

Complementary uses of Free Electron Lasers and High-Harmonic Sources



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SCIENCE@FELs Villigen, 17-9-2014

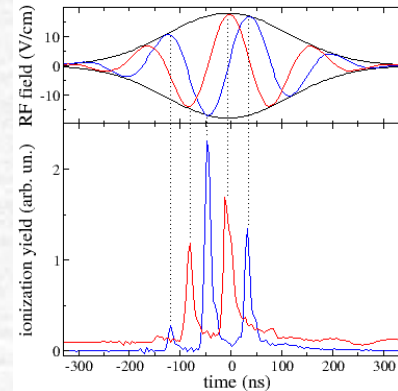
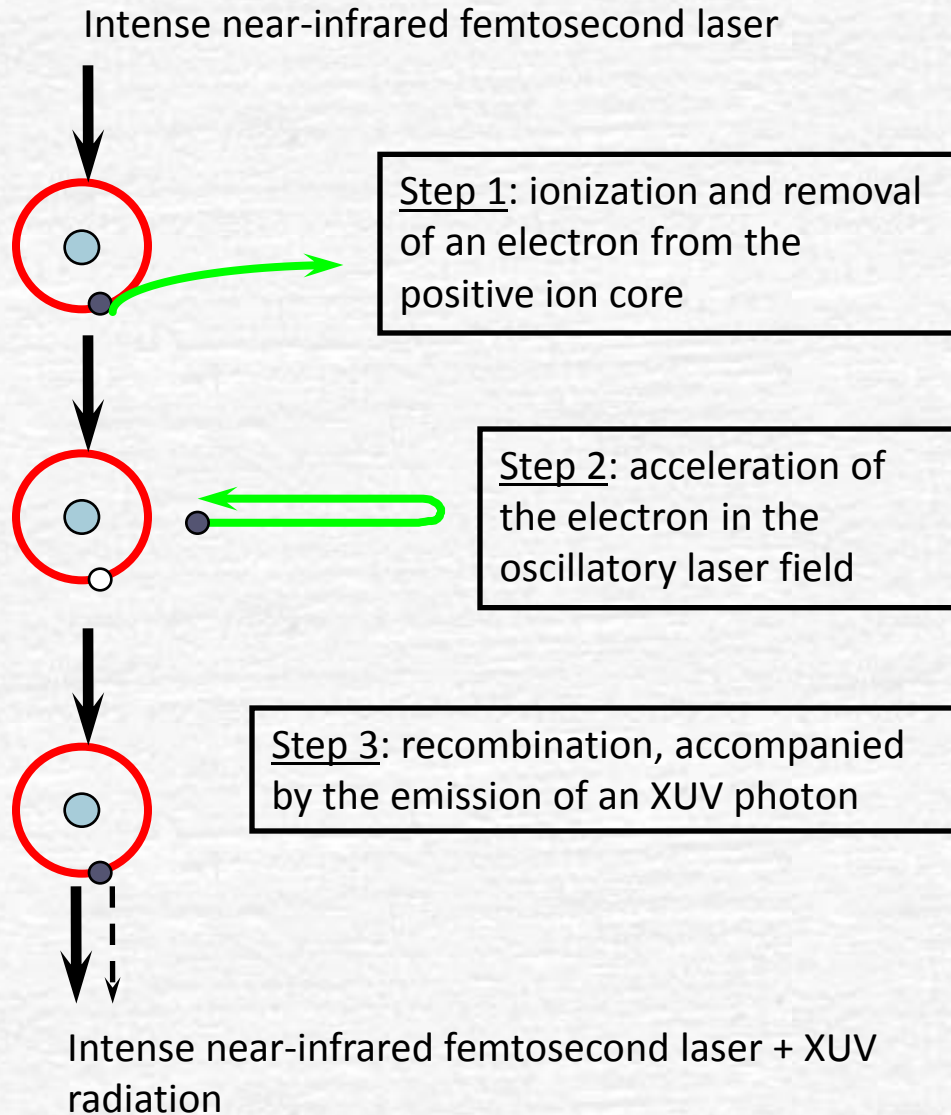
Source parameter comparison FEL vs. HHG

Parameters	Flash FEL (Hamburg 2005)	Fermi (Trieste, 2010)	LCLS (Stanford, 2009)	SACLA (Hyogo, 2011)	EU XFEL (Hamburg, 2015)
E_e	1.25 GeV	1.24 GeV	13.6 GeV	8 GeV	17.5 GeV
γ	2,450	2,300	26,600	15,700	35,000
$\hat{I}$	1.3 kA	300 A	3.4 kA	3 kA	5 kA
λ_u	27.3 mm	55 mm	30 mm	18 mm	35.6 mm
N	989	216	3733	4986	4000
L_u	27 m	14 m	112 m	90 m	200 m
$\hbar\omega$	30-300 eV (4-40 nm)	20-60 eV (20-60 nm)	250 eV - 12 keV (1-50 Å)	4.5-15 keV (0.8-2.8 Å)	4-12 keV (1-3 Å)
$\lambda/\Delta\lambda_{FWHM}$	100	1000	200-500	200-400	1000
$\Delta\tau_{FWHM}$	25 fsec	85 fsec	70 fsec	30 fsec	100 fsec
$\dot{f}$ (ph/pulse)	3×10^{12}	5×10^{12}	2×10^{12}	5×10^{11}	10^{12}
rep rate	5 Hz	10 Hz	120 Hz	60 Hz	5 Hz/27 kHz
$\hat{P}$	1 GW	1 GW	25 GW	30 GW	20 GW
L	260 m	200 m	2 km	710 m	3.4 km
Polarization	linear	variable	linear	linear	variable (?)
Mode	SASE	Seeded (3 ω Ti: sapphire)	SASE	SASE	SASE

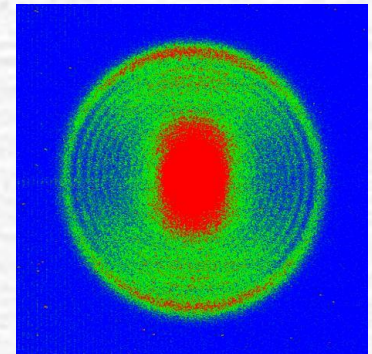
Flash II, Fermi II, SLS FEL, LCLS II,

FreeElectronLasersChart_June2014.ai

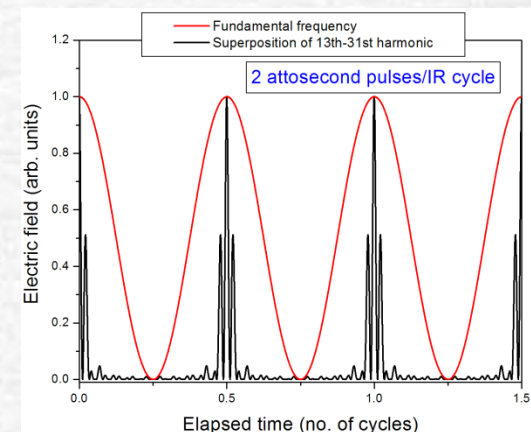
XUV/X-ray pulses from High Harmonic Generation



Guertler et.al., Phys. Rev. Lett. 92, 063901 (2004)



Nicole et.al., Phys. Rev. Lett. 88, 133001 (2002)

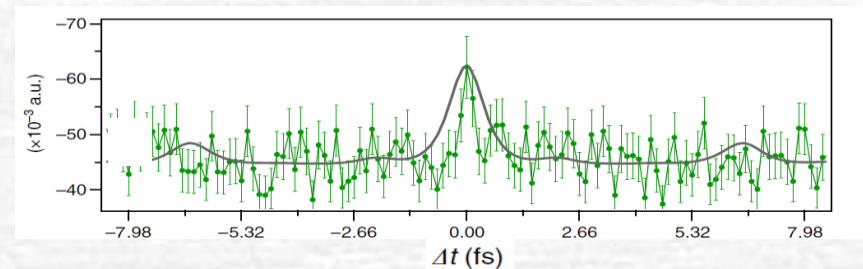


Source parameter comparison FEL vs. HHG

A sampling of what is possible using HHG

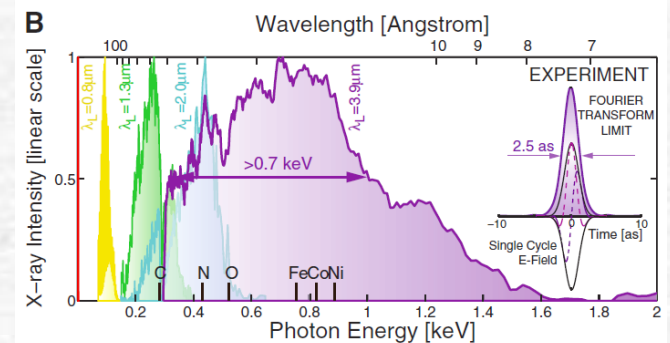
**K. Midorikawa and co-workers,
Nat. Comm. 4, 2691 (2013)**

9 mJ, 800 nm, 30 fs + 2.5 mJ, 1300 nm, 35 fs
→ 1.3 μ J pulse in 28-35 eV, 500 as



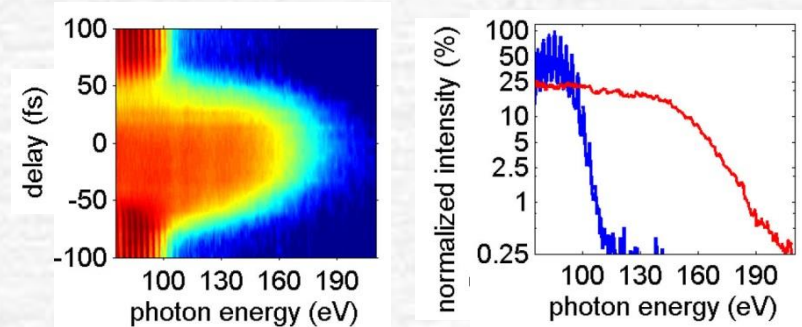
**A. Baltuska and co-workers,
Science 336, 6086 (2012)**

10 mJ, 3.9 μ m, 80 fs → 10^5 photons per shot
in a fractional bandwidth of 1% at 1 keV



**Our own work (B. Schütte et al., in
preparation)**

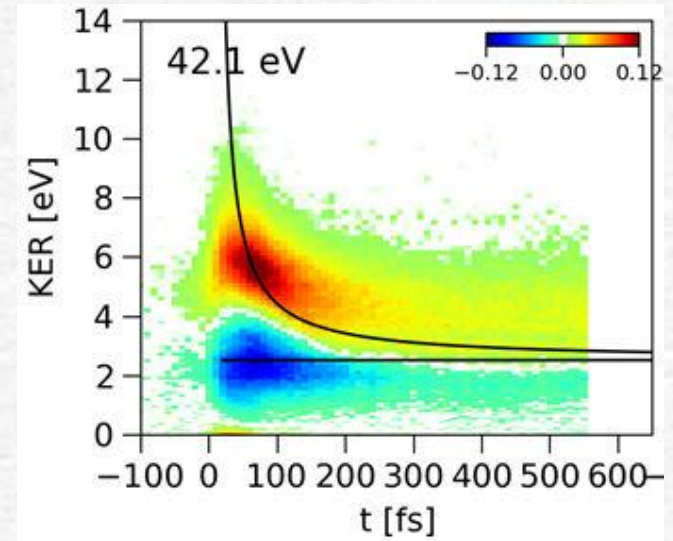
4×10^{14} W/cm², 800 nm, 50 fs + 9×10^{14} W/cm²
1300 nm, 50 fs → 10^7 photons per shot
between 70 and 160 eV



Contents of this talk

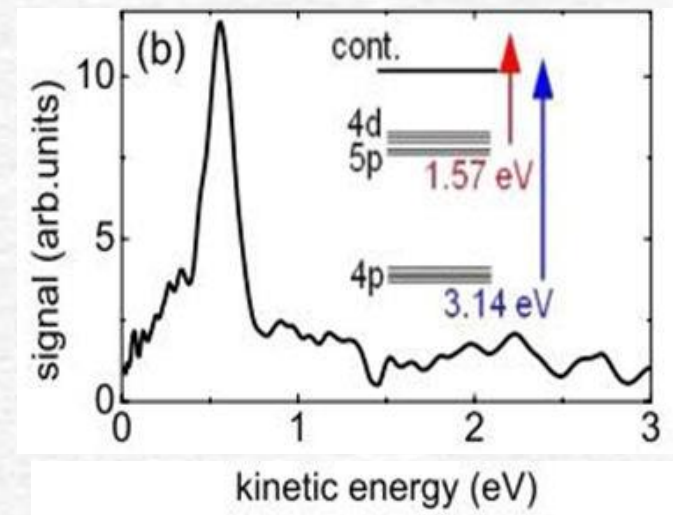
➤ Pump-probe spectroscopy on molecules

- Using photoelectron holography for time-resolved measurements of molecular structure
- XUV-induced molecular dynamics

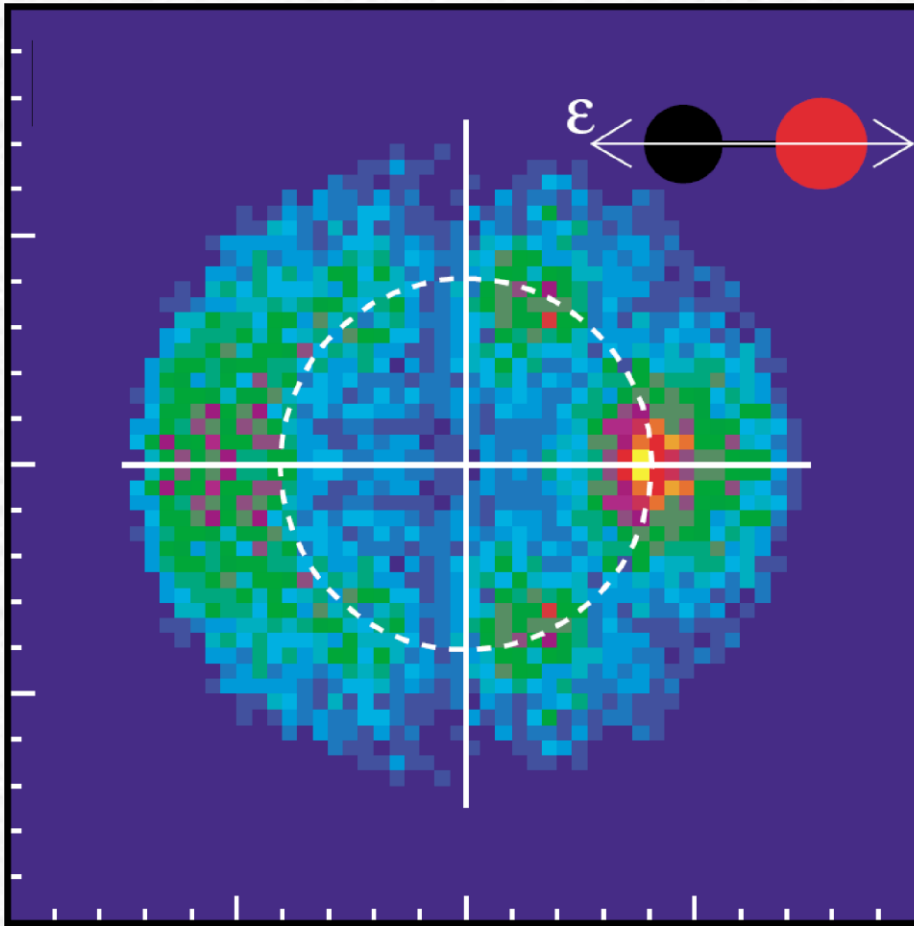


➤ Pump-probe spectroscopy on clusters

- Probing charging and dissociation of highly-excited rare gas clusters



Holography in XUV/X-ray photoionization

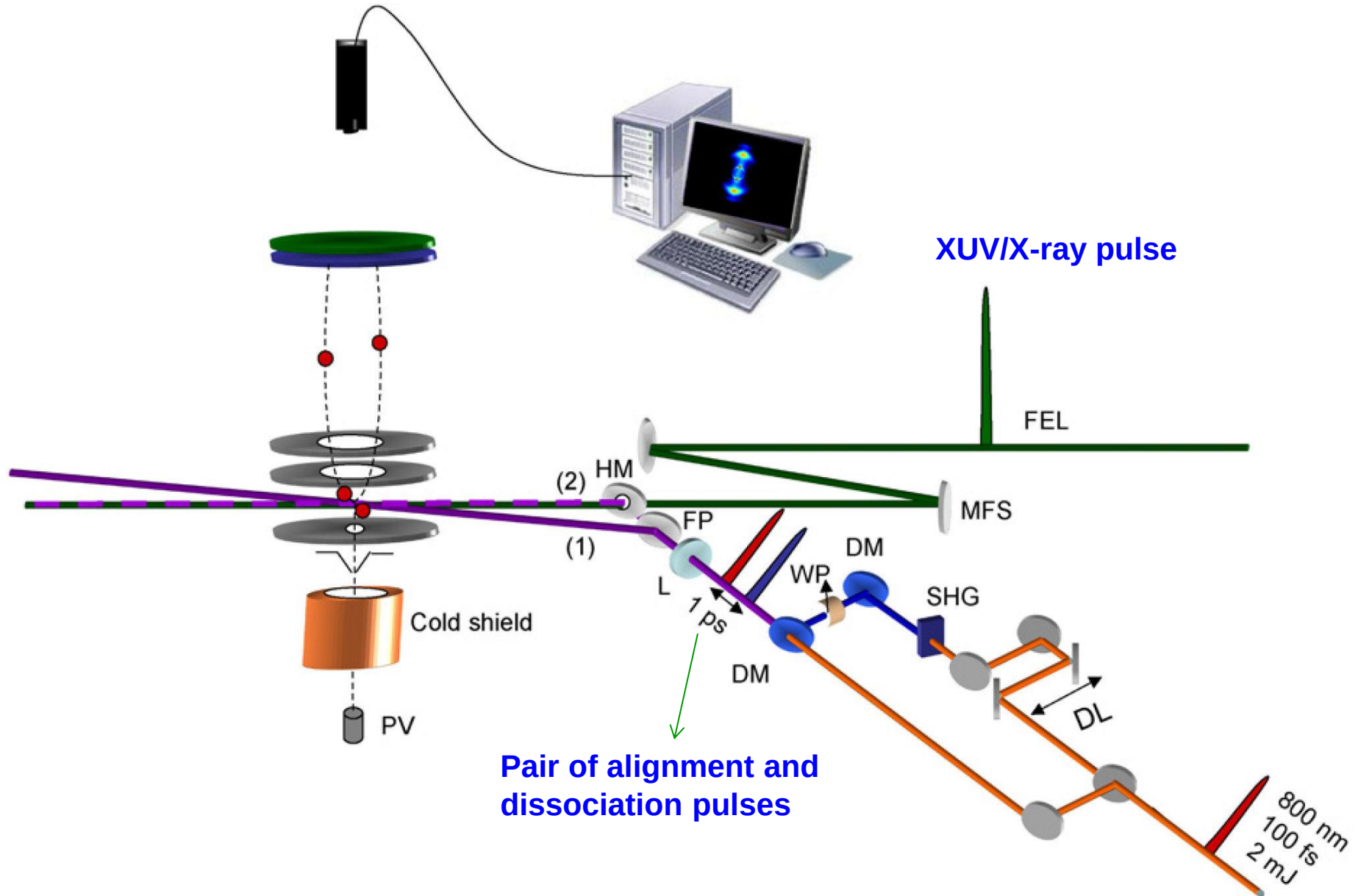


C(1s) core-shell photo-emission from CO using 294 to 326 eV radiation.

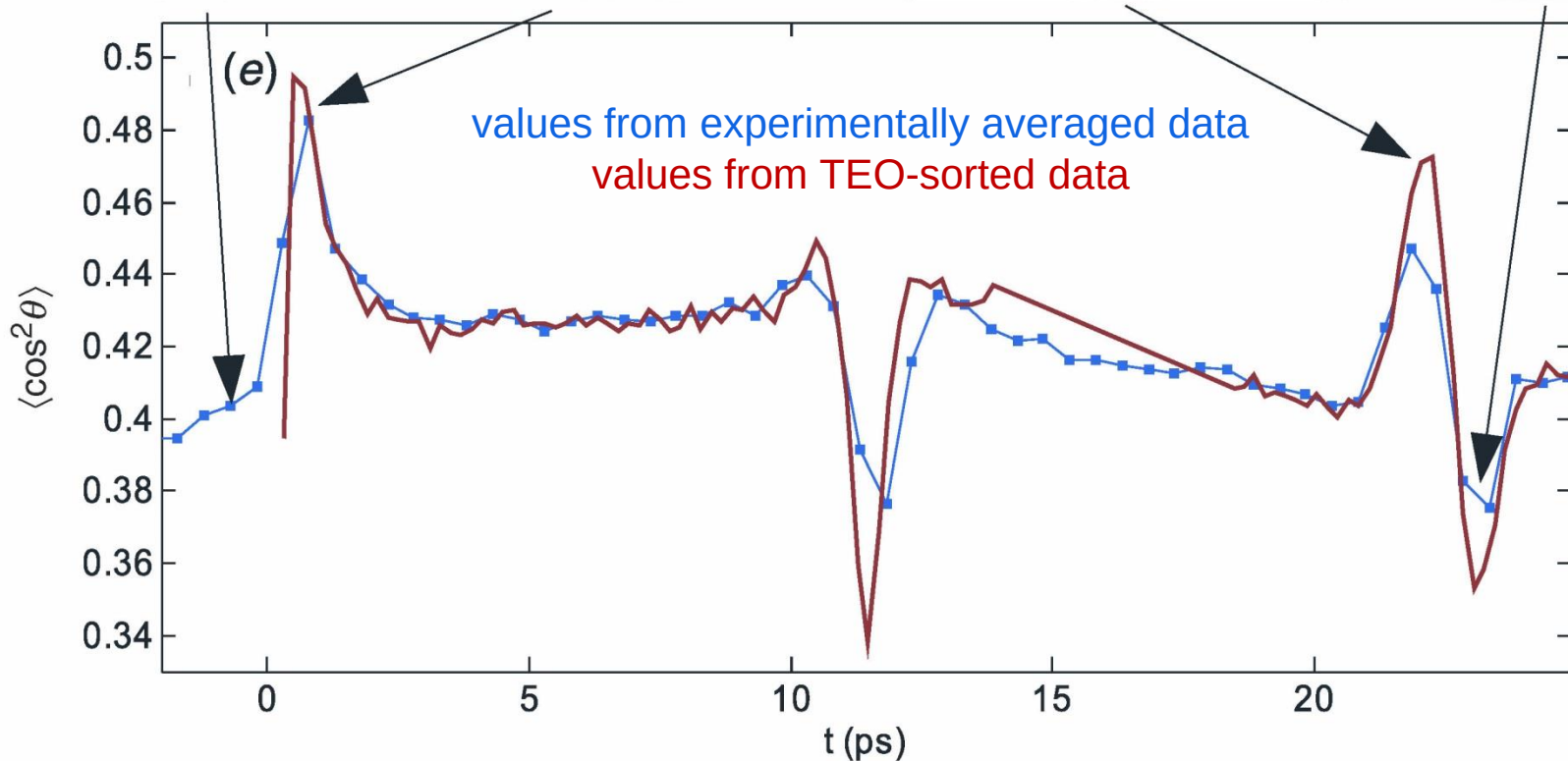
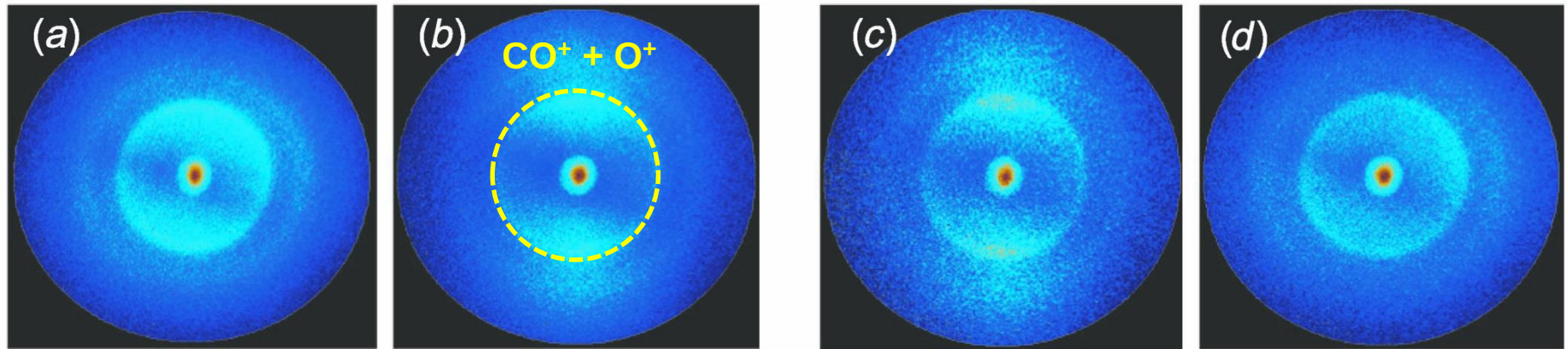
Away from the Carbon atom (black) the angular distribution is relatively unstructured.

In the direction of the Oxygen atom (red) a holographic structure is observed.

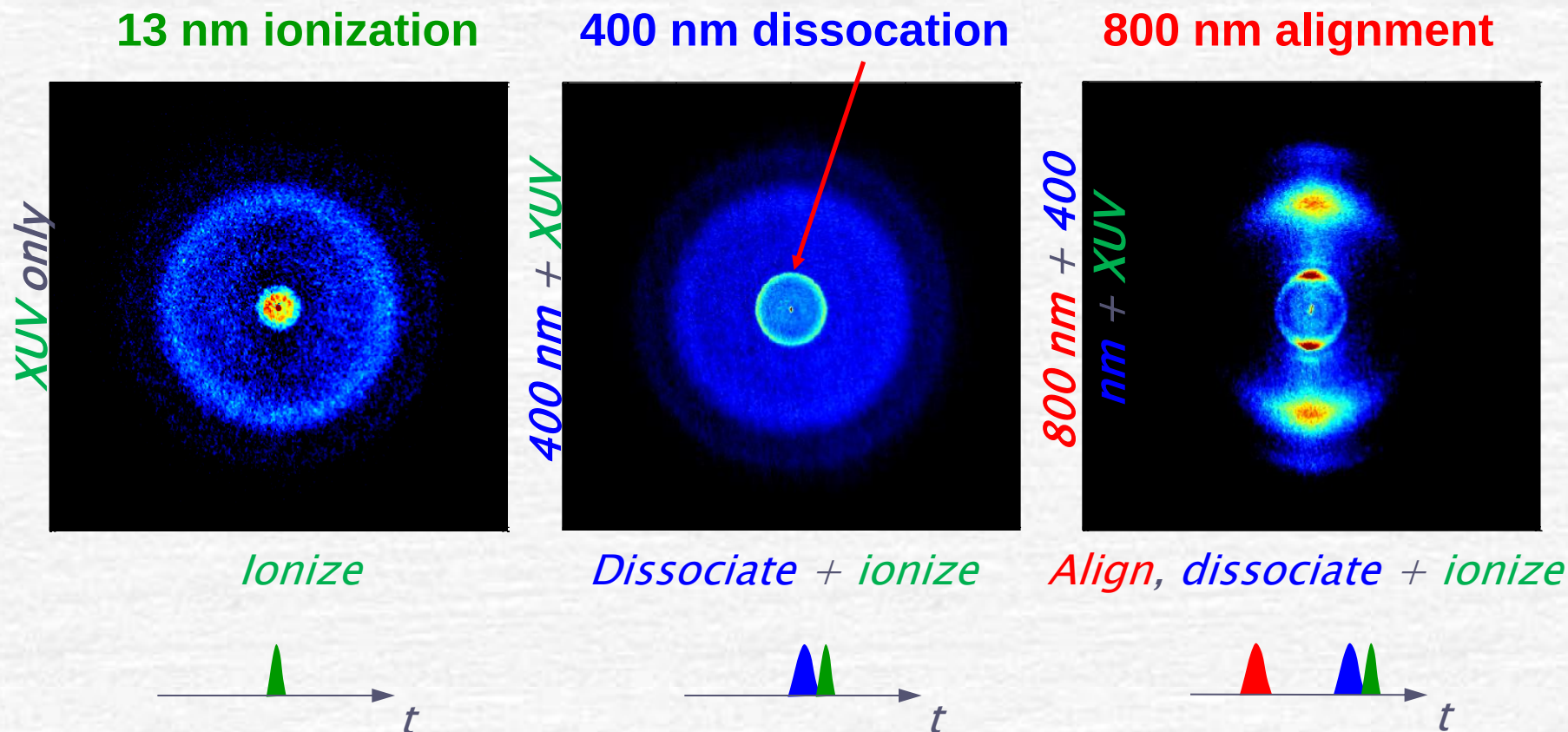
Pump-probe experiments at the FLASH FEL



Dynamic alignment of CO₂ at the FLASH FEL



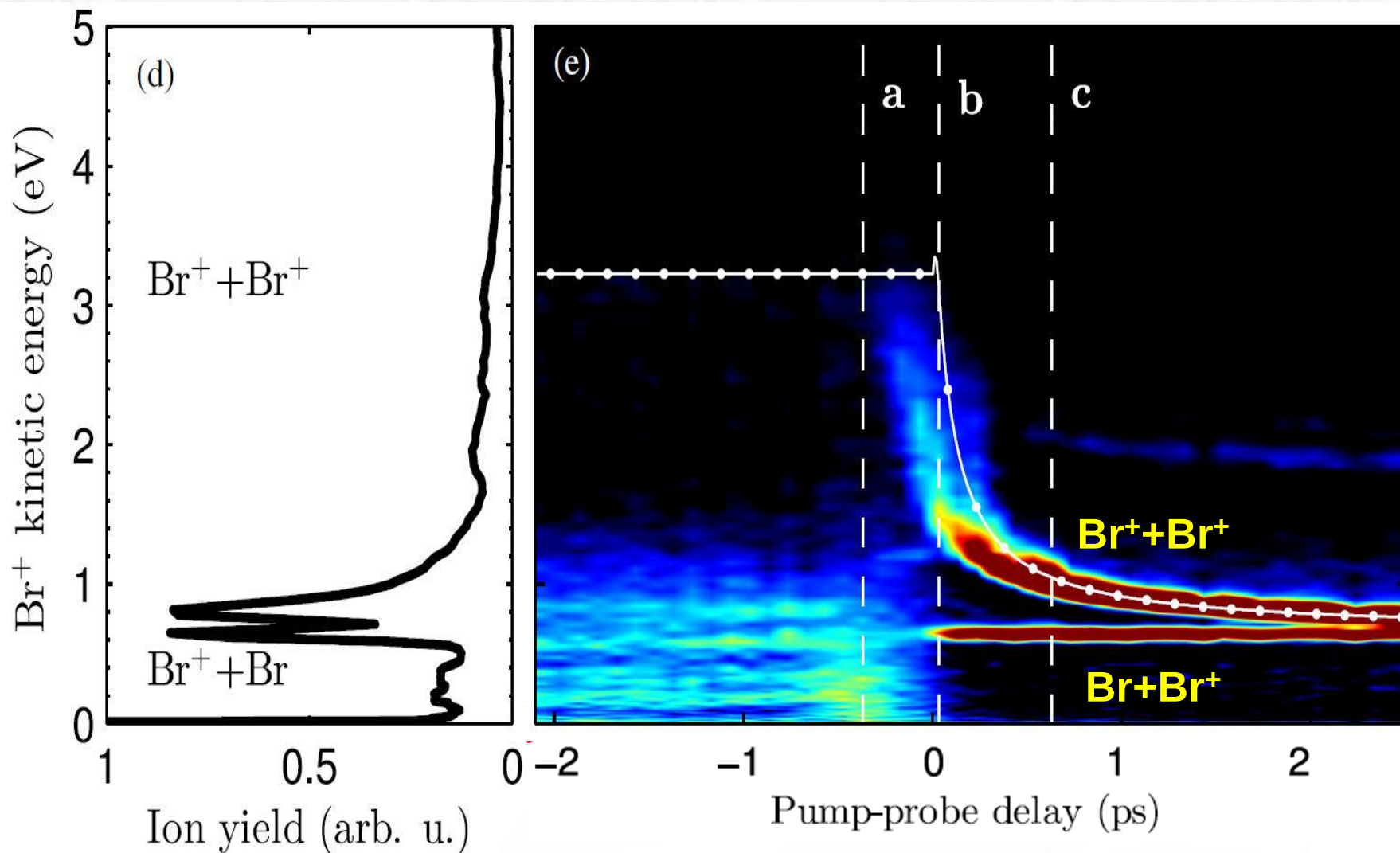
Dissociation of aligned Br_2 at the FLASH FEL



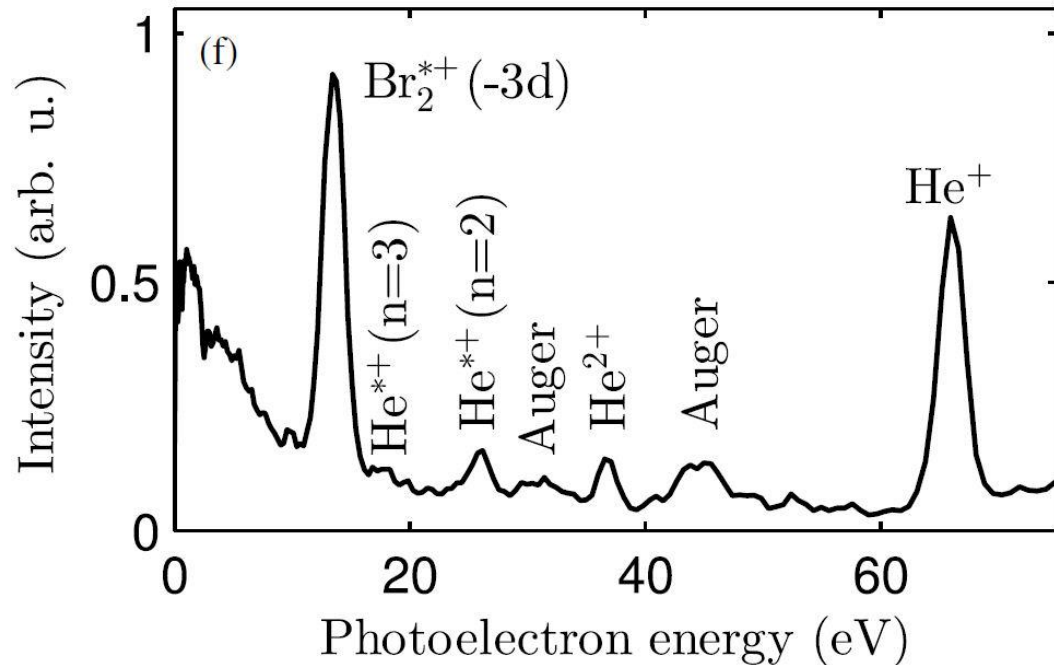
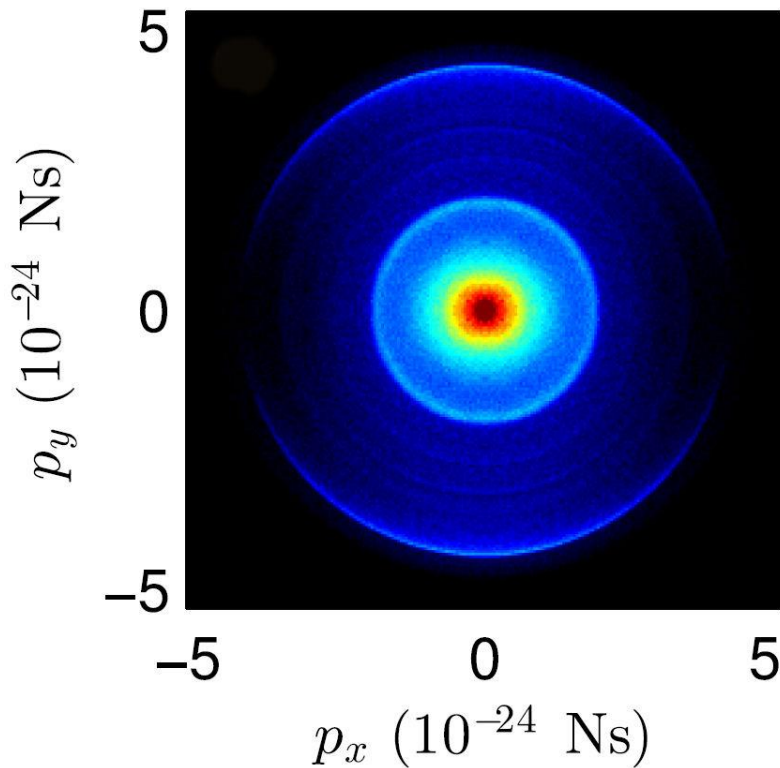
Implementation of protocol for three-color alignment-pump-probe experiments

FEL-IR/UV overlap on the basis of dissociative ionization of H_2

Dissociation of aligned Br_2 at the FLASH FEL



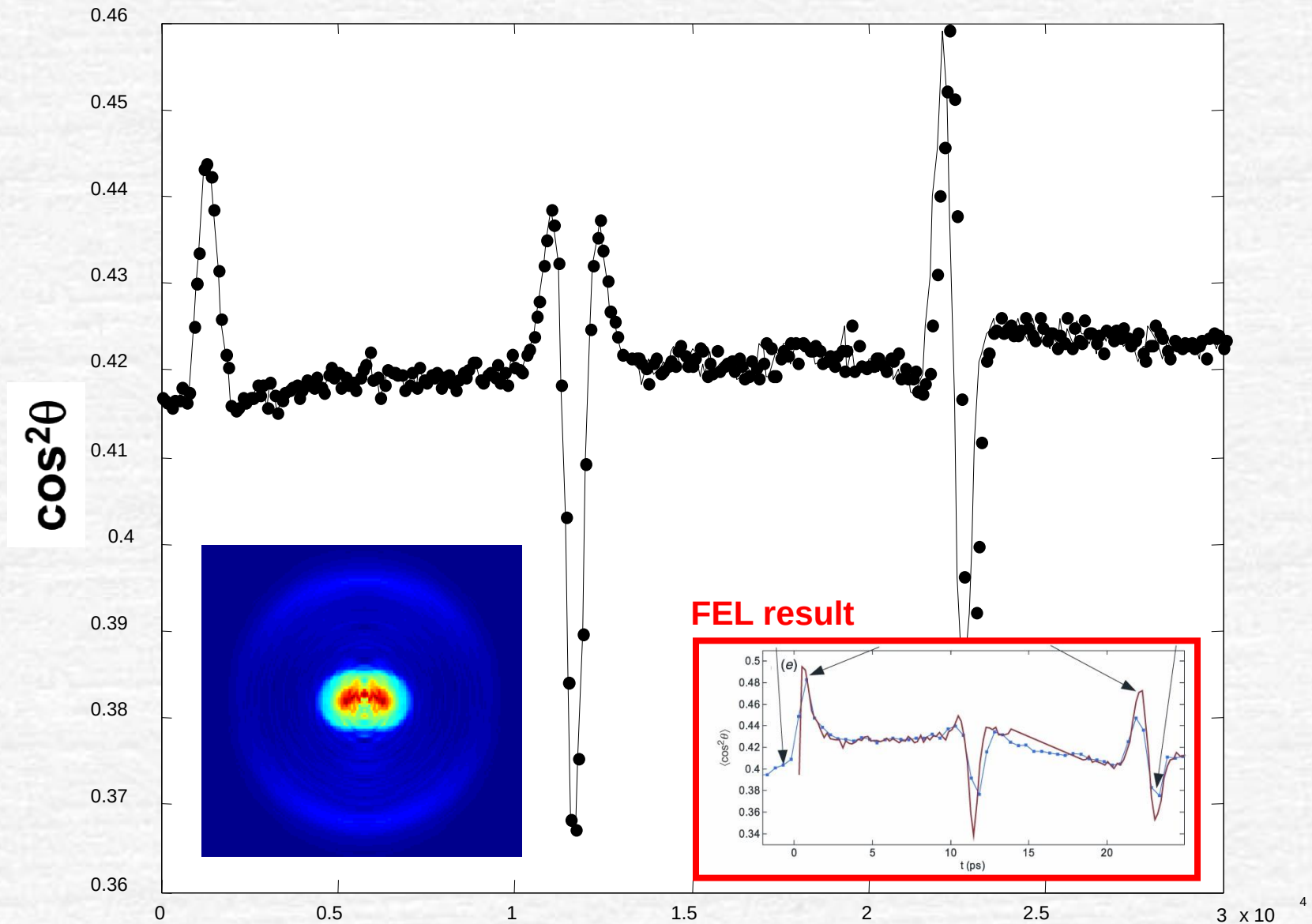
Dissociation of aligned Br₂ at the FLASH FEL



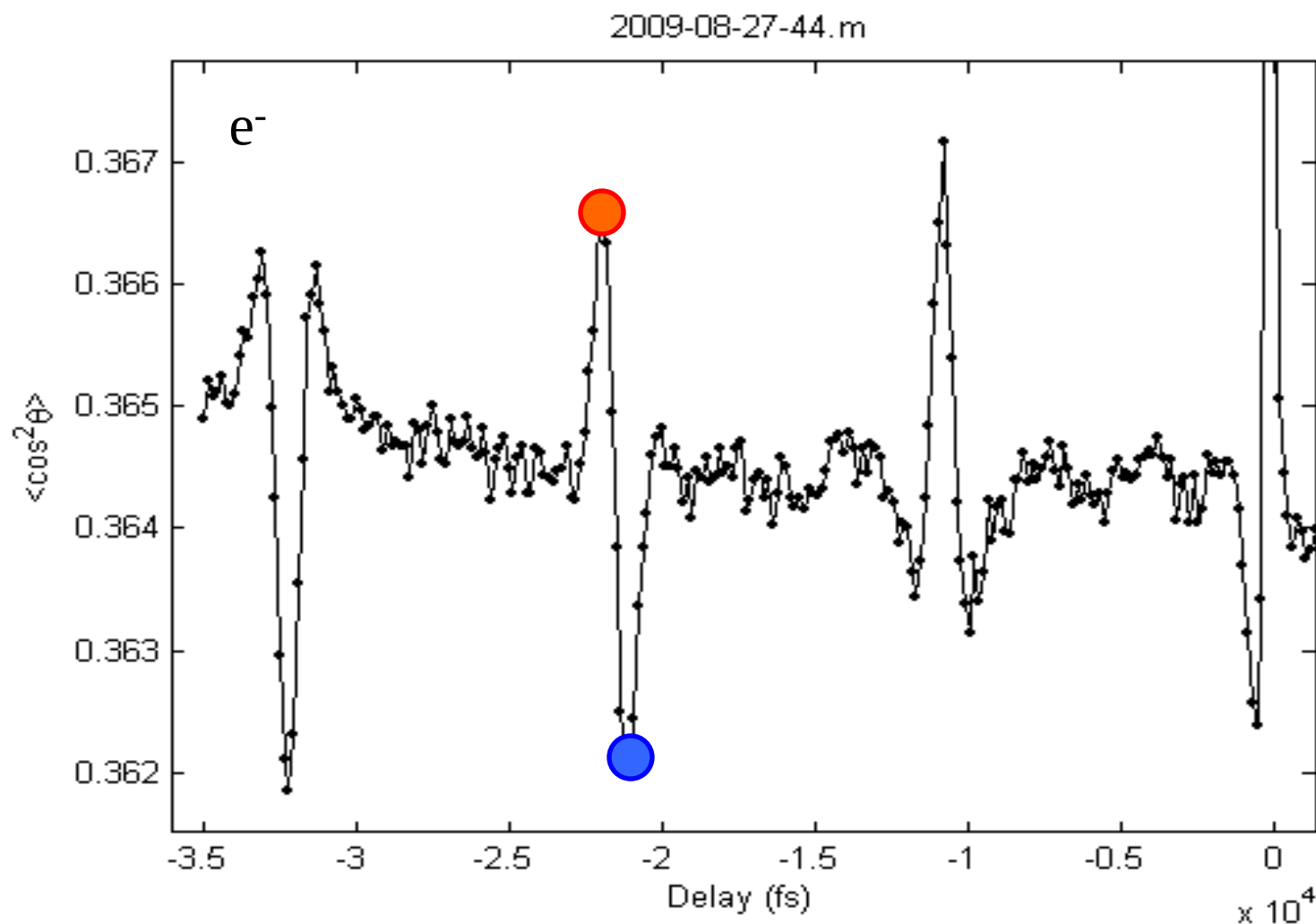
Demonstrated all the essential ingredients required for time-dependent photoelectron holography experiment: alignment, dissociation, photoelectron imaging, three-pulse pump-pump-probe scans

But: no successful photoelectron delay scans yet

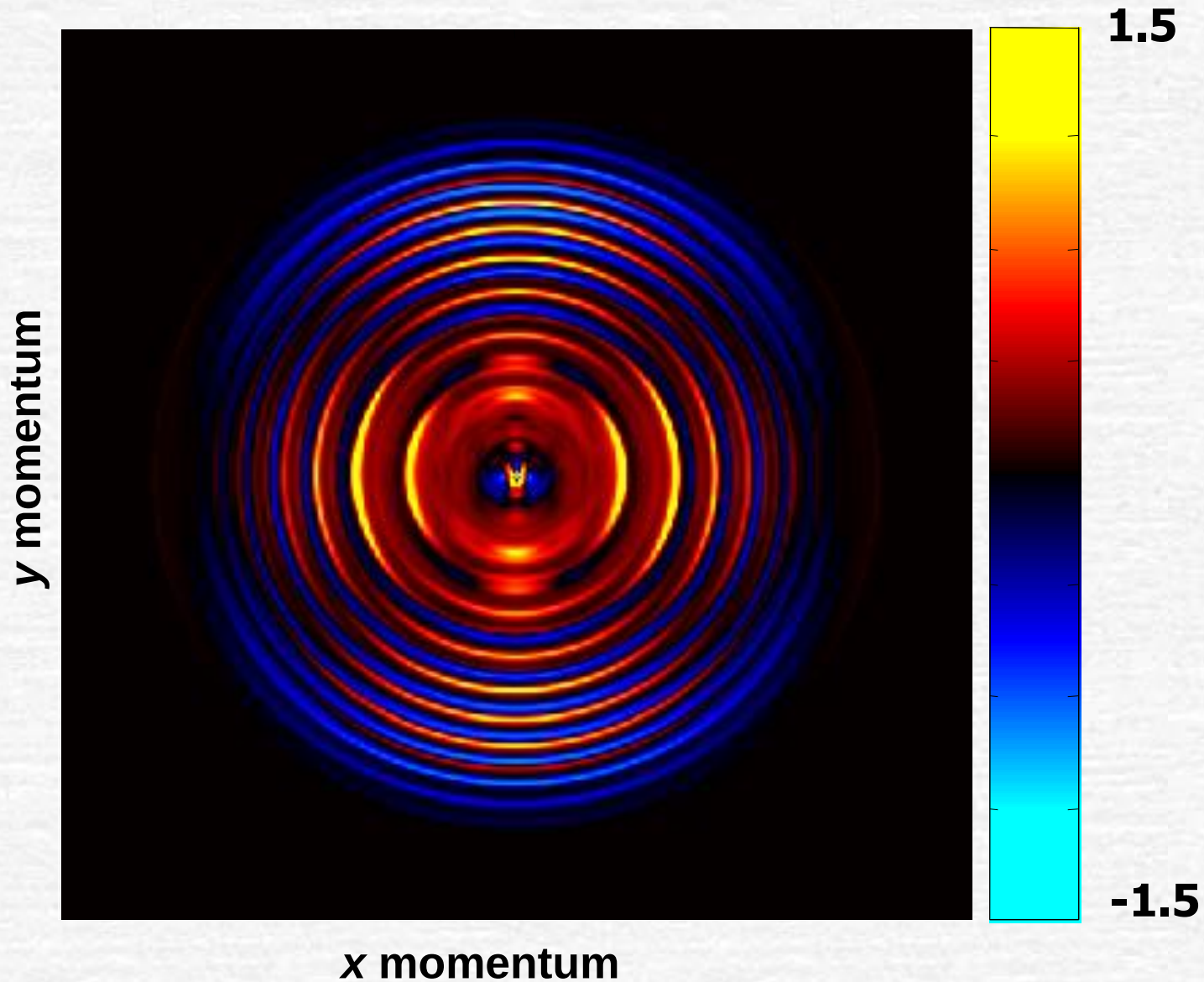
Dynamic alignment of CO₂ probed with HHG source



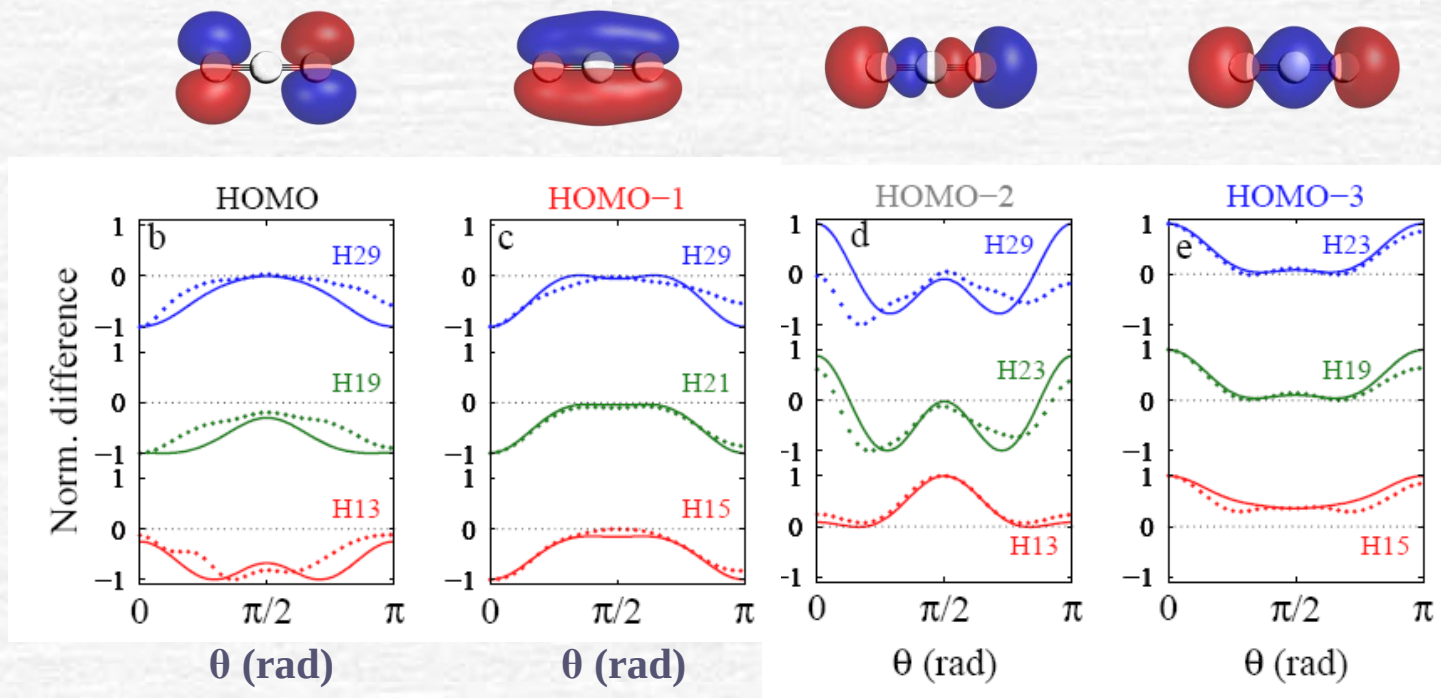
Photoelectron angular distributions from aligned CO_2 molecules



Differential PAD: alignment – anti-alignment

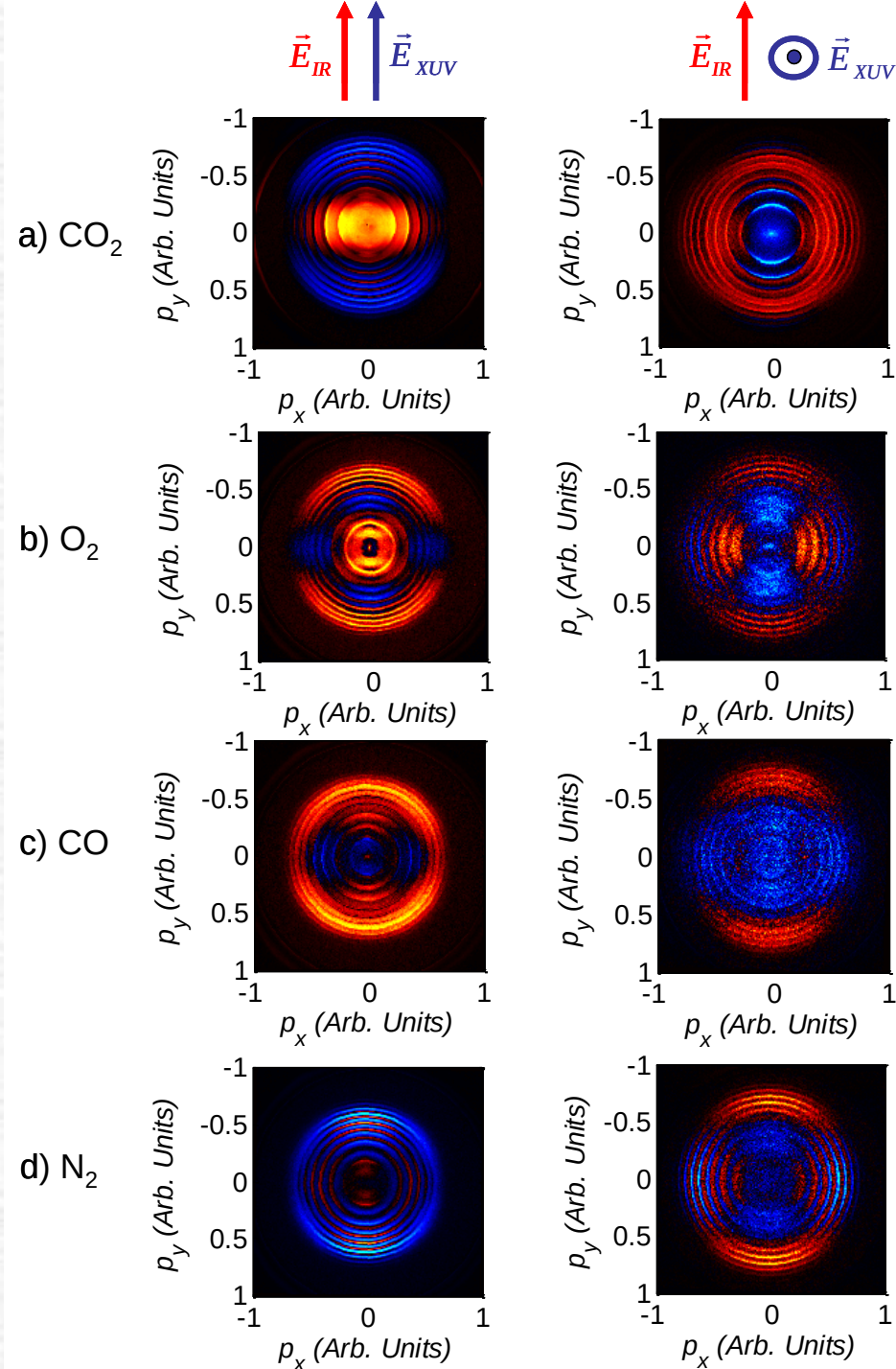


Angular distributions evolve with energy

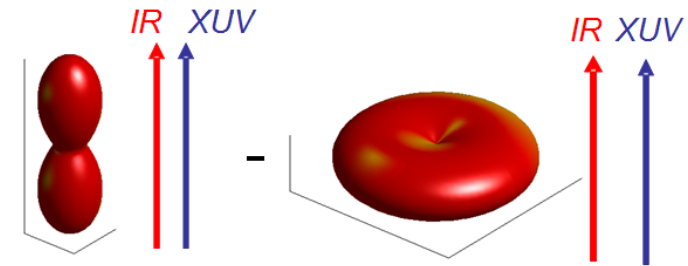


1. Angular distribution sensitive to electronic structure
2. Energy dependence from interaction with molecular Coulomb field $\rightarrow$ molecular structure

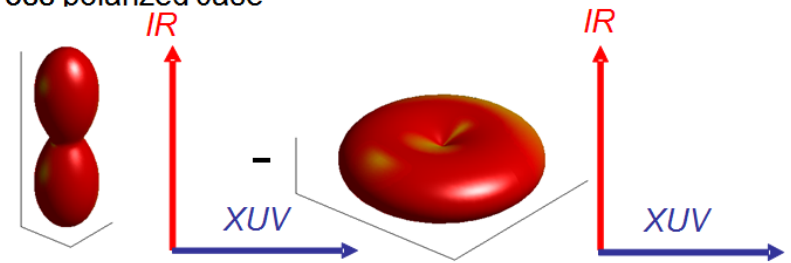
Going from co-polarized to cross-polarized pulses



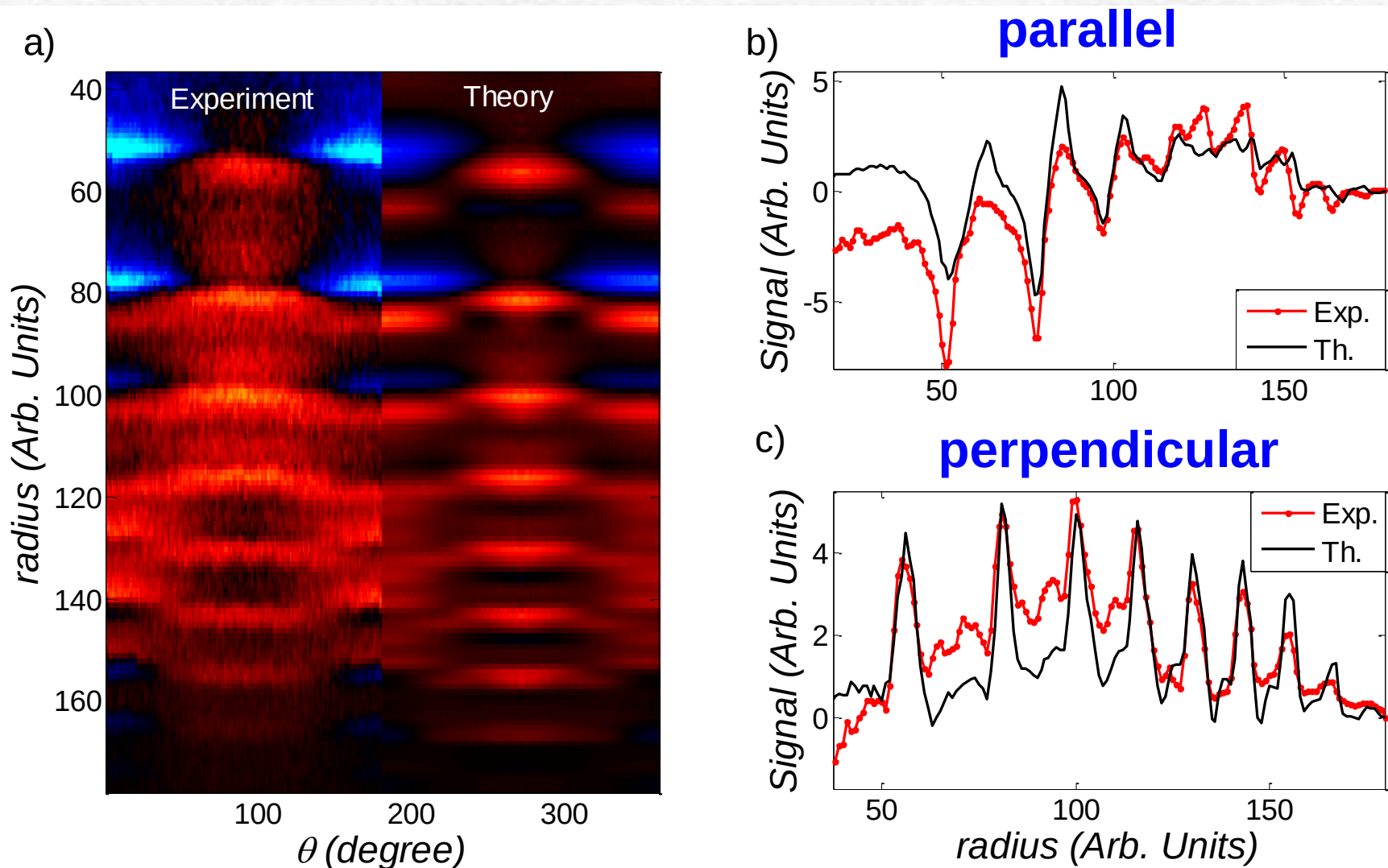
a) co-polarized case



b) cross-polarized case



Comparison with R-matrix calculations for CO₂



Predictions for better alignment

$\langle \cos^2\theta \rangle = 0.71$

$\langle \cos^2\theta \rangle = 0.78$

$\langle \cos^2\theta \rangle = 0.84$

$\langle \cos^2\theta \rangle = 0.91$

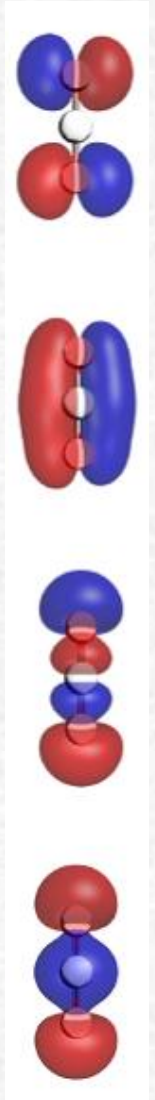
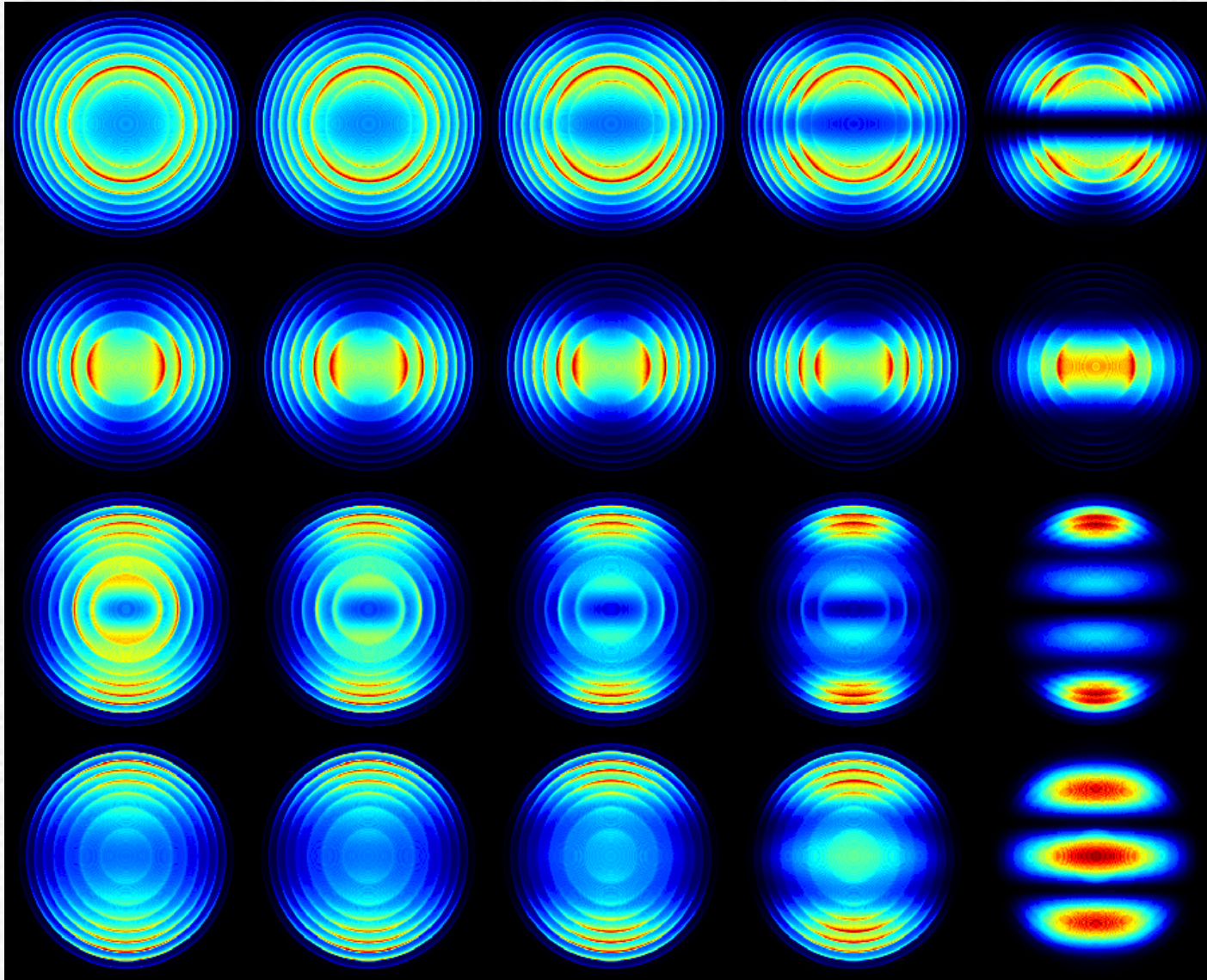
$\langle \cos^2\theta \rangle = 1$

X channel

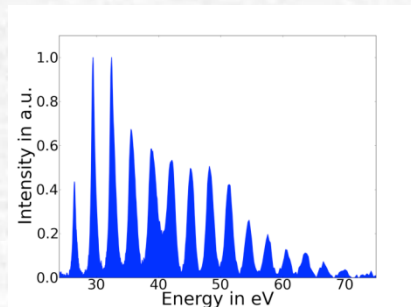
A channel

B channel

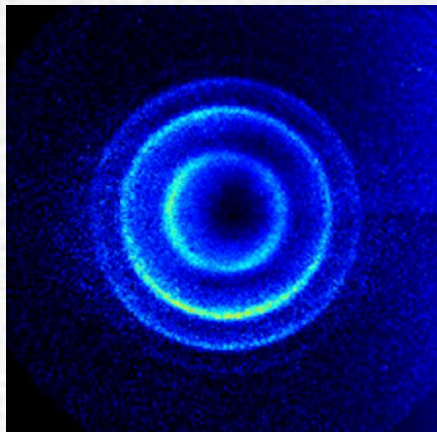
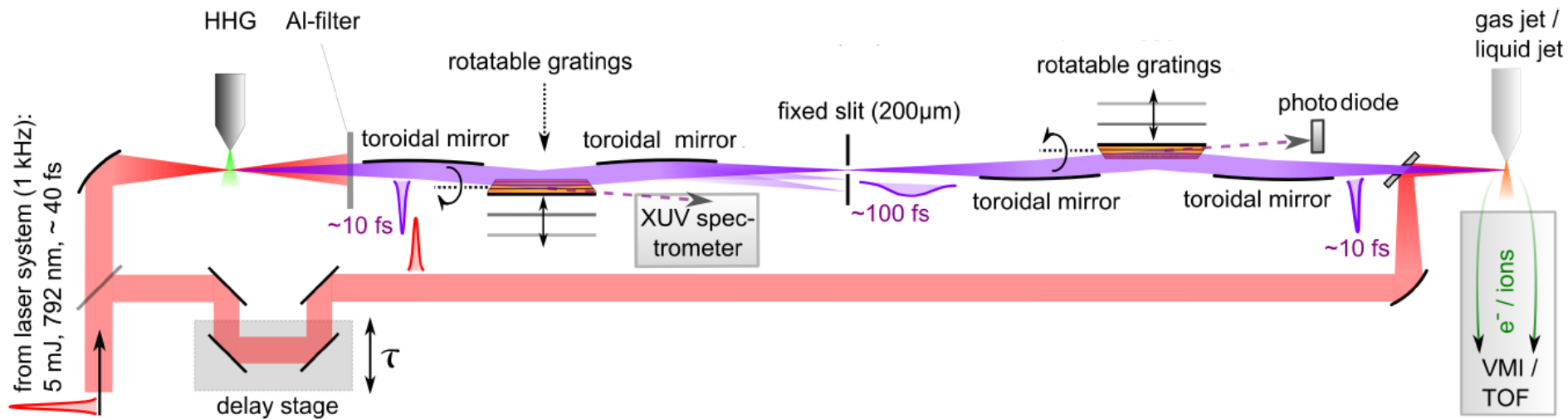
C channel



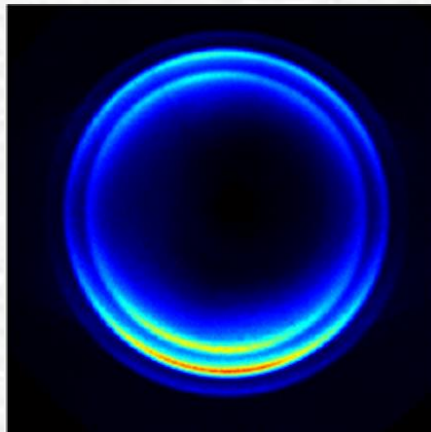
Selecting single harmonics with a time-compensating monochromator



Time-compensating XUV monochromator



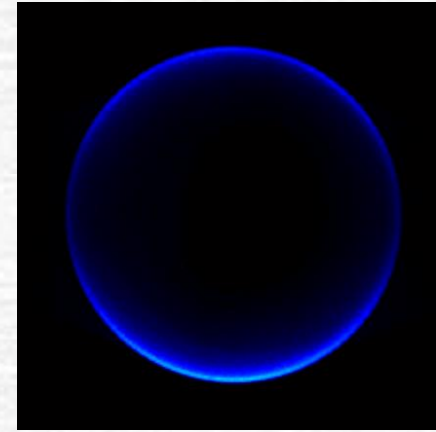
zero-th order



first order

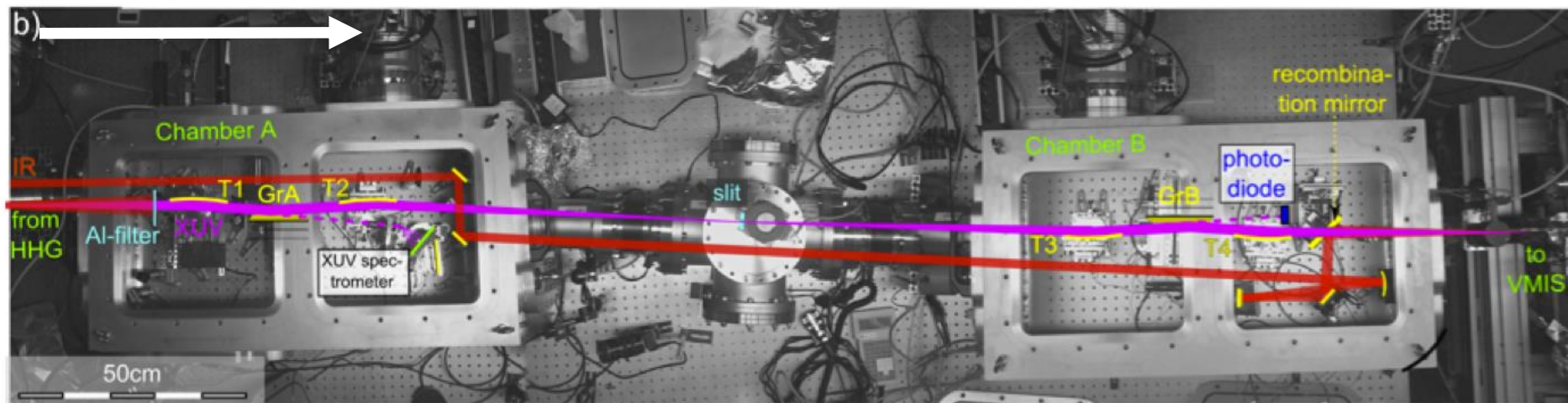


H17



H19

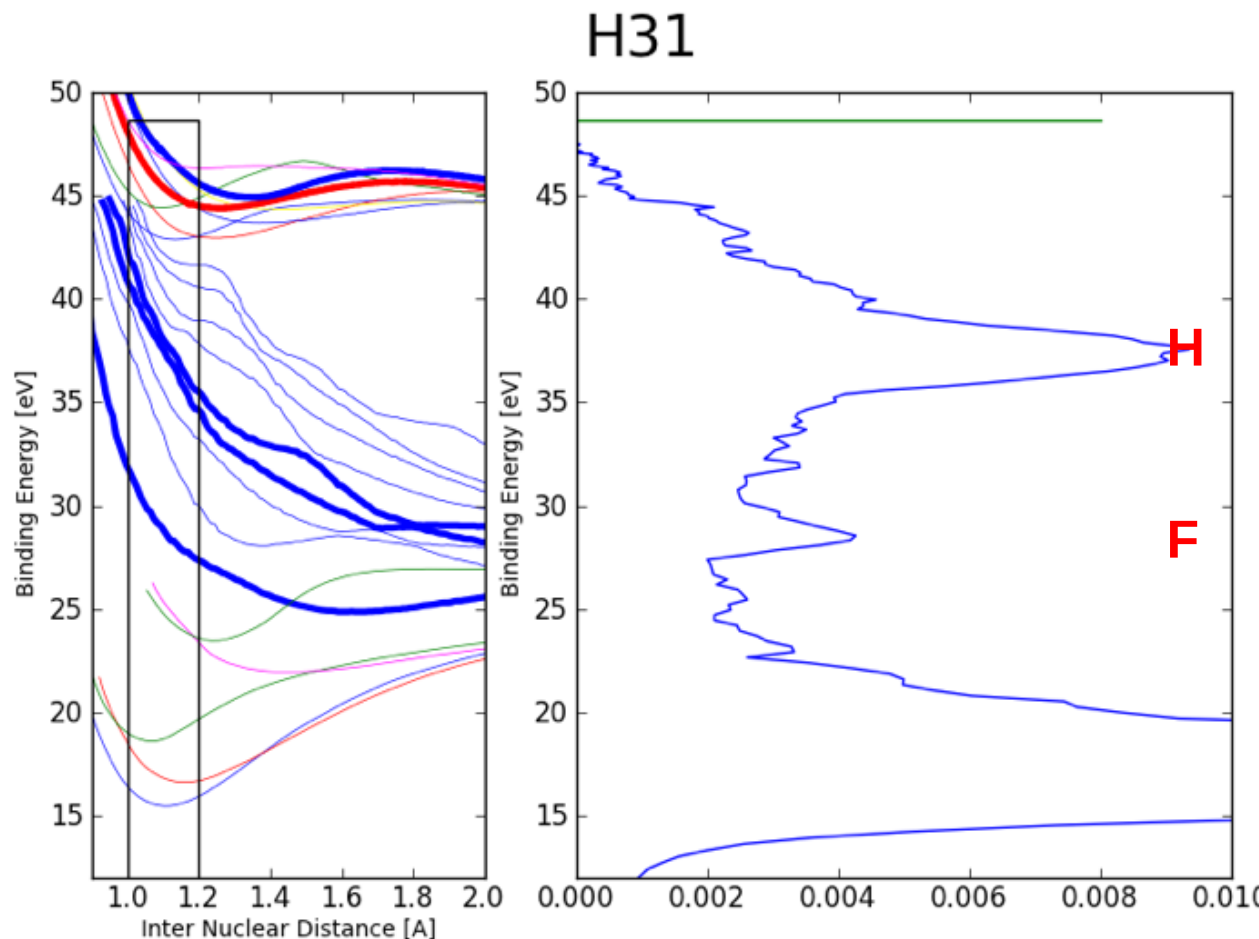
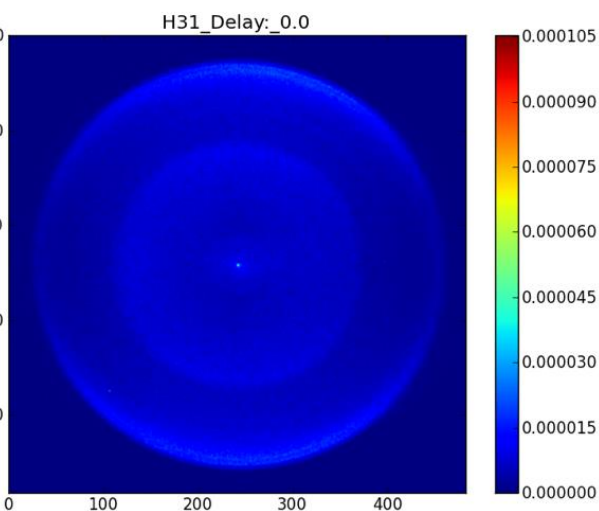
Beamline design and parameters



- Time resolution: <25 fs (cross-correlation between XUV and IR)
- Spectral resolution: <500 meV (deduced from Ar photoelectron spectra)
- Transmission: 3%-16% for 3 eV - 50 eV (calibrated XUV photodiode)
- Continuous automated scans in the complete spectral range
- Long term stability: scan durations >36 hrs
- Up to 10^7 photons per pulse at 1kHz repetition rate

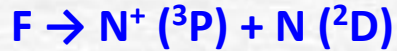
Proof of concept: Photochemistry of N₂

VMI Image



XUV-only photoelectron images revealing the participation of successive ionic states

Extracting dissociation thresholds for the F- and H-state



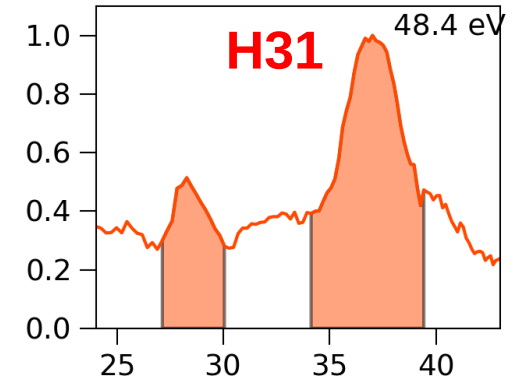
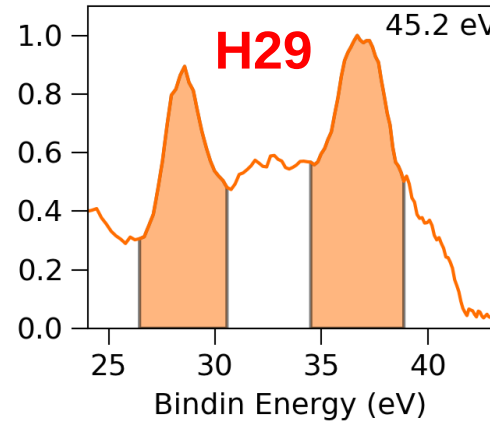
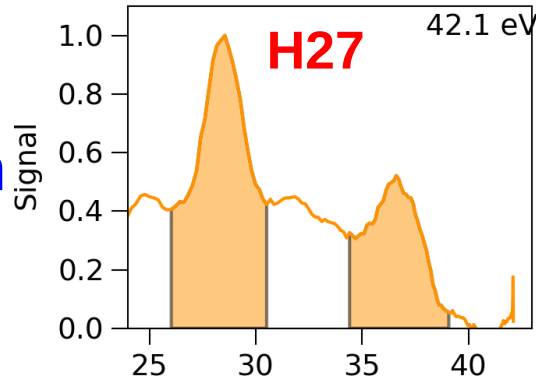
$E = 26.676 \text{ eV (L3)}$



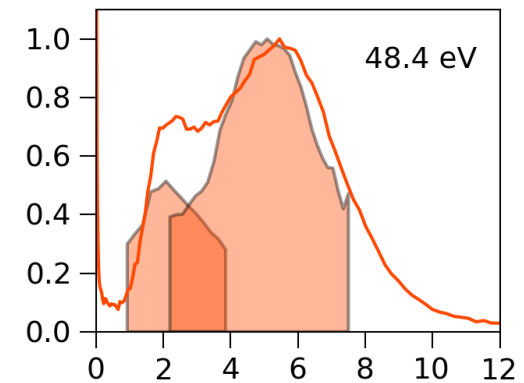
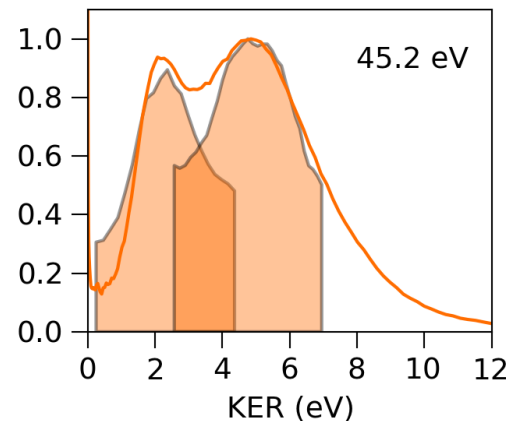
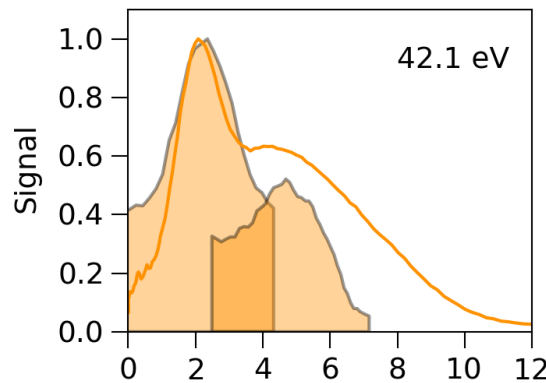
$E = 31.921 \text{ eV (L9)}$

X conventional wisdom!

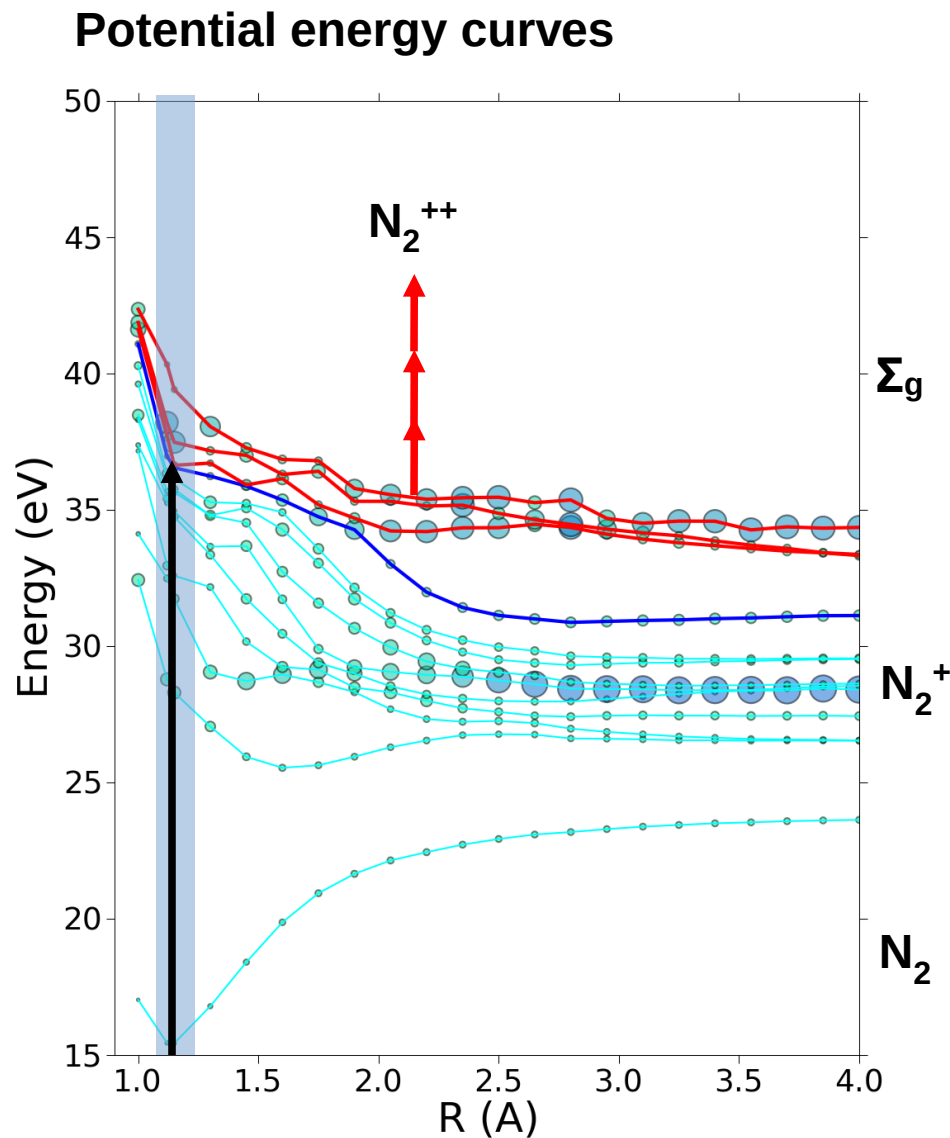
Electron



Ion



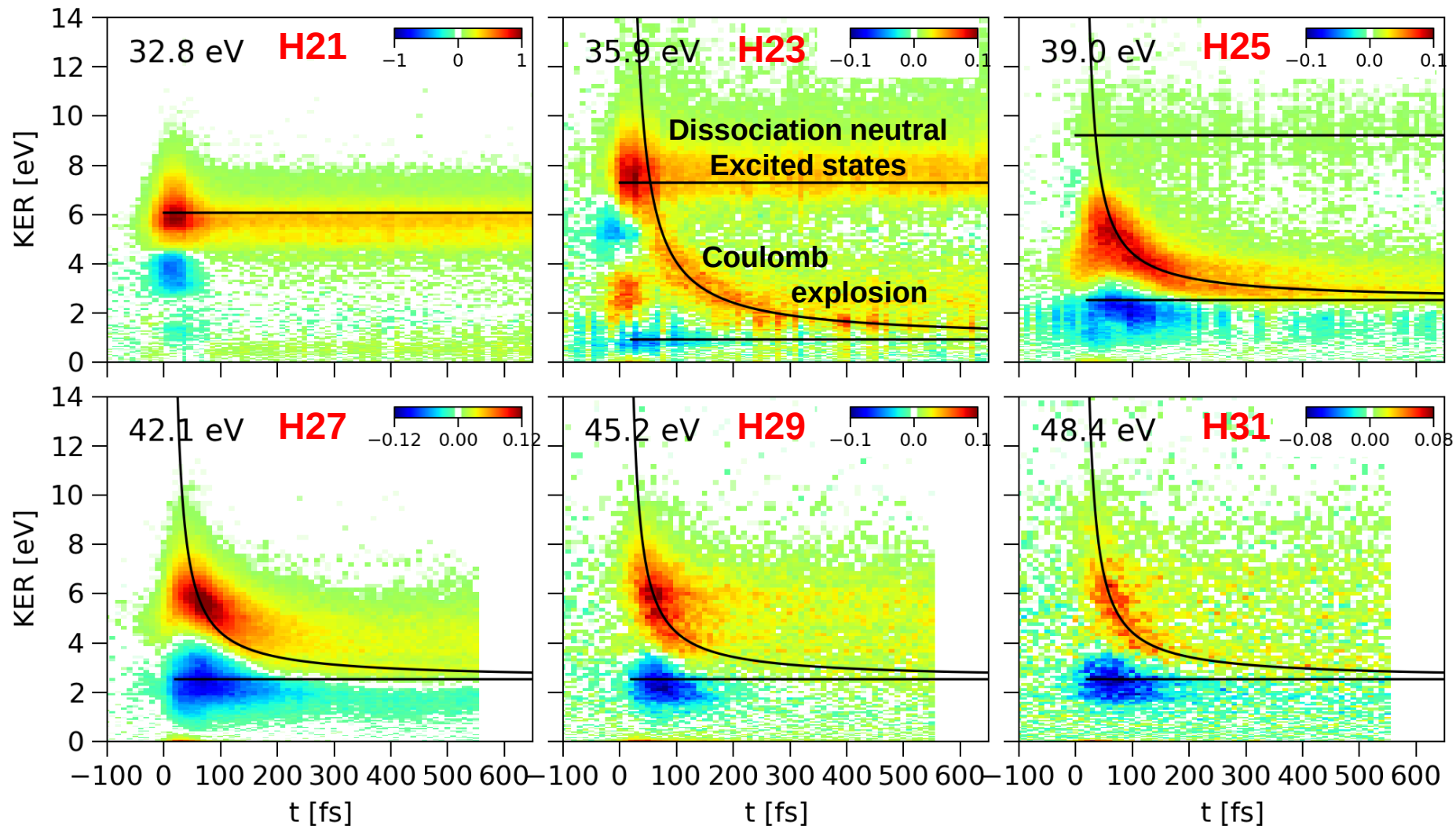
Time-dependent dynamics: XUV + IR



Removal of an electron from the $2\sigma_g$ orbital excites the N_2^+ ion in the H-band and leads to dissociation, dominantly to the L9 dissociation threshold, and with smaller contributions from higher-lying states.

The contributions from higher-lying states can be further ionized using 3-photon excitation by the IR laser, leading to Coulomb explosion of the molecule

Time-dependent dynamics: XUV + IR Differential Ion KER distributions

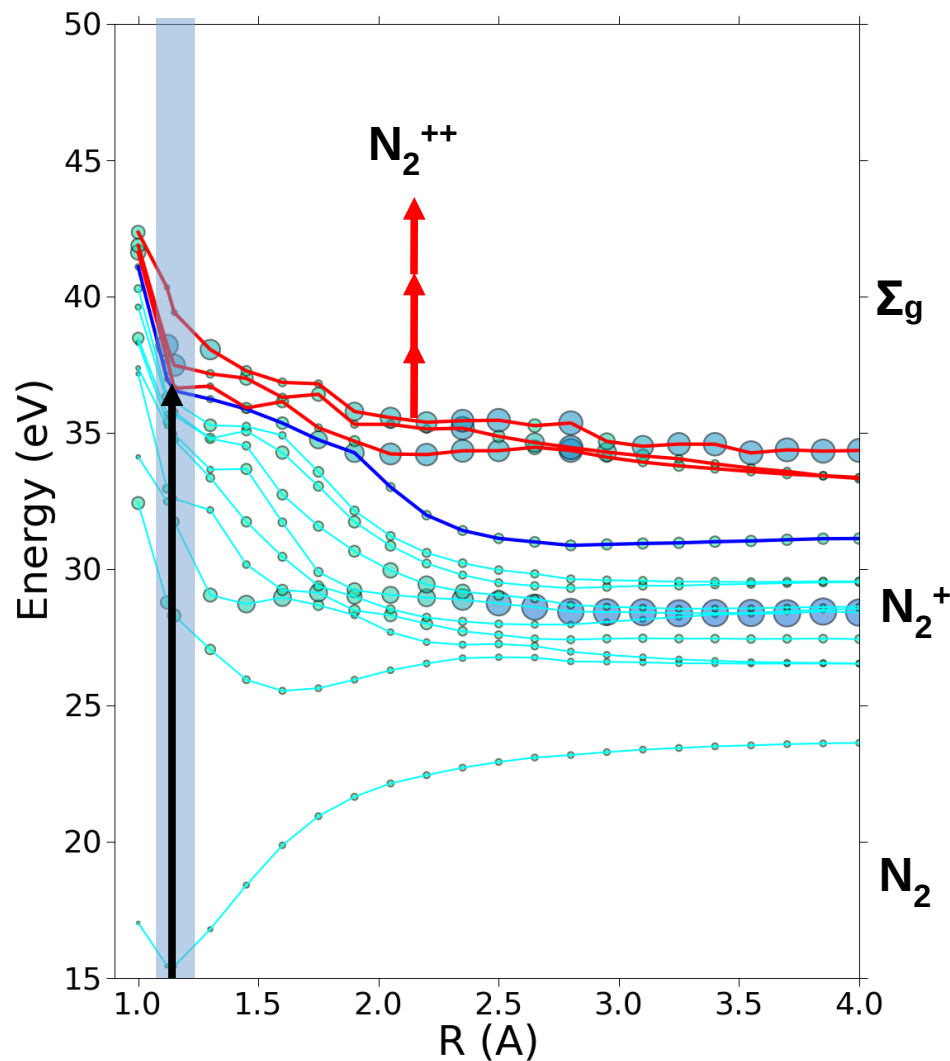


Enhancement

Depletion

Time-dependent dynamics: XUV + IR

Potential energy curves

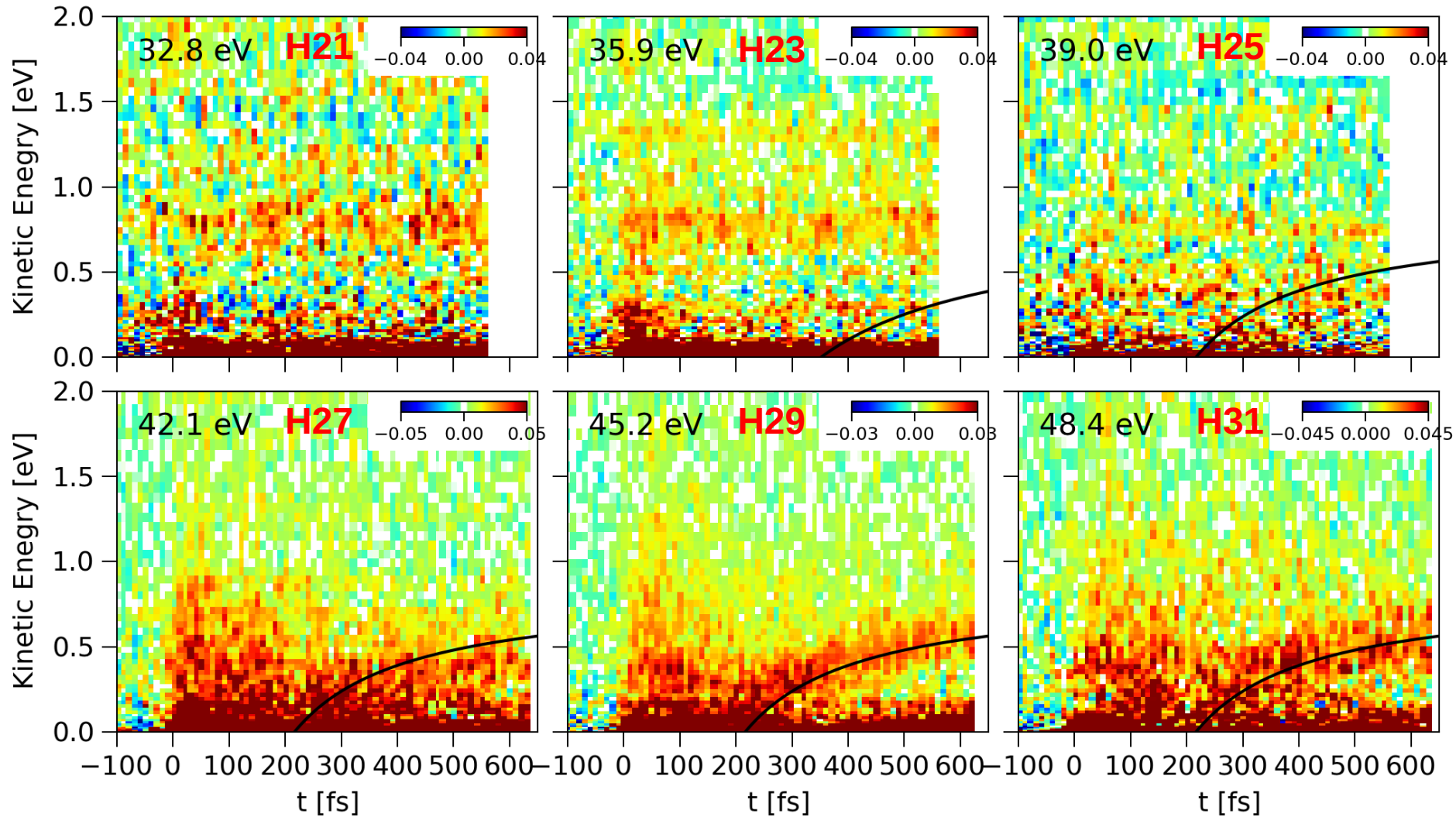


Removal of an electron from the $2\sigma_g$ orbital excites the N_2^+ ion in the H-band and leads to dissociation, dominantly to the L9 dissociation threshold, and with smaller contributions from higher-lying states.

The contributions from higher-lying states can be further ionized using 3-photon excitation by the IR laser, leading to Coulomb explosion of the molecule

The Coulomb explosion gives a characteristic contribution in the photoelectron spectrum

Time-dependent dynamics: XUV + IR Differential Electron KE distributions



Enhancement

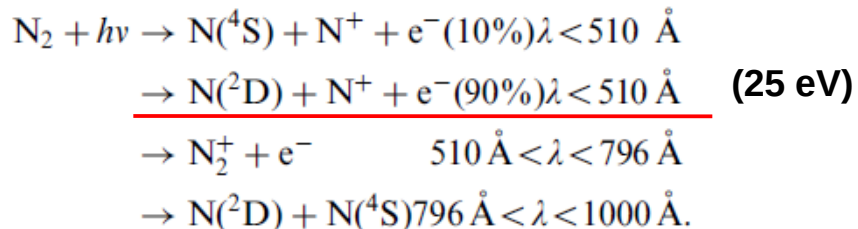
Depletion

Conclusions and Implications

- Commissioning XUV time-compensating monochromator beamline
- Major H-band dissociation channel → L9: $\text{N}^+ (^1\text{S}) + \text{N} (^2\text{P})$
- Main pump-probe response comes from secondary channel → L11
- Strong impact on atmospheric models: ^2P atoms are 2 orders of magnitude less reactive than ^2D atoms!!!

2.3.1. Photodissociation

The complex chemistry of Titan's atmosphere is initiated with the photodissociation/photoionization of molecular nitrogen and methane (in the current calculations the contributions of energetic particles are not included and ionization is considered only for N_2). Photons in the EUV region of the solar radiation spectrum, depending on their energy, are able to ionize and dissociate N_2 according to the following pathways (Banks and Kockarts, 1973; Nicolas et al., 2003):



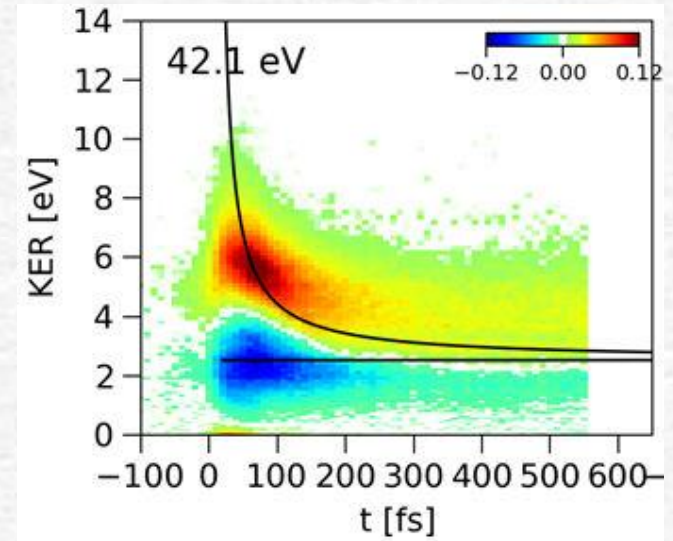
Huygens-Cassini mission to Saturn



Contents of this talk

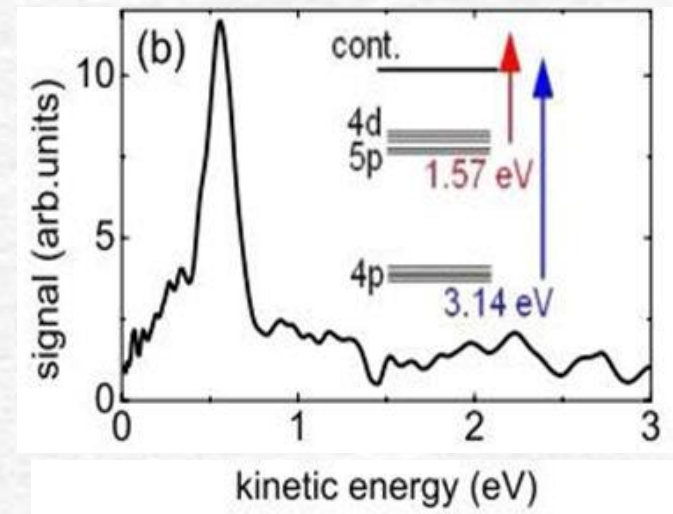
➤ Pump-probe spectroscopy on molecules

- Using photoelectron holography for time-resolved measurements of molecular structure
- XUV-induced molecular dynamics

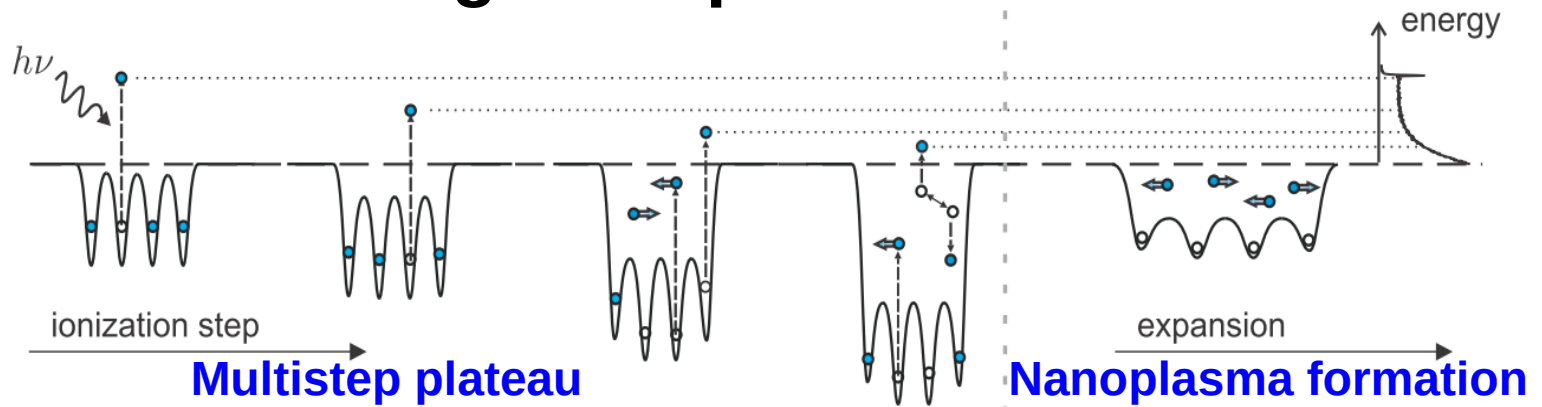


➤ Pump-probe spectroscopy on clusters

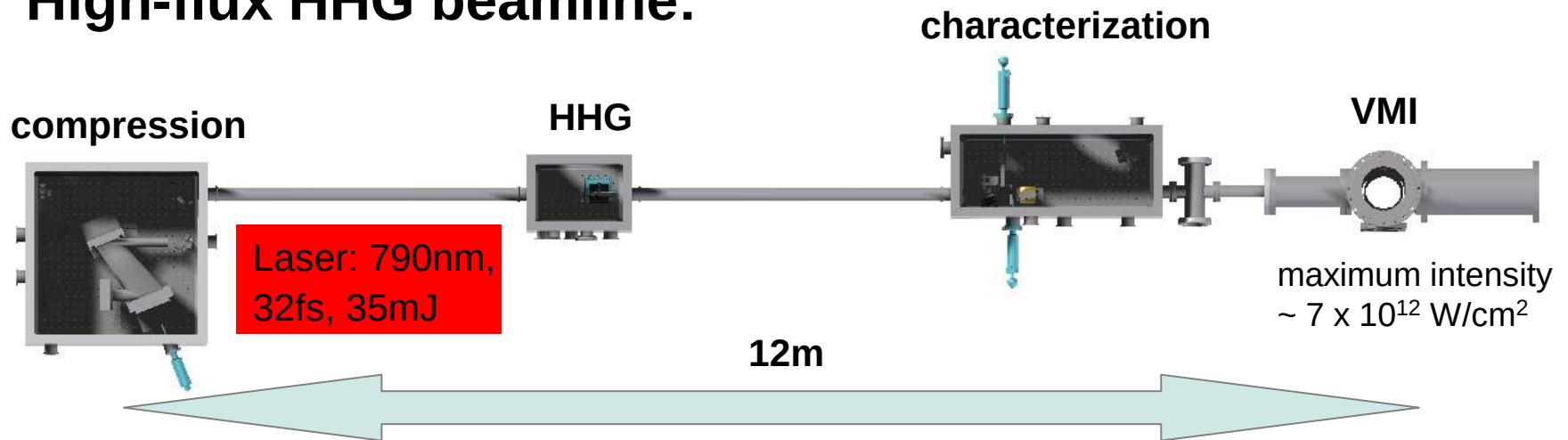
- Probing charging and dissociation of highly-excited rare gas clusters



Understanding from previous FEL research:

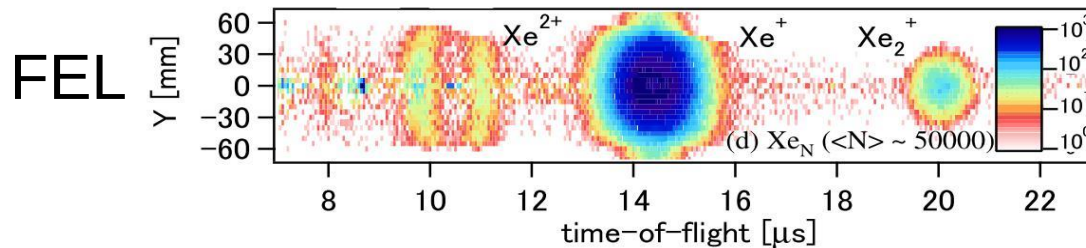


High-flux HHG beamline:



HHG very promising for improved understanding of cluster dynamics

Clusters in HHG / FEL pulses

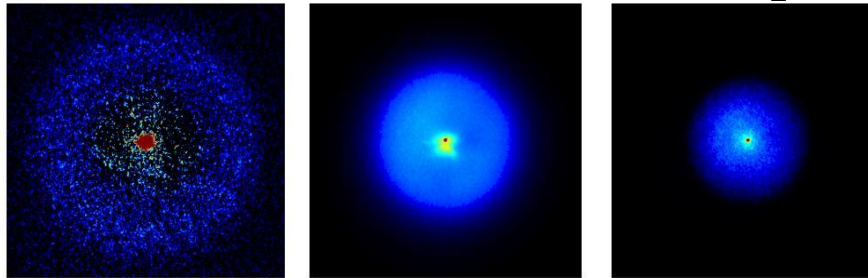


H. Iwayama *et al.*, J. Phys. B **43**, 161001 (2010).

Xe^{2+}

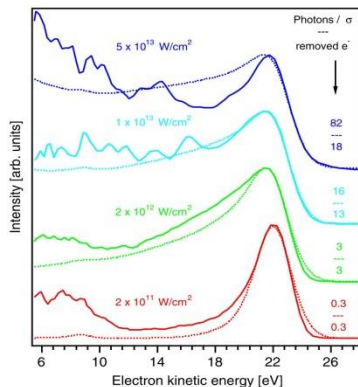
Xe^+

Xe_2^+

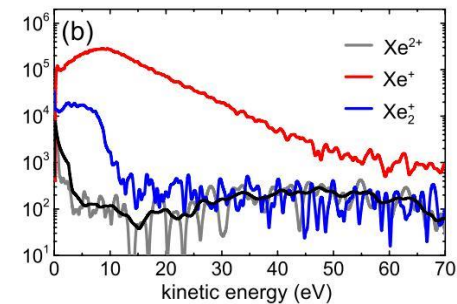
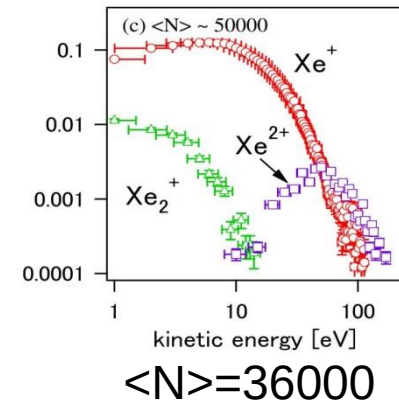


B. Schütte *et al.*, PRL **112**, 073003 (2014).

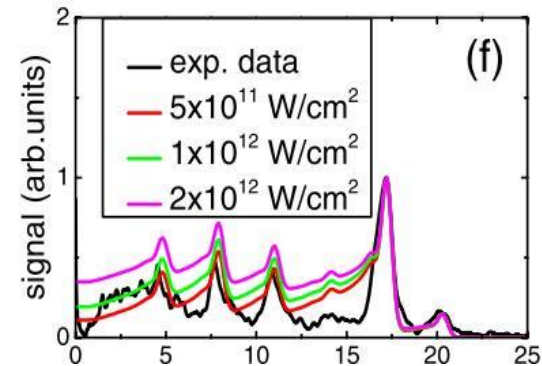
FEL



C. Bostedt *et al.*, PRL **100**, 133401 (2008).



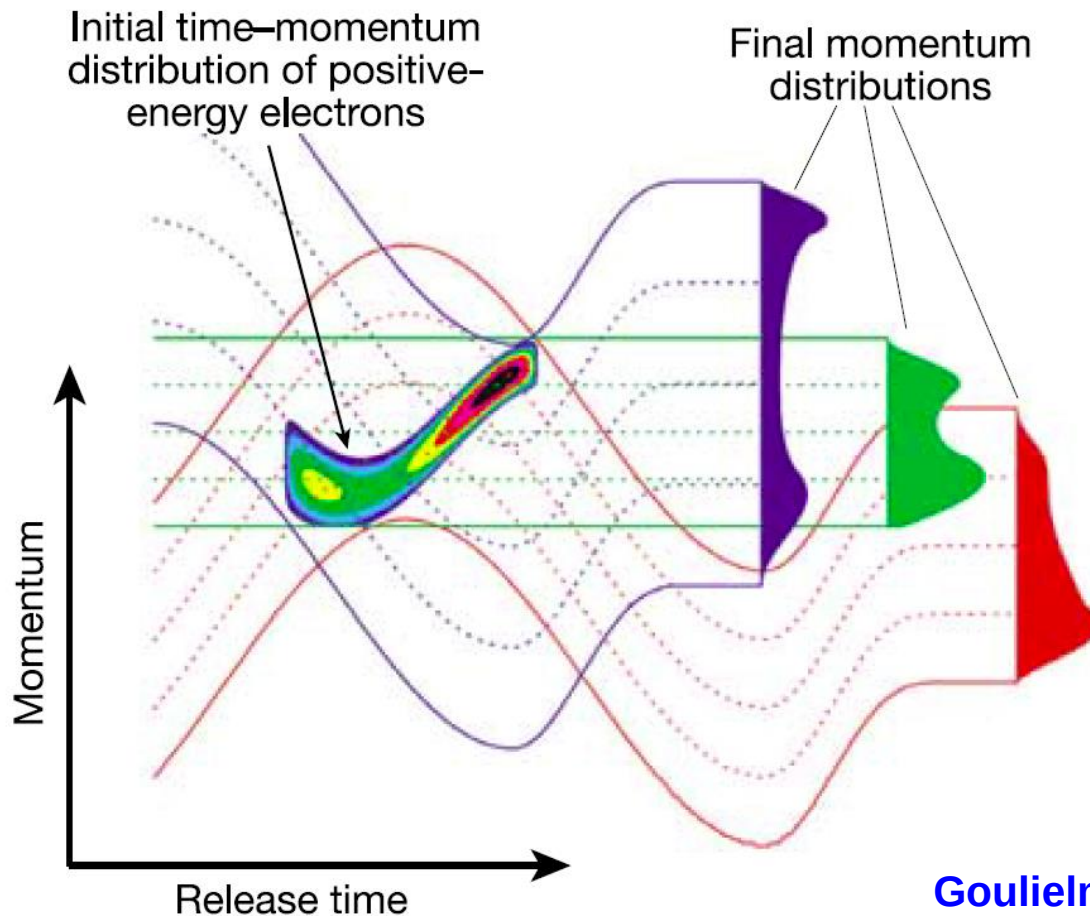
HHG



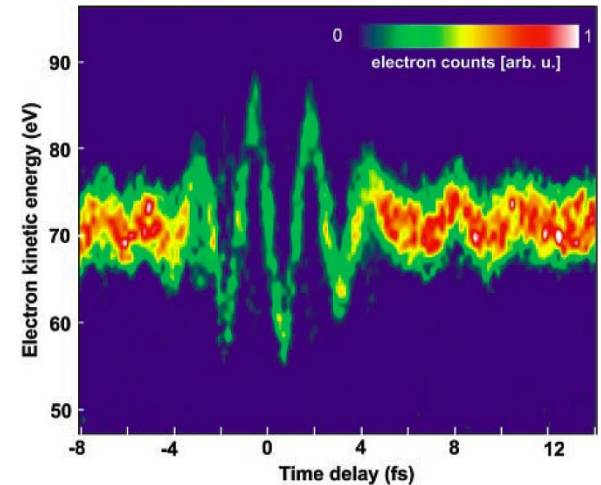
B. Schütte *et al.*, PRL **112**, 073003 (2014).

FEL results can be qualitatively reproduced!

Look at charging dynamics using THz streaking



Blue and red:
Lines of constant canonical momentum $p=k(t)+qA(t)$ give different tomographic projections



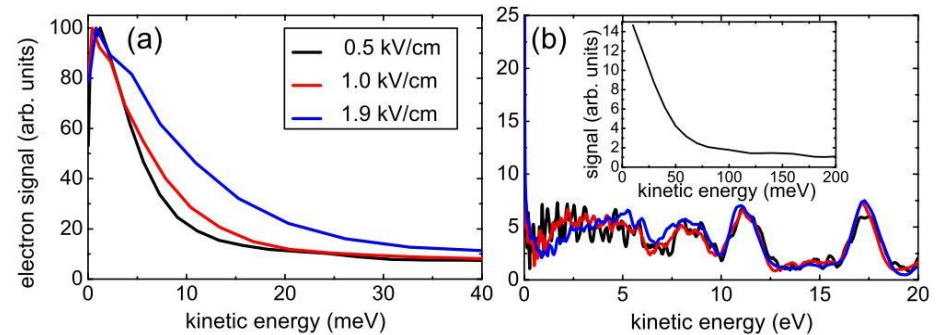
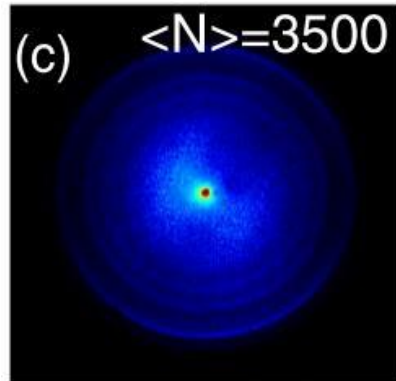
Goulielmakis et al., Science 305, 1267 (2004)

Lesson from attosecond science: time-dependent photoionization can be revealed by a streaking measurement

Kienberger et. al. Nature 427, 817 (2004)

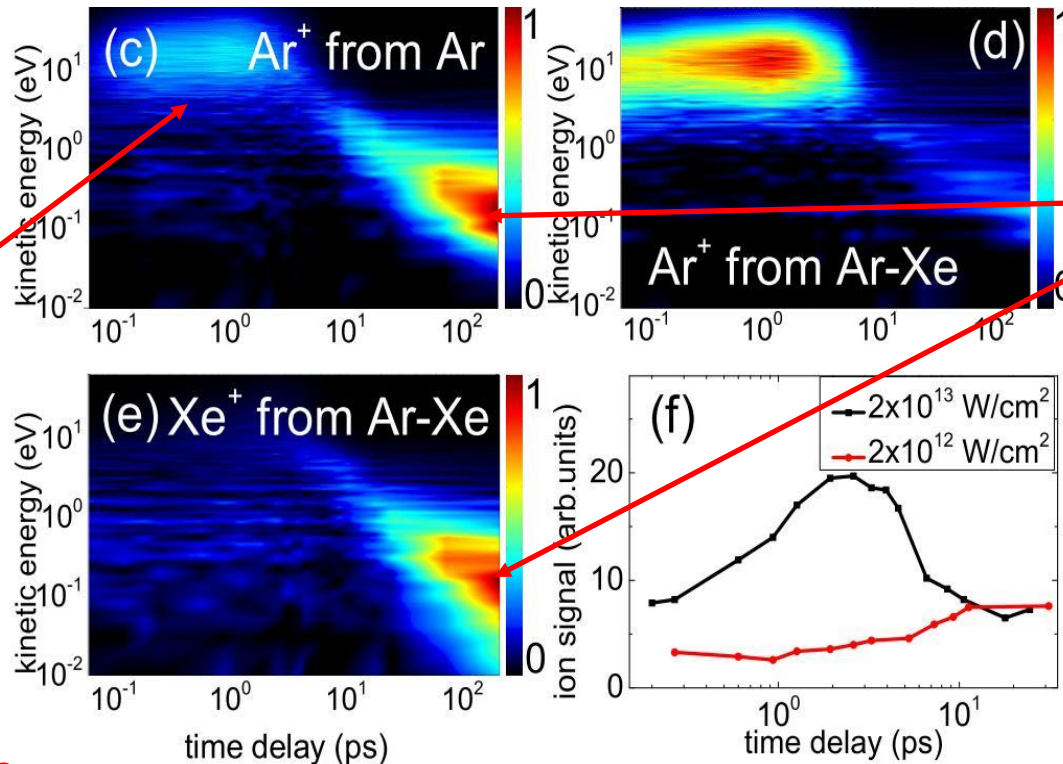
Fragmentation and Recombination dynamics

Frustrated
Recombination:



B. Schütte *et al.*, PRL 112, 073003 (2014)

Reionization of
Excited Atoms from
Recombination
(REAR)



Long time delays:
Recombination

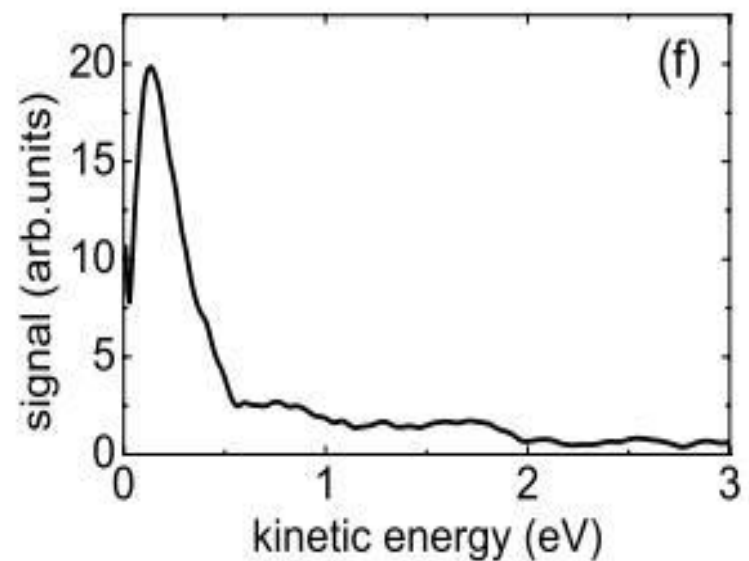
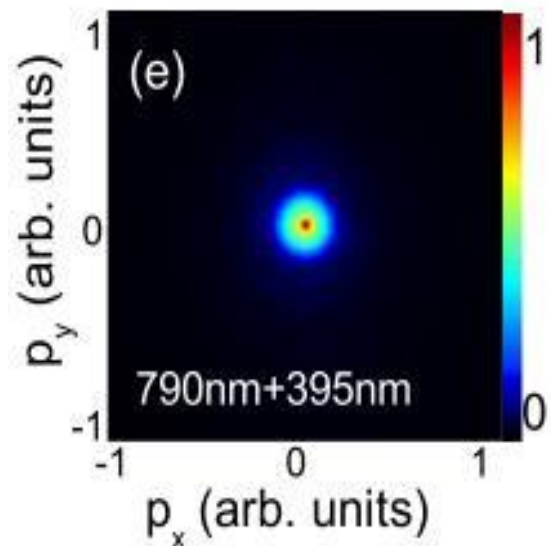
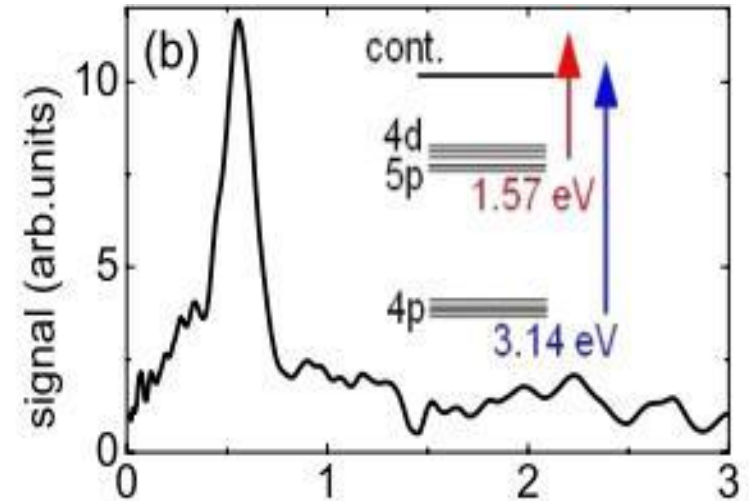
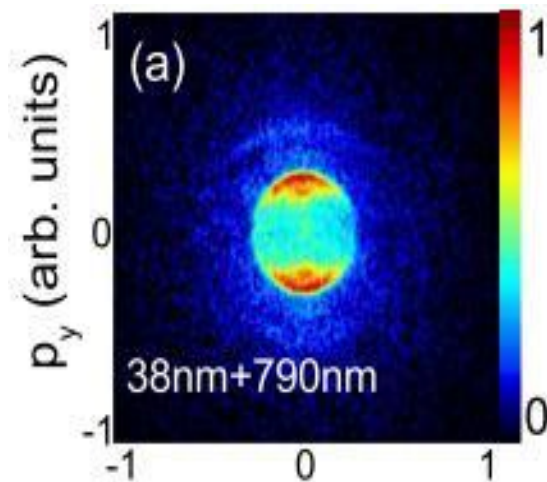
Short time delays:
Resonance effects

NIR induced signals

B. Schütte *et al.*, PRL 112, 253401 (2014)

Fragmentation and Recombination dynamics

Reionization of
Excited Atoms from
Recombination
(REAR)



Summary and outlook

Free electron lasers offer unparalleled performance when very high photon energies (> 1 keV) or very high XUV/X-ray non-linearities are required

High harmonic generation (HHG) can serve as a complementary source under conditions where the required photon energies are somewhat lower (< 1 keV) or under conditions where the dependence on the XUV/X-ray intensity is linear or of modest order

The easy (full-time) access and the ready availability of synchronized secondary sources (IR, UV, THz) is a major advantage

Anticipated technical improvements:

- Broadband continua up to and possibly beyond the O K-edge
- Monochromatized XUV-IR/UV pump-probe spectroscopy with ≤ 10 fs resolution
- Focusability of single harmonics to $\geq 5 \times 10^{13}$ W/cm²
- XUV pump-XUV probe spectroscopy with attosecond time resolution

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