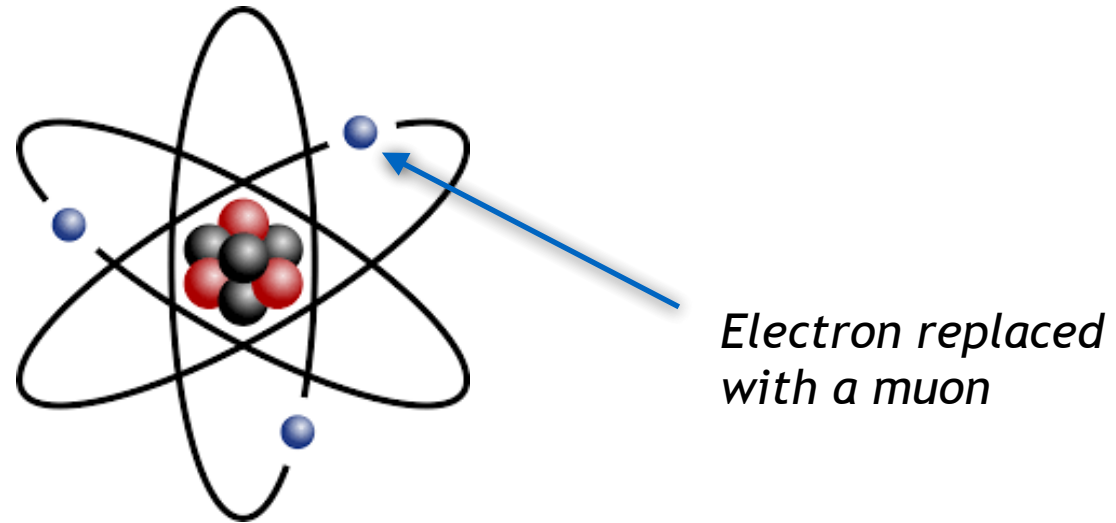


*muon more sensitive to the nucleus*



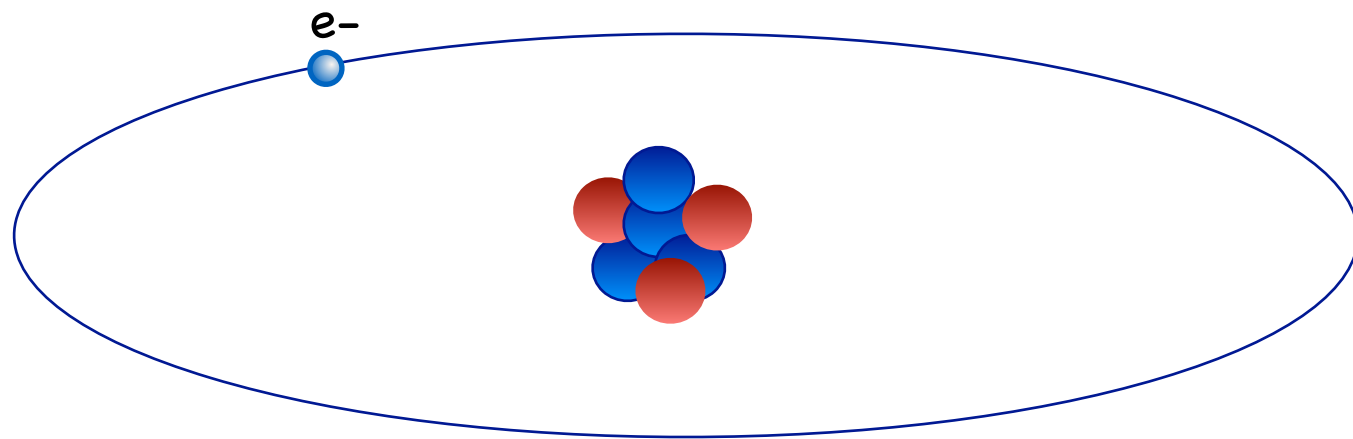
# What are muonic atoms?

*Exotic atoms*

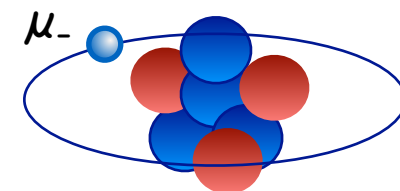


## Hydrogen-like systems

*Ordinary atoms*



*Muonic atoms*



*muon more sensitive to the nucleus*

**Can be used as a precision probe for the nucleus**

# Proton Radius Puzzle

The proton charge radius is measured from:

● electron-proton interactions:  $0.8770 \pm 0.0045$  fm

eH spectroscopy

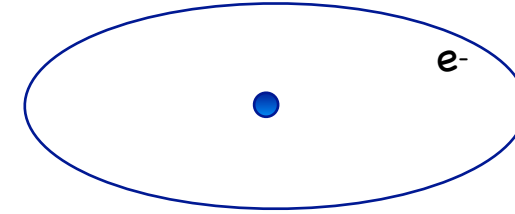
e-p scattering

● muonic -proton interactions:  $0.8409 \pm 0.0004$  fm

$\mu$ H Lamb-shift

Pohl *et al.*, Nature (2010)

Antognini *et al.*, Science (2013)



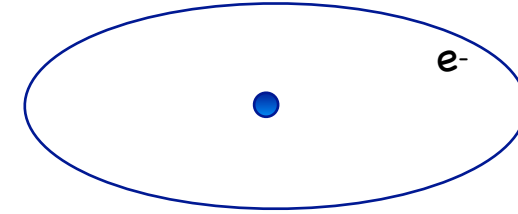
# Proton Radius Puzzle

The proton charge radius is measured from:

● electron-proton interactions:  $0.8770 \pm 0.0045$  fm

eH spectroscopy

e-p scattering



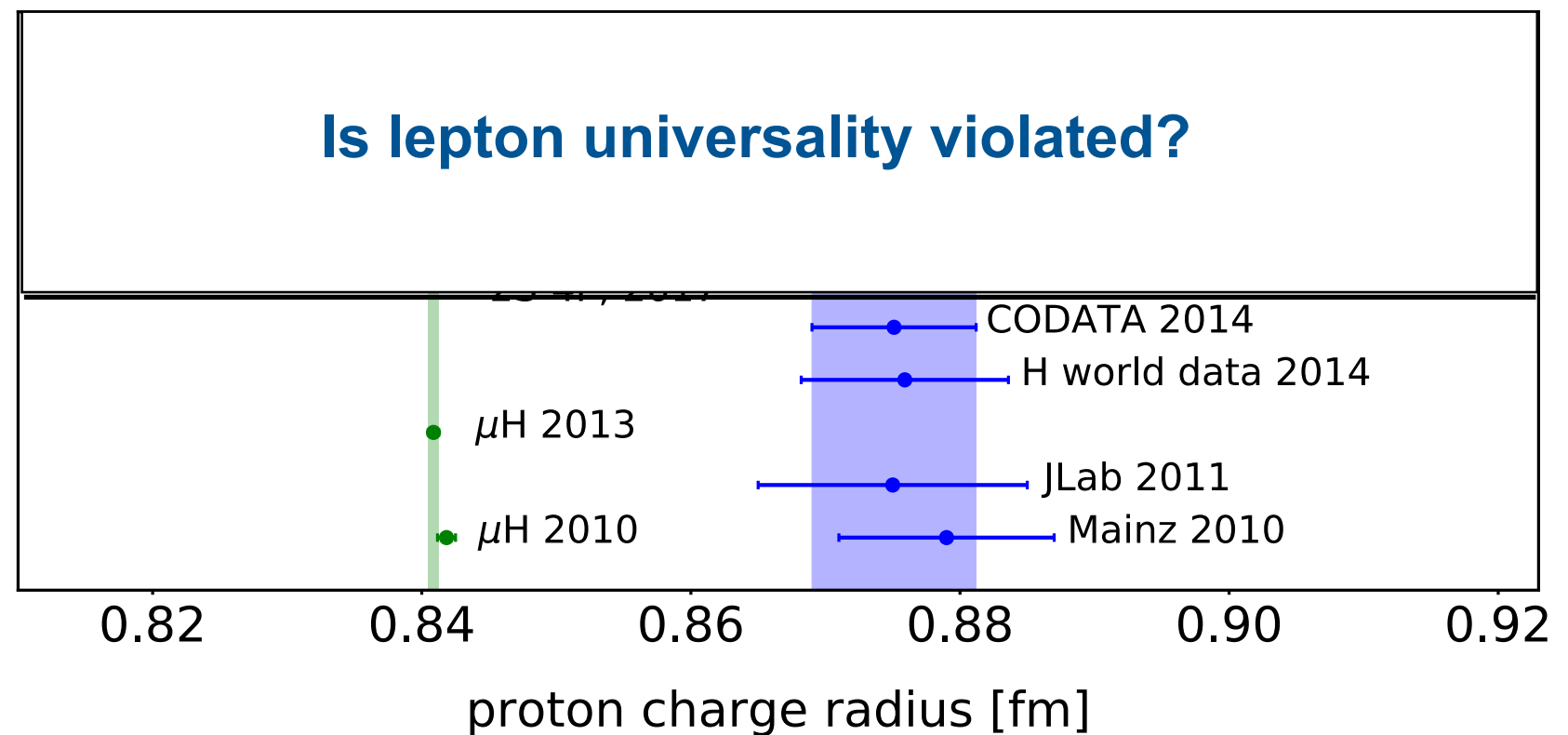
● muonic -proton interactions:  $0.8409 \pm 0.0004$  fm

$\mu$ H Lamb-shift



Pohl *et al.*, Nature (2010)

Antognini *et al.*, Science (2013)



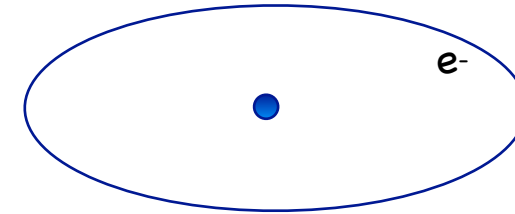
# Proton Radius Puzzle

The proton charge radius is measured from:

● electron-proton interactions:  $0.8770 \pm 0.0045$  fm

eH spectroscopy

e-p scattering



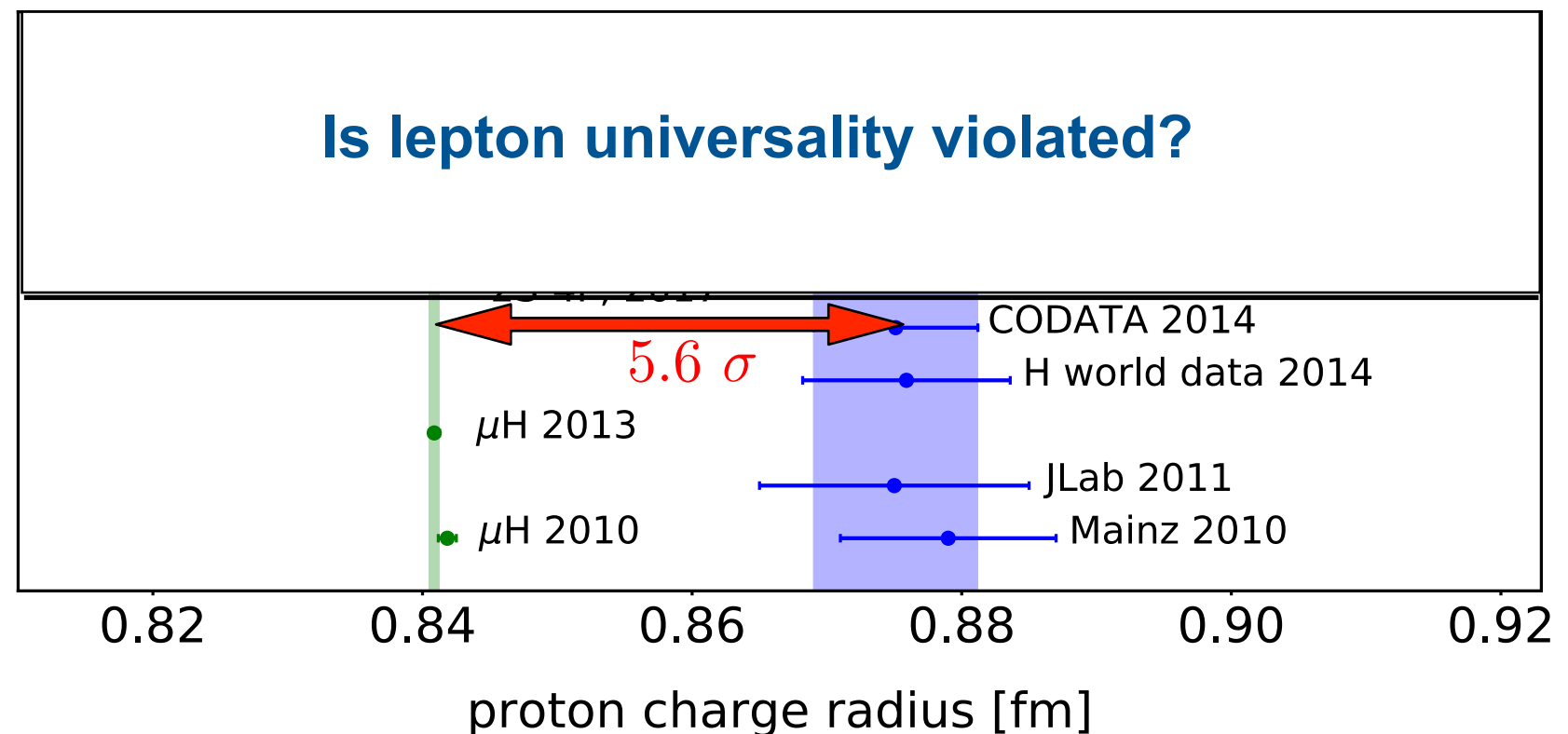
● muonic -proton interactions:  $0.8409 \pm 0.0004$  fm

$\mu$ H Lamb-shift



Pohl *et al.*, Nature (2010)

Antognini *et al.*, Science (2013)





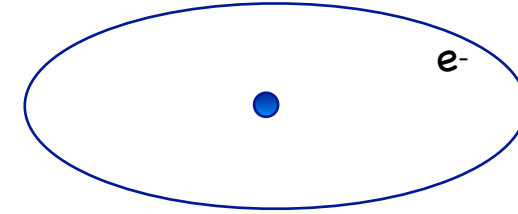
# Proton Radius Puzzle

The proton charge radius is measured from:

● electron-proton interactions:  $0.8770 \pm 0.0045$  fm

eH spectroscopy

e-p scattering



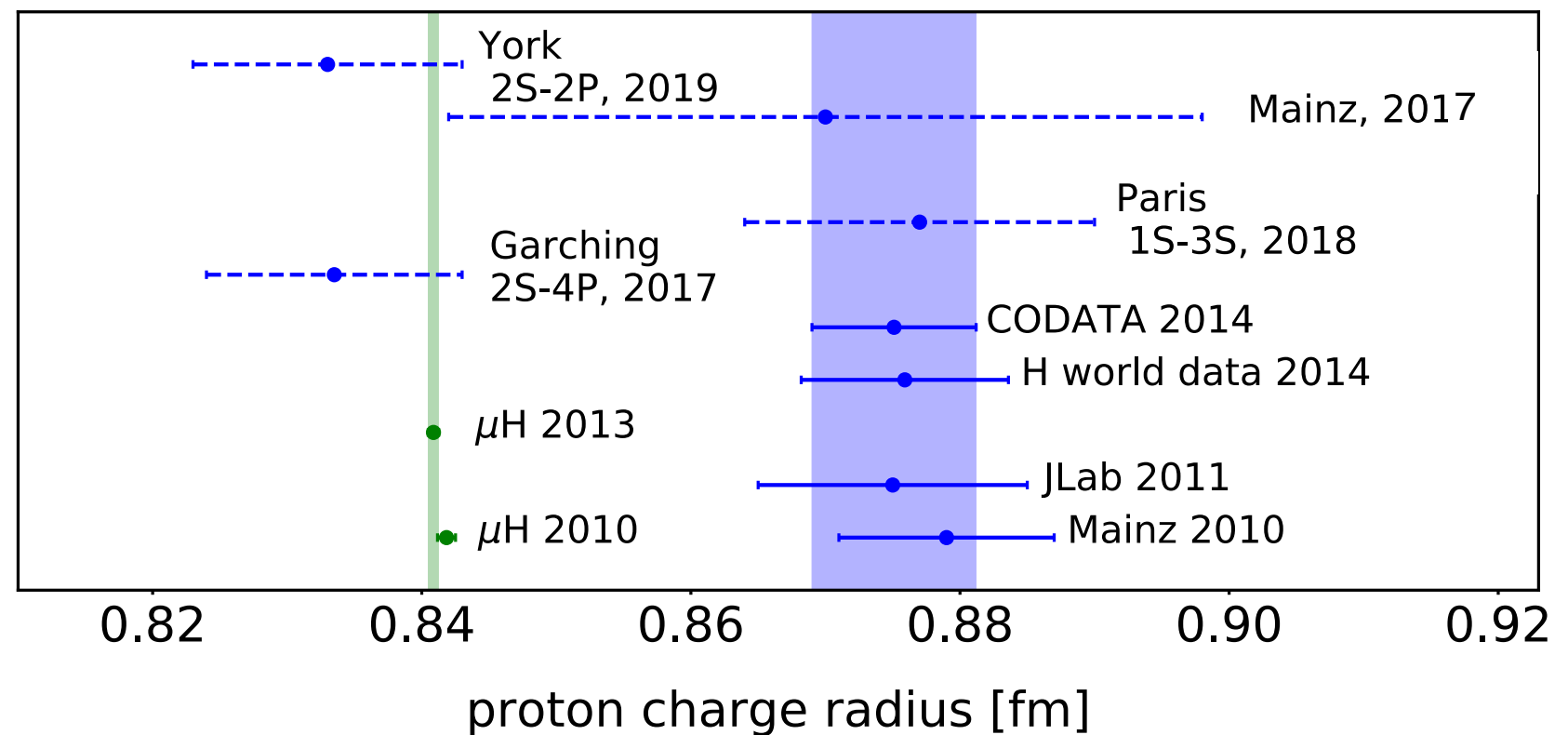
● muonic -proton interactions:  $0.8409 \pm 0.0004$  fm

$\mu$ H Lamb-shift



Pohl *et al.*, Nature (2010)

Antognini *et al.*, Science (2013)



# Experimental efforts

## 🌟 Higher precision electron scattering experiments

$Q^2$  from  $10^{-4} \text{ GeV}^2$  to  $10^{-2} \text{ GeV}^2$

Jefferson Lab

ISR measurement, not competitive, [Phys.Lett. B 771 \(2017\) 194-198](#)

Repeat traditional measurement with windowless gas jet target (2019/20)



# Experimental efforts

## 🌟 Higher precision electron scattering experiments

$Q^2$  from  $10^{-4} \text{ GeV}^2$  to  $10^{-2} \text{ GeV}^2$

Jefferson Lab

ISR measurement, not competitive, [Phys.Lett. B 771 \(2017\) 194-198](#)

Repeat traditional measurement with windowless gas jet target (2019/20)



## 🌟 MUSE collaboration

Measure both  $e^+/e^-$  and  $\mu^+/\mu^-$  to reduce uncertainties



# Experimental efforts

## Higher precision electron scattering experiments

$Q^2$  from  $10^{-4} \text{ GeV}^2$  to  $10^{-2} \text{ GeV}^2$

Jefferson Lab

ISR measurement, not competitive, *Phys.Lett. B* **771** (2017) 194-198

Repeat traditional measurement with windowless gas jet target (2019/20)



## MUSE collaboration

Measure both  $e^+/e^-$  and  $\mu^+/\mu^-$  to reduce uncertainties



## CREMA collaboration

CREMA collaboration currently measuring Lamb shift in light muonic atoms: Deuterium, Helions, etc.



# Experimental efforts

## Higher precision electron scattering experiments

$Q^2$  from  $10^{-4} \text{ GeV}^2$  to  $10^{-2} \text{ GeV}^2$

Jefferson Lab

ISR measurement, not competitive, *Phys.Lett. B* **771** (2017) 194-198

Repeat traditional measurement with windowless gas jet target (2019/20)



## MUSE collaboration

Measure both  $e^+/e^-$  and  $\mu^+/\mu^-$  to reduce uncertainties



## CREMA collaboration

CREMA collaboration currently measuring Lamb shift in light muonic atoms: Deuterium, Helions, etc.

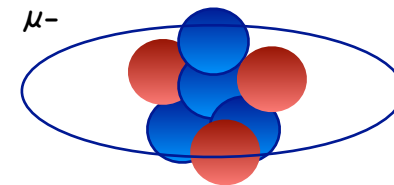




# Other light muonic atoms

Strong experimental program at PSI from the CREMA collaboration to unravel the mystery by studying other muonic atoms:

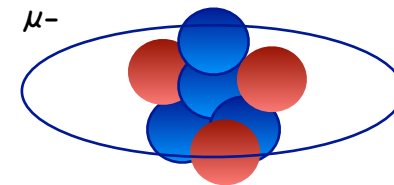
- $\mu\text{D}$  (results released)
- $\mu^4\text{He}^+$  (analyzing data)
- $\mu^3\text{He}^+$  (analyzing data)
- $\mu^3\text{H}$  (impossible)
- $\mu^6\text{Li}^{2+}$ ,  $\mu^7\text{Li}^{2+}$  (future)



# Other light muonic atoms

Strong experimental program at PSI from the CREMA collaboration to unravel the mystery by studying other muonic atoms:

- $\mu\text{D}$  (results released)
- $\mu^4\text{He}^+$  (analyzing data)
- $\mu^3\text{He}^+$  (analyzing data)
- $\mu^3\text{H}$  (impossible)
- $\mu^6\text{Li}^{2+}$ ,  $\mu^7\text{Li}^{2+}$  (future)

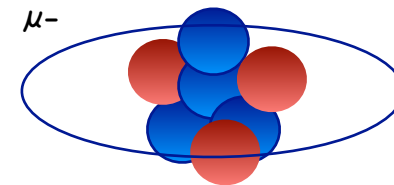


$$\Delta E_{2S-2P} = \delta_{\text{QED}} + \mathcal{A}_{\text{OPE}} \langle r_c^2 \rangle + \delta_{\text{TPE}}$$

# Other light muonic atoms

Strong experimental program at PSI from the CREMA collaboration to unravel the mystery by studying other muonic atoms:

- $\mu\text{D}$  (results released)
- $\mu^4\text{He}^+$  (analyzing data)
- $\mu^3\text{He}^+$  (analyzing data)
- $\mu^3\text{H}$  (impossible)
- $\mu^6\text{Li}^{2+}$ ,  $\mu^7\text{Li}^{2+}$  (future)



$$\Delta E_{2S-2P} = \delta_{\text{QED}} + \mathcal{A}_{\text{OPE}} \langle r_c^2 \rangle + \delta_{\text{TPE}}$$

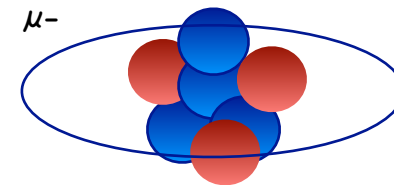


well known

# Other light muonic atoms

Strong experimental program at PSI from the CREMA collaboration to unravel the mystery by studying other muonic atoms:

- $\mu\text{D}$  (results released)
- $\mu^4\text{He}^+$  (analyzing data)
- $\mu^3\text{He}^+$  (analyzing data)
- $\mu^3\text{H}$  (impossible)
- $\mu^6\text{Li}^{2+}$ ,  $\mu^7\text{Li}^{2+}$  (future)



$$\Delta E_{2S-2P} = \delta_{\text{QED}} + \mathcal{A}_{\text{OPE}} \langle r_c^2 \rangle + \delta_{\text{TPE}}$$

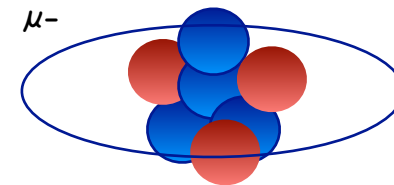


well known

# Other light muonic atoms

Strong experimental program at PSI from the CREMA collaboration to unravel the mystery by studying other muonic atoms:

- $\mu\text{D}$  (results released)
- $\mu^4\text{He}^+$  (analyzing data)
- $\mu^3\text{He}^+$  (analyzing data)
- $\mu^3\text{H}$  (impossible)
- $\mu^6\text{Li}^{2+}$ ,  $\mu^7\text{Li}^{2+}$  (future)



$$\Delta E_{2S-2P} = \delta_{\text{QED}} + \mathcal{A}_{\text{OPE}} \langle r_c^2 \rangle + \delta_{\text{TPE}}$$

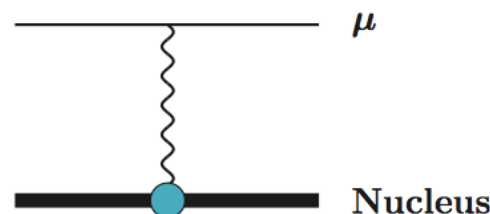


well known



$\propto$

$$\frac{m_r^4 (Z\alpha)^4}{12}$$

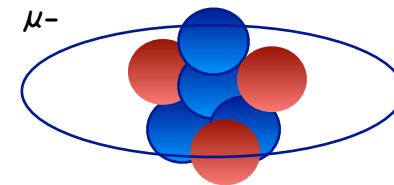




# Other light muonic atoms

Strong experimental program at PSI from the CREMA collaboration to unravel the mystery by studying other muonic atoms:

- $\mu\text{D}$  (results released)
- $\mu^4\text{He}^+$  (analyzing data)
- $\mu^3\text{He}^+$  (analyzing data)
- $\mu^3\text{H}$  (impossible)
- $\mu^6\text{Li}^{2+}$ ,  $\mu^7\text{Li}^{2+}$  (future)



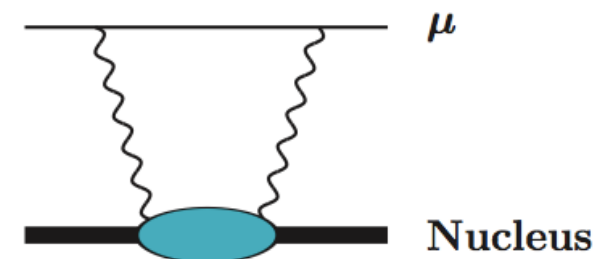
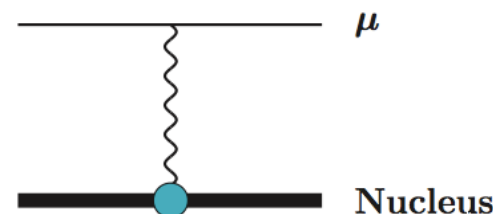
$$\Delta E_{2S-2P} = \delta_{\text{QED}} + \mathcal{A}_{\text{OPE}} \langle r_c^2 \rangle + \delta_{\text{TPE}}$$

well known

$\propto$

$$\frac{m_r^4 (Z\alpha)^4}{12}$$

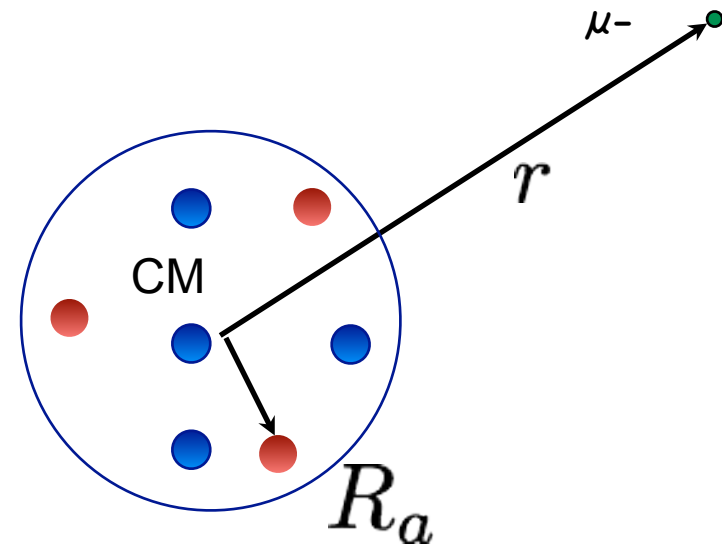
unknown  
NS corrections



# Theoretical derivation of TPE

$$H = H_N + H_\mu + \Delta V$$

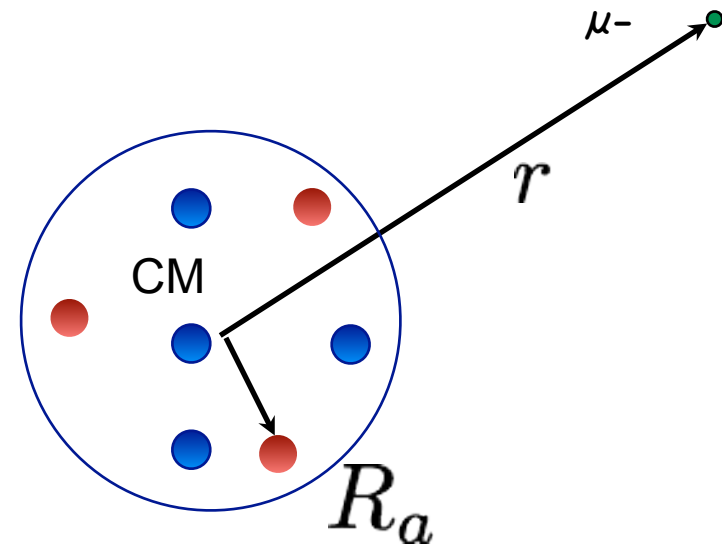
$$H_\mu = \frac{p^2}{2m_r} - \frac{Z\alpha}{r}$$



# Theoretical derivation of TPE

$$H = H_N + H_\mu + \Delta V$$

$$H_\mu = \frac{p^2}{2m_r} - \frac{Z\alpha}{r}$$



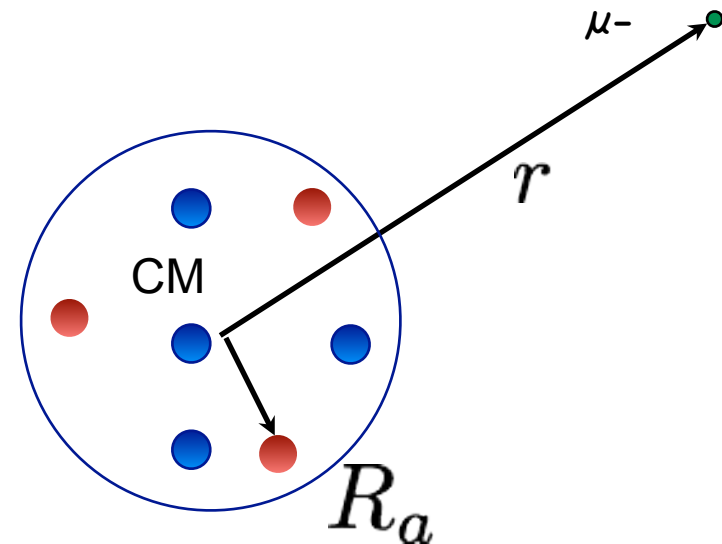
Perturbative potential: correction to the bulk Coulomb

$$\Delta V = \sum_a^Z \alpha \left( \frac{1}{r} - \frac{1}{|\vec{r} - \vec{R}_a|} \right)$$

# Theoretical derivation of TPE

$$H = H_N + H_\mu + \Delta V$$

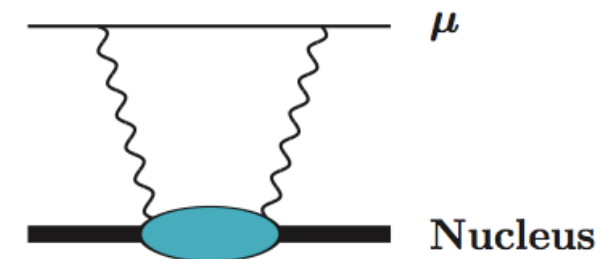
$$H_\mu = \frac{p^2}{2m_r} - \frac{Z\alpha}{r}$$



Perturbative potential: correction to the bulk Coulomb

$$\Delta V = \sum_a^Z \alpha \left( \frac{1}{r} - \frac{1}{|\vec{r} - \vec{R}_a|} \right)$$

Using perturbation theory at second order one obtains the expression for TPE up to order  $(Z\alpha)^5$



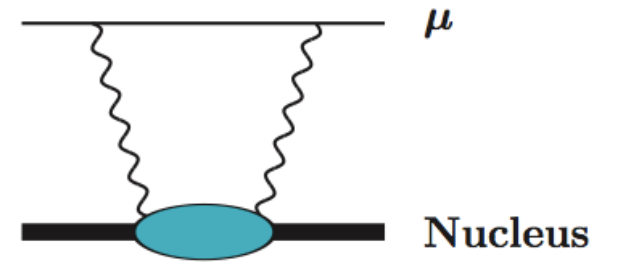
# Theoretical derivation of TPE

## ● Non relativistic term

Take non-relativistic kinetic energy in muon propagator

Neglect Coulomb force in the intermediate state

Expand the muon matrix elements in powers of  $\eta = \sqrt{\frac{m_r}{m_N}} \sim 0.17$





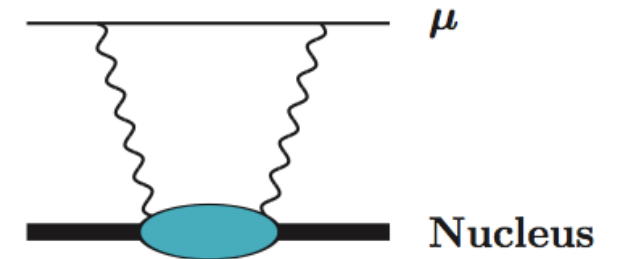
# Theoretical derivation of TPE

## Non relativistic term

Take non-relativistic kinetic energy in muon propagator

Neglect Coulomb force in the intermediate state

Expand the muon matrix elements in powers of  $\eta = \sqrt{\frac{m_r}{m_N}} \sim 0.17$



★ **Leading-order term**, related to the energy-weighted integral of the **dipole response function**

$$\delta_{D1}^{(0)} = -\frac{2\pi m_r^3}{9} (Z\alpha)^5 \int_{\omega_{th}}^{\infty} d\omega \sqrt{\frac{2m_r}{\omega}} S_{D1}(\omega)$$

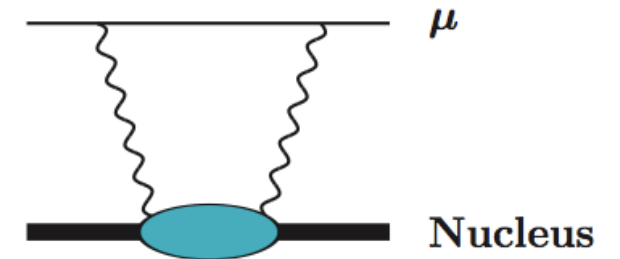
# Theoretical derivation of TPE

## Non relativistic term

Take non-relativistic kinetic energy in muon propagator

Neglect Coulomb force in the intermediate state

Expand the muon matrix elements in powers of  $\eta = \sqrt{\frac{m_r}{m_N}} \sim 0.17$



★ **Leading-order term**, related to the energy-weighted integral of the **dipole response function**

$$\delta_{D1}^{(0)} = -\frac{2\pi m_r^3}{9} (Z\alpha)^5 \int_{\omega_{th}}^{\infty} d\omega \sqrt{\frac{2m_r}{\omega}} S_{D1}(\omega)$$

★ **Next to leading-order term**, related to Zemach moment elastic contribution

$$\delta_{Z3}^{(1)} = \frac{\pi}{3} m_r (Z\alpha)^2 \phi^2(0) \iint d^3R d^3R' |\mathbf{R} - \mathbf{R}'|^3 \rho_0^p(\mathbf{R}) \rho_0^p(\mathbf{R}')$$

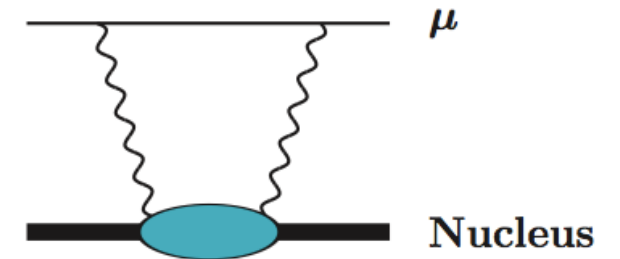
# Theoretical derivation of TPE

## Non relativistic term

Take non-relativistic kinetic energy in muon propagator

Neglect Coulomb force in the intermediate state

Expand the muon matrix elements in powers of  $\eta = \sqrt{\frac{m_r}{m_N}} \sim 0.17$



★ **Leading-order term**, related to the energy-weighted integral of the **dipole response function**

$$\delta_{D1}^{(0)} = -\frac{2\pi m_r^3}{9} (Z\alpha)^5 \int_{\omega_{th}}^{\infty} d\omega \sqrt{\frac{2m_r}{\omega}} S_{D1}(\omega)$$

★ **Next to leading-order term**, related to Zemach moment elastic contribution

$$\delta_{Z3}^{(1)} = \frac{\pi}{3} m_r (Z\alpha)^2 \phi^2(0) \iint d^3R d^3R' |\mathbf{R} - \mathbf{R}'|^3 \rho_0^p(\mathbf{R}) \rho_0^p(\mathbf{R}')$$

★ **Next-to-next to leading-order term**, related to **monopole and quadrupole response functions...**

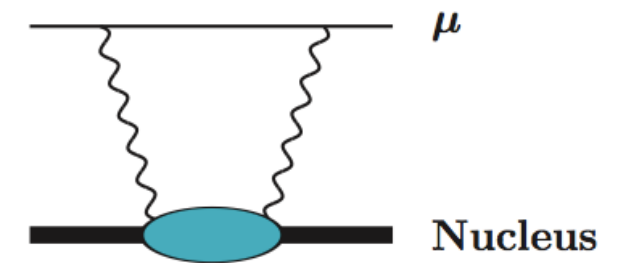
$$S_{R^2}(\omega), S_Q(\omega), S_{D1D3}(\omega)$$

# Theoretical derivation of TPE

## 🌟 Coulomb term

Consider the Coulomb force in the intermediate states  
Naively  $\delta_C^{(0)} \sim (Z\alpha)^6$ , actually logarithmically enhanced  
 $\delta_C^{(0)} \sim (Z\alpha)^5 \log(Z\alpha)$  [Friar \(1977\)](#), [Pachucki \(2011\)](#)

Related to the **dipole response function**

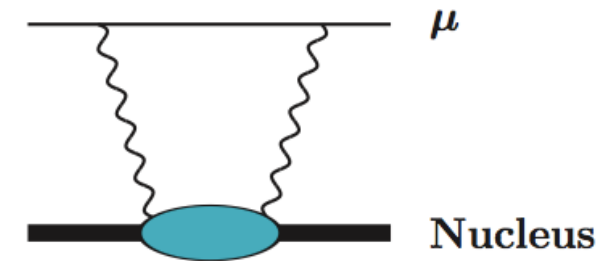


# Theoretical derivation of TPE

## ● Coulomb term

Consider the Coulomb force in the intermediate states  
Naively  $\delta_C^{(0)} \sim (Z\alpha)^6$ , actually logarithmically enhanced  
 $\delta_C^{(0)} \sim (Z\alpha)^5 \log(Z\alpha)$  Friar (1977), Pachucki (2011)

Related to the **dipole response function**



## ● Relativistic terms

Take the relativistic kinetic energy in muon propagator

Related to the **dipole response function**

$$\delta_{L(T)}^{(0)} = \frac{2m_r^3}{9} (Z\alpha)^5 \int_{\omega_{\text{th}}}^{\infty} d\omega K_{L(T)} \left( \frac{\omega}{m_r} \right) S_{D_1}(\omega)$$

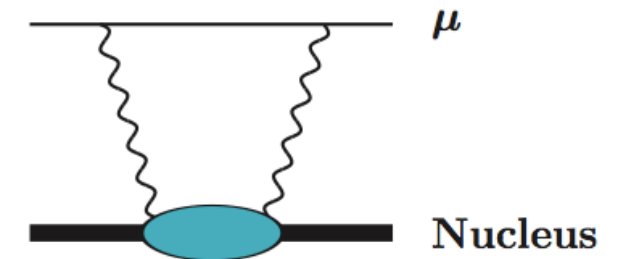


# Theoretical derivation of TPE

## Coulomb term

Consider the Coulomb force in the intermediate states  
Naively  $\delta_C^{(0)} \sim (Z\alpha)^6$ , actually logarithmically enhanced  
 $\delta_C^{(0)} \sim (Z\alpha)^5 \log(Z\alpha)$  Friar (1977), Pachucki (2011)

Related to the **dipole response function**



## Relativistic terms

Take the relativistic kinetic energy in muon propagator

Related to the **dipole response function**

$$\delta_{L(T)}^{(0)} = \frac{2m_r^3}{9} (Z\alpha)^5 \int_{\omega_{\text{th}}}^{\infty} d\omega K_{L(T)} \left( \frac{\omega}{m_r} \right) S_{D_1}(\omega)$$

## Finite nucleon-size corrections

Consider finite nucleon-size by including their charge distributions and obtain terms, e.g.,

$$\delta_{R1}^{(1)} = -8\pi m_r (Z\alpha)^2 \phi^2(0) \int \int d^3 R d^3 R' |\mathbf{R} - \mathbf{R}'| \left[ \frac{2}{\beta^2} \rho_0^{pp}(\mathbf{R}, \mathbf{R}') - \lambda \rho_0^{np}(\mathbf{R}, \mathbf{R}') \right]$$

# Theoretical derivation of TPE

$$\delta_{\text{TPE}} = \delta_{\text{Zem}}^A + \delta_{\text{Zem}}^N + \delta_{\text{pol}}^A + \delta_{\text{pol}}^N$$

# Theoretical derivation of TPE

$$\delta_{\text{TPE}} = \delta_{\text{Zem}}^A + \delta_{\text{Zem}}^N + \delta_{\text{pol}}^A + \delta_{\text{pol}}^N$$

$$\begin{aligned} \delta_{\text{pol}}^A = & \delta_{D1}^{(0)} + \delta_{R3}^{(1)} + \cancel{\delta_{Z3}^{(1)}} + \delta_{R^2}^{(2)} + \delta_Q^{(2)} + \delta_{D1D3}^{(2)} + \delta_C^{(0)} \\ & + \delta_L^{(0)} + \delta_T^{(0)} + \delta_M^{(0)} + \delta_{R1}^{(1)} + \cancel{\delta_{Z1}^{(1)}} + \delta_{NS}^{(2)} \end{aligned}$$

$$\delta_{\text{Zem}}^A = -\cancel{\delta_{Z3}^{(1)}} - \cancel{\delta_{Z1}^{(1)}} \quad \text{Friar and Payne ('97)}$$

Need to calculate  $\delta_{\text{TPE}}$  and related uncertainties.

# A matter of precision

The uncertainty of the extracted radius depends on the precision of the TPE

$$\Delta E_{2S-2P} = \delta_{\text{QED}} + \mathcal{A}_{\text{OPE}} \langle r_c^2 \rangle + \delta_{\text{TPE}}$$

Roughly:      95%                      4%                      1%

TPE needs to be known precisely, in order to exploit the experimental precision.

# A matter of precision

The uncertainty of the extracted radius depends on the precision of the TPE

$$\Delta E_{2S-2P} = \delta_{\text{QED}} + \mathcal{A}_{\text{OPE}} \langle r_c^2 \rangle + \delta_{\text{TPE}}$$

Roughly:      95%                      4%                      1%

TPE needs to be known precisely, in order to exploit the experimental precision.

## Uncertainties comparison

Atom	Exp uncertainty on $\Delta E_{2S-2P}$	Uncertainty on TPE prior to the discovery of the proton radius puzzle
$\mu^2\text{H}$	0.003 meV	0.03 meV*
$\mu^3\text{He}^+$	0.08 meV	1 meV
$\mu^4\text{He}^+$	0.06 meV	0.6 meV
$\mu^{6,7}\text{Li}^{++}$	0.7 meV	4 meV

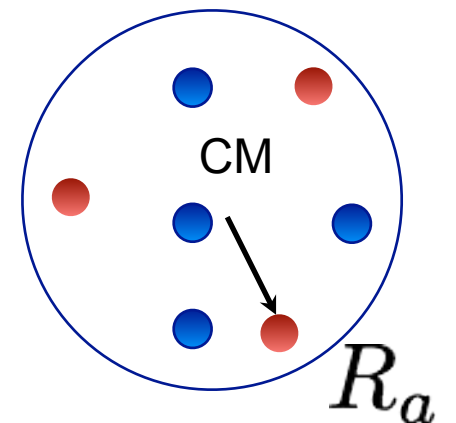
\*Leidemann, Rosenfelder '95 using few-body methods

# Ab Initio Nuclear Theory

- Solve the Schrödinger equation for few-nucleons

$$H_N|\psi_i\rangle = E_i|\psi_i\rangle$$

$$H_N = T + V_{NN}(\Lambda) + V_{3N}(\Lambda) + \dots$$

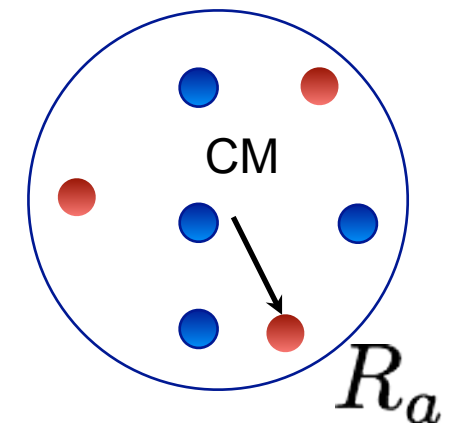


# Ab Initio Nuclear Theory

- Solve the Schrödinger equation for few-nucleons

$$H_N |\psi_i\rangle = E_i |\psi_i\rangle$$

$$H_N = T + V_{NN}(\Lambda) + V_{3N}(\Lambda) + \dots$$



Hyper-spherical harmonics expansions and Lorenz integral transform method for  $A=3,4,6,7$

Barnea, Leidemann, Orlandini PRC **61** (2000) 054001

Efros, *et al.*, JPG.: Nucl.Part.Phys. **34** (2007) R459

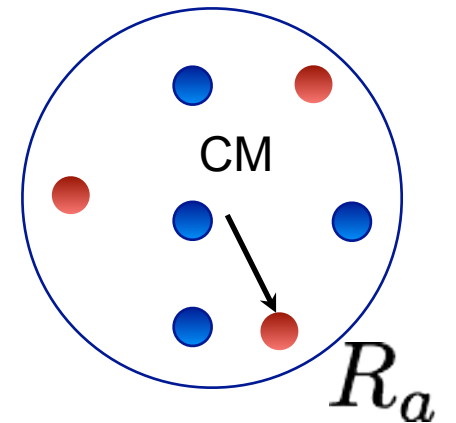
For  $A=2$  we use an harmonic oscillator expansion

# Ab Initio Nuclear Theory

- Solve the Schrödinger equation for few-nucleons

$$H_N |\psi_i\rangle = E_i |\psi_i\rangle$$

$$H_N = T + V_{NN}(\Lambda) + V_{3N}(\Lambda) + \dots$$



Hyper-spherical harmonics expansions and Lorenz integral transform method for  $A=3,4,6,7$

Barnea, Leidemann, Orlandini PRC **61** (2000) 054001

Efros, *et al.*, JPG.: Nucl.Part.Phys. **34** (2007) R459

For  $A=2$  we use an harmonic oscillator expansion

- We will use nuclear interactions derived from traditional potentials (AV18+UIX) and from chiral effective field theory (at various orders)


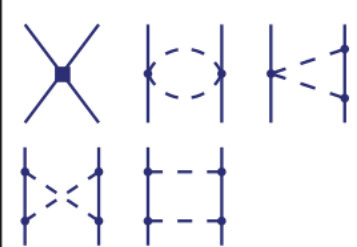
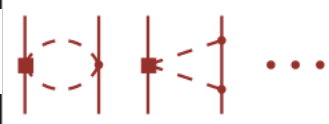

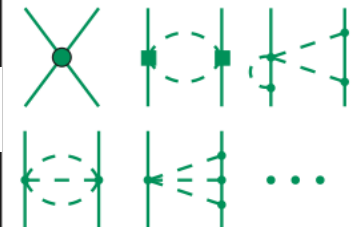

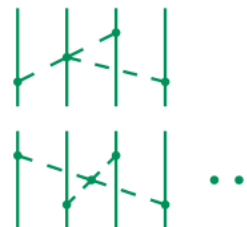


# Ab Initio Nuclear Theory

Chiral effective field theory

Systematic expansion

$$\mathcal{L} = \sum_{\nu} c_{\nu} \left( \frac{Q}{\Lambda_b} \right)^{\nu}$$

	2N force	3N force	4N force
LO $\nu = 0$			
NLO $\nu = 2$			
N2LO $\nu = 3$			
N3LO $\nu = 4$			

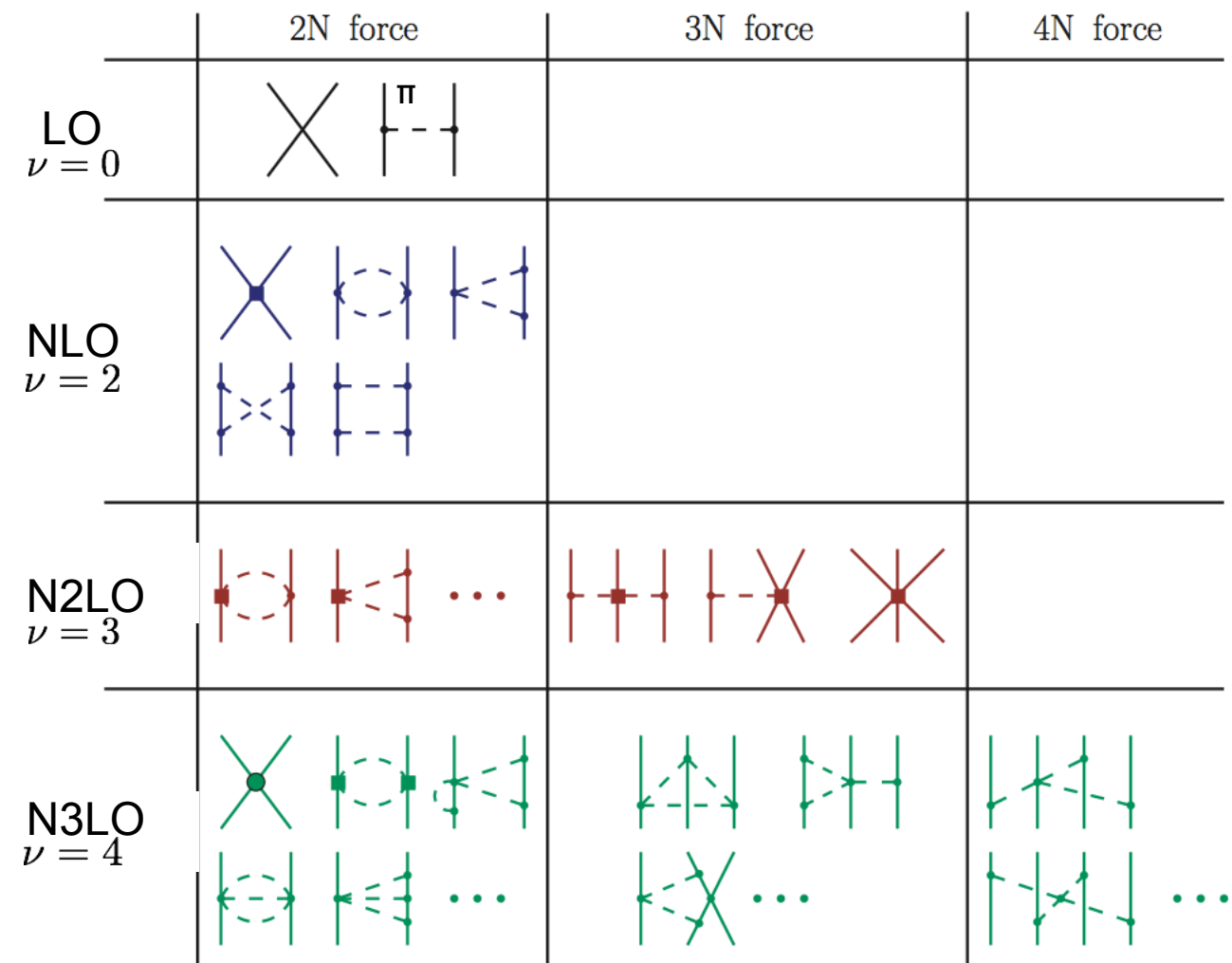
# Ab Initio Nuclear Theory

Chiral effective field theory

Systematic expansion

$$\mathcal{L} = \sum_{\nu} c_{\nu} \left( \frac{Q}{\Lambda_b} \right)^{\nu}$$

Details of short distance physics not resolved, but captured in **low energy constants (LEC)**



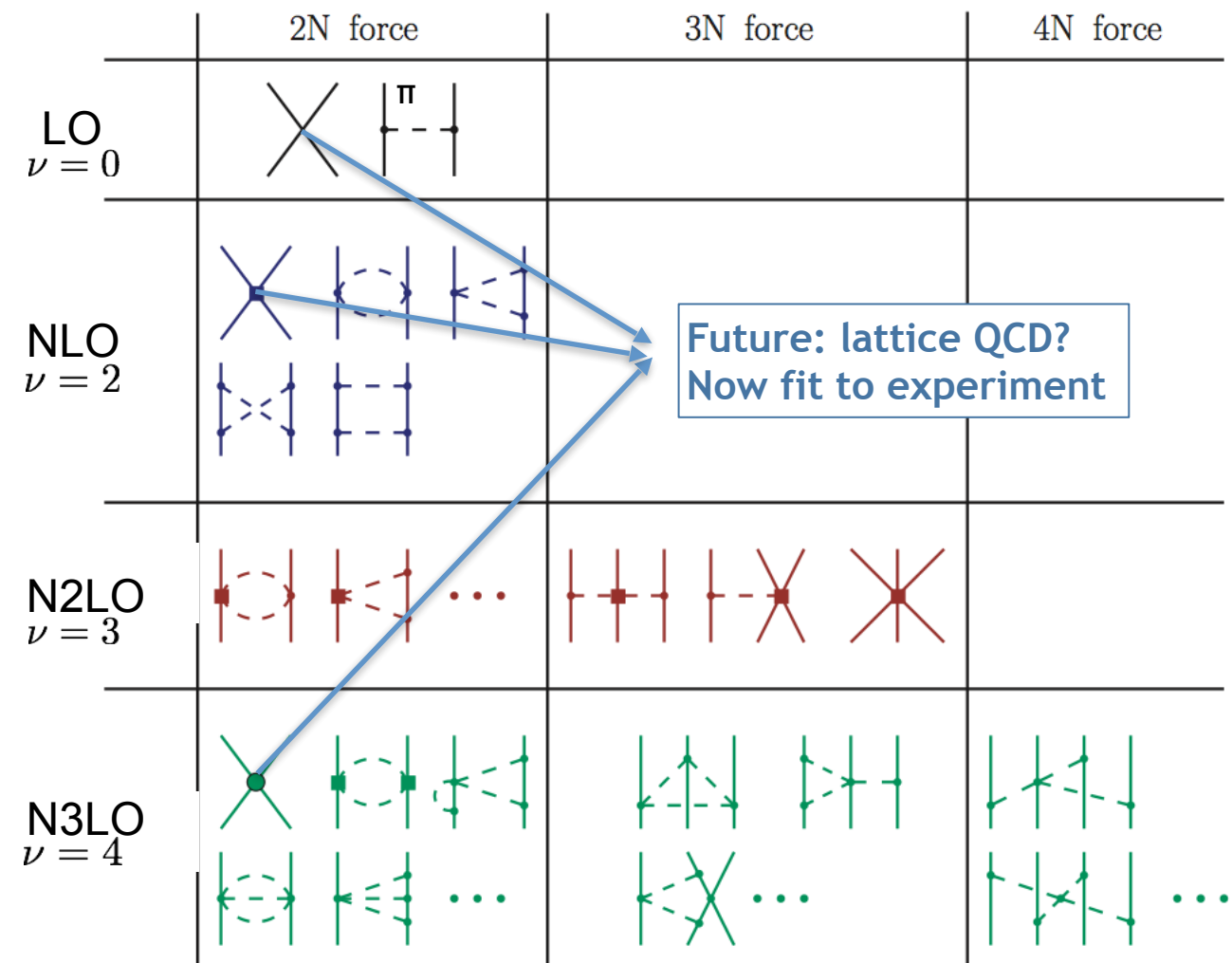
# Ab Initio Nuclear Theory

Chiral effective field theory

Systematic expansion

$$\mathcal{L} = \sum_{\nu} c_{\nu} \left( \frac{Q}{\Lambda_b} \right)^{\nu}$$

Details of short distance physics not resolved, but captured in **low energy constants (LEC)**



# Ab Initio Nuclear Theory

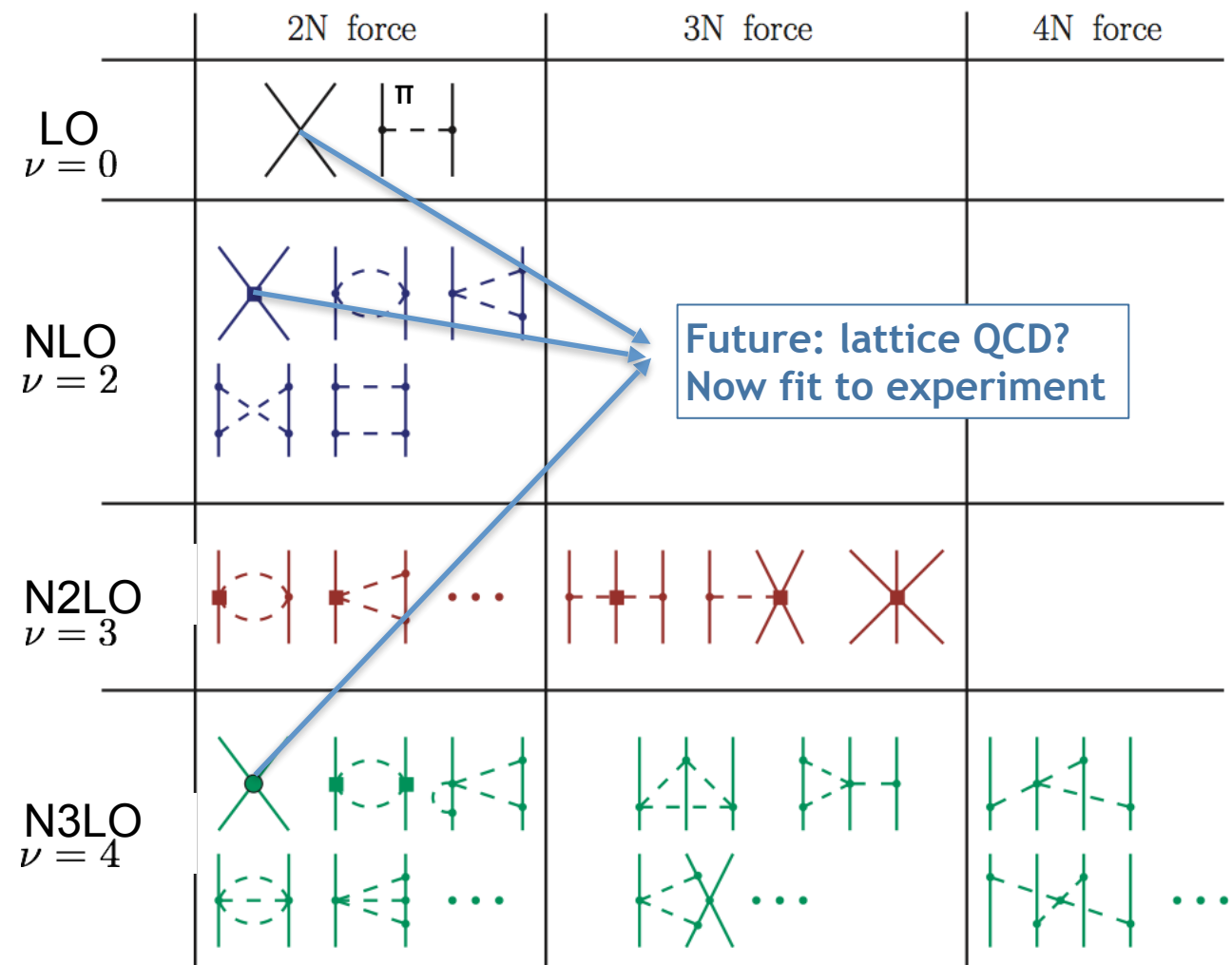
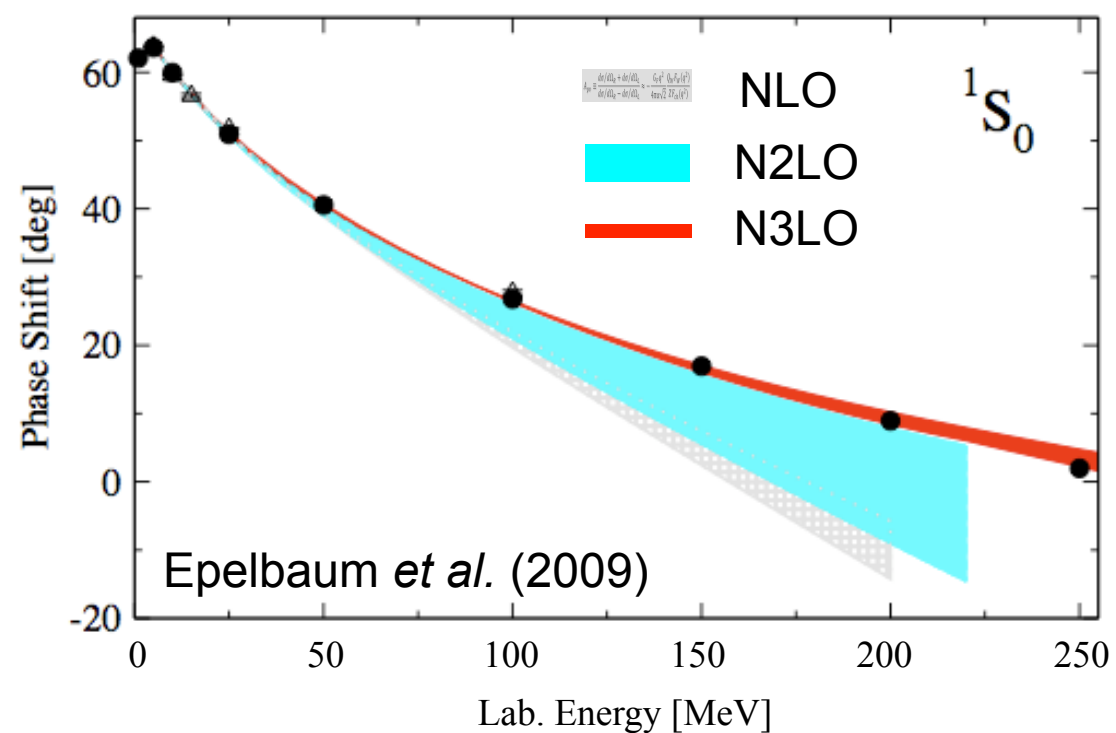
Chiral effective field theory

Systematic expansion

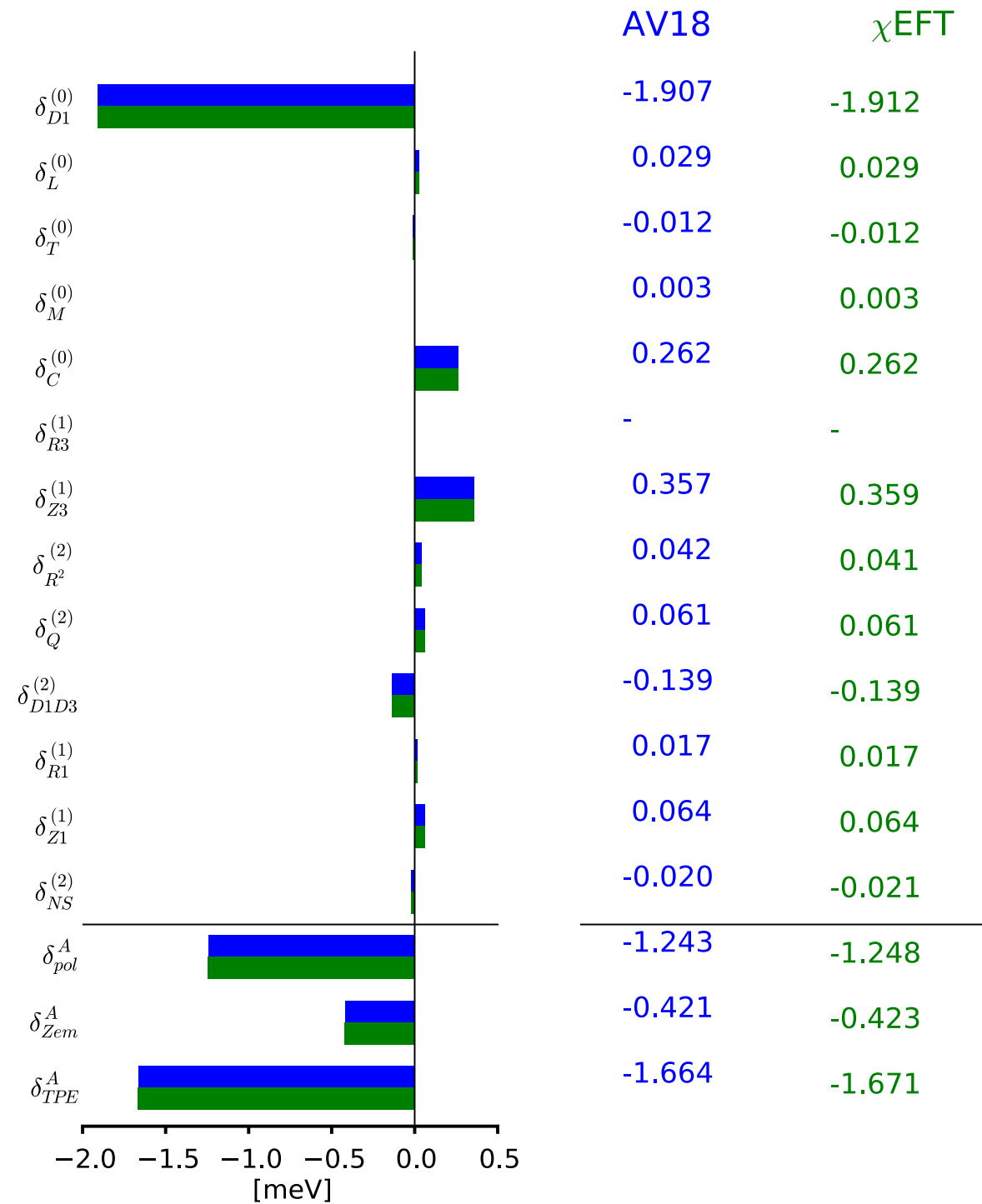
$$\mathcal{L} = \sum_{\nu} c_{\nu} \left( \frac{Q}{\Lambda_b} \right)^{\nu}$$

Details of short distance physics not resolved, but captured in **low energy constants (LEC)**

**LEC** fit to experiment - NN sector -



# Muonic Deuterium

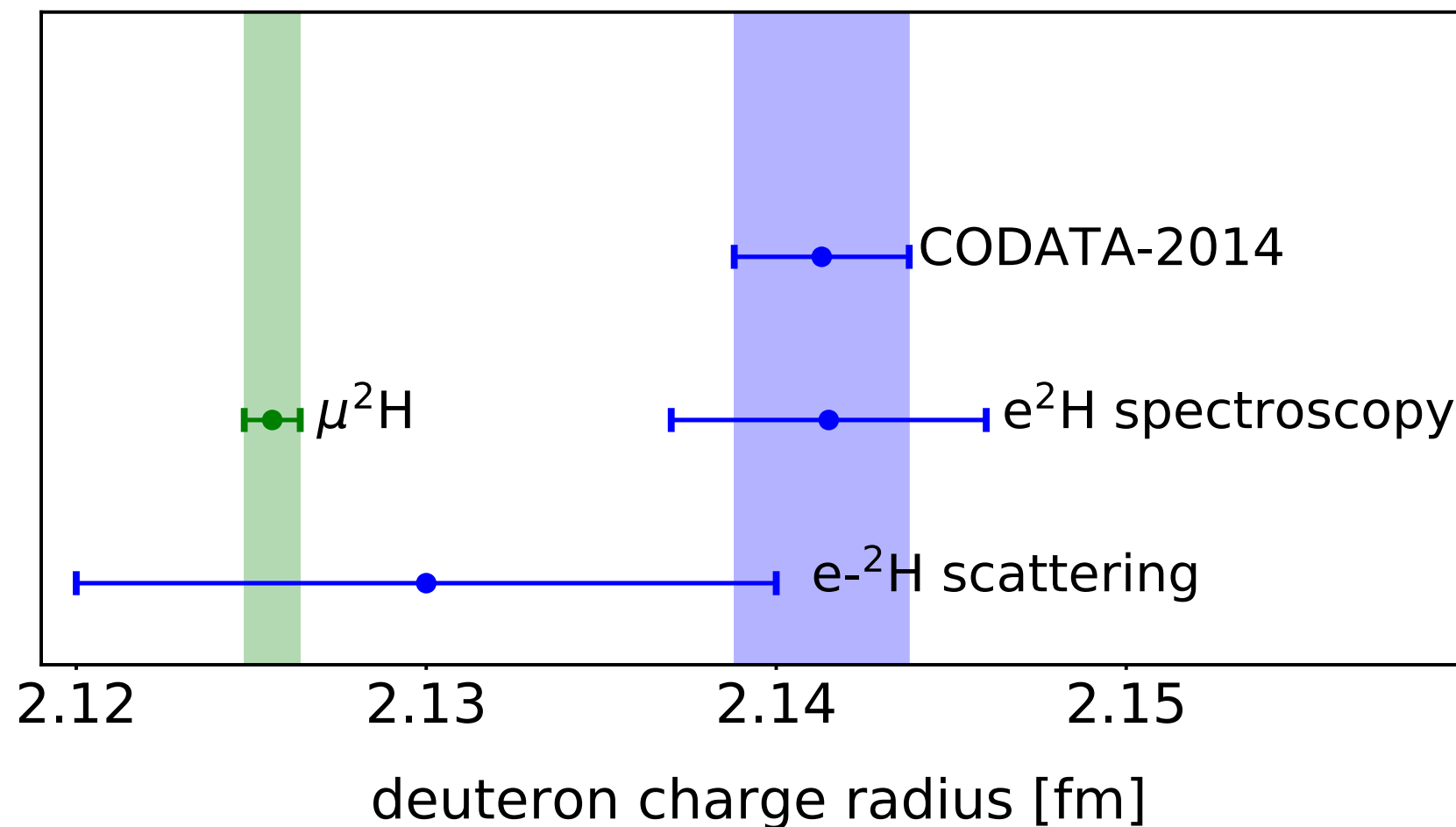


J. Hernandez et al, Phys. Lett. B **736**, 344 (2014)

AV18 in agreement with Pachucki (2011)+ Pachucki, Wienczek (2015)

# Deuteron radius puzzle

Pohl et al., Science **353**, 669 (2016)



$$\Delta E_{2S-2P} = \delta_{\text{QED}} + \mathcal{A}_{\text{OPE}} \langle r_c^2 \rangle + \delta_{\text{TPE}}$$

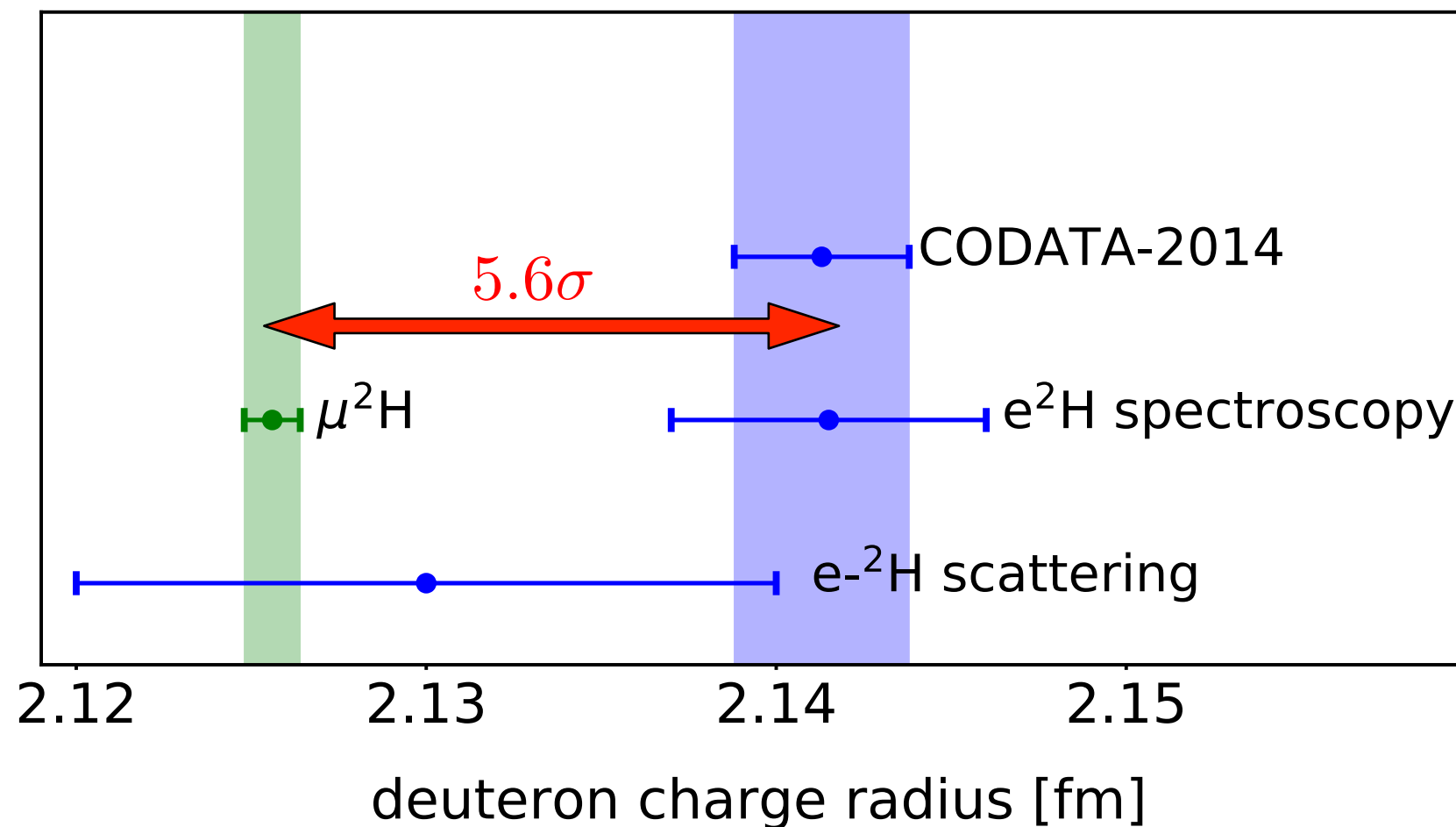


Hernandez et al., PLB **736**, 334 (2014)

Pachucki (2011)+ Pachucki, Wienczek (2015)

# Deuteron radius puzzle

Pohl et al., Science **353**, 669 (2016)



$$\Delta E_{2S-2P} = \delta_{\text{QED}} + \mathcal{A}_{\text{OPE}} \langle r_c^2 \rangle + \delta_{\text{TPE}}$$

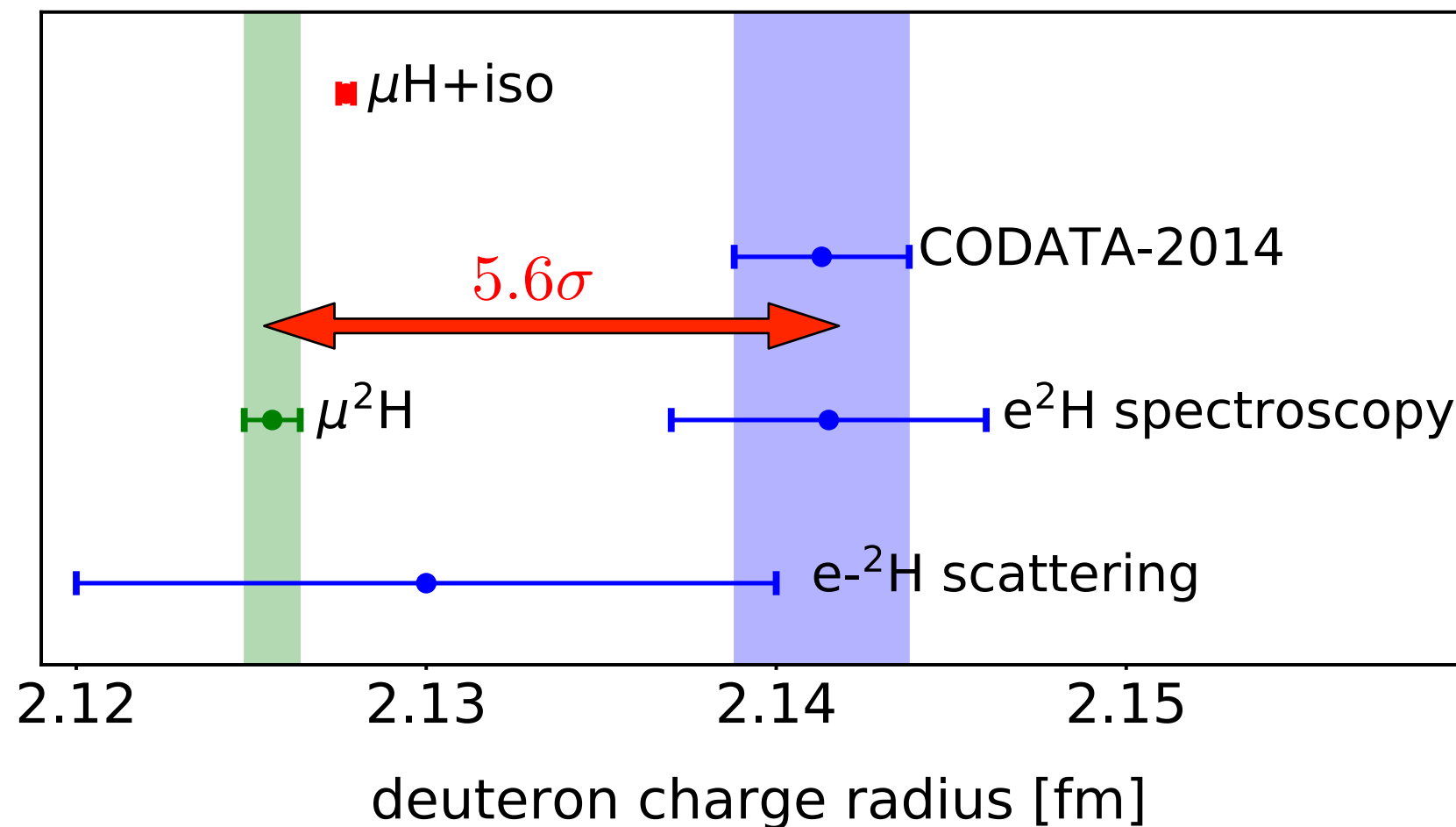


Hernandez et al., PLB **736**, 334 (2014)

Pachucki (2011)+ Pachucki, Wienczek (2015)

# Deuteron radius puzzle

Pohl et al., Science **353**, 669 (2016)



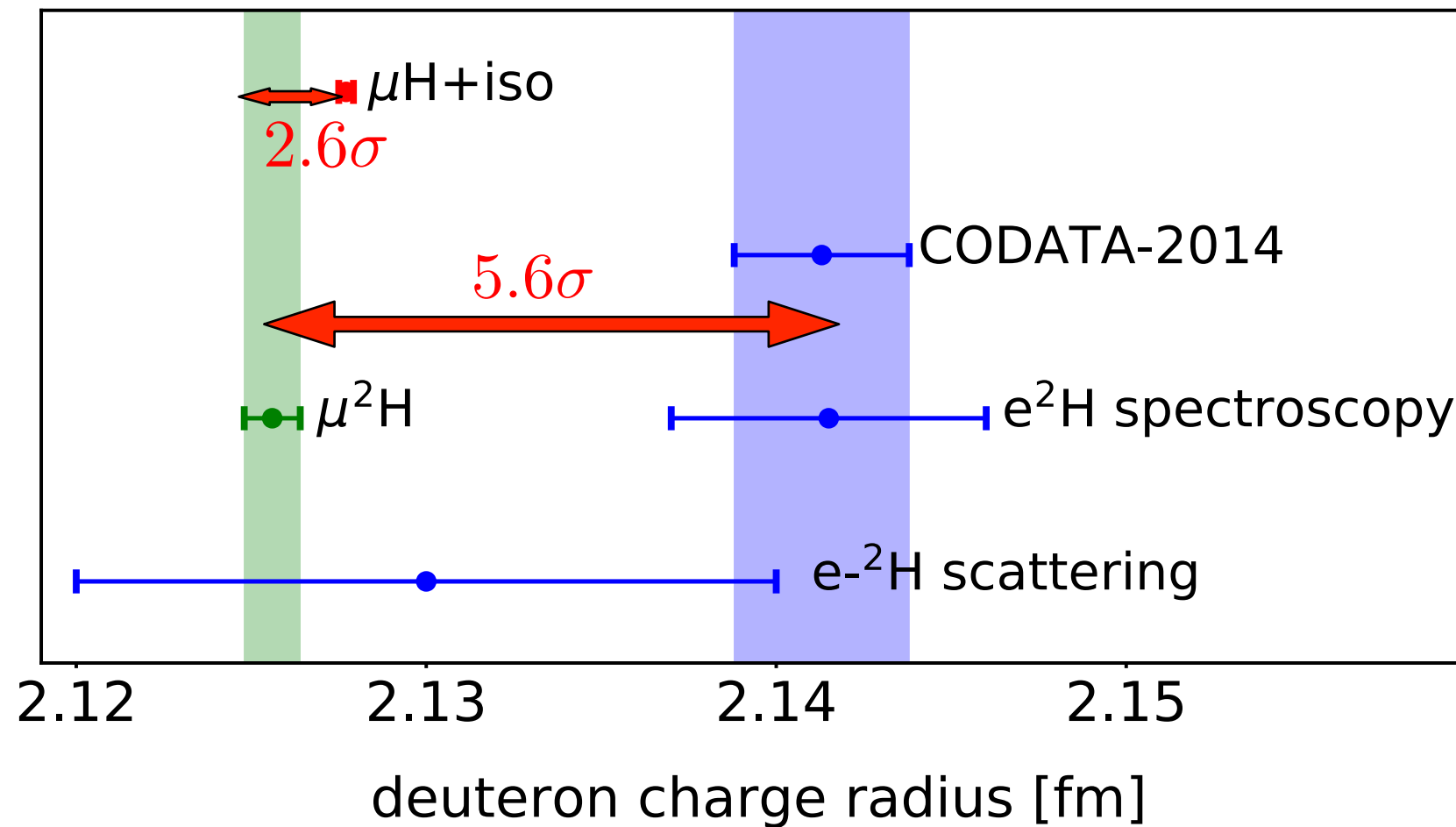
$$\Delta E_{2S-2P} = \delta_{\text{QED}} + \mathcal{A}_{\text{OPE}} \langle r_c^2 \rangle + \delta_{\text{TPE}}$$

$\mu\text{H}+\text{iso}$ :  $r_p$  from  $\mu\text{H}$  and deuterium isotopic shift  $r_d^2 - r_p^2$ : Parthey et al., PRL **104** 233001 (2010)



# Deuteron radius puzzle

Pohl et al., Science **353**, 669 (2016)



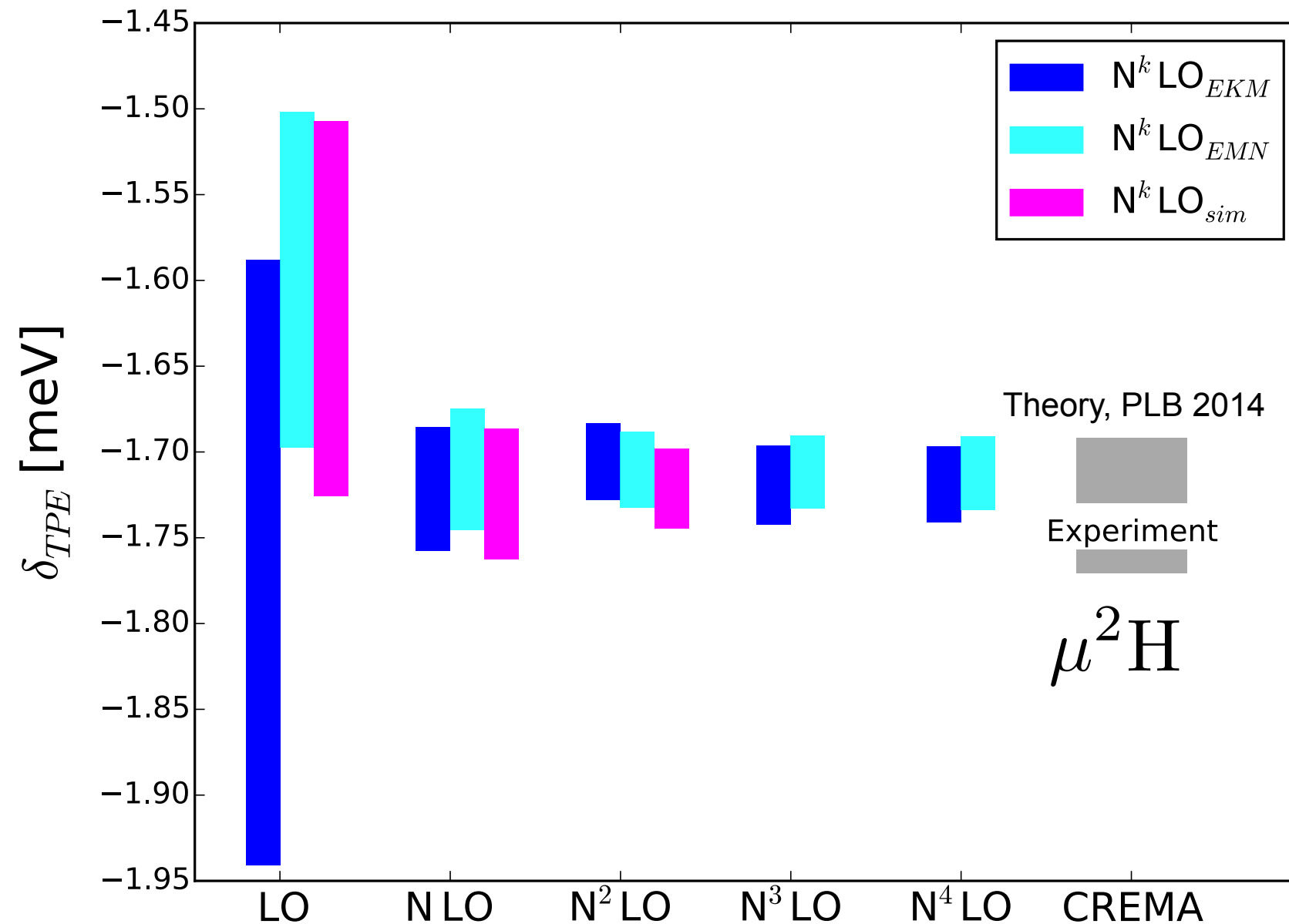
$$\Delta E_{2S-2P} = \delta_{\text{QED}} + \mathcal{A}_{\text{OPE}} \langle r_c^2 \rangle + \delta_{\text{TPE}}$$

$\mu\text{H}+\text{iso}$ :  $r_p$  from  $\mu\text{H}$  and deuterium isotopic shift  $r_d^2 - r_p^2$ : Parthey et al., PRL **104** 233001 (2010)

# Order-by-order chiral expansion

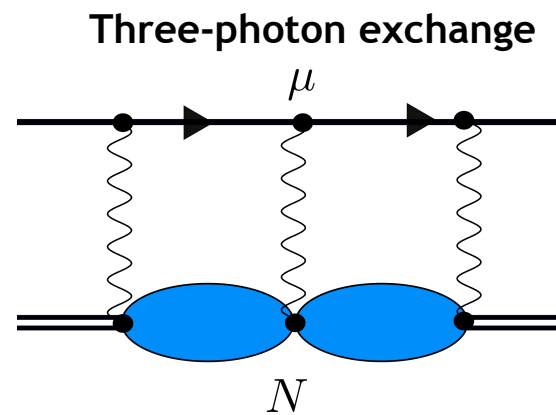
Statistical and systematic uncertainty analysis

J. Hernandez et al., Phys. Lett. B **778**, 377 (2018)



Only slightly mitigate the “small” proton radius puzzle (2.6 to  $2\sigma$ )

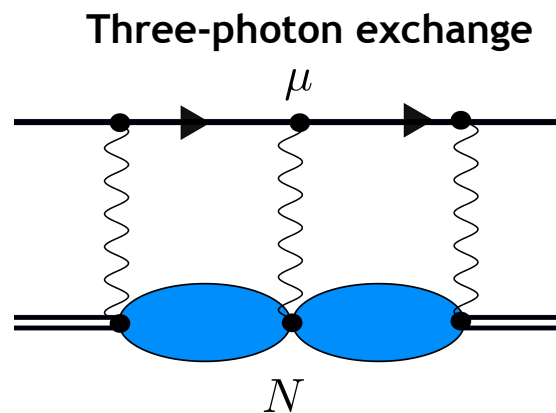
# Higher order corrections in $\alpha$



Pachucki et al., Phys. Rev. A **97** 062511 (2018)

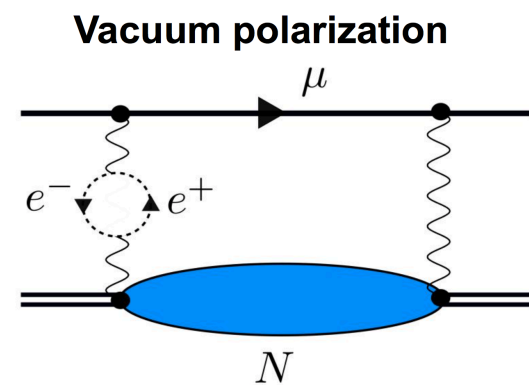
$(Z\alpha)^6$  correction, negligible

# Higher order corrections in $\alpha$



Pachucki et al., Phys. Rev. A **97** 062511 (2018)

$(Z\alpha)^6$  correction, negligible



One the many  $\alpha^6$  corrections, supposedly the largest

Kalinowski, Phys. Rev. A **99** 030501 (2019)

$$\delta_{\text{TPE}} = -1.750_{-16}^{+14} \text{ meV} \quad \text{Theory}$$

$$\delta_{\text{TPE}} = -1.7638(68) \text{ meV} \quad \text{Exp}$$

Consistent within  $1\sigma$

solves the small deuteron-radius puzzle

Large deuteron-radius puzzle still unsolved

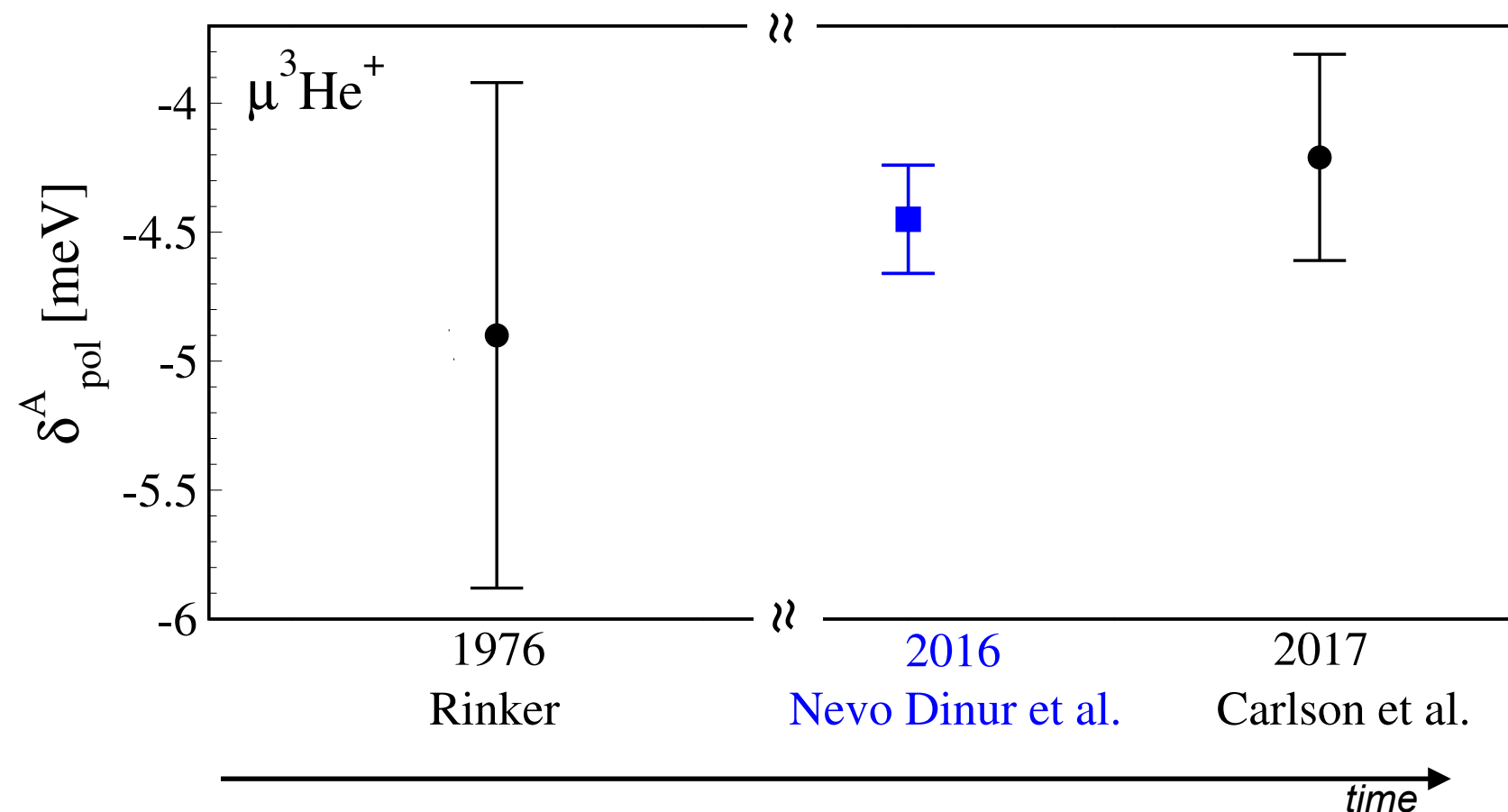
# Reduction of Uncertainties

Atom	Exp uncertainty on $\Delta E_{2S-2P}$	Uncertainty on TPE prior to the discovery of the puzzle	Uncertainty on TPE: <i>ab initio</i>
$\mu^2\text{H}$	0.003 meV	0.03 meV	0.02 meV
$\mu^3\text{He}^+$	0.08 meV	1 meV	0.3 meV
$\mu^4\text{He}^+$	0.06 meV	0.6 meV	0.4 meV
$\mu^{6,7}\text{Li}^{++}$	0.7 meV	4 meV	

# Reduction of Uncertainties

Atom	Exp uncertainty on $\Delta E_{2S-2P}$	Uncertainty on TPE prior to the discovery of the puzzle	Uncertainty on TPE: <i>ab initio</i>
$\mu^2\text{H}$	0.003 meV	0.03 meV	0.02 meV
$\mu^3\text{He}^+$	0.08 meV	1 meV	0.3 meV
$\mu^4\text{He}^+$	0.06 meV	0.6 meV	0.4 meV
$\mu^{6,7}\text{Li}^{++}$	0.7 meV	4 meV	

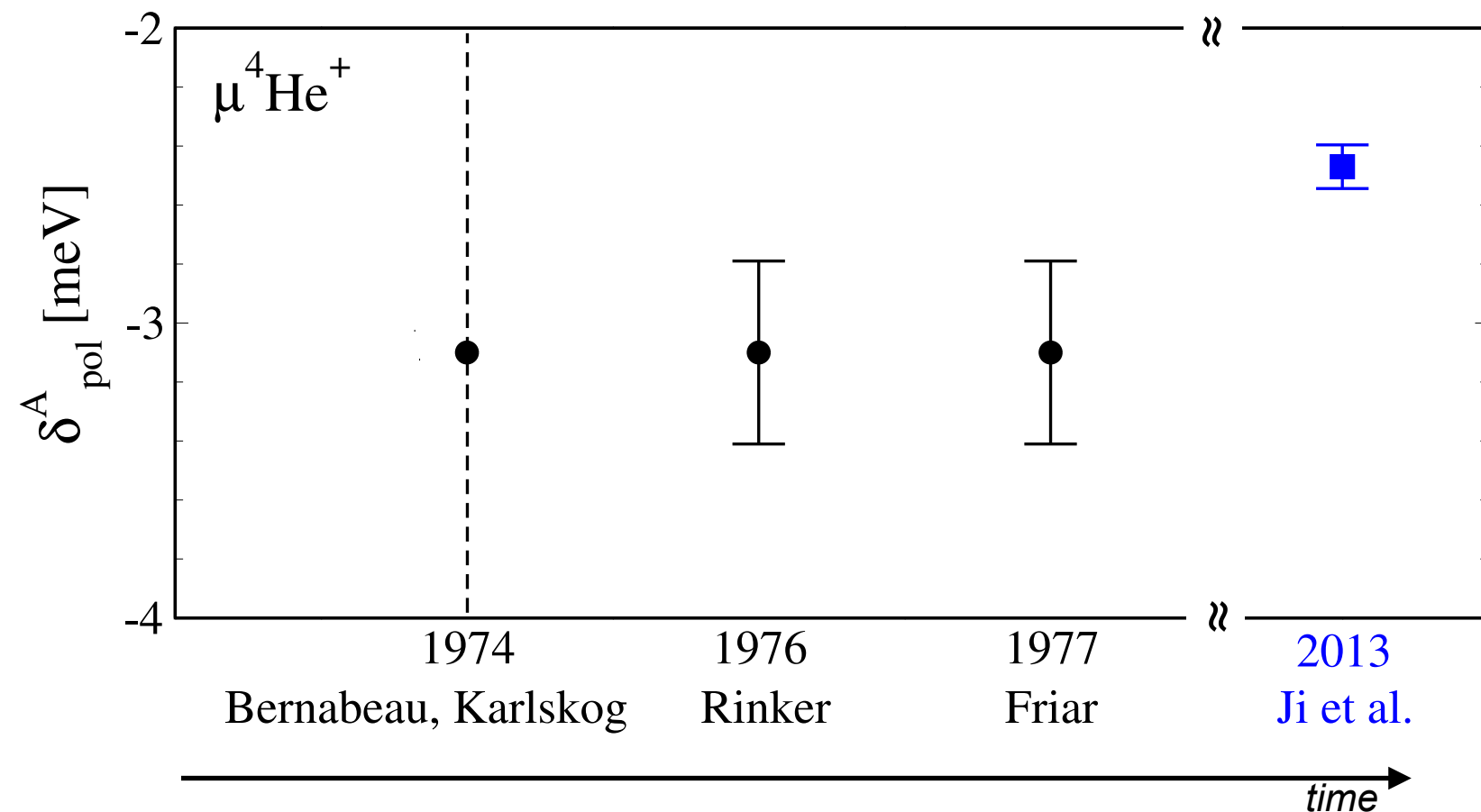
C.Ji et al., JPG: Part. Nucl. 45, 093002 (2018)



# Reduction of Uncertainties

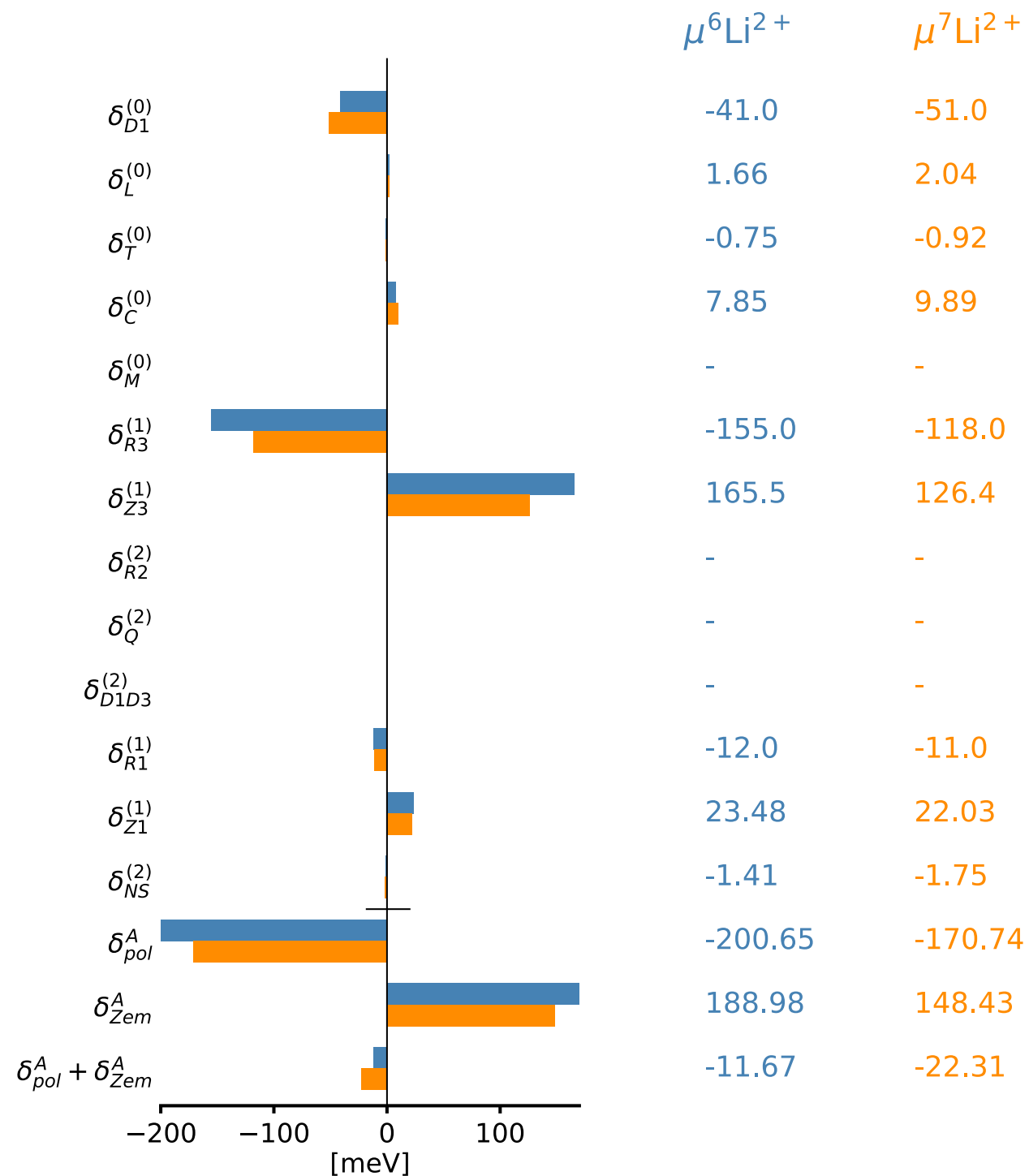
Atom	Exp uncertainty on $\Delta E_{2S-2P}$	Uncertainty on TPE prior to the discovery of the puzzle	Uncertainty on TPE: <i>ab initio</i>
$\mu^2\text{H}$	0.003 meV	0.03 meV	0.02 meV
$\mu^3\text{He}^+$	0.08 meV	1 meV	0.3 meV
$\mu^4\text{He}^+$	0.06 meV	0.6 meV	0.4 meV
$\mu^{6,7}\text{Li}^{++}$	0.7 meV	4 meV	

C.Ji et al., JPG: Part. Nucl. 45, 093002 (2018)



# Muonic Lithiums, unpublished

Preliminary



S. Li Muli et al.





# Summary and Outlook

- Ab initio calculations have allowed to microscopically compute TPE and to substantially reduce uncertainties
- Independently on the nature of the puzzle, these calculations are needed to support any spectroscopic measurement with muonic atoms
- In the future we will investigate the hyperfine splitting of muonic deuterium and the Lamb shift in muonic Lithium and Beryllium atoms

# Summary and Outlook

- Ab initio calculations have allowed to microscopically compute TPE and to substantially reduce uncertainties
- Independently on the nature of the puzzle, these calculations are needed to support any spectroscopic measurement with muonic atoms
- In the future we will investigate the hyperfine splitting of muonic deuterium and the Lamb shift in muonic Lithium and Beryllium atoms

Thanks to my collaborators

N.Barnea, O.J. Hernandez, C.Ji, S.Li Muli, N.Nevo Dinur, A. Poggialini

# Summary and Outlook

- Ab initio calculations have allowed to microscopically compute TPE and to substantially reduce uncertainties
- Independently on the nature of the puzzle, these calculations are needed to support any spectroscopic measurement with muonic atoms
- In the future we will investigate the hyperfine splitting of muonic deuterium and the Lamb shift in muonic Lithium and Beryllium atoms

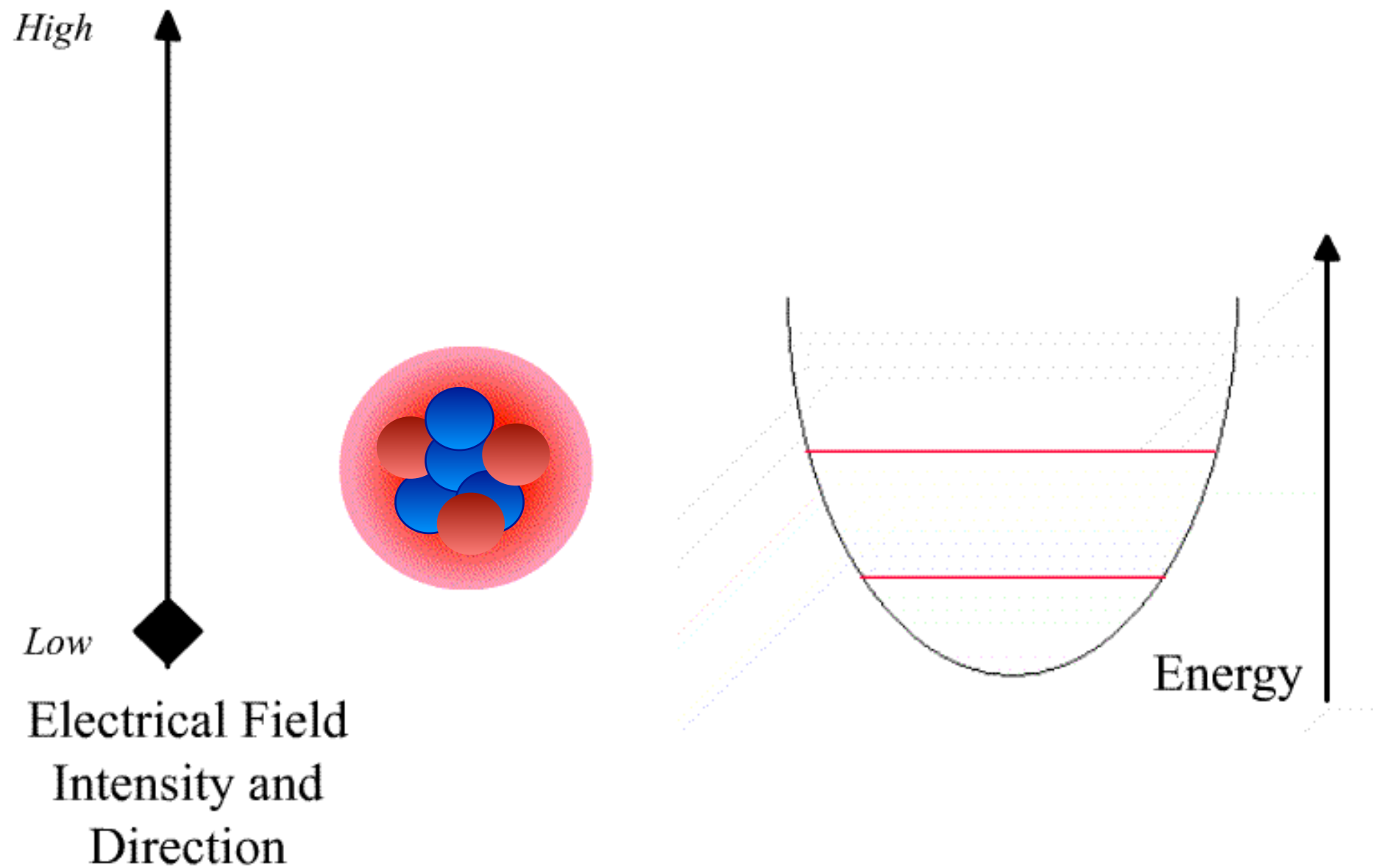
Thanks to my collaborators

N.Barnea, O.J. Hernandez, C.Ji, S.Li Muli, N.Nevo Dinur, A. Poggialini

Thank you for your attention!

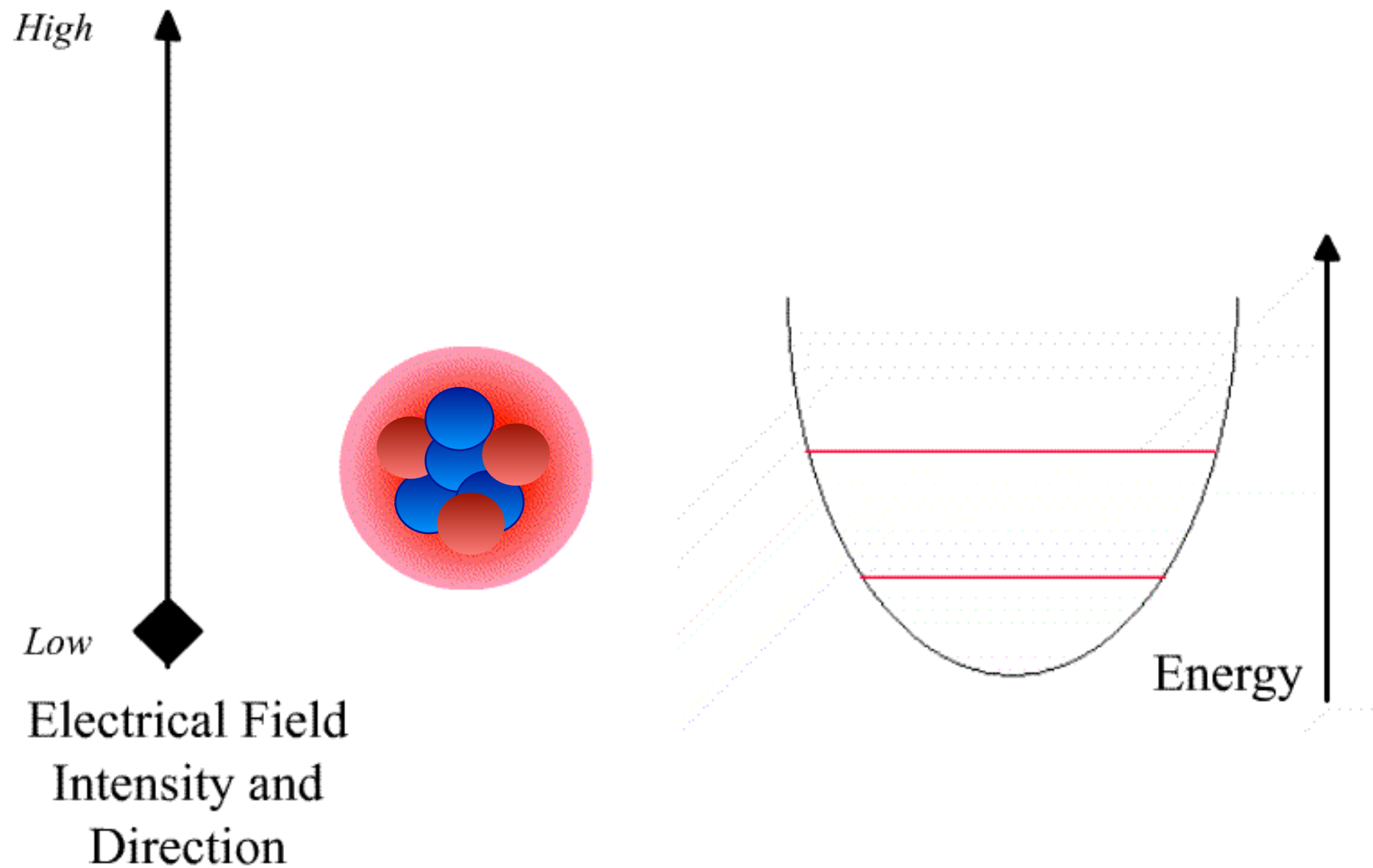
# Backup Slides

# Concept of polarizability



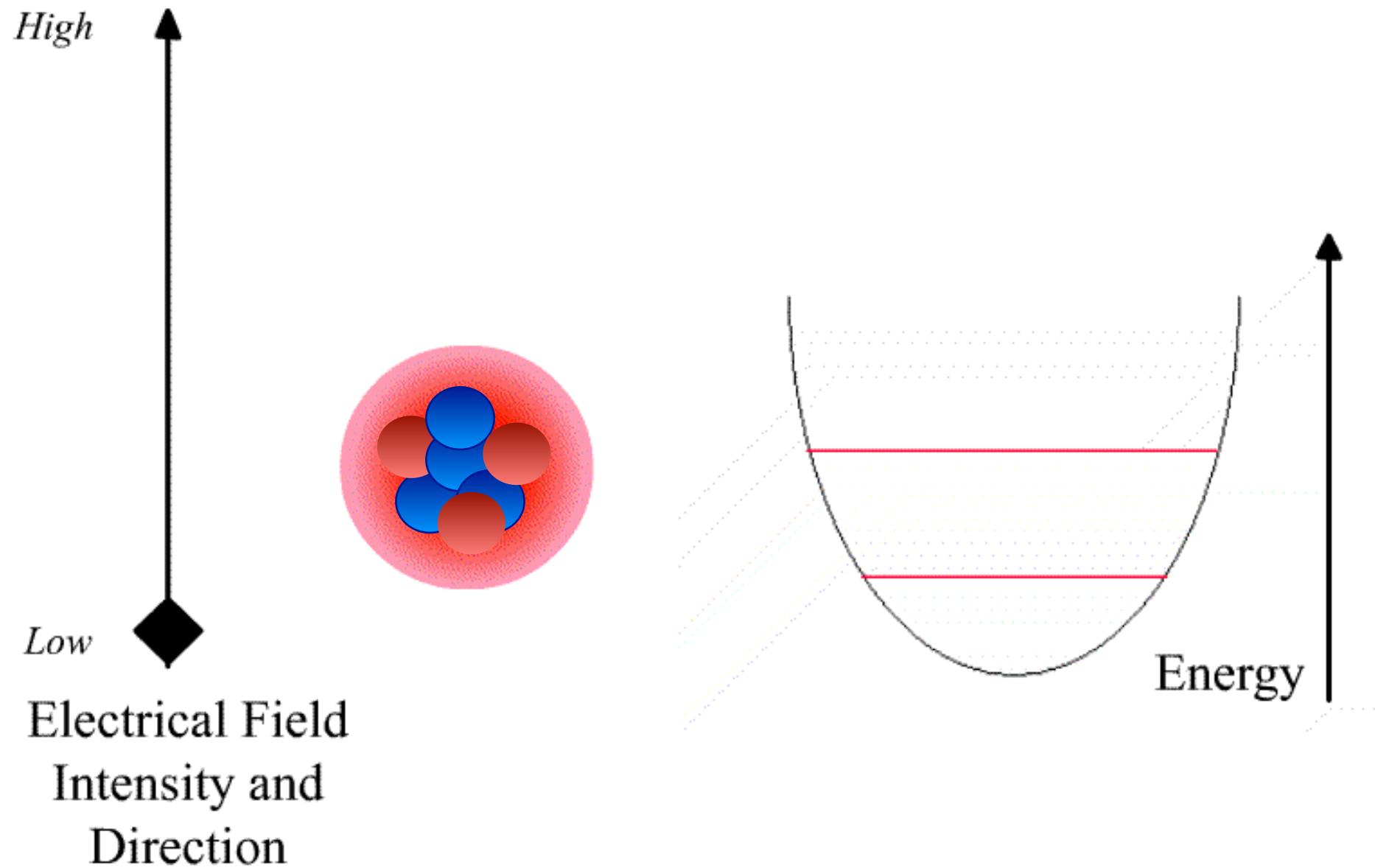
$$\mathbf{P} = \alpha_{\mathbf{E}1} \mathbf{E}$$

# Concept of polarizability



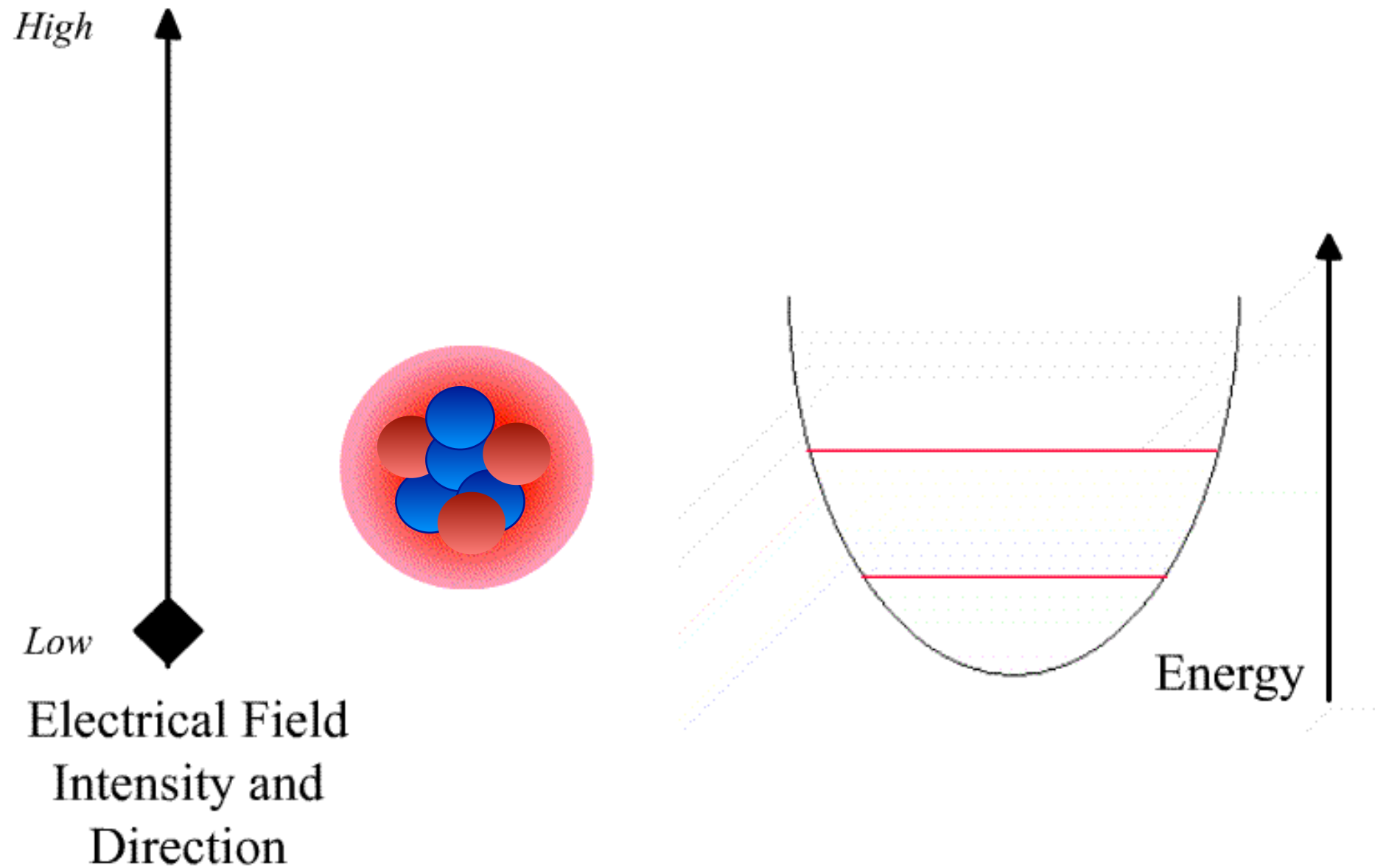
$$\mathbf{P} = \alpha_{\mathbf{E}1} \mathbf{E}$$

# Concept of polarizability



$$\mathbf{P} = \alpha_{\mathbf{E}1} \mathbf{E}$$

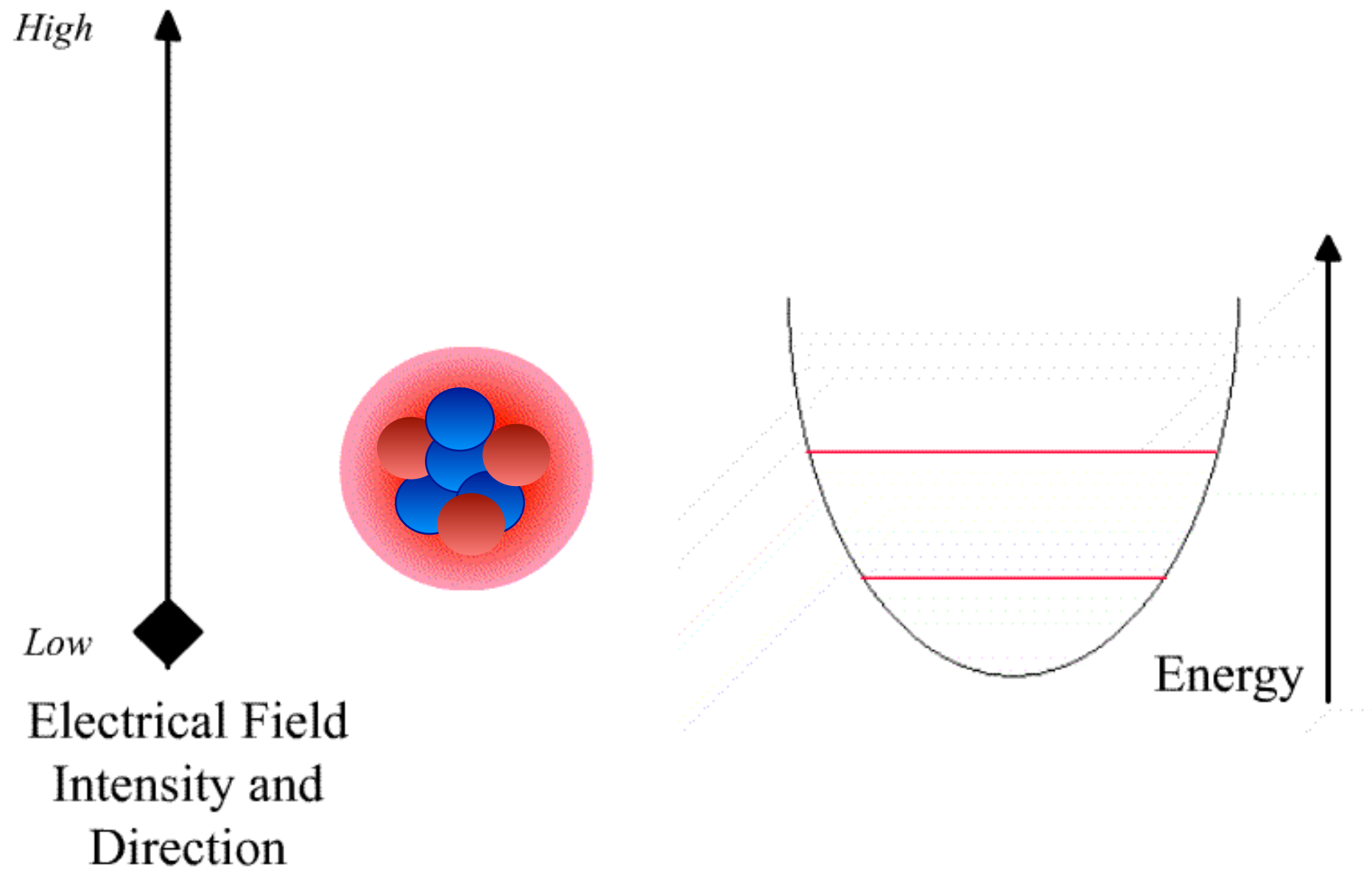
# Concept of polarizability



$$\mathbf{P} = \alpha_{\mathbf{E}1} \mathbf{E}$$



# Concept of polarizability



$$\mathbf{P} = \alpha_{\mathbf{E}1} \mathbf{E}$$

# Uncertainties quantifications

## Uncertainties sources

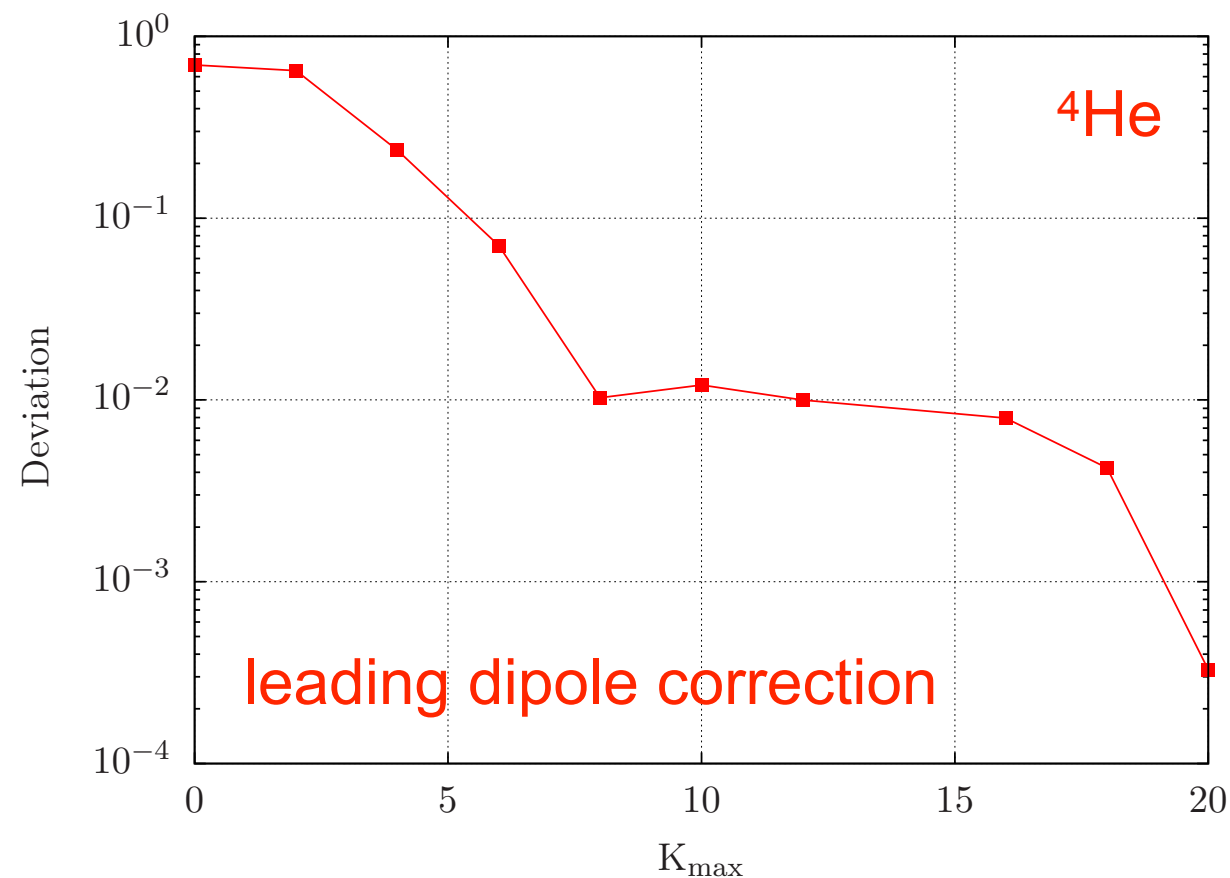
- Numerical
- Nuclear model
- Isospin symmetry breaking
- Nucleon-size
- Truncation of multiples
- $\eta$ -expansion
- expansion in  $Z\alpha$

# Uncertainties quantifications

## Uncertainties sources

- Numerical
- Nuclear model
- Isospin symmetry breaking
- Nucleon-size
- Truncation of multiples
- $\eta$ -expansion
- expansion in  $Z\alpha$

C.Ji et al., JPG: Part. Nucl. **45**, 093002 (2018)



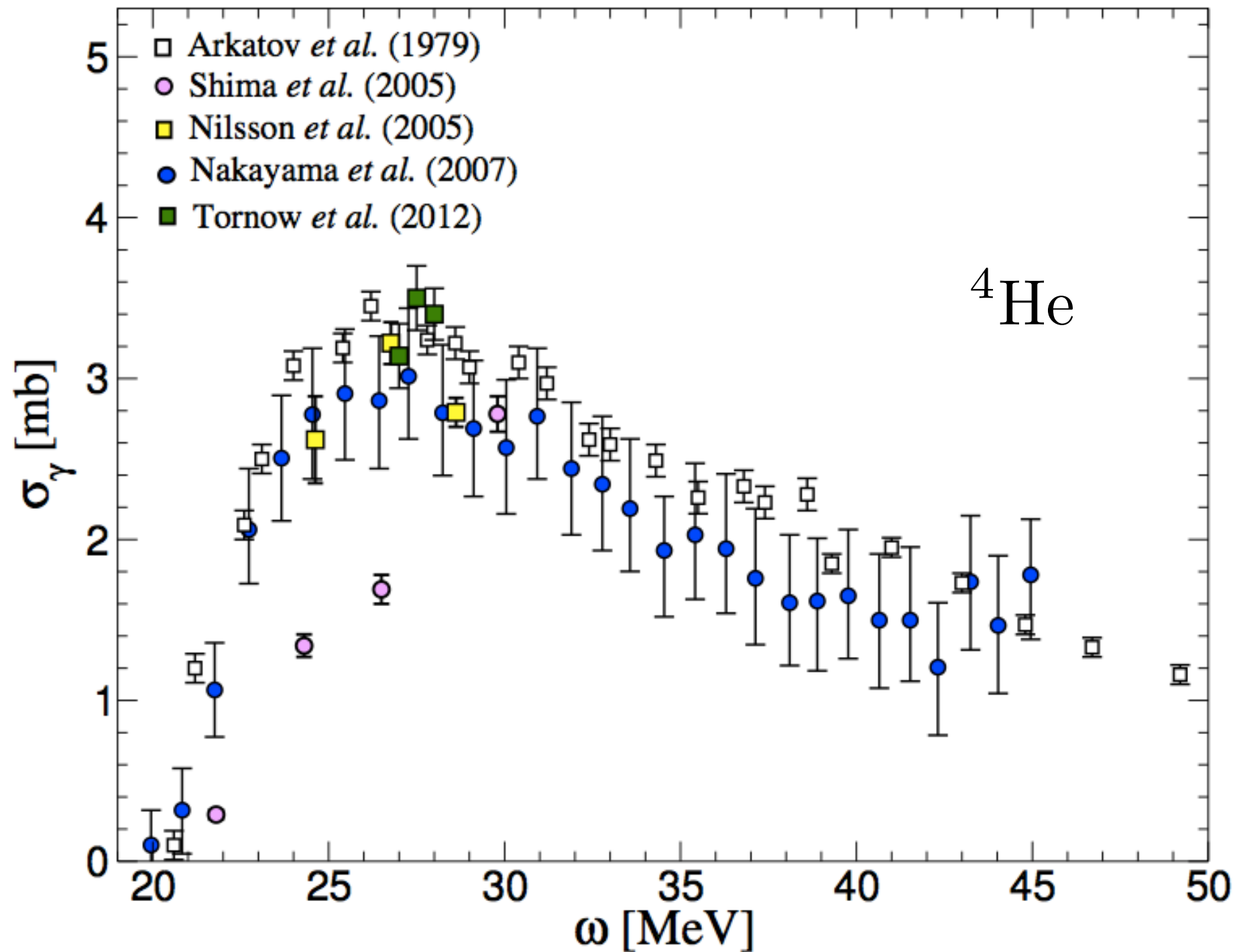
# Uncertainties quantifications

- Uncertainties sources
- Numerical
  - Nuclear model
  - Isospin symmetry breaking
  - Nucleon-size
  - Truncation of multiples
  - $\eta$ -expansion
  - expansion in  $Z\alpha$

C.Ji et al., JPG: Part. Nucl. 45, 093002 (2018)

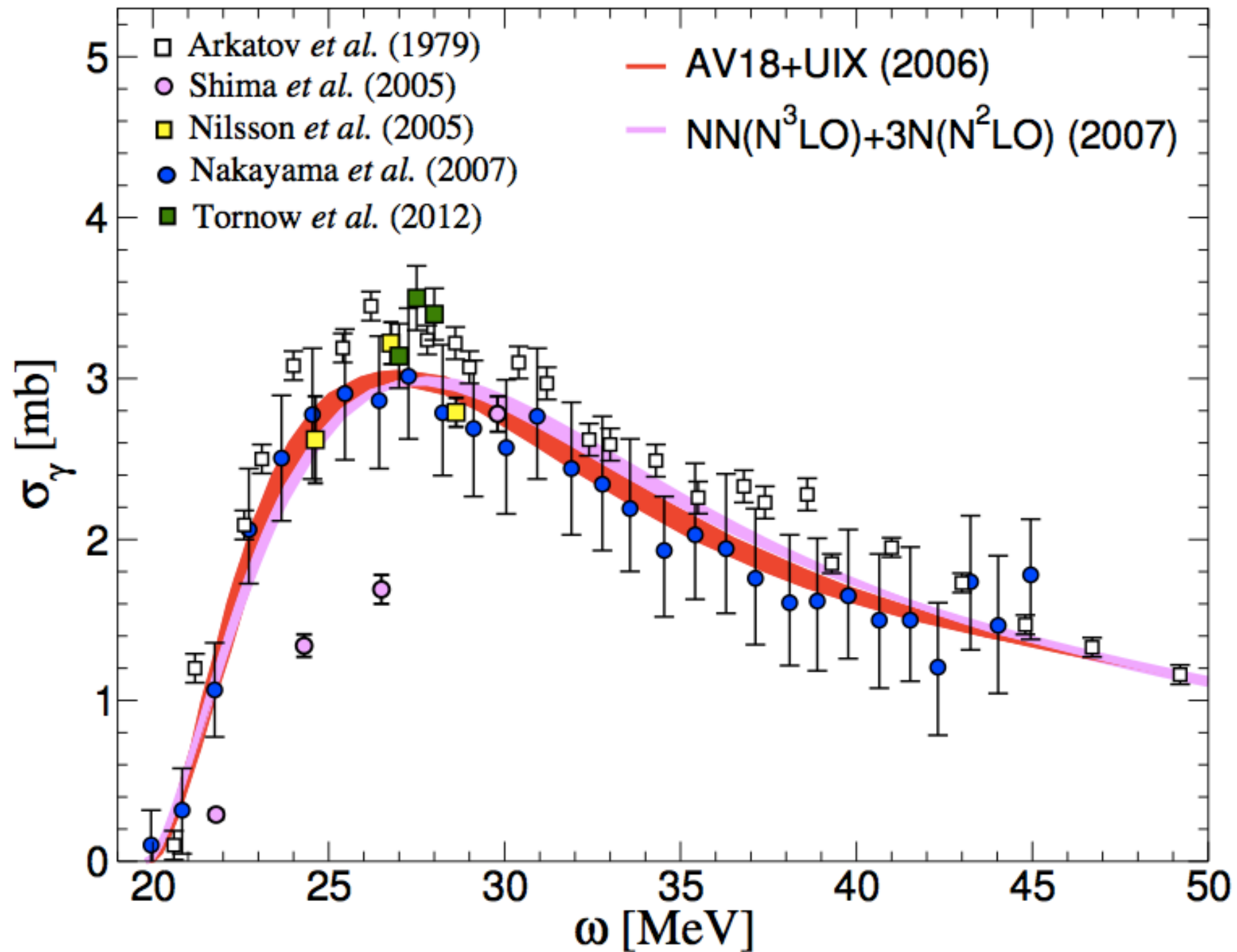
	$\mu^2\text{H}$			$\mu^3\text{H}$			$\mu^3\text{He}^+$			$\mu^4\text{He}^+$		
	$\delta_{\text{pol}}^A$	$\delta_{\text{Zem}}^A$	$\delta_{\text{TPE}}^A$	$\delta_{\text{pol}}^A$	$\delta_{\text{Zem}}^A$	$\delta_{\text{TPE}}^A$	$\delta_{\text{pol}}^A$	$\delta_{\text{Zem}}^A$	$\delta_{\text{TPE}}^A$	$\delta_{\text{pol}}^A$	$\delta_{\text{Zem}}^A$	$\delta_{\text{TPE}}^A$
Numerical	0.0	0.0	0.0	0.1	0.0	0.1	0.4	0.1	0.1	0.4	0.3	0.4
Nuclear-model	0.3	0.5	0.4	1.3	2.4	1.7	0.7	1.8	1.5	3.9	4.6	4.4
ISB	0.2	0.2	0.2	0.7	0.2	0.5	1.8	0.2	0.5	2.2	0.5	0.5
Nucleon-size	0.3	0.8	0.0	0.6	0.9	0.2	1.2	1.3	0.9	2.7	2.0	1.2
Relativistic	0.0	—	0.0	0.1	—	0.1	0.4	—	0.1	0.1	—	0.0
Coulomb	0.4	—	0.3	0.5	—	0.3	3.0	—	0.9	0.4	—	0.1
$\eta$ -expansion	0.4	—	0.3	1.3	—	0.9	1.1	—	0.3	0.8	—	0.2
Higher $Z\alpha$	0.7	—	0.5	0.7	—	0.5	1.5	—	0.4	1.5	—	0.4
Total	1.0	0.9	0.8	2.3	2.2	2.0	4.2	2.2	2.1	5.5	5.1	4.6

# An example



S.B. and Saori Pastore, Journal of Physics G.: Nucl. Part. Phys. **41**, 123002 (2014)

# An example



S.B. and Saori Pastore, Journal of Physics G.: Nucl. Part. Phys. **41**, 123002 (2014)