

On compact finite differences for the Poisson equation

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Outline

Motivation

1D Poisson problems

2D Poisson problems

3D Poisson problems

Introduction: Purpose of the talk

- Poisson equation on rectangular domains often solved by finite differences (5-point stencil).
Ditto in 3D with the 7-point stencil.
- These methods converge with $\mathcal{O}(h^2)$ in the mesh width h
- Higher orders of accuracy requires bigger stencils or more brain.
- Higher orders of accuracy lead to (much) smaller linear systems of equations for the same accuracy.
- We discuss how to get fourth order compact finite difference schemes.
- Emphasis on fast Poisson solvers: the solution is obtained by the FFT.

References

1. L. Collatz. *The Numerical Treatment of Differential Equations*. Springer, 3rd ed., 1960. (→ *Mehrstellenmethode*)
2. R. J. LeVeque. *Finite Difference Methods for Ordinary and Partial Differential Equations*. SIAM, 2007.
3. W. F. Spitz and G. F. Carey. A high-order compact formulation for the 3D Poisson equation. *Numer. Methods Partial Differ. Equations*, 12:235–243, 1996.
4. S. O. Settle, C. C. Douglas, I. Kim, and D. Sheen. On the derivation of highest-order compact finite difference schemes for the one- and two-dimensional Poisson equation with Dirichlet boundary conditions. *SINUM*, 51:2470–2490, 2013.

The 1D case: problem statement

- Interval $I = (0, a)$
- Poisson equation:

$$-u''(x) = f(x), \quad 0 < x < a, \quad u(0) = u(a) = 0.$$

- Equidistant mesh $0 = x_0 < x_1 < \dots < x_n < x_{n+1} = a$.
- Mesh width $h = x_j - x_{j-1} = a/(n+1)$.
- Approximation $u_j \approx u(x_j)$.
- Approximate Poisson equation by

$$\frac{-u_{j-1} + 2u_j - u_{j+1}}{h^2} = f(x_j), \quad 1 \leq j \leq n. \quad (1)$$

The 1D case: linear system

The n equations in (1) can be collected in matrix equation

$$\frac{1}{h^2} \mathbf{T}_n \mathbf{u} = \frac{1}{h^2} \begin{pmatrix} 2 & -1 & & & \\ -1 & 2 & -1 & & \\ & \ddots & \ddots & \ddots & \\ & & -1 & 2 & -1 \\ & & & -1 & 2 \end{pmatrix} \begin{bmatrix} u_1 \\ u_2 \\ \vdots \\ u_{n-1} \\ u_n \end{bmatrix} = \begin{bmatrix} f(x_1) \\ f(x_2) \\ \vdots \\ f(x_{n-1}) \\ f(x_n) \end{bmatrix} = \mathbf{f}.$$

$\mathbf{T}_n \in \mathbb{R}^{n \times n}$ has the spectral decomposition

$$\mathbf{T}_n = \mathbf{Q}_n \mathbf{\Lambda}_n \mathbf{Q}_n^T, \quad (2)$$

with diagonal $\mathbf{\Lambda}_n$

$$\mathbf{\Lambda}_n = \text{diag}(\lambda_1^{(n)}, \dots, \lambda_n^{(n)}), \quad \lambda_k^{(n)} = 4 \sin^2 \frac{k\pi}{2(n+1)}. \quad (3)$$

The 1D case: linear system (cont.)

\mathbf{Q}_n is orthogonal, i.e., $\mathbf{Q}_n^{-1} = \mathbf{Q}_n^T$, with elements

$$q_{jk} = \left(\frac{2}{n+1} \right)^{1/2} \sin \frac{jk\pi}{n+1}.$$

Multiplying with \mathbf{Q}_n or \mathbf{Q}_n^T is related to the Fourier transform.

If n is chosen properly then the Fast Sine Transform (\sim Fast Fourier Transform) can be employed to solve (1).

This does not make sense in the 1D case, though.

The 1D case: local truncation error

The **local truncation error** is obtained by plugging the exact solution in the FD formula,

$$\frac{-u(x-h) + 2u(x) - u(x+h)}{h^2} - f(x) = \tau(x; h)$$

The 1D case: local truncation error (cont.)

Using the Taylor series expansion

$$u(x \pm h) = u(x) \pm hu'(x) + \frac{h^2}{2}u''(x) \pm \frac{h^3}{6}u'''(x) + \frac{h^4}{24}u''''(x) + \mathcal{O}(h^5)$$

we obtain

$$\begin{aligned} & u(x_{j-1}) - 2u(x_j) + u(x_{j+1}) \\ &= u(x_j) - hu'(x_j) + \frac{h^2}{2}u''(x_j) - \frac{h^3}{6}u'''(x_j) + \frac{h^4}{24}u''''(x_j) + \dots \\ & \quad - 2u(x_j) \\ & \quad + u(x_j) + hu'(x_j) + \frac{h^2}{2}u''(x_j) + \frac{h^3}{6}u'''(x_j) + \frac{h^4}{24}u''''(x_j) + \dots \\ &= h^2u''(x_j) + \frac{h^4}{12}u''''(x_j) + \mathcal{O}(h^6) \end{aligned}$$

The 1D case: local truncation error (cont.)

Using the Taylor series expansion

$$u(x \pm h) = u(x) \pm hu'(x) + \frac{h^2}{2}u''(x) \pm \frac{h^3}{6}u'''(x) + \frac{h^4}{24}u''''(x) + \mathcal{O}(h^5)$$

we obtain

$$u(x_{j-1}) - 2u(x_j) + u(x_{j+1}) = h^2u''(x_j) + \frac{h^4}{12}u''''(x_j) + \mathcal{O}(h^6)$$

or, using $-u''(x) = f(x)$,

$$\frac{-u(x_{j-1}) + 2u(x_j) - u(x_{j+1}))}{h^2} = f(x_j) + \underbrace{\frac{h^2}{12}u''''(x_j) + \mathcal{O}(h^4)}_{\tau(x_j)} \quad (4)$$

The 1D case: local truncation error (cont.)

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$$u(x \pm h) = u(x) \pm hu'(x) + \frac{h^2}{2}u''(x) \pm \frac{h^3}{6}u'''(x) + \frac{h^4}{24}u''''(x) + \mathcal{O}(h^5)$$

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$$u(x_{j-1}) - 2u(x_j) + u(x_{j+1}) = h^2u''(x_j) + \frac{h^4}{12}u''''(x_j) + \mathcal{O}(h^6)$$

or, using $-u''(x) = f(x)$,

$$\frac{-u(x_{j-1}) + 2u(x_j) - u(x_{j+1}))}{h^2} = f(x_j) - \underbrace{\frac{h^2}{12}u''''(x_j) + \mathcal{O}(h^4)}_{\tau(x_j)} \quad (4)$$

The 1D case: global error

$$\frac{-u_{j-1} + 2u_j - u_{j+1}}{h^2} = f(x_j). \quad (5)$$

$$\frac{-u(x_{j-1}) + 2u(x_j) - u(x_{j+1}))}{h^2} = -u''(x_j) + \tau(x_j). \quad (6)$$

Subtracting (5) from (6) we get for the **error** $e(x_j) = u(x_j) - u_j$

$$h^{-2} \mathbf{T}_n \mathbf{e} = \boldsymbol{\tau}.$$

So, the L_2 -error behaves like the local truncation error since

$$\|h^2 \mathbf{T}_n^{-1}\|_2 < C \quad \text{for all } h \text{ (or } n).$$

C^{-1} is a lower bound for the smallest eigenvalue of $h^{-2} \mathbf{T}_n$.

The 1D case: Improving accuracy

1. Use longer stencil

$$\frac{1}{12h^2}(-u_{j-2} + 16u_{j-1} - 30u_j + 16u_{j+1} - u_{j+2}) = f(x_j)$$

- 2.

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2. Closer look at truncation error

$$\frac{-u(x_{j-1}) + 2u(x_j) - u(x_{j+1}))}{h^2} = -u''(x_j) + \tau(x_j).$$

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Replace finite difference stencil by

$$\frac{-u_{j-1} + 2u_j - u_{j+1}}{h^2} = f(x_j) + \frac{h^2}{12}f''(x_j) \quad (7)$$

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Replace finite difference stencil by

$$\frac{-u_{j-1} + 2u_j - u_{j+1}}{h^2} = f(x_j) + \frac{h^2}{12}f''(x_j) \quad (7)$$

or

$$\frac{-u_{j-1} + 2u_j - u_{j+1}}{h^2} = f(x_j) + \frac{1}{12}(f(x_{j-1}) - 2f(x_j) + f(x_{j+1}))) \quad (8)$$

The 1D case: Matlab demo

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generate_convergence_plot1D
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The 2D case: problem statement

- Rectangle $\Omega = (0, a_x) \times (0, a_y)$
- Poisson equation

$$-\nabla^2 u(x, y) = f(x, y) \quad \text{in } \Omega, \quad u = 0 \quad \text{on } \partial\Omega. \quad (9)$$

- Rectangular mesh: $n_x + 2 \times n_y + 2$ grid points (incl. boundary)
- Mesh widths: $h_x = a_x / (n_x + 1)$ and $h_y = a_y / (n_y + 1)$
- **5-point stencil** is most used approximation of the Laplacian
- Approximation $u_{ij} \approx u(x_i, y_j)$
- Approximate Poisson equation by

$$\frac{-u_{i-1,j} + 2u_{ij} - u_{i+1,j}}{h_x^2} + \frac{-u_{i,j-1} + 2u_{ij} - u_{i,j+1}}{h_y^2} = f(x_i, y_j) \quad (10)$$

for $0 < i \leq n_x, 0 < j \leq n_y$.

The 2D case: stencil

Often, the discretized Poisson equation is displayed as a *stencil*

$$-\nabla_{\frac{2}{5}}^2 u(x, y) = \begin{array}{c} -\frac{1}{h_y^2} \\ \bullet \\ | \\ -\frac{1}{h_x^2} \quad \frac{2}{h_x^2} \quad \frac{2}{h_y^2} \quad -\frac{1}{h_x^2} \\ \bullet \quad \blacksquare \quad \bullet \\ | \\ -\frac{1}{h_y^2} \\ \bullet \end{array} \odot u(x, y) = f(x, y)$$

which shows nicely the five involved grid points with their weights.

The 2D case: linear system

Collect the $u_{ij}/f(x_i, y_j)$ in a vector $\mathbf{u}, \mathbf{f} \in \mathbb{R}^{n_x n_y}$.

The $n_x n_y$ equations in (10) can be collected in matrix form

$$\left(\frac{1}{h_x^2} \mathbf{I}_{n_y} \otimes \mathbf{T}_{n_x} + \frac{1}{h_y^2} \mathbf{T}_{n_y} \otimes \mathbf{I}_{n_x} \right) \mathbf{u} = \mathbf{f}, \quad (11)$$

where \otimes denotes **Kronecker product**. Then, (11) can be written as

$$(\mathbf{Q}_{n_y} \otimes \mathbf{Q}_{n_x}) \left(\frac{1}{h_x^2} \mathbf{I}_{n_y} \otimes \mathbf{\Lambda}_{n_x} + \frac{1}{h_y^2} \mathbf{\Lambda}_{n_y} \otimes \mathbf{I}_{n_x} \right) (\mathbf{Q}_{n_y}^T \otimes \mathbf{Q}_{n_x}^T) \mathbf{u} = \mathbf{f}. \quad (12)$$

Matrix in the middle is **diagonal**.

With $n = n_x n_y$, (12) can be solved with $\mathcal{O}(n \log n)$ flops, if FFT is applicable.

The 2D case: truncation error

Local truncation error for 5-point stencil is

$$-\nabla_5^2 u(x, y) - f(x, y) = -\frac{h_x^2}{12} \partial_x^4 u(x, y) - \frac{h_y^2}{12} \partial_y^4 u(x, y) + \mathcal{O}(h_x^4 + h_y^4).$$

Can we do better in 2D as well?

The 2D case: improving accuracy

Define a 9-point (**compact**) stencil

$$\begin{aligned} \nabla_9^2 u_{i,j} \equiv & \nabla_5^2 u_{i,j} + \frac{1}{12} \left(4u_{i,j} - 2(u_{i+1,j} + u_{i-1,j} + u_{i,j+1} + u_{i,j-1}) \right. \\ & \left. + u_{i+1,j+1} + u_{i-1,j+1} + u_{i+1,j-1} + u_{i-1,j-1} \right) \left(\frac{1}{h_x^2} + \frac{1}{h_y^2} \right). \end{aligned}$$

For the local truncation error of the Poisson equation we get

$$\begin{aligned} -\nabla_9^2 u(x, y) - f(x, y) = & -\frac{h_x^2}{12} (\partial_x^4 u(x, y) + \partial_x^2 \partial_y^2 u(x, y)) \\ & -\frac{h_y^2}{12} (\partial_x^2 \partial_y^2 u(x, y) + \partial_y^4 u(x, y)) + \mathcal{O}((h_x^2 + h_y^2)^2), \end{aligned}$$

which does not look like an improvement w.r.t. the 5-pt stencil.

The 2D case: improving accuracy (cont.)

BUT

$$\begin{aligned} -\nabla_9^2 u(x, y) - f(x, y) &= -\frac{h_x^2}{12} (\partial_x^4 u(x, y) + \partial_x^2 \partial_y^2 u(x, y)) \\ &\quad -\frac{h_y^2}{12} (\partial_x^2 \partial_y^2 u(x, y) + \partial_y^4 u(x, y)) + \mathcal{O}((h_x^2 + h_y^2)^2) \end{aligned}$$

The 2D case: improving accuracy (cont.)

BUT

$$\begin{aligned} -\nabla_9^2 u(x, y) - f(x, y) &= -\frac{h_x^2}{12} (\partial_x^4 u(x, y) + \partial_x^2 \partial_y^2 u(x, y)) \\ &\quad - \frac{h_y^2}{12} (\partial_x^2 \partial_y^2 u(x, y) + \partial_y^4 u(x, y)) + \mathcal{O}((h_x^2 + h_y^2)^2) \\ &= -\frac{h_x^2}{12} (\partial_x^2 (\partial_x^2 u(x, y) + \partial_y^2 u(x, y))) \\ &\quad - \frac{h_y^2}{12} (\partial_y^2 (\partial_x^2 u(x, y) + \partial_y^2 u(x, y))) + \dots \end{aligned}$$

The 2D case: improving accuracy (cont.)

BUT

$$\begin{aligned} -\nabla_9^2 u(x, y) - f(x, y) &= -\frac{h_x^2}{12} (\partial_x^4 u(x, y) + \partial_x^2 \partial_y^2 u(x, y)) \\ &\quad - \frac{h_y^2}{12} (\partial_x^2 \partial_y^2 u(x, y) + \partial_y^4 u(x, y)) + \mathcal{O}((h_x^2 + h_y^2)^2) \\ &= -\frac{h_x^2}{12} (\partial_x^2 (\partial_x^2 u(x, y) + \partial_y^2 u(x, y))) \\ &\quad - \frac{h_y^2}{12} (\partial_y^2 (\partial_x^2 u(x, y) + \partial_y^2 u(x, y))) + \dots \\ &= -\frac{h_x^2}{12} \partial_x^2 \nabla^2 u(x, y) - \frac{h_y^2}{12} \partial_y^2 \nabla^2 u(x, y) + \dots \end{aligned}$$

The 2D case: improving accuracy (cont.)

BUT

$$\begin{aligned} -\nabla_9^2 u(x, y) - f(x, y) &= -\frac{h_x^2}{12} (\partial_x^4 u(x, y) + \partial_x^2 \partial_y^2 u(x, y)) \\ &\quad - \frac{h_y^2}{12} (\partial_x^2 \partial_y^2 u(x, y) + \partial_y^4 u(x, y)) + \mathcal{O}((h_x^2 + h_y^2)^2) \\ &= -\frac{h_x^2}{12} (\partial_x^2 (\partial_x^2 u(x, y) + \partial_y^2 u(x, y))) \\ &\quad - \frac{h_y^2}{12} (\partial_y^2 (\partial_x^2 u(x, y) + \partial_y^2 u(x, y))) + \dots \\ &= -\frac{h_x^2}{12} \partial_x^2 \nabla^2 u(x, y) - \frac{h_y^2}{12} \partial_y^2 \nabla^2 u(x, y) + \dots \\ &= \frac{h_x^2}{12} \partial_x^2 f(x, y) + \frac{h_y^2}{12} \partial_y^2 f(x, y) + \mathcal{O}((h_x^2 + h_y^2)^2) \end{aligned}$$

The 2D case: improving accuracy (cont.)

If the second derivatives of f not available or too expensive to compute, replace them by finite differences:

$$\frac{1}{12} \times \begin{array}{ccccc} -\frac{1}{h_x^2} & -\frac{1}{h_y^2} & \frac{2}{h_x^2} & -\frac{10}{h_y^2} & -\frac{1}{h_x^2} & -\frac{1}{h_y^2} \\ \bullet & & \bullet & & \bullet & \\ \frac{2}{h_y^2} & \frac{10}{h_x^2} & \frac{20}{h_x^2} & \frac{20}{h_y^2} & \frac{2}{h_y^2} & \frac{10}{h_x^2} \\ \bullet & & \bullet & & \bullet & \\ -\frac{1}{h_x^2} & -\frac{1}{h_y^2} & \frac{2}{h_x^2} & \frac{10}{h_y^2} & -\frac{1}{h_x^2} & -\frac{1}{h_y^2} \\ \bullet & & \bullet & & \bullet & \end{array} \odot u(x, y) = \frac{1}{12} \times \begin{array}{ccc} & 1 & \\ \bullet & & \bullet \\ & 8 & \\ \bullet & & \bullet \\ & 1 & \end{array} \odot f(x, y).$$

A fourth order local truncation error is the best one can get in 2D by compact FD (Settle et al. SINUM 2013).

The 2D case: linear system for compact FD

The matrix form of the stencil before is

$$\begin{aligned} & \left(\frac{1}{h_x^2} \mathbf{I}_{n_y} \otimes \mathbf{T}_{n_x} + \frac{1}{h_y^2} \mathbf{T}_{n_y} \otimes \mathbf{I}_{n_x} - \frac{1}{12} \left(\frac{1}{h_x^2} + \frac{1}{h_y^2} \right) \mathbf{T}_{n_y} \otimes \mathbf{T}_{n_x} \right) \mathbf{u} \\ & = \left(\mathbf{I} - \frac{1}{12} (\mathbf{I}_{n_y} \otimes \mathbf{T}_{n_x} + \mathbf{T}_{n_y} \otimes \mathbf{I}_{n_x}) \right) \mathbf{f} \end{aligned}$$

Using the spectral decompositions of the matrices \mathbf{T}_{n_x} , \mathbf{T}_{n_y} gives

$$\begin{aligned} \mathbf{u} & = (\mathbf{Q}_{n_y} \otimes \mathbf{Q}_{n_x}) \left(\mathbf{I}_{n_y} \otimes \mathbf{\Lambda}_{n_x} + \mathbf{\Lambda}_{n_y} \otimes \mathbf{I}_{n_x} - \frac{h_x^2 + h_y^2}{12} \mathbf{\Lambda}_{n_y} \otimes \mathbf{\Lambda}_{n_x} \right)^{-1} \times \\ & \quad \times h_x^2 h_y^2 \left(\mathbf{I} - \frac{1}{12} (\mathbf{I}_{n_y} \otimes \mathbf{\Lambda}_{n_x} + \mathbf{\Lambda}_{n_y} \otimes \mathbf{I}_{n_x}) \right) (\mathbf{Q}_{n_y}^T \otimes \mathbf{Q}_{n_x}^T) \mathbf{f} \end{aligned}$$

In the middle there is again a diagonal matrix.

The 2D case: Matlab demo

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generate_convergence_plot2D
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The 3D case: problem statement

- Cuboid $\Omega = (0, a_x) \times (0, a_y) \times (0, a_z)$
- Poisson equation

$$-\nabla^2 u(x, y, z) = f(x, y, z) \quad \text{in } \Omega, \quad u = 0 \text{ on } \partial\Omega. \quad (13)$$

- Rectangular mesh: $n_x + 2 \times n_y + 2 \times n_z + 2$ grid points
- Mesh widths: h_x, h_y, h_z
- **7-point stencil** is standard approximation of the Laplacian
- Approximation $u_{ij} \approx u(x_i, y_j)$
- In interior $n_x n_y n_z$ grid points approximate Poisson eq. by

$$\begin{aligned} & \frac{-u_{i-1,j,k} + 2u_{ijk} - u_{i+1,j,k}}{h_x^2} + \frac{-u_{i,j-1,k} + 2u_{ijk} - u_{i,j+1,k}}{h_y^2} \\ & + \frac{-u_{i,j,k-1} + 2u_{ijk} - u_{i,j,k+1}}{h_z^2} = f(x_i, y_j, z_k) \end{aligned}$$

The 3D case: linear system for the 7-point stencil

Collect values u_{ijk} , $f(x_i, y_j, z_k)$ in vectors \mathbf{u} , $\mathbf{f} \in \mathbb{R}^{n_x n_y n_z}$, similarly as in the 2D case. Then, the matrix form of above equations is

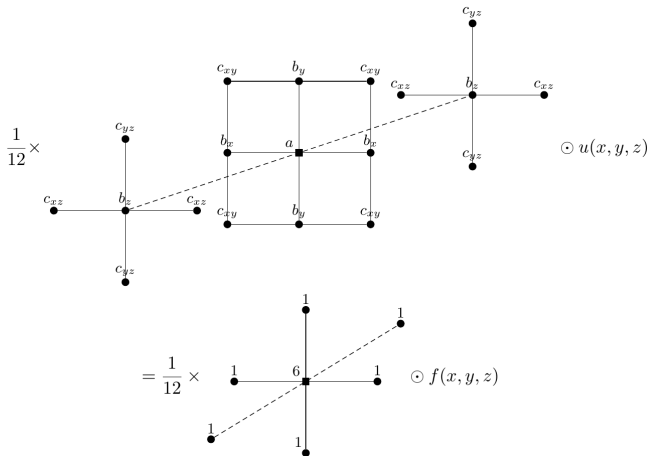
$$\left(\frac{1}{h_x^2} \mathbf{I}_{n_z} \otimes \mathbf{I}_{n_y} \otimes \mathbf{T}_{n_x} + \frac{1}{h_y^2} \mathbf{I}_{n_z} \otimes \mathbf{T}_{n_y} \otimes \mathbf{I}_{n_x} + \frac{1}{h_z^2} \mathbf{T}_{n_z} \otimes \mathbf{I}_{n_y} \otimes \mathbf{I}_{n_x} \right) \mathbf{u} = \mathbf{f}.$$

Using the spectral decomposition of the \mathbf{T} 's this becomes

$$\begin{aligned} & (\mathbf{Q}_{n_z} \otimes \mathbf{Q}_{n_y} \otimes \mathbf{Q}_{n_x}) \\ & \left(\frac{1}{h_x^2} \mathbf{I}_{n_z} \otimes \mathbf{I}_{n_y} \otimes \mathbf{\Lambda}_{n_x} + \frac{1}{h_y^2} \mathbf{I}_{n_z} \otimes \mathbf{\Lambda}_{n_y} \otimes \mathbf{I}_{n_x} + \frac{1}{h_z^2} \mathbf{\Lambda}_{n_z} \otimes \mathbf{I}_{n_y} \otimes \mathbf{I}_{n_x} \right) \\ & (\mathbf{Q}_{n_z}^T \otimes \mathbf{Q}_{n_y}^T \otimes \mathbf{Q}_{n_x}^T) \mathbf{u} = \mathbf{f}. \end{aligned}$$

The diagonal matrix in the middle can be precomputed.

The 3D case: linear system for 4th order 19-point stencil



Cf. Spetz&Carey

The 3D case: linear system for 4th order 19-point stencil (cont.)

The matrix form of this stencil is

$$\begin{aligned} & \left(\frac{1}{h_x^2} I_{n_z} \otimes I_{n_y} \otimes T_{n_x} + \frac{1}{h_y^2} I_{n_z} \otimes T_{n_y} \otimes I_{n_x} + \frac{1}{h_z^2} T_{n_z} \otimes I_{n_y} \otimes I_{n_x} \right. \\ & - \frac{1}{12} \left(\frac{1}{h_x^2} + \frac{1}{h_y^2} \right) I_{n_z} \otimes T_{n_y} \otimes T_{n_x} - \frac{1}{12} \left(\frac{1}{h_x^2} + \frac{1}{h_z^2} \right) T_{n_z} \otimes I_{n_y} \otimes T_{n_x} \\ & \quad \left. - \frac{1}{12} \left(\frac{1}{h_y^2} + \frac{1}{h_z^2} \right) T_{n_z} \otimes T_{n_y} \otimes I_{n_x} \right) \mathbf{u} \\ & = \left(I - \frac{1}{12} (I_{n_z} \otimes I_{n_y} \otimes T_{n_x} + I_{n_z} \otimes T_{n_y} \otimes I_{n_x} + T_{n_z} \otimes I_{n_y} \otimes I_{n_x}) \right) \mathbf{f}. \end{aligned}$$

Remark: Spetz&Carey also give a $\mathcal{O}(h^6)$ 27-pt stencil for the Laplacian that does not lead to a compact stencil for the Poisson equation, though.

The 3D case: numerical example

$$a_x = 1.1, \quad a_y = 1.0, \quad a_z = 0.9,$$

$$f(x, y, z) = \pi^2 \left(\frac{1}{a_x^2} + \frac{9}{a_y^2} + \frac{25}{a_z^2} \right) \sin\left(\frac{\pi x}{a_x}\right) \sin\left(\frac{3\pi y}{a_y}\right) \sin\left(\frac{5\pi z}{a_z}\right).$$

$$u(x, y, z) = \sin(\pi x/a_x) \sin(3\pi y/a_y) \sin(5\pi z/a_z).$$

in the Matlab code the approximation error is plotted versus the mesh width $h \sim 1/n$. The norm of the error is computed as

$$\|\mathbf{e}\| = \sqrt{\frac{1}{n_x n_y n_z} \sum_{i=1}^{n_x} \sum_{j=1}^{n_y} \sum_{k=1}^{n_z} |u_{i,j,k} - u(x_i, y_j, z_k)|^2}.$$

In this example we have $n = n_x = n_y = n_z$.

The 3D case: Matlab demo

```
generate_convergence_plot3D
```

Conclusions

- High-order methods can generate accurate solutions on coarse grids
- Solutions have to be smooth enough
- Matrices get denser as order increases, but we use its spectral decomposition and FFT
- Class of operators is limited, but Laplacian is fine
- In 3D 6th order is possible but the stencil for the right-hand side is not compact anymore
- To use compact FD inside other software, the (input) data has to be accurate